

1 An Updated Parameterization of the Unstable Atmospheric Surface 2 Layer in the WRF Modeling System

3 Prabhakar Namdev¹, Maithili Sharan¹, Piyush Srivastava², Saroj K. Mishra¹

4 ¹Centre for Atmospheric Sciences, Indian Institute of Technology Delhi, New Delhi, 110016, India

5 ²Centre of Excellence in Disaster Mitigation and Management, Indian Institute of Technology Roorkee, Roorkee, 247667,
6 India

7 *Correspondence to:* Prabhakar Namdev (Prabhakarnmdv587@gmail.com)

8 **Abstract.** Accurate parameterization of atmospheric surface layer processes is crucial for weather forecasts using numerical
9 weather prediction models. Here, an attempt has been made to improve the surface layer parameterization in the Weather
10 Research and Forecasting Model (WRFv4.2.2) by implementing similarity functions proposed by Kader and Yaglom (1990)
11 to make it consistent in producing the transfer coefficient for momentum observed over tropical region (Srivastava and Sharan
12 2015). The surface layer module in WRFv4.2.2 is modified in such a way that it contains the commonly used similarity
13 functions for momentum (φ_m) and heat (φ_h) under convective conditions instead of the existing single functional form. The
14 updated module has various alternatives of φ_m and φ_h , which can be controlled by a flag introduced in the input file. The
15 impacts of utilizing different functional forms have been evaluated using the bulk flux algorithm as well as real-case
16 simulations with the WRFv4.2.2 model. The model-simulated variables have been evaluated with observational data from a
17 flux tower at Ranchi (23.412N, 85.440E; India) and the ERA5-Land reanalysis dataset. The transfer coefficient for momentum
18 simulated using the implemented scheme is found to agree well with its observed non-monotonic behaviour in convective
19 conditions (Srivastava and Sharan, 2021). The study suggests that the updated surface layer scheme performs well in simulating
20 the surface transfer coefficients and could be potentially utilized for parameterization of surface fluxes under convective
21 conditions in the WRF model.

22 1 Introduction

23 Inadequate representation of near-surface turbulent processes adds significant uncertainty in both climate projections and
24 seasonal weather forecasts obtained from atmospheric models (Bourassa et al., 2013). Most of the numerical weather prediction
25 and general circulation models utilize Monin-Obukhov similarity theory (MOST; Monin and Obukhov 1954) to parameterize
26 surface turbulent fluxes. To estimate these fluxes and near-surface atmospheric variables, the theory utilizes similarity
27 functions of momentum (φ_m) and heat (φ_h) often prescribed as functions of ζ (stability parameter). However, the exact
28 functional forms for these functions have not been provided by MOST, rather it suggests some asymptotic predictions under
29 near-neutral to very stable and unstable conditions. Over the years, researchers have developed many functional forms for
30 these functions based on the different experiments, conducted over different locations and have separate expressions for stable

31 and unstable stratifications (Webb, 1970; Businger, 1971; Carl et al., 1973; Dyer, 1974; Hicks, 1976; Holtslag and De Bruin,
32 1988; Brutsaert, 1992; Bruin, 1999; Wilson, 2001; Cheng & Brutsaert, 2005; Grachev et al., 2007; Gryanik et al. 2020;
33 Srivastava et al. 2020).

34 In most of the atmospheric models, the commonly used similarity functions under convective conditions are those
35 proposed by Businger (1966) and A. J. Dyer [1965, unpublished work; see Businger (1988)] and referred to as Businger-Dyer
36 (BD) functions. However, these functional forms are unable to follow the classical free convection limit. The study by Rao et
37 al. (1996) suggests that the MOST using Businger relations is unable to define transfer coefficient for momentum (C_D)
38 consistent with its observed behaviour, specifically at low wind convective conditions, indicating that MOST needs to be
39 modified in the (nearly) windless free convection limits. As a result, a revised scaling of heat flux for weakly forced convection
40 in the atmosphere has been proposed by Rao et al. (2006). Later, the issues of using BD functions in the surface layer scheme
41 based on the fifth-generation Pennsylvania State University-National Centre for Atmospheric Research Mesoscale Model
42 (MM5) of a regional scale model (Weather Research and Forecasting; WRF) have been reported in a study by Jimenez et al.
43 (2012). They implemented the new scheme (referred to as revised MM5 scheme; Jimenez et al., 2012) in the WRF modeling
44 system and replaced the BD functions by those proposed by Fairall et al. (1996) (F96) under convective conditions. F96
45 functions are the combination of BD functions, and the functions suggested by Carl et al. (1973) and are valid for the entire
46 range of atmospheric instability. Note that the most recent version of the WRF model still utilizes F96 functions under
47 convective conditions.

48 Srivastava and Sharan (2015) analyzed the observed behaviour of C_D over an Indian land surface and suggested that
49 the observed C_D shows non-monotonic behaviour with $-\zeta$, unlike the behaviour of predicted C_D from MOST based
50 parameterization using commonly used φ_m and φ_h (Businger et al., 1971; Carl et al., 1973; Fairall et al., 1996). Later, a
51 theoretical study by Srivastava and Sharan (2021) revealed that the three-sublayer model based on Kader and Yaglom (1990)
52 is able to predict C_D consistent with its observed non-monotonic behaviour. Note that the three-sublayer model has not yet
53 been newly installed and evaluated in the WRF modeling framework. However, it is already being operational in the surface
54 layer scheme (Community Land Model; CLM) of National Centre for Atmospheric Research Community Atmosphere Model
55 version 5 (NCAR-CAM5) as well as Regional Climate Model (RegCM).

56 The study by Srivastava and Sharan (2021) also analyzed the possible uncertainties associated with the use of different
57 functional forms of φ_m and φ_h under convective conditions. To quantify the impacts of different functional forms, they
58 classified available φ_m and φ_h in four classes based on the exponents appearing in the expressions of φ_m and φ_h as (1)
59 functional forms having the exponents of φ_m and φ_h as $-1/4$ and $-1/2$, respectively (Businger et al. 1971; Hogstrom 1996).
60 (2) functional forms having the exponent of φ_m and φ_h as $-1/3$ (Carl et al. 1973). (3) functional forms having the exponent
61 of φ_m and φ_h as $-1/4$ and $-1/2$, respectively in near-neutral conditions while $-1/3$ in very unstable conditions (Fairall et
62 al. 1996; Grachev et al. 2000; Fairall et al. 2003). (4) functional forms having the exponent of φ_m and φ_h as $-1/4$ and $-1/2$,
63 respectively in near-neutral conditions however, $1/3$ for φ_m and $-1/3$ for φ_h in strong unstable conditions (Kader and
64 Yaglom 1990; Zeng et al. 1998). This study concluded that utilizing different functional forms of similarity functions in the

65 bulk flux algorithm results in a large deviation in the values of estimated fluxes. The detailed description of different functional
 66 forms for φ_m and φ_h considered in different classes are given in Appendix A. We wish to highlight that all available functional
 67 forms for φ_m and φ_h under convective conditions fall in one of the classes stated above.

68 The revised MM5 surface layer scheme of the WRF model version 4.2.2 (WRFv4.2.2) employed φ_m and φ_h based
 69 on Fairall et al. (1996), which belong to class 3. As a result, this scheme is not appropriate in producing C_D consistent with its
 70 observed behaviour, specifically over the Indian land as stated above. Recently Namdev et al. (2023) argue that the
 71 performance of NWP models varies a lot over different seasons and surface types depending upon the functional behaviour of
 72 φ_m and φ_h . Thus, to enhance the potential applicability of the WRF modeling framework, this study attempted to incorporate
 73 all the commonly used similarity functions under convective conditions along with KY90 as well as existing functional forms
 74 in the revised MM5 surface layer scheme of WRFv4.2.2. A namelist flag has been introduced in WRF model to choose between
 75 various φ_m and φ_h in the modified scheme. The modified surface layer scheme proposed in this study has been evaluated
 76 using offline simulations with bulk flux algorithm as well as the real-case simulations with WRFv4.2.2 during the pre-monsoon
 77 season (March-April-May) of 2009 over a domain centered around the location of the flux tower installed at Ranchi (23.412N,
 78 85.440E), India.

79 **2 Methodology and data**

80 **2.1 Surface flux computation in the WRF modeling system**

81 The Monin-Obukhov similarity theory serves as the foundation for the surface layer parameterization (revised MM5 scheme)
 82 in the WRF model, and the surface turbulent fluxes are calculated based on the bulk approach using bulk transfer coefficients
 83 for momentum (C_D) and heat (C_H) (Namdev et al., 2024; Srivastava et al., 2021; Srivastava and Sharan, 2021). Following
 84 MOST they are formulated as follows:

$$85 \quad C_D = k^2 \left[\ln \left(\frac{z + z_0}{z_0} \right) - \left\{ \psi_m \left(\frac{z + z_0}{L} \right) - \psi_m \left(\frac{z_0}{L} \right) \right\} \right]^{-2} \quad (1)$$

$$86 \quad C_H = k^2 \left[\ln \left(\frac{z + z_0}{z_0} \right) - \left\{ \psi_m \left(\frac{z + z_0}{L} \right) - \psi_m \left(\frac{z_0}{L} \right) \right\} \right]^{-1} \left[\ln \left(\frac{z + z_h}{z_h} \right) - \left\{ \psi_h \left(\frac{z + z_h}{L} \right) - \psi_h \left(\frac{z_h}{L} \right) \right\} \right]^{-1} \quad (2)$$

87 in which k is the von Karmann constant; z_0 and z_h are the roughness lengths for momentum and heat, respectively; ψ_m and
 88 ψ_h are the integrated similarity functions for momentum and heat, respectively; and L is the Obukhov length scale.

89 Their determination based on MOST using integrated forms of the similarity functions is explained in Appendix B. In the
 90 following, the default similarity functions used in WRF are explained and other functions are introduced in Section 2.2.

91 The default version of the revised MM5 scheme in the WRF model utilizes similarity functions suggested by Cheng
 92 and Brutsaert (2005) under stable atmospheric conditions ($\zeta > 0$), which are developed using the CASES-99 dataset. The
 93 integrated forms of functions proposed by Cheng and Brutsaert are

$$94 \quad \psi_m(\zeta) = -a \ln(\zeta + [1 + \zeta^b]^{1/b}), \quad \zeta > 0 \quad (3)$$

$$95 \quad \psi_h(\zeta) = -c \ln(\zeta + [1 + \zeta^d]^{1/d}), \quad \zeta > 0 \quad (4)$$

96 where $d = 1.1$, $c = 5.3$, $b = 2.5$ and $d = 6.1$.

97 On the other hand, the similarity functions for unstable atmospheric surface layer ($\zeta < 0$) are those proposed by
98 Fairall et al. (1996; F96). The corresponding integrated functional forms ψ_m and ψ_h are defined as:

$$99 \quad \psi_\alpha(\zeta) = \frac{\psi_{\alpha BD}(\zeta) + \zeta^2 \psi_{\alpha conv}(\zeta)}{1 + \zeta^2}, \quad \alpha = m, h. \quad (5)$$

100 where $\psi_{\alpha BD}$ and $\psi_{\alpha conv}$ denote the integrated functional forms based on Businger and Dyer, and Carl et al. (1973),
101 respectively. The expressions for $\psi_{\alpha BD}$ and $\psi_{\alpha conv}$ are

$$102 \quad \psi_{mBD}(\zeta) = 2 \ln\left(\frac{1+x}{2}\right) + \ln\left(\frac{1+x^2}{2}\right) - 2 \tan^{-1} x + \frac{\pi}{2}, \quad (6)$$

$$103 \quad \psi_{hBD}(\zeta) = 2 \ln\left(\frac{1+x^2}{2}\right), \quad (7)$$

104 in which $x = (1 - 16\zeta)^{1/4}$ and

$$105 \quad \psi_{\alpha conv} = \frac{3}{2} \ln(y^2 + y + 1/3) - \sqrt{3} \tan^{-1}(2y + 1/\sqrt{3}) + \frac{\pi}{\sqrt{3}} \quad (8)$$

106 with $y = [1 - \beta_{m,h}\zeta]^{1/3}$. The values of the constants β_m and β_h are taken as 10 and 34 based on Grachev et al. (2000).

107 2.2 Implementation of different similarity functions

108 In this section, we briefly describe the implementation of different similarity functions for unstable stratification in the surface
109 layer parameterization of WRFv4.2.2. Note that two sets of functional forms, namely those suggested by Carl et al. (1973) and
110 the three sub-layer model proposed by Kader and Yaglom (1990) for convective conditions have not been included and tested
111 in the surface layer scheme of the WRF modeling framework.

112 2.2.1 Functions by Businger et al. (1971) (BD71)

113 Similarity functions suggested by Businger et al. (1971) are based on the KANSAS dataset (Izumi, 1971). These functions do
114 not satisfy the classical free convection limit as predicted by the MOST. They are already implemented in the old version of

115 the MM5 surface layer scheme (Grell et al., 1994) in the WRF model. The integrated functional forms (ψ_m and ψ_h) for φ_m
 116 and φ_h stated in Eqns. (A1) and (A2) (Appendix A) are given in Eqns. (6) and (7).

117 2.2.2 Functions by Carl et al. (1973) (CL73)

118 Carl et al. (1973) proposed an expression of similarity functions φ_m and φ_h valid for the stability range $-10 \leq \zeta \leq 0$. The
 119 expressions for φ_m and φ_h are given in Eqns. (A3) and (A4) (Appendix A). The similarity functions proposed by Carl et al.
 120 (1973) have not been analyzed in the surface layer scheme of the WRF model. The integrated forms (ψ_m and ψ_h) of similarity
 121 functions φ_m and φ_h are given by Eqn. (8).

122 2.2.3 Functions by Kader and Yaglom (1990) (KY90)

123 Kader and Yaglom (1990) introduced a three-sublayer model for convective conditions. The three sublayers are categorized
 124 based on ζ values as (1) the dynamic sublayer which corresponds to near-neutral conditions, (2) the dynamic convective
 125 sublayer which corresponds to moderately unstable conditions and (3) the free convective conditions. The present study
 126 utilized φ_m and φ_h expressions given in Eqns. (A9), and (A10) (Appendix A) that are being used in the surface layer scheme
 127 (CLM4.0; Zeng et al. 1998) of NCAR-CAM5 model. The corresponding integrated forms for φ_m and φ_h are

$$128 \quad \psi_m(\zeta) = \begin{cases} \psi_{m1}(\zeta_m) + \ln \frac{\zeta}{\zeta_m} - 1.14 [(-\zeta)^{1/3} - (-\zeta_m)^{1/3}], & \zeta \leq -1.574 (= \zeta_m) \\ \psi_{m1}(\zeta) = 2 \ln \left(\frac{1+x}{2} \right) + \ln \left(\frac{1+x^2}{2} \right) - 2 \tan^{-1} x + \frac{\pi}{2}, & -1.574 < \zeta < 0 \end{cases} \quad (9)$$

$$129 \quad \psi_h(\zeta) = \begin{cases} \psi_{h1}(\zeta_h) + \ln \frac{\zeta}{\zeta_h} - 0.8 [(-\zeta)^{-1/3} - (-\zeta_h)^{-1/3}], & \zeta \leq -0.465 (= \zeta_h) \\ \psi_{h1}(\zeta) = 2 \ln \left(\frac{1+x^2}{2} \right), & -0.465 < \zeta < 0 \end{cases} \quad (10)$$

130 where $x = (1 - 16\zeta)^{1/4}$.

131 Note that all the functions stated above have been newly installed in the revised MM5 surface layer scheme of
 132 WRFv4.2.2 and can be used in place of F96 functions already employed in the model. Here, we have introduced a new surface
 133 layer module where different options for φ_m and φ_h can be controlled using an appropriate value of namelist parameter
 134 (psimhu_opt). The parameter psimhu_opt is added under the physics section of the namelist file. The variable psimhu_opt can
 135 have values 0, 1, 2, and 3 for different options for functions F96 (default), BD71, CL73, and KY90, respectively. A brief
 136 structure and different choices for psimhu_opt based on newly installed and default functional forms of φ_m and φ_h in the
 137 default and modified revised MM5 scheme are shown in Figure 1.

138 2.3 Characteristics of default and newly installed similarity functions

139 The expressions of φ_m and φ_h for different functional forms utilized in this study are stated in Appendix A. Figure S1
140 (supplementary material) shows the variation of different (a) φ_m and (b) φ_h under moderately to strongly unstable conditions.
141 It is evident from Figure S1 that all the different functional forms provide similar values of φ_m and φ_h in near-neutral to
142 moderately unstable conditions (up to $\zeta = -0.1$ approximately). However, at higher instabilities one can expect noticeable
143 differences between different functional forms of φ_m and φ_h . Note that the functional forms for φ_m corresponding to BD71
144 and CL73 decrease continuously on increasing instability; however, φ_m corresponding to KY90 functional forms show
145 decreasing behaviour in near-neutral to moderately unstable conditions and attain a minimum at $\zeta = -1.574$, and, as the
146 instability further increases, it starts increasing with $-\zeta$ (Figure S1a). This implies that φ_m based on class 4 functions shows
147 non-monotonic behaviour which contradicts the classical MOST prediction. On the other hand, in case of φ_h , all the functional
148 forms provide continuously decreasing behaviour of φ_h from near-neutral to moderately unstable conditions (Figure S1b).

149 Figure 2 illustrates the variation of default (F96) and newly installed integrated similarity functions ψ_m and ψ_h
150 (BD71, CL73, and KY90) with respect to $-\zeta$. One can see from Figure 2a that ψ_m corresponding to F96, BD71, and CL73
151 functional forms increases continuously with $-\zeta$ in moderately to strongly unstable conditions. However, a non-monotonic
152 behaviour has been found for ψ_m corresponding to the KY90 functions implying it first increases with $-\zeta$ and reaches a
153 maximum at $\zeta = -1.574$ and then starts decreasing as instability further grows. On the other hand, ψ_h corresponding to all
154 the considered functional forms increases continuously in near-neutral to strong unstable conditions. However, the rate of
155 increase is slightly higher for F96 in comparison to the other three functions (BD71, CL73, and KY90), whose results are very
156 similar to each other (Fig. 2b).

157 2.4 Observational data for model evaluation

158 For the evaluation of different simulations corresponding to newly installed similarity functions, observational data derived
159 from the micrometeorological tower installed at Ranchi (India) has been utilized (Srivastava and Sharan, 2019; Srivastava et
160 al., 2020; 2021). The dataset (Ranchi data) is derived from an instrument mounted on a 32-m tall tower at the Birla Institute
161 of Technology Mesra in Ranchi, India (Dwivedi et al., 2014) with an average elevation of 609 m above sea level in a tropical
162 region. The site has a few buildings in between east and northwest; agriculture land in between northwest and west; and
163 residential area, and dense trees in between southeast and east. The site also has a relatively flat area in between southeast and
164 west which is free from any obstacle (Srivastava and Sharan, 2015). A fast response sensor (CSAT3 Sonic Anemometer) at a
165 height of 10 m with an average elevation 609 m above sea level provides the temperature and the three components of wind at
166 a 10 Hz frequency. The eddy covariance technique (Stull 1988) is used to estimate heat and momentum fluxes at one-hour
167 time resolution, however the hourly temperature at 2-m is determined by averaging temperature observations available at a
168 temporal scale of 1 minute from the slow response sensors located at logarithmic heights on the same tower. We have utilized

169 hourly data for considered variables. Apart from this we have also utilized the ERA5-Land reanalysis dataset available at
170 $0.10^\circ \times 0.10^\circ$ spatial resolution to evaluate the spatial distribution of the model simulated near surface atmospheric variables.
171 For consistency, we have regridded the model output to the same grid resolution of reanalysis/observed dataset.

172 **3 Numerical simulations**

173 To analyze the impacts of newly installed similarity functions together with the existing functional forms in the surface layer
174 scheme of WRFv4.2.2, the performance of the default and newly installed similarity functions is investigated in two steps. The
175 first one is independent of the WRF model. Namely, we apply Eqn. (B8) (Appendix B) to iteratively determine C_D and C_H as
176 a function of ζ by prescribing the bulk Richardson number (Ri_B) and surface roughness parameters for momentum (z_0) and
177 heat (z_h). The value of ζ is estimated by calculating the root of least magnitude of Eqn. (B8) for a given value of Ri_B . Once ζ
178 is calculated then utilizing it in Eqns. (B9) and (B10), the values of C_D and C_H can be estimated. We call this in the following
179 offline simulation. For the computation, z is taken as 10 m and Ri_B is in the range $-2 \leq Ri_B \leq 0$. The offline simulations are
180 carried out over three different surface types by considering surface roughness (z_0) to be 0.01 m (smooth surface), 0.1 m
181 (transition surface) and 1 m (rough surface) to analyze the impact of roughness of underlying surface on the simulation of ζ ,
182 C_D and C_H .

183 The second step is to apply all the parameterizations of the similarity functions in the WRF model version 4.2.2 over
184 an Indian land site whose output is compared then with the observations during the pre-monsoon (March-April-May; MAM)
185 season of the year 2009. The simulations have been conducted over a nested domain centred around the location of a
186 micrometeorological tower installed at Ranchi (23.412°N , 85.44°E), India (Figure 3). Domain d01 ($6 \times 6 \text{ km}^2$) consists of 233
187 east-west and 210 north-south grid points and domain d02 ($2 \times 2 \text{ km}^2$) consists of 223 east-west and 196 north-south grid
188 points which covers $1398 \times 1260 \text{ km}^2$ and $446 \times 392 \text{ km}^2$ spatial area around the centre point, respectively. Each domain
189 was configured with 50 vertical eta levels from surface to top of the atmosphere. We kept five vertical levels below 100 m
190 height. Initial and boundary conditions were taken from ERA5 global atmospheric reanalysis dataset at a resolution of
191 $0.25^\circ \times 0.25^\circ$ and boundary conditions were forced every 6 hours. For land use and land cover (LU/LC) information, we
192 have used dataset from MODIS (Moderate Resolution Imaging Spectroradiometer; Friedl et al., 2002). Various physical
193 parameterizations utilized in the simulations are listed in Appendix C. In this study, four sets of simulations were carried out,
194 as given in Table 1.

195 Note that the revised MM5 surface layer scheme has lower limits on the values of \mathbf{u}_* ($> 0.001 \text{ m s}^{-1}$) and \mathbf{U} (> 0.1
196 m s^{-1}) that allow nocturnal values of \mathbf{u}_* at night and control Ri_B values to be inordinately high, respectively (Jimenez et al.,
197 2012). However, the stability parameter ζ or Ri_B is not restricted in the revised MM5 surface layer scheme, which gives
198 complete freedom to the WRF model to show its sensitivity to the tested similarity functions (Jimenez et al., 2012). Moreover,
199 some of the LES studies reported in the literature suggest that the friction velocity cannot be zero when the mean wind drops

200 to zero, i.e., there should be a minimum friction velocity that is proportional to the \mathbf{w}_* (Schumann, 1980). For this purpose,
201 the existing version of the revised MM5 scheme sets 0.001 m s^{-1} as the minimum value of \mathbf{u}_* based on the recommendations
202 by Jimenez et al. (2012). Thus, to avoid the complexity that arises when mean wind drops to zero, the updated revised MM5
203 scheme proposed in the present study also utilizes a minimum value of \mathbf{u}_* ($> 0.001 \text{ m s}^{-1}$) as suggested by Jimenez et al.
204 (2012) in the existing version of the revised MM5 scheme. Moreover, the scheme uses constant values of \mathbf{z}_0 , while the values
205 of \mathbf{z}_h are calculated from the expression suggested by Brutsaert (1982).

206 The whole simulation period is divided into segments of 4 days with 24 h overlapping time between different
207 segments to ensure continuity. The model is initialized at 0000 UTC of the first day of each simulation and runs for 96 hours.
208 In order to avoid the potential spin-up problems at the beginning of the simulation, we discard the first day of each simulation
209 as spin up time and consider the last three days for the analysis (Jimenez et al., 2010; 2012).

210 For the evaluation of the real-case simulations, different statistical parameters such as mean absolute error (MAE),
211 root mean square error (RMSE), mean bias (MB), index of agreement (IOA), different measures of correlation coefficient
212 (CC), mean bias (%) (bias), and standard deviation of the model predicted output normalized by that of the observations are
213 used. A brief description of the performance indicators for validation utilized in the present study is given in Appendix C.

214 **4 Results**

215 **4.1 Offline simulations**

216 To analyze the functional dependence of ζ , C_D and C_H on the utilized forms of similarity functions, the offline simulations
217 independent of the WRF model have been conducted utilizing newly installed functions (BD71, CL73, and KY90) together
218 with F96 functions existing in the default version of the surface layer scheme of the WRF model for three different roughness
219 lengths for momentum (z_0), which are representative of smooth ($z_0 = 0.01 \text{ m}$), transition ($z_0 = 0.1 \text{ m}$), and rough
220 ($z_0 = 1.0 \text{ m}$) surfaces. The results for ζ (Figure 4a, b, and c) with R_{iB} , C_D (Figure 4d, e, and f) and C_H (Figure 4g, h, and i)
221 with ζ across various surface types and sublayers have been analyzed. The different sublayers associated with convective
222 stratification include dynamic (DNS), dynamic-dynamic convective transition (DNS-DCS), dynamic convective (DCS),
223 dynamic convective-free convective transition (DCS-FCS), and free convective (FCS) (Srivastava and Sharan, 2021). Note
224 that the sublayers DNS ($-0.04 \leq \zeta \leq 0$) and DNS-DCS transition ($-0.12 \leq \zeta < -0.04$) are corresponding to weakly to
225 moderately unstable conditions, while sublayers DCS ($-1.20 \leq \zeta < -0.12$), DCS-FCS ($-2.0 \leq \zeta < -1.20$), and FCS ($\zeta <$
226 -2.0) belong to moderately to strongly convective conditions (Srivastava and Sharan, 2015).

227 It is found that the simulated values of ζ at smaller values of R_{iB} (i.e., in DNS to DCS) from different forms of
228 similarity functions are found to be almost identical to the F96 functional forms (Figure 4a-c). Moreover, results from the
229 BD71, CL73, and F96 functions are even similar at higher instabilities (i.e., the whole range of ζ values), while they differ

230 strongly from values obtained using the KY90 functions (Figure 4a-c). Notably, BD71, CL73, and F96 functional forms predict
231 relatively smaller absolute values of ζ for a given value of Ri_B . However, KY90 functions are found to produce a relatively
232 larger magnitude of ζ for a given value of Ri_B . This behaviour is found to be consistent for all ratios z/z_0 (Figures 4a-c)
233 representative of smooth, transition, and rough surfaces. A relatively larger magnitude of ζ for a given value of Ri_B and the
234 smaller values of ψ_m and ψ_h (Figure 2) in KY90 functional forms implies that the momentum and heat fluxes predicted using
235 KY90 functions will be smaller than those anticipated in BD71, CL73, and F96 functional forms.

236 Figure 4d-f shows the variation of C_D with ζ estimated using BD71, CL73, KY90, and F96 functional forms over
237 different surfaces. Notice that the C_D values calculated from BD71, CL73, and F96 forms of functions are relatively higher
238 than those produced by KY90 functional forms and continue to rise as instability progresses from DCS to FCS. It is important
239 to highlight that C_D estimated using KY90 functions shows a non-monotonic behaviour, which is consistent with the observed
240 behaviour of C_D over the Indian region reported in the literature (Srivastava and Sharan, 2019; 2021). Note that this non-
241 monotonic behaviour is consistent for all three cases of different roughness lengths (Figure 4d-f).

242 On the other hand, across all three surfaces, one can see that the values of C_H estimated from all four functional forms
243 increase with increasing instability (Figure 4g-i), while the rate of increase of C_H in KY90 functions is relatively slower.
244 Moreover, BD71, CL73, and F96 functions predict almost similar values over all three types of surfaces. Noticeably, C_H
245 estimated using KY90 functions also exhibits non-monotonic behaviour with ζ over rough surfaces, which contradicts the
246 predictions of the other three functional forms. In addition, it is important to note that C_D and C_H predicted by KY90 functional
247 forms are found to bound by twice their near-neutral values, while the other functional forms predict continuously increasing
248 values of C_D and C_H on increasing instability.

249 Note that the error caused by different values of z_0 can be so large that the stability dependence of using different
250 forms of similarity functions is less important in the computation of C_D and C_H . As a result, three different values of z_0 have
251 been chosen, similar to a recent study by Srivastava and Sharan (2021), which are representative of smooth ($z_0 = 0.01$ m),
252 transition ($z_0 = 0.1$ m), and rough ($z_0 = 1.0$ m) surfaces to account for the impacts of using different z_0 on the estimation of
253 C_D and C_H from different functional forms of similarity functions in offline simulations.

254 Moreover, Figure 4 depicts the offline simulations with equal values of z_0 and z_h . While in the revised MM5 surface
255 layer scheme available in the WRF model, the values of z_0 and z_h are not the same. Thus, we have also attempted to discuss
256 the results from the offline simulations with different values of z_h , assuming $z_0 = 0.1$ m. Figure S2 (supplementary material)
257 shows the variation of ζ with Ri_B , C_D , and C_H with ζ calculated from the bulk flux algorithm using similarity functions
258 corresponding to BD71, CL73, KY90, and F96 with different values of z_h while z_0 is fixed. The values of z_h are taken such
259 that the ratio $\ln(z_0/z_h)$ assumes 0.1, 1, 2, 3, and 4. Figure S2 clearly shows that the estimated values of ζ are similar in near-
260 neutral to moderately unstable conditions for all values of z_h ; however, relatively smaller values have been found as the ratio
261 $\ln(z_0/z_h)$ increases for each form of similarity function. Since the computation of C_D does not involve the values of z_h (Eqn.
262 B9), the estimated values of C_D for each form of similarity function are found to be approximately the same for different values

263 of z_h . However, in the case of C_H , differences are clearly visible if one uses different values of z_h . The estimated C_H using
264 various similarity functions behaves similarly for different values of z_h , while the magnitude decreases as the ratio $\ln(z_0/z_h)$
265 increases.

266 Hence, it is evident that the BD71, CL73, and F96 functional forms predict values of ζ , C_D , and C_H that are almost
267 same over all the three different surface types. However, using KY90 functions compared to other commonly used φ_m and
268 φ_h , one can expect a significant reduction in the estimated values of transfer coefficients in moderately to strongly unstable
269 stratification.

270 4.2 Results of the WRF model using different sets of integrated similarity functions

271 In this section, observational and reanalysis datasets have been used to analyze the simulations performed with WRFv4.2.2
272 utilizing newly installed and default φ_m and φ_h . The model simulated output has been extracted at the location of the flux
273 tower and compared against the observations derived from the flux tower installed at Ranchi (23.412N, 85.440E), India. The
274 mean spatial patterns of certain variables averaged over daytime (04:00-12:00 UTC) have been compared against the ERA5-
275 Land reanalysis dataset. Further, to access the effects of newly installed functions under free convective conditions, the mean
276 spatial patterns of considered variables averaged across strong convective conditions (hours in which $\zeta < -10$ over most of
277 the domain) have been analyzed against respective hours of ERA5-Land reanalysis data. Bilinear interpolation has been used
278 to interpolate the model output to the same grid resolution as the ERA5-Land data in order to allow a consistent comparison.

279 4.2.1 Evaluation against observations derived from flux tower installed at Ranchi (India)

280 Figure 5 depicts the variation of (a) ζ with Ri_B , (b) C_D , and (c) C_H at the first model level with ζ from different experiments
281 (Exp1, Exp2, and Exp3) and CTRL simulation. Although the absolute values of the parameters differ from each other due to
282 the different prescribed roughnesses, the variation of ζ with Ri_B , C_D and C_H with ζ is very similar to the offline results. Note
283 that, at the moment, due to the inaccessibility of long-term data on detailed surface properties such as vegetation structure
284 needed to quantify the roughness length, we do not have an access to the precise values of z_0 and z_h at the Ranchi station.
285 Moreover, the values of z_0 and z_h do not directly involve in the estimation of C_D , C_H , and the surface fluxes from the
286 observational data, while they are important in computing these variables using the MOST framework. Thus, the default value
287 of z_0 is used in the revised MM5 surface layer scheme available in the WRF model, which is found to be approximately in the
288 range 0.1 – 0.2 m at the Ranchi station. We wish to highlight that the z_0 used in the WRF model simulations at the Ranchi
289 station is nearly similar to the case of $z_0 = 0.1$ m presented in Figure 4, and the offline simulations also indicate that the
290 behaviour of the estimated C_D and C_H with ζ remains almost the same for different values of z_0 with slightly varying
291 magnitudes. Thus, one can interpret the results of C_D and C_H shown in Figures 4 and 5 from the offline simulations and the
292 WRF model, respectively, and can compare the WRF model simulated C_D with the observed one at the Ranchi station.

293 Although the model simulations and observed data may have a different z_0 , the comparison of model simulated variables with
294 the Ranchi data allows for an impression of the structural behaviour of model results as a function of stratification compared
295 with measurements.

296 It is clear from Figure 5 that the values of simulated variables are found to be almost identical in DNS to DCS
297 sublayers for all the experiments. Moreover, in FCS, the results obtained from Exp1, 2 and CTRL simulation are found to be
298 nearly similar however, relatively strong differences have been found in results from Exp 3 (Figure 5a, b, and c). Simulated ζ
299 for a given Ri_B in Exp2 and CTRL simulation are similar and found to be relatively smaller in magnitude than Exp1 and Exp3
300 in FCS. However, the absolute values of ζ in Exp3 (KY90 functions) are relatively higher in FCS than in all other experiments.

301 Figure 5b shows the variation of simulated C_D with ζ from different experiments. Purple circles denote the variation
302 of observed C_D with ζ at the location of flux tower (Figure 5b). It is found that the observed C_D increases as the instability
303 increases from DNS to DCS and has the maximum value in the DCS (at $\zeta = -0.1$ approx.) and then starts to decrease as
304 instability grows further from DCS to FCS. It is evident that C_D simulated using φ_m and φ_h based on class 4 functions (Exp3)
305 exhibits non-monotonic behaviour (Figure 5b), which is consistent with the observed behaviour of C_D (Srivastava and Sharan,
306 2015; 2021). The magnitude of C_D predicted in Exp3 is significantly smaller than that simulated from other experiments as
307 well as CTRL simulation, specifically in FCS. This may be due to the large differences between the KY90 functional forms of
308 ψ_m and ψ_h and other forms of functions. On the other hand, C_D simulated using φ_m and φ_h based on the first three classes
309 (Exp1, Exp2, and CTRL simulation) increases continuously as instability grows from DNS to FCS (Figure 5b).

310 However, it is found that the C_D predicted from the original forms of class 4 functions (Exp3) show large disagreement
311 with its observed behaviour, as the predicted C_D starts decreasing at ζ lying in FCS, which is different from that observed, i.e.,
312 ζ lying in DCS. As a result, the study also highlighted the necessity of fine-tuning the original KY90 functional forms and
313 evaluating their performance in the WRF model with additional observational datasets from various land sites and seasons.

314 Note that Srivastava and Sharan (2021) tuned the original forms of class 4 functions by enforcing the matching of the
315 point at which both observed and model predicted C_D attain their maximum value. However, more studies in terms of predicting
316 the observed variation of the non-dimensional vertical gradients of mean wind speed and temperature with ζ are essential to
317 further tune the original KY90 functions for the Indian region using observed data from various locations under different
318 seasons.

319 Further, we would like to point out that currently no observational datasets are available which show a better
320 agreement with the KY90 functions over Indian land. However, it is desirable to further validate these functional forms over
321 Indian land once such observational datasets become available.

322 We wish to highlight that utilizing KY90 (Exp3) functions in the revised MM5 scheme of the WRF model makes it
323 consistent in predicting C_D with its observed non-monotonic behaviour over the Indian region.

324 The variation of simulated C_H with ζ from different experiments is shown in Figure 5c. C_H simulated from Exp1-3 as
325 well as CTRL simulation shows continuously increasing behaviour with ζ . The magnitude of simulated C_H from CTRL

326 simulation and Exp1-2 is relatively higher than that of Exp3 in FCS beyond $\zeta < -10$ (approximately). It is also evident that
327 at higher instabilities, even C_H shows non-monotonic behaviour with ζ (Figure 5c). We wish to point out that a relatively larger
328 scatter has been found in the values of C_H than C_D . As the WRF model utilizes constant values for z_0 , while z_h is calculated
329 using expression suggested by Brutsaert (1982). The relatively large scatter in the values of C_H simulated from the WRF model
330 can be due to the parameterization of the ratio of momentum and scalar roughness lengths in the model.

331 Note that the transfer coefficients C_D and C_H shown in Figure 5 are at the reference height corresponding to the lowest
332 model grid level, which is ~ 12 m in the present study. However, we have also analyzed the behaviour of C_D and C_H at 10 m
333 height with ζ and found that they behave similarly to those presented in Figure 5.

334 The analysis presented here indicates that the KY90 functions in the revised MM5 surface layer scheme are found to
335 be appropriate in producing non-monotonic behaviour of C_D consistent with its observed nature. However, all other functional
336 forms of φ_m and φ_h produce C_D , which increases continuously with ζ from DNS to FCS.

337 To quantify the uncertainties involved in the simulated surface fluxes and certain near-surface variables using KY90
338 (Exp3) as well as other functional forms (Exp1-2 and CTRL simulation), model simulations have been compared against the
339 observations. Figure 6 compares the model-simulated (a) u_*^2 ($\text{m}^2 \text{s}^{-2}$) (representative of momentum flux), (b) SHF (W m^{-2})
340 (sensible heat flux), (c) U_{10} (m s^{-1}) (10-m wind speed), and (d) T_{2m} (K) (2-m temperature) with the observed data obtained
341 from the flux tower at Ranchi (23.412N, 85.440E), India. The model output was extracted at a single grid point closest to the
342 flux tower to allow a consistent comparison. In Figure 7, a Taylor diagram is displayed along with the normalized standard
343 deviations and correlations of considered variables. Figure 8 shows the scatter plot between CC vs. RMSE for considered
344 variables simulated using different experiments. In case of u_*^2 , Exp1 and Exp2 are found to be comparable to the CTRL
345 simulation, while Exp3 considerably improved the simulation of u_*^2 (Figures 6a, 7 and 8). Exp3 reduced MAE (RMSE) from
346 0.09 (0.16) $\text{m}^2 \text{s}^{-2}$ to 0.08 (0.14) $\text{m}^2 \text{s}^{-2}$ (Table 2; Figures 7 and 8) and improved the CC (0.74) and IOA (0.84) for u_*^2 (Table 2).
347 A Q-Q plot is shown in Figure S3a (supplementary material) suggesting that Exp3 (KY90 functions) is found to be slightly
348 better than all other experiments and CTRL simulation for u_*^2 . For SHF, all the experiments are comparable to the CTRL
349 simulation; however, Exp3 shows less scatter than other experiments (Figure 6a).

350 In case of U_{10} , Exp3 shows less scatter and appears to be closer to the observations than other experiments (Figure
351 6c). Exp3 noticeably improved the simulation of U_{10} by reducing MAE (RMSE) from 1.20 (1.54) m s^{-1} to 1.16 (1.47) m s^{-1} and
352 MB up to 5 % (Figures 6c, and 7; Table 2). It considerably improved the CC (IOA) for U_{10} from 0.66 (0.73) to 0.68 (0.75)
353 (Figure 7 and Table 2). A Q-Q plot (Figure S3b: supplementary material) reveals that Exp3 is found to be better than all other
354 experiments and CTRL simulation for U_{10} . Thus, the KY90 functions in the surface layer scheme of the WRF model
355 considerably improve the model in simulating U_{10} (Figures. 6c, 7, 8, and S2b) at the location of the flux tower. Further, in case
356 of T_{2m} , Figures 7 and 8 exhibit that all the experiments are found to be comparable with the CTRL simulation.

357 Note that earlier studies, especially the ones done in the GABLS model intercomparison projects, have studied the
358 impacts of the similarity functions on the modelled profiles and fluxes (though mostly for stable conditions). However, they
359 learnt that applying different stability functions in the surface and boundary layer parameterizations may trigger unnatural

360 kinks in the model simulated wind speed and temperature profiles. Here, we have analyzed the profiles of U_{10} and T_{2m} simulated
361 from WRF model using different similarity functions in the surface layer scheme for the occurrence of unnatural kinks in their
362 values. One can see that the U_{10} predicted from CTRL simulation, as well as different experiments corresponding to different
363 similarity functions at certain hours goes higher than that of its observed maximum value (approx. 8 m s^{-1}) (Figure S4:
364 supplementary material). These relatively higher magnitudes may be linked with some localised weather phenomenon
365 characterized by rapid changes in weather including strong wind, lightning and thunderstorms and are justifiable. However,
366 the simulated T_{2m} from different similarity functions are found to be in line with the observed values across the whole
367 simulation period (Figure S5: supplementary material). This suggests that the values of U_{10} and T_{2m} predicted from WRF
368 model are found to be in justifiable range and no unnatural kinks have been found.

369 4.2.2 Evaluation of mean spatial distribution of simulated variables against ERA5-Land reanalysis data during daytime

370 In this section, mean spatial distribution of simulated variables from different experiments as well as CTRL simulation
371 averaged during daytime (04:00-12:00 UTC) for entire simulation period, is compared with the ERA5-Land reanalysis data.
372 Figure 9 depicts the mean spatial patterns of simulated ζ ($= \frac{z}{L}$) ($a1 - 4$), C_D ($c1-c4$), and C_H ($e1-4$) from CTRL simulation
373 and other experiments, as well as their differences with respect to CTRL simulation. It is found that the absolute value of ζ
374 simulated in Exp3 (KY90 functions) is relatively smaller than CTRL simulation (Figure 9b3) across the whole domain, which
375 is consistent with Figure 5a, and offline simulations presented in Figure 4(a-c). This could be because the magnitude of KY90
376 functions (φ_m and φ_h) (Figure S1: supplementary material) is relatively smaller than the functions employed in default scheme
377 (CTRL simulation).

378 On the other hand, Exp1 also provides slightly smaller absolute values of ζ (Figure 9b1), while Exp2 is almost
379 comparable to the CTRL simulation (Figure 9b2). Model simulated C_D is found to be relatively smaller in Exp3 than CTRL
380 simulation (Figure 9d3), while Exp1 and Exp2 provide comparable values of C_D to CTRL simulation (Figure 9d1-2). In the
381 case of C_H , the simulated values from different experiments are comparable to the CTRL simulation over the whole study
382 domain (Figure 9f1-3). Note that simulated C_H is found to be comparable in all the experiments while one can see slight
383 differences in C_D in Exp3 than all other experiments which may be related to the fact that only φ_m functions are involved in
384 the computation of C_D (Eqn. 1), and the differences between φ_m corresponding to Exp3 are relatively more than φ_h , so are
385 the differences in C_D . The hatched regions in Figure 9 shows the differences between simulated variables from different
386 experiments with respect to CTRL simulation are statistically significant at 95% confidence level.

387 The slight differences in C_D in Exp3 reflected further in the simulated $u_*^2 \text{ m}^2 \text{ s}^{-2}$ (a measure of momentum flux) (Figure
388 10b3). A slight reduction has been found in simulated u_*^2 in Exp3 compared to the CTRL simulation over some parts of the
389 domain (Figure 10b3), while in Exp1 and Exp2 values are comparable with the CTRL simulation (Figure 10b1-2). In case of
390 SHF and LHF, the mean spatial distribution from all the experiments is found to be consistent with the ERA5-Land reanalysis

391 data, and the magnitude of differences between model simulation and ERA5-Land data is comparable for all the experiments
392 (Table S1: supplementary material).

393 For T_{2m} (upper panel of Figure 11), T_s (middle panel of Figure 11), and U_{10} (lower panel of Figure 11), mean spatial
394 distribution from different experiments and CTRL simulation agreed well with slightly varying magnitude to the ERA5-Land
395 reanalysis data. One can see a warm bias up to 2 K (3 K) for T_{2m} (T_s) simulated from different experiments and CTRL
396 simulation over most of the domain. For T_{2m} , bias, RMSE, and PCC between different experiments together with CTRL
397 simulation and ERA5-Land reanalysis data are found to be comparable (Table S1: supplementary material). However, Exp3
398 slightly improved the PCC from 0.50 to 0.51 for T_s (Table S1: supplementary material). Further, in the case of U_{10} , all the
399 simulations exhibit overprediction over the whole domain (lower panel of Figure 11: b1-4) and Exp3 is found to be slightly
400 better than all other experiments as well as CTRL simulation as it reduced bias% (RMSE) from 32.28 (0.54) $m\ s^{-2}$ to 32.06
401 (0.53) $m\ s^{-2}$ and improved the PCC from 0.89 to 0.91 (Table S1: supplementary material).

402 **4.2.3 Evaluation of newly installed functions during strong unstable conditions with respect to ERA5-Land reanalysis** 403 **data**

404 This section describes the impacts of utilizing different similarity functions (φ_m and φ_h) on simulated variables during highly
405 convective regime (i.e., $\zeta < -10$) with respect to the ERA5-Land reanalysis dataset. Since the functional forms of ψ_m and ψ_h
406 are almost identical in near-neutral to moderately unstable conditions, however, in strong unstable conditions, the differences
407 between different functional forms are more pronounced. Thus, the corresponding differences in the simulated values of
408 considered variables are expected to be more pronounced during highly convective regimes. For this purpose, the model output
409 has been extracted for those hours in daytime which show ζ smaller than -10 over most of the domain and compared with the
410 respective hours of ERA5-Land reanalysis data.

411 Figure S6 (Supplementary material) depicts the mean spatial distribution of ζ (a1-4), C_D (c1-4), and C_H (e1-4) as well
412 as their deviations from CTRL simulation. Notice that the magnitude of differences for all variables (ζ , C_D , and C_H) in this
413 case are found to be larger than the case of mean spatial patterns averaged during the whole daytime (section 6.2.2). It is
414 evident from Figure S6b3 (supplementary material) that Exp3 produce large absolute values of ζ and smaller values of C_D and
415 C_H (Figures S6b3, d3 and f3: supplementary material) than all other experiments and the CTRL simulation. While Exp1 and
416 Exp2 are found to be comparable to the CTRL simulation for both C_D and C_H (Figures S6d1-2 and f1-2).

417 The model simulations for T_{2m} and T_s do not capture the spatial patterns well in comparison to ERA5-Land data
418 (Figures S7a1-5 and S8a1-5: supplementary material). All experiments, as well as the CTRL simulation, exhibit overprediction
419 across the whole domain (Figures S7b1-4 and S8b1-4). We wish to highlight that the differences between various experiments
420 and CTRL simulation are seen up to 0.5 K for T_{2m} (Figure S7c1-3: supplementary material) as well as T_s (Figure S8c1-3)
421 which is slightly higher than the case of mean spatial patterns averaged over whole daytime (upper and middle panels of Figure
422 11). For T_{2m} , it is evident from Figure S7 (supplementary material) and Table 3 that Exp3 noticeably reduced the bias%

423 (RMSE) from 0.64 (2.13) K to 0.62 (2.10) K and improved the PCC from 0.43 to 0.46 (approximately 6%). In case of T_s as
424 well, Exp3 slightly improved the PCC and reduced the bias% (RMSE) from 1.25 (4.01) K to 1.24 (3.97) K (Table 3 and Figure
425 12).

426 For U_{10} , the mean spatial patterns simulated using different experiments agreed well with the ERA5-Land reanalysis
427 data (Figure S9a1-5: supplementary material) and the magnitude of biases is found to be up to 1 m s^{-1} . Exp3 outperformed all
428 other experiments and the CTRL simulation by lowering the bias% from -4.96 to -0.28 m s^{-1} and improved the PCC from 0.34
429 to 0.36 with comparable RMSE values (Figures S9 and 12; Table 3).

430 The results presented so far suggest that the changes corresponding to different functional forms of similarity
431 functions in the surface layer parameterization of the WRF model are more pronounced in convective conditions during
432 daytime hours. For the number of grid points over the study domain that are being affected by the changed similarity functions,
433 no fixed pattern was found; however, the changes depend on the considered variable and similarity functions. Furthermore,
434 we observe that the changes are more pronounced in grids that experience strong instability during the daytime.

435 **5 Summary and concluding remarks**

436 In the present study, the revised MM5 surface layer scheme of the WRFv4.2.2 model has been modified to incorporate φ_m
437 and φ_h suggested by Kader and Yaglom (1990) to make it consistent in producing the transfer coefficient for momentum (C_D)
438 in line with its observed behaviour. The revised MM5 scheme is modified in such a way that it contains all commonly used
439 φ_m and φ_h under convective conditions instead of a single functional form. Various alternatives of φ_m and φ_h in the modified
440 scheme can be controlled by a flag (psimhu_opt) that has been introduced in the physics section of the namelist file. The
441 impacts of utilizing different functional forms of φ_m and φ_h in the proposed scheme have been evaluated using offline
442 simulations (with bulk flux algorithm) as well as real-case simulations with WRFv4.2.2 model. The model-simulated surface
443 turbulent fluxes and certain near-surface variables have been compared with observational data from a flux tower at Ranchi
444 (23.412N, 85.440E; India), and the spatial patterns have been evaluated with the ERA5-Land reanalysis dataset.

445 Offline simulations indicate that at nearly neutral to moderately unstable conditions, ζ simulated using various
446 functional forms of φ_m and φ_h is comparable, and as the instability grows (free convective conditions), the differences
447 between different experiments become more pronounced. This might be connected to the corresponding variations between
448 different functional forms of similarity functions in the respective regimes. Similarly, for simulated C_D , Exp3 (KY90 functions)
449 demonstrates nonmonotonic behaviour with $-\zeta$ across all three surface types (representing smooth, transition, and rough
450 surfaces), which is consistent with its observed behaviour. However, all other experiments and CTRL simulation indicate
451 continuously increasing C_D with $-\zeta$ from near-neutral to free convective conditions over all three surface types, which is
452 inconsistent with its observed behaviour over the study domain. The non-monotonic behaviour of C_D in Exp3 (KY90 functions)
453 may be associated to the analogous non-monotonic behaviour of the corresponding ψ_m in the respective regime.

454 In real-case simulations, the model simulated ζ , C_D and C_H are found to be consistent with the offline simulations.
 455 One can see that the variation of C_D in Exp3 (KY90 functions) with $-\zeta$ is nonmonotonic, as reported in offline simulations
 456 and found to be consistent with its observed behaviour. This indicates that the KY90 functions in the surface layer scheme of
 457 the WRF model make it compatible in producing C_D consistent with its observed behaviour over Indian land. As compared
 458 with the observations over Ranchi (India), the simulations using KY90 (Exp3) functions are found to perform better for most
 459 of the considered variables compared to all other experiments. Further, in the mean spatial distribution averaged during daytime
 460 (04:00–12:00 UTC) over the entire simulation period, the significant increase in absolute value of ζ from Exp3 resulted in a
 461 noticeable reduction in the values of C_D and C_H , which further impacted the simulated values of T_s , T_{2m} , and U_{10} . When
 462 compared with the ERA5-Land reanalysis data, the spatial patterns for T_{2m} , T_s , and U_{10} from Exp3 (KY90 functions) provided
 463 more consistent results. A reduction has been found in bias (%) and RMSE values for T_s , and U_{10} . Moreover, in case of highly
 464 convective regime ($\zeta < -10$), Exp3 (KY90 functions) slightly improved the performance of the model by reducing the bias
 465 (%) and RMSE for T_{2m} , T_s , and U_{10} and increasing the correlation to some extent.

466 Thus, it is concluded that the similarity functions proposed by Kader and Yaglom (1990) (KY90 functions; Exp3) are
 467 found to be more appropriate for use in the WRF model as they can simulate C_D consistent with its observed behaviour and
 468 improve the simulation for most of the considered variables over the study domain. However, due to the limited spatial
 469 coverage of the domain considered in this study and the limited availability of observational data, KY90 functional forms need
 470 to be further evaluated in the WRF modeling framework utilizing observations from other sites. The modified surface layer
 471 scheme proposed in this study could enhance the potential applicability of the WRF modeling framework for the community
 472 in investigating the role of different functional forms of similarity functions under convective conditions for selected
 473 events/case studies such as extreme weather events, heat waves during summer, cyclonic storms, and fog predictions using the
 474 WRF model.

475 **Appendix A**

476 Here, the detailed description of commonly used functions (φ_m and φ_h) in numerical models under convective conditons is
 477 provided.

478 Based on Businger (1966) and A. J. Dyer [1965, unpublished work; see Businger (1988) for details] the expressions
 479 for φ_m and φ_h are as follows:

$$480 \varphi_m = (1 - \gamma_m \zeta)^{-\frac{1}{4}} \quad (A1)$$

$$481 \varphi_h = Pr_t (1 - \gamma_h \zeta)^{-\frac{1}{2}} \quad (A2)$$

482 in which $\gamma_m = 15$, $\gamma_h = 9$, and $Pr_t = 0.74$ is the turbulent Prandtl number. Note that in case of Dyer (1974) the values of
 483 $\gamma_m = \gamma_h = 16$ and $Pr_t = 1.0$. These functions commonly known as Businger-Dyer similarity (BD) functions and do not
 484 satisfy the classical free convection limit (Srivastava et al. 2021).

485 The similarity functions proposed by Carl et al. (1973) under convective conditions are applicable for the range
 486 $-10 \leq \zeta \leq 0$. The expressions for φ_m and φ_h suggested by Carl et al. (1973) are:

$$487 \quad \varphi_m = (1 - \beta_m \zeta)^{-\frac{1}{3}} \quad (\text{A3})$$

$$488 \quad \varphi_h = (1 - \beta_h \zeta)^{-\frac{1}{3}} \quad (\text{A4})$$

489 in which $\beta_m = \beta_h = 15$. However, based on various studies reported in the literature β_m and β_h can take different values. For
 490 example, Delage and Girard (1992) proposed $\beta_m = \beta_h = 40$, on the other hand, Fairall et al. (1996) suggested that $\beta_m = \beta_h =$
 491 12.87.

492 Fairall et al. (1996, 2003) proposed an interpolation function applicable for the entire range of atmospheric instability,
 493 which was based on BD functions and functions suggested by Carl et al. (1973). This interpolation function does not have the
 494 gradient form (φ_m and φ_h), as they have interpolated the integrated forms of the functions. We wish to highlight that the
 495 revised MM5 surface layer scheme of Weather Research and Forecasting Model version 4.2.2 utilized the interpolation
 496 functions suggested by Fairall et al. (1996).

497 Kader and Yaglom (1990) proposed a three-sublayer model under convective conditions. The dynamic sublayer
 498 corresponds to near-neutral conditions in which $\varphi_m = 1$ and $\varphi_h = \text{Pr}_t$. Further, in the dynamic convective sublayer,
 499 mechanical energy is in the x direction, while buoyancy-induced energy is in the z direction. Thus, in this sublayer, the
 500 functional forms for similarity functions, as determined by dimensional analysis, are

$$501 \quad \varphi_m(\zeta) = A_u(-\zeta)^{-\frac{1}{3}} \quad (\text{A5})$$

$$502 \quad \varphi_h(\zeta) = A_T(-\zeta)^{-\frac{1}{3}} \quad (\text{A6})$$

503 in which A_u and A_T are constants.

504 Moreover, in the free-convective sublayer, buoyancy dominates the mechanical production of energy, and the
 505 pressure redistribution term feeds the buoyant energy in the vertical direction into the horizontal direction (Kader and Yaglom,
 506 1990). Thus, in this case, the dimensional analysis suggests

$$507 \quad \varphi_m(\zeta) = B_u(-\zeta)^{\frac{1}{3}} \quad (\text{A7})$$

$$508 \quad \varphi_h(\zeta) = B_T(-\zeta)^{\frac{1}{3}} \quad (\text{A8})$$

509 in which B_u and B_T are constants.

510 Thus, under unstable conditions, φ_m exhibits a nonmonotonic behaviour with respect to $-\zeta$ as the three sublayer
 511 theory suggested that for sufficiently large values of $-\zeta$, φ_m varies as the $+1/3$ power of ζ , in contrast to the case of the free
 512 convection limit, where both φ_m and φ_h follow the $-1/3$ power law. In the literature, various expressions for φ_m and φ_h are
 513 available based on the Kader and Yaglom (1990) three-sublayer model. However, the present study employs φ_m and φ_h based
 514 on the expressions implemented in the surface layer scheme (CLM4.0) of NCAR-CAM5 (Zeng et al., 1998) model. The
 515 expressions for φ_m and φ_h utilized in this study are as follows:

$$516 \quad \varphi_m = \begin{cases} (1 - 16\zeta)^{-\frac{1}{4}}, & -1.574 \leq \zeta \leq 0 \\ 0.7k^{\frac{2}{3}}(-\zeta)^{\frac{1}{3}}, & \zeta \leq -1.574 \end{cases} \quad (A9)$$

517 and

$$518 \quad \varphi_h = \begin{cases} (1 - 16\zeta)^{-\frac{1}{2}}, & -0.465 \leq \zeta \leq 0 \\ 0.9k^{\frac{4}{3}}(-\zeta)^{-\frac{1}{3}}, & \zeta \leq -0.465 \end{cases} \quad (A10)$$

519

520 Srivastava and Sharan (2021) classified these commonly used similarity functions stated above into four different classes based
521 on the exponents appearing in the expressions of φ_m and φ_h . The classification is as follows:

522

523 **Class 1.** This class consists of functions having the exponents of φ_m and φ_h as $-1/4$ and $-1/2$ (as in Eqns. A1 and A2),
524 respectively from near-neutral to strong unstable conditions. φ_m and φ_h proposed by Businger (1971) and Hogstrom (1996)
525 are the examples of class 1 functions.

526

527 **Class 2.** In this class, the similarity functions (φ_m and φ_h) having exponents of φ_m and φ_h as $-1/3$ for the entire range from
528 near-neutral to moderately unstable conditions (as in Eqns. A3 and A4), respectively are included. The functional forms
529 suggested by Carl et al. (1973) are the example of class 2 functions.

530

531 **Class 3.** φ_m and φ_h having exponents as $-1/4$ and $-1/2$, respectively in near-neutral conditions while $-1/3$ in strong
532 unstable conditions are included in this class. φ_m and φ_h based on Fairall et al. (1996), Grachev et al. (2000) and Fairall et al.
533 2003 are some examples of class 3 functions.

534

535 **Class 4.** Functional forms of φ_m and φ_h having the exponents as $-1/4$ and $-1/2$, respectively in near-neutral conditions
536 however, $1/3$ for φ_m and $-1/3$ for φ_h in strong unstable conditions are classified in this class (as in Eqns. A9 and A10). The
537 three-sublayer model for φ_m and φ_h suggested by Kader and Yaglom (1990) (Zeng et al. 1998) is one of the examples of
538 functions in this class.

539 Appendix B

540 This section consists of a brief description of the computation of surface turbulent fluxes in the revised MM5 surface layer
541 scheme. In a homogeneous surface layer, the dimensionless wind and temperature gradients are defined as

$$542 \quad \frac{kz \partial U}{u_* \partial z} = \varphi_m(\zeta), \quad (B1)$$

$$543 \quad \frac{kz \partial \theta}{\theta_* \partial z} = \varphi_h(\zeta). \quad (B2)$$

544 where L denotes the Obukhov length scale and U is the wind speed at height z ; k represents the von Karman constant and its
 545 value is taken as 0.4. Integrating Eqns. (B1) and (B2) with respect to z leads to

$$546 \quad U = \frac{u_*}{k} \left[\ln \left(\frac{z}{z_0} \right) - \left\{ \psi_m(\zeta) - \psi_m \left(\frac{z_0}{L} \right) \right\} \right], \quad (\text{B3})$$

$$547 \quad (\theta_a - \theta_g) = \frac{\theta_*}{k} \left[\ln \left(\frac{z}{z_h} \right) - \left\{ \psi_h(\zeta) - \psi_h \left(\frac{z_h}{L} \right) \right\} \right] \quad (\text{B4})$$

548 in which ψ_m and ψ_h denote the integrated form of similarity functions φ_m and φ_h . The roughness lengths for momentum and
 549 heat are denoted by z_0 and z_h , respectively. The ground and surface air potential temperature are denoted by θ_g and θ_a ,
 550 respectively. $\zeta (= \frac{z}{L})$ is the stability parameter and is defined as

$$551 \quad \zeta = \frac{kgz \theta_*}{\theta_a u_*^2} \quad (\text{B5})$$

552 ψ_m and ψ_h can be calculated from the following expression (e.g., Panofsky, 1963):

$$553 \quad \psi_m(\zeta) = \psi_h(\zeta) = \int_0^\zeta \frac{1 - \varphi_{m,h,q}(\zeta')}{\zeta'} d\zeta' \quad (\text{B6})$$

554 The bulk Richardson number (Ri_B) is given by:

$$555 \quad Ri_B = \frac{g(\theta_a - \theta_g)(z - z_0)^2}{\bar{\theta} U^2 (z - z_h)} \quad (\text{B7})$$

556 Substituting the values of U and $(\theta_a - \theta_g)$ from Eqns. (B3) and (B4) in Eqn. (B7), one gets

$$557 \quad Ri_B = \zeta \left[\frac{\left(1 - \frac{z_0}{z}\right)^2}{\left(1 - \frac{z_h}{z}\right)} \right] \frac{\left[\ln \left(\frac{z}{z_h} \right) - \left\{ \psi_h(\zeta) - \psi_h \left(\zeta \frac{z_h}{z} \right) \right\} \right]}{\left[\ln \left(\frac{z}{z_0} \right) - \left\{ \psi_m(\zeta) - \psi_m \left(\zeta \frac{z_0}{z} \right) \right\} \right]^2} \quad (\text{B8})$$

558 Note that Eqn. (B8) is a transcendental equation, and for a given value of Ri_B , the corresponding ζ value can be calculated
 559 using any iterative method.

560 The bulk transfer coefficient for momentum (C_D) and heat (C_H) are defined as:

$$561 \quad C_D = k^2 \left[\ln \left(\frac{z + z_0}{z_0} \right) - \left\{ \psi_m \left(\frac{z + z_0}{L} \right) - \psi_m \left(\frac{z_0}{L} \right) \right\} \right]^{-2} \quad (\text{B9})$$

$$562 \quad C_H = k^2 \left[\ln \left(\frac{z + z_0}{z_0} \right) - \left\{ \psi_m \left(\frac{z + z_0}{L} \right) - \psi_m \left(\frac{z_0}{L} \right) \right\} \right]^{-1} \left[\ln \left(\frac{z + z_h}{z_h} \right) - \left\{ \psi_h \left(\frac{z + z_h}{L} \right) - \psi_h \left(\frac{z_h}{L} \right) \right\} \right]^{-1} \quad (\text{B10})$$

563 Once we get C_D and C_H , then the momentum (τ), and sensible heat (H) fluxes are calculated using the following expressions:

$$564 \quad \tau = \rho C_D U^2 \quad (\text{B11})$$

$$565 \quad H = -\rho c_p C_H U (\theta_a - \theta_g), \quad (\text{B12})$$

566 **Appendix C**

567 In this section, the details of various physical parameterizations utilized in the real-case simulations using WRFv4.2.2 model
 568 and the different statistical indicators used for model evaluation.

569 The real-case simulations with the WRFv4.2.2 model utilised the Purdue Lin microphysics scheme (Lin et al., 1983);
 570 YSU (Hong, Noh, and Dudhia, 2006) PBL scheme; Kain-Fritsch (Kain and John, 2004) cumulus scheme; Dudhia (Dudhia,
 571 1989) shortwave scheme; RRTM (Mlawer et al., 1997) longwave scheme; Noah-MP land surface model (Niu et al., 2011);
 572 and revised MM5 surface layer scheme (Jimenez et al., 2012).

573 In the present study, different statistical indicators have been used for the model evaluation with respect to
 574 observations/reanalysis datasets. Statistical parameters such as mean absolute error (MAE), root mean square error (RMSE),
 575 mean bias (MB), index of agreement (IOA), and correlation coefficient (CC) are defined as:

- 576 1. Mean absolute error:

577
$$MAE = \frac{\sum_{i=1}^n |p_i - o_i|}{n}$$

- 578 2. Root mean square error:

579
$$RMSE = \sqrt{\frac{\sum_{i=1}^n (p_i - o_i)^2}{n}}$$

- 580 3. Mean bias

581
$$MB = \overline{(p_i - o_i)}$$

- 582 4. Index of agreement

583
$$IOA = 1 - \frac{\sum_{i=1}^n (o_i - p_i)^2}{\sum_{i=1}^n (|p_i - \bar{o}| + |o_i - \bar{o}|)^2}$$

- 584 5. Correlation coefficient

585
$$CC = \frac{\sum_{i=1}^n (p_i - \bar{p})(o_i - \bar{o})}{\sqrt{\sum_{i=1}^n (p_i - \bar{p})^2} \sqrt{\sum_{i=1}^n (o_i - \bar{o})^2}}$$

586 in which p_i and o_i represent the predicted and observed time series, respectively, while \bar{p} and \bar{o} are the predicted
 587 and observed mean for a considered variable, respectively.

- 588 6. Taylor diagram: It exhibits how well patterns match each other in terms of their correlation, ratio of their variances,
 589 and root mean square differences (Taylor, 2001).

- 590 7. Q-Q plot: It is a graphical technique used to compare the overall distribution of predicted and observed values for a
 591 variable (Venkatram, 1999)

592
 593 The error or deviation between observed and simulated values is measured by MAE, RMSE, and MB. On the other
 594 hand, IOA is used to assess the trend relationship, or how closely the magnitudes and signs of the observed values are related

595 to the projected values (Schlunzen and Sokhi 2008). In order to evaluate the spatial patterns with ERA5-Land reanalysis
596 dataset, statistical metrics such as mean bias (%), RMSE, and pattern correlation (PCC) have been used.

597 **Code and data availability:** Weather Research and Forecasting Model version 4.2.2 (WRFv4.2.2) is an open source model
598 and can be downloaded from https://www2.mmm.ucar.edu/wrf/users/download/get_source.html. The model output at the
599 location of the flux tower at Ranchi (23.412N, 85.440E), India is openly available at <https://doi.org/10.5281/zenodo.10435513>.
600 The raw observational data derived from the flux tower at Ranchi (23.412N, 85.440E; India) utilized in the present study can
601 be obtained from the Indian National Centre for Ocean Information Service upon request
602 (<http://www.incois.gov.in/portal/datainfo/ctczdata.jsp>). Hourly ERA5-Land reanalysis data utilized in this study can be found
603 in its official website <https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-land?tab=form>.

604 **Author contribution:** All authors contributed to the design of the study, analysis, and writing of the manuscript. PN carried
605 out the computations as well as the analysis of the model output.

606 **Competing interests:** The authors have declared that they have no conflict of interest.

607 **Acknowledgements:**

608 We would like to thank Dr. Manoj Kumar for providing observational data at Ranchi. The authors acknowledge the use of
609 NCAR-NCL and ERA5-Land reanalysis dataset for this study. The use of supercomputing facility (HPC) provided by IIT
610 Delhi is gratefully acknowledged. This work is partially supported by INSA, DST, DST-INSPIRE, and YES Foundation. We
611 wish to thank the reviewers for their helpful comments and suggestions, which have significantly enhanced the quality of this
612 paper.

613 **References:**

614 Bruin, H. A. R. de.: A Note on Businger's Derivation of Nondimensional Wind and Temperature Profiles under Unstable
615 Conditions, *J. Appl. Meteor. Climatol.*, 38, 626–28, [https://doi.org/10.1175/1520-0450\(1999\)038<0626:ANOBSD>2.0.CO;2](https://doi.org/10.1175/1520-0450(1999)038<0626:ANOBSD>2.0.CO;2),
616 1999.
617 Brutsaert, W.: *Evaporation into the Atmosphere: Theory, History, and Applications*. Springer, Dordrecht, 299.
618 <http://dx.doi.org/10.1007/978-94-017-1497-6>, 1982.
619 Brutsaert, W.: Stability Correction Functions for the Mean Wind Speed and Temperature in the Unstable Surface Layer,
620 *Geophys. Res. Lett.*, 19, 469–72. <https://doi.org/10.1029/92GL00084>, 1992.

621 Businger, J. A., Wyngaard, J. C., Izumi, Y., & Bradley, E. F. 1971. "Flux-Profile Relationships in the Atmospheric Surface
622 Layer in the Atmospheric Surface Layer". *J. Atmos. Sci.*, 28(2), 181-189. [https://doi.org/10.1175/1520-0469\(1971\)028<0181:FPRITA>2.0.CO;2](https://doi.org/10.1175/1520-0469(1971)028<0181:FPRITA>2.0.CO;2), 1971.

624 Carl, D. M., Tarbell, T. C., and Panofsky, H. A.: Profiles of Wind and Temperature from Towers over Homogeneous Terrain,
625 *J. Atmos. Sci.*, 30, 788-794, [http://dx.doi.org/10.1175/1520-0469\(1973\)030<0788:POWATF>2.0.CO;2](http://dx.doi.org/10.1175/1520-0469(1973)030<0788:POWATF>2.0.CO;2), 1973.

626 Cheng, Y., and Brutsaert, W.: Flux-Profile Relationships for Wind Speed and Temperature in the Stable Atmospheric
627 Boundary Layer, *Boundary-Layer Meteorol.*, 114, 519–38. <https://doi.org/10.1007/s10546-004-1425-4>, 2005.

628 Dwivedi, A. K., Chandra, S., Kumar, M., Kumar, S., and Kumar, N. V. P. K.: Spectral Analysis of Wind and Temperature
629 Components during Lightning in Pre-Monsoon Season over Ranchi, *Meteorol. Atmos. Phys.*, 127, 95–105,
630 <https://doi.org/10.1007/s00703-014-0346-0>, 2015.

631 Dyer, A. J.: A Review of Flux-Profile Relationships, *Boundary-Layer Meteorol.*, 7, 363–372.
632 <https://doi.org/10.1007/BF00240838>, 1974.

633 Fairall, C. W., Bradley, E. F., Hare, J. E., Grachev, A. A., and Edson, J. B.: Bulk Parameterization of Air–Sea Fluxes: Updates
634 and Verification for the COARE Algorithm, *J. Climate*, 16, 571–591, [https://doi.org/10.1175/1520-0442\(2003\)016<0571:BPOASF>2.0.CO;2](https://doi.org/10.1175/1520-0442(2003)016<0571:BPOASF>2.0.CO;2), 2003.

636 Fairall, C. W., Bradley, E. F., Rogers, D. P., Edson, J. B., and Young, G. S.: Bulk Parameterization of Air-Sea Fluxes for
637 Tropical Ocean global Atmosphere Coupled-Ocean Atmosphere Response Experiment, *J. Geophys. Res.*, 101, 3747–3764,
638 doi:10.1029/95JC03205, 1996.

639 Friedl, M. A., McIver, D. K., Hodges, J. C. F., Zhang, X. Y., Muchoney, D., Strahler, A. H., Woodcock, C. E., Gopal, S.,
640 Schneider, A., Cooper, A., Baccini, A., Gao, F., and Schaaf, C.: Global Land Cover Mapping from MODIS: Algorithms and
641 Early Results, *Remote Sens. Environ.* 83, 287–302, [https://doi.org/10.1016/S0034-4257\(02\)00078-0](https://doi.org/10.1016/S0034-4257(02)00078-0), 2002.

642 Giorgi, F., Coppola, E., Solmon, F., Mariotti, L., Sylla, M. B., Bi, X., Elguindi, N., et al.: RegCM4: Model Description and
643 Preliminary Tests over Multiple CORDEX Domains, *Clim. Res.*, 52, 7–29, <http://doi.org/10.3354/cr01018>, 2012.

644 Grachev, A. A., Fairall, C. W., and Bradley, E. F.: Convective Profile Constants Revisited, *Boundary-Layer Meteorol.*, 94,
645 495–515, <https://doi.org/10.1023/A:1002452529672>, 2000.

646 Grachev, A. A., Andreas, E. L., Fairall, C. W., Guest, P. S., and Persson, P. O. G.: SHEBA Flux-Profile Relationships in the
647 Stable Atmospheric Boundary Layer, *Boundary-Layer Meteorol.*, 124, 315–33, <https://doi.org/10.1007/s10546-007-9177-6>,
648 2007.

649 Grell, G. A., Dudhia, J., & Stauffer, D.: A description of the fifth-generation Penn State/NCAR Mesoscale Model (MM5) (No.
650 NCAR/TN-398+STR), University Corporation for Atmospheric Research. <http://doi:10.5065/D60Z716B>, 1994.

651 Hicks, B. B.: Wind Profile Relationships from the 'Wangara' Experiment, *Q. J. R. Meteorol. Soc.*, 102, 535–51,
652 <https://doi.org/10.1002/qj.49710243304>, 1976.

653 Hogstrom, U.: Review of Some Basic Characteristics of the Atmospheric Surface Layer, *Boundary-Layer Meteorol.*, 78, 215–
654 246, <https://doi.org/10.1007/BF00120937>, 1996.

655 Holtslag, A. A. M., and De Bruin, H. A. R.: Applied Modeling of the Night-time Surface Energy Balance over Land, *J. Appl.*
656 *Meteor. Climatol.*, 27, 689-704, [https://doi.org/10.1175/1520-0450\(1988\)027<0689:AMOTNS>2.0.CO;2](https://doi.org/10.1175/1520-0450(1988)027<0689:AMOTNS>2.0.CO;2), 1988.

657 Izumi, Y.: Kansas 1968 Field Program Data Report. Bedford, MA, Air Force Cambridge Research Papers, No. 379, 79 pp,
658 1971.

659 Jiménez, P. A., González-Rouco, J. F., García-Bustamante, E., Navarro, J., Montávez, J. P., de Arellano, J. V., Dudhia, J., &
660 Muñoz-Roldan, A.: Surface Wind Regionalization over Complex Terrain: Evaluation and Analysis of a High-Resolution WRF
661 Simulation, *J. Appl. Meteor. Climatol.*, 49, 268-287, <https://doi.org/10.1175/2009JAMC2175.1>, 2010.

662 Jiménez, P. A., Dudhia, J., González-Rouco, J. F., Navarro, J., Montávez, J. P., & García-Bustamante, E.: A Revised Scheme
663 for the WRF Surface Layer Formulation, *Mon. Wea. Rev.*, 140, 898-918. <https://doi.org/10.1175/MWR-D-11-00056.1>, 2012.

664 Kader, B. A., and Yaglom, A. M.: Mean Fields and Fluctuation Moments in Unstably Stratified Turbulent Boundary Layers,
665 *Journal of Fluid Mechanics*, 212, 637–662, <https://doi.org/10.1017/S0022112090002129>, 1990.

666 Monin, A. S., and Obukhov, A. M.: Basic Laws of Turbulent Mixing in the Surface Layer of the Atmosphere, *Tr. Akad. Nauk*
667 *SSSR Geophys. Inst* 24(151), 163–187, 1954.

668 Namdev, P., Sharan, M. & Mishra, S.K.: Impact of the similarity functions of surface layer parametrization in a climate model
669 over the Indian region, *Q. J. R. Meteorol. Soc.*, 149, 152–170, <https://doi.org/10.1002/qj.4400>, 2023.

670 Namdev, P., Srivastava, P., Sharan, M., & Mishra, S. K.: An Update to WRF Surface Layer Parameterization over an Indian
671 Region, *Dynam. Atmos. Ocean*, 105, 101414, <https://doi.org/10.1016/j.dynatmoce.2023.101414>, 2024.

672 Panofsky, H. and Dutton, J.: Atmospheric Turbulence, John Wiley & Sons, New York, 397 p, 1984.

673 Rao, K. G., Narasimha, R., and Prabhu, A.: Estimation of Drag Coefficient at Low Wind Speeds over the Monsoon Trough
674 Land Region during MONTBLEX-90, *Geophys. Res. Lett.*, 23, 2617–2620, <https://doi.org/10.1029/96GL02368>, 1996.

675 Rao, K. G., and Narasimha, R.: Heat-Flux Scaling for Weakly Forced Turbulent Convection in the Atmosphere, *Journal of*
676 *Fluid Mechanics*, 547, 115–135, <https://doi.org/10.1017/S0022112005007251>, 2006.

677 Sharan, M., and Srivastava, P.: Characteristics of the Heat Flux in the Unstable Atmospheric Surface Layer, *J. Atmos. Sci.*,
678 73, 4519–4529, <https://doi.org/10.1175/JAS-D-15-0291.1>, 2016.

679 Srivastava, P., and Sharan, M.: Characteristics of the Drag Coefficient over a Tropical Environment in Convective Conditions,
680 *J. Atmos. Sci.*, 72, 4903–4913, <https://doi.org/10.1175/JAS-D-14-0383.1>, 2015.

681 Srivastava, P., and Sharan, M.: Analysis of Dual Nature of Heat Flux Predicted by Monin-Obukhov Similarity Theory: An
682 Impact of Empirical Forms of Stability Correction Functions, *J. Geophys. Res. Atmos.*, 124, 3627–3646,
683 <https://doi.org/10.1029/2018JD029740>, 2019.

684 Srivastava, P., and Sharan, M.: Uncertainty in the Parameterization of Surface Fluxes under Unstable Conditions, *J. Atmos.*
685 *Sci.*, 78, 2237–2247, <https://doi.org/10.1175/JAS-D-20-0350.1>, 2021.

686 Srivastava, P., Sharan, M., and Kumar, M.: A Note on Surface Layer Parameterizations in the Weather Research and Forecast
687 Model, *Dynam. Atmos. Ocean*, 96, 101259, <https://doi.org/10.1016/j.dynatmoce.2021.101259>, 2021.

688 Srivastava, P., Sharan, M., Kumar, M., and Dhuria, A. K.: On Stability Correction Functions over the Indian Region under
689 Stable Conditions, *Meteorol. Appl.*, 27:e1880, <https://doi.org/10.1002/met.1880>, 2020.

690 Stull, R. B.: *An Introduction to Boundary Layer Meteorology*, Kluwer Academic Publishers, Dordrecht, The Netherlands, 13,
691 670 pp, <https://doi.org/10.1007/978-94-009-3027-8>, 1988.

692 Venkatram, A.: Applying a framework for evaluating the performance of air quality models, in: *Proceedings of the sixth*
693 *International Conference on Harmonisation within Atmospheric Dispersion modeling for Regulatory Applications*, Rouen,
694 France, 11 – 14 October, 1999, 11 – 14, 1999.

695 Webb, E. K.: Profile Relationships: The Log-linear Range, and Extension to Strong Stability, *Q. J. R. Meteorol. Soc.*, 96, 67–
696 90, <https://doi.org/10.1002/qj.49709640708>, 1970.

697 Wilson, D. K.: An Alternative Function for the Wind and Temperature Gradients in Unstable Surface Layers, *Boundary-Layer*
698 *Meteorol.*, 99, 151–158, <https://doi.org/10.1023/A:1018718707419>, 2001.

699 Zeng, X., Zhao, M., and Dickinson, R. E.: Intercomparison of Bulk Aerodynamic Algorithms for the Computation of Sea
700 Surface Fluxes Using TOGA COARE and TAO Data, *J. Climate*, 11, 2628–2644, [https://doi.org/10.1175/1520-](https://doi.org/10.1175/1520-0442(1998)011<2628:IOBAAF>2.0.CO;2)
701 [0442\(1998\)011<2628:IOBAAF>2.0.CO;2](https://doi.org/10.1175/1520-0442(1998)011<2628:IOBAAF>2.0.CO;2), 1998.

702

703

704

705

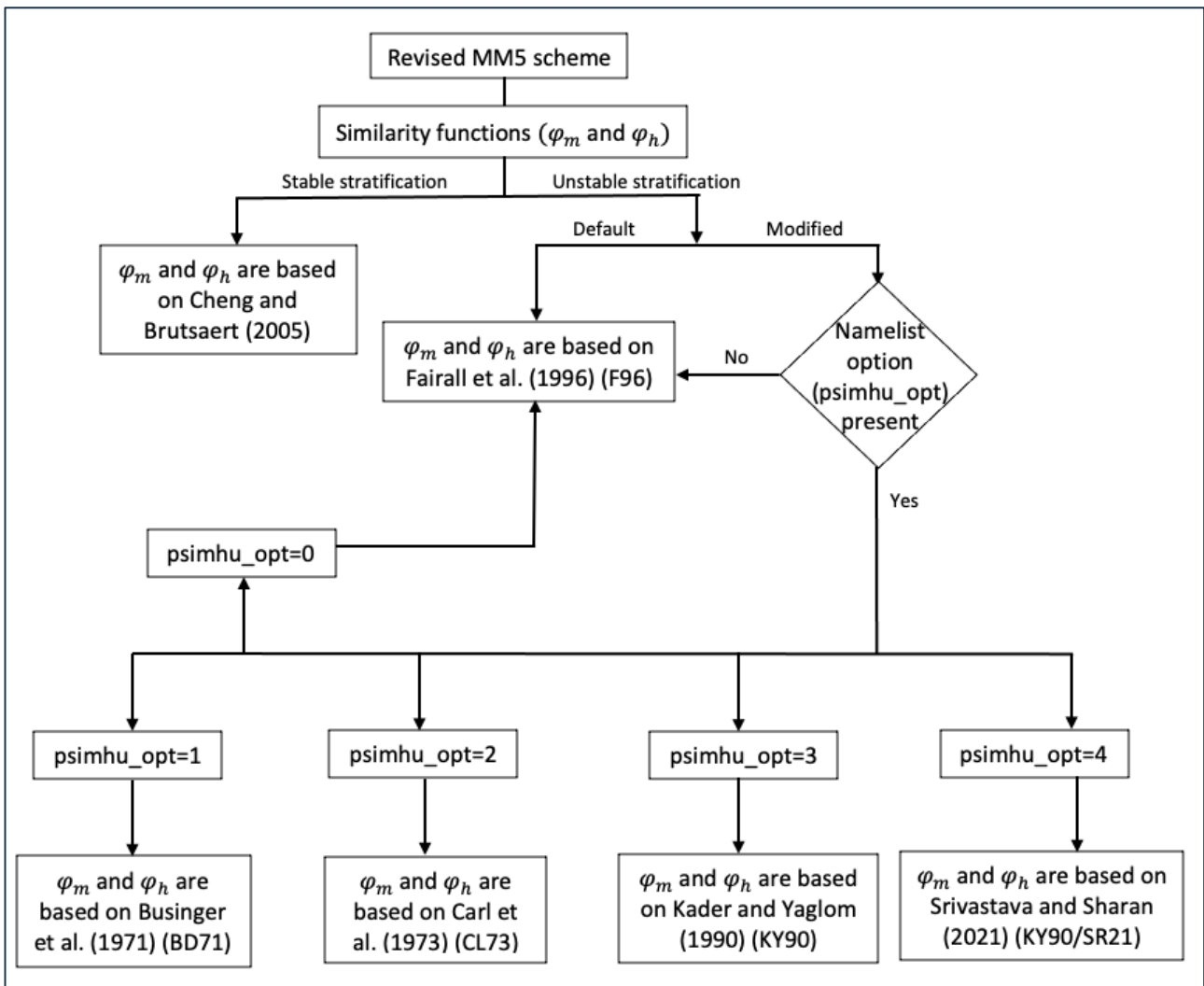
706

707

708

709

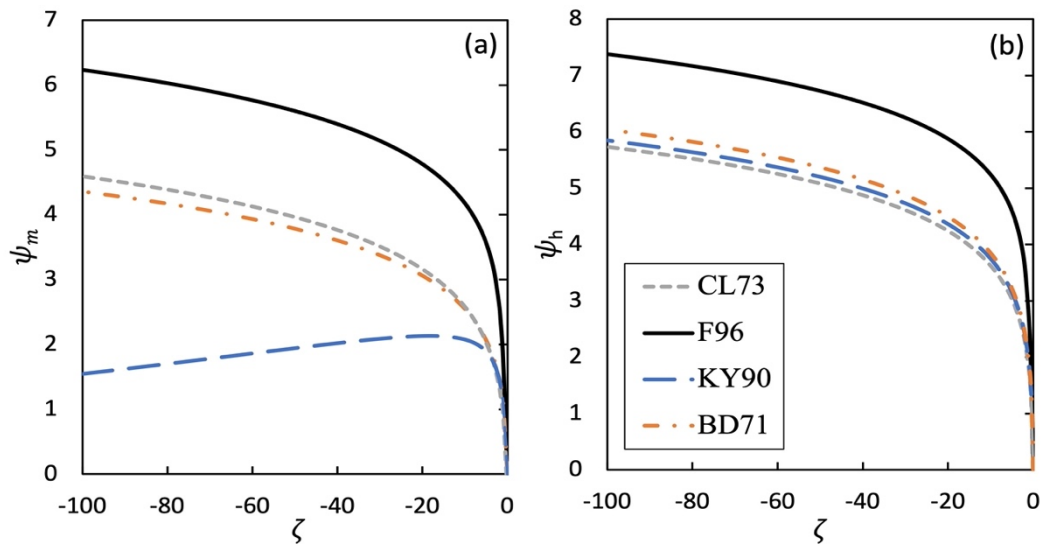
710



711

712 **Figure 1: Flowchart to provide a brief description of different options for similarity functions in the modified surface layer scheme**
 713 **that can be controlled by namelist variable psimhu_opt.**

714



716 **Figure 2: Integrated similarity functions $\psi_{m,h}(\zeta)$ for momentum and heat for default (F96; black line) and newly installed (BD71,**
 717 **CL73, and KY90; orange, grey and blue lines, respectively) functions for unstable atmospheric surface layer.**

718

719

720

721

722

723

724

725

726

727

728

729

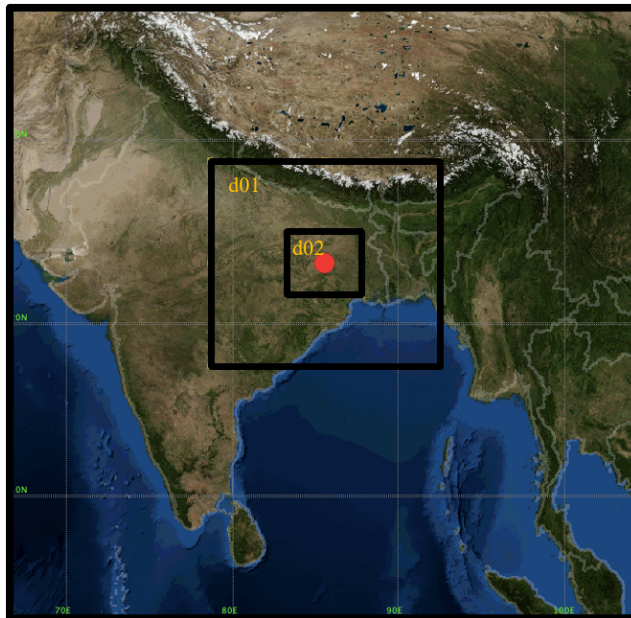
730

731

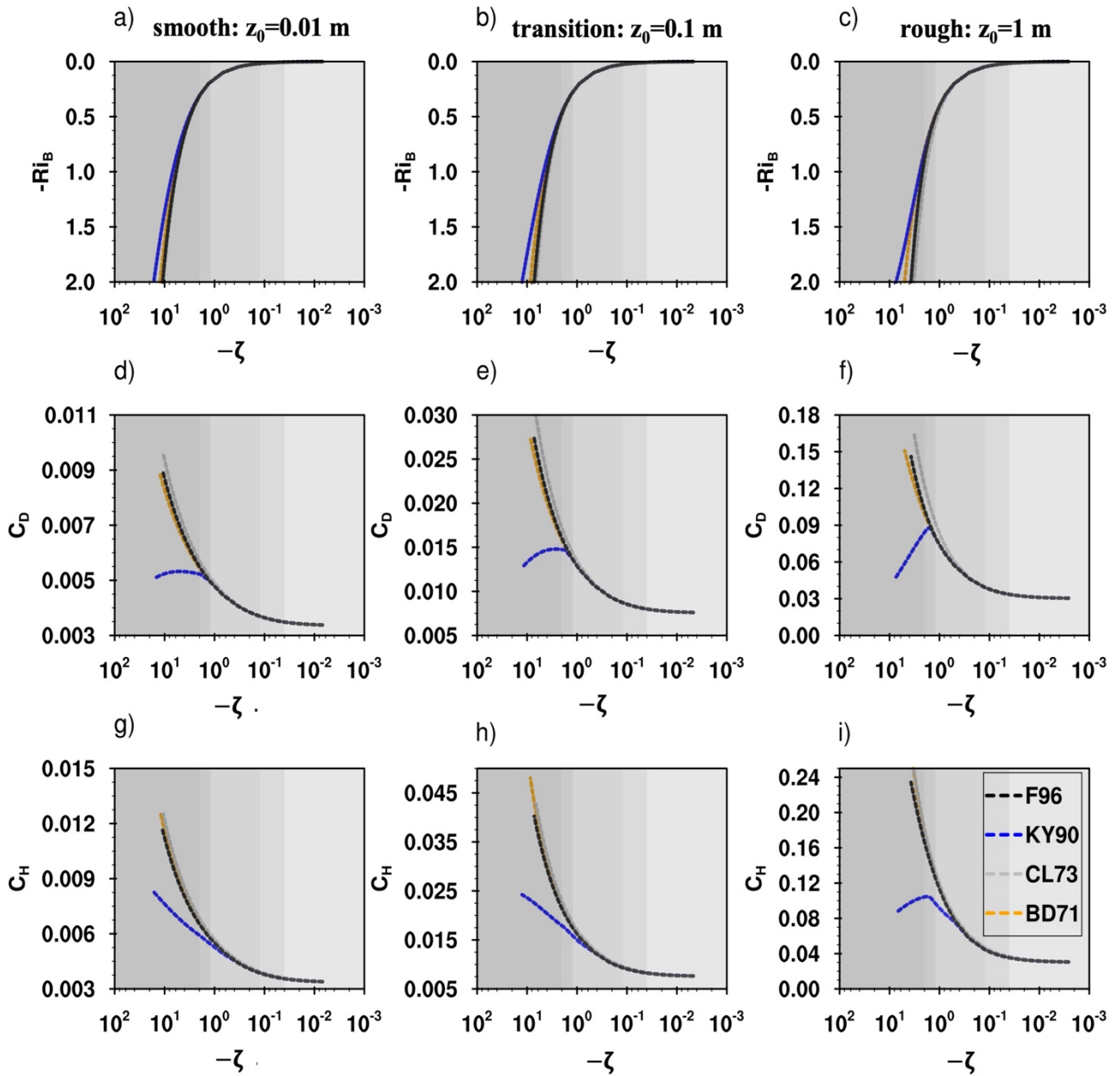
732

733

734



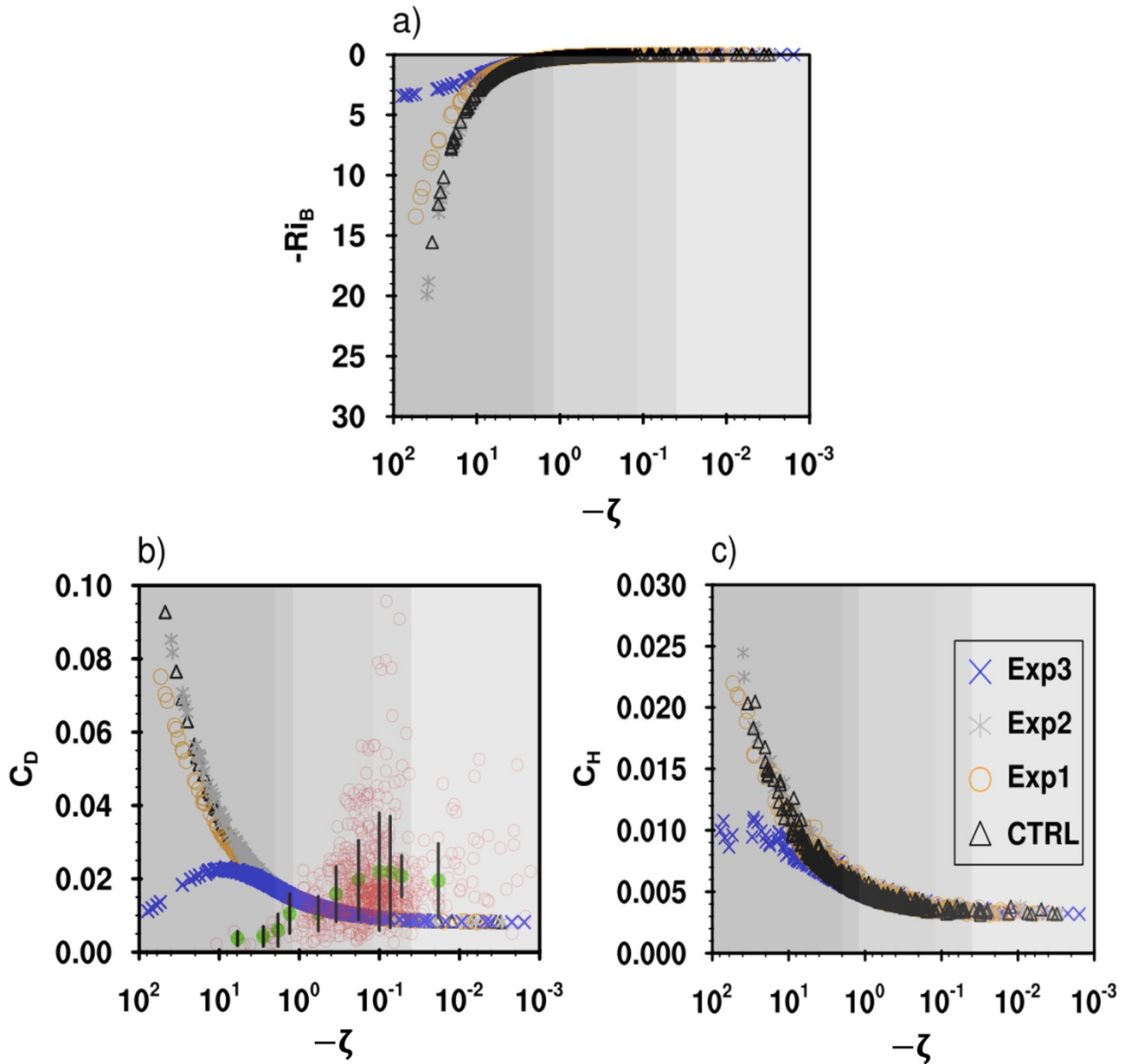
735 **Figure 3: Spatial distribution of domain used for the simulations using the WRF model. The spatial resolution for domains d01 and**
 736 **d02 is 6×6 km and 2×2 km, respectively. The domain d02 covers 446×392 km² area around the centre point.**



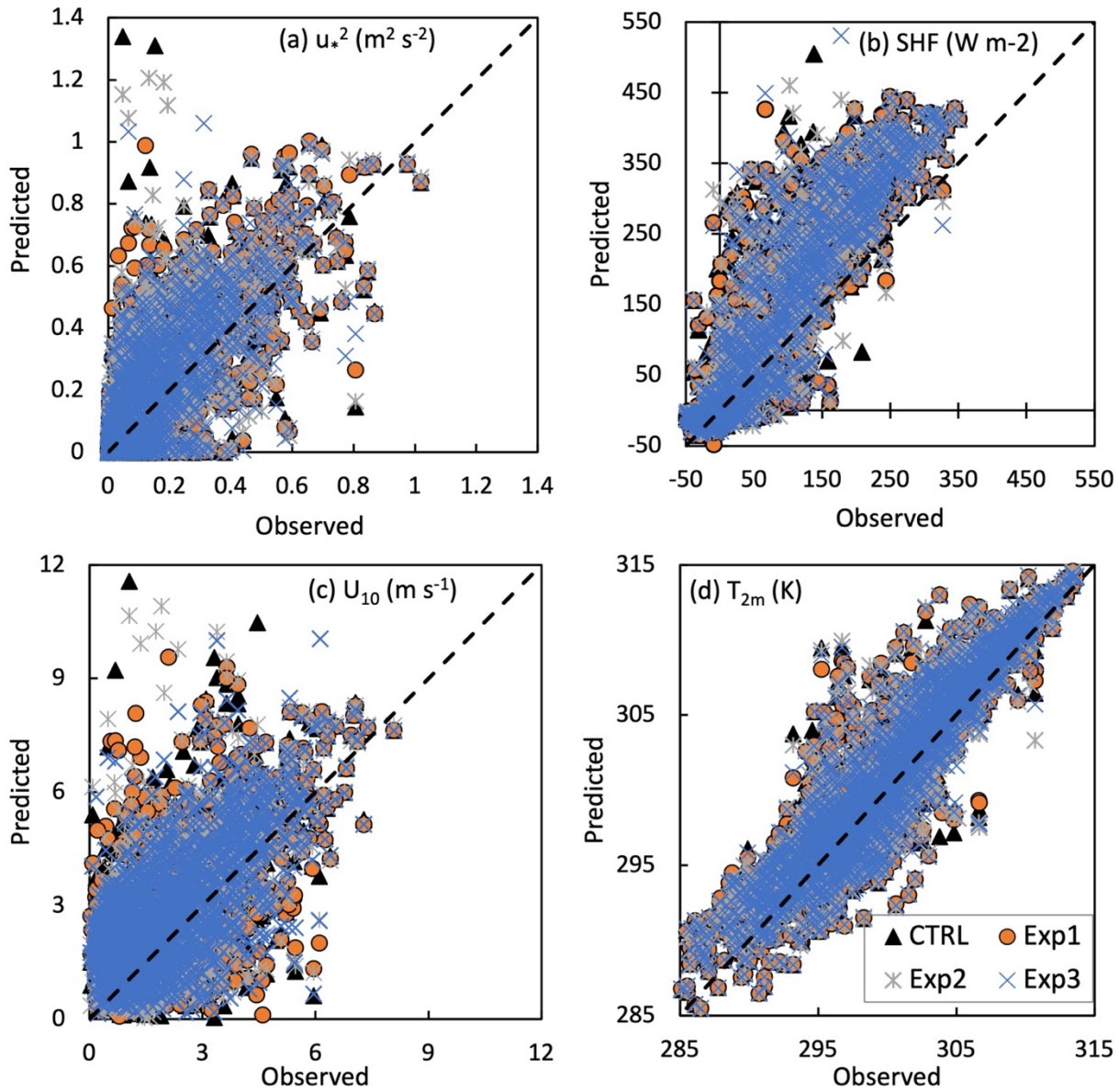
738

739 **Figure 4: Variation of ζ with Ri_B (upper panel), C_D (middle panel) and C_H (lower panel) with ζ calculated from bulk flux algorithm**
 740 **(offline simulation) for different functional forms of ψ_m and ψ_h corresponding to BD71, CL73, KY90, and F96 forms for smooth**
 741 **($z_0 = 0.01$ m; 1st column), transition ($z_0 = 0.1$ m; 2nd column), and rough ($z_0 = 1.0$ m; 3rd column) surfaces. The background**
 742 **colour corresponds to different sublayers in convective conditions (Kader and Yaglom 1990), from the dynamic sublayer ($0 \geq \zeta >$**
 743 **-0.04 ; light grey) to the free convective sublayer ($\zeta < -2$; dark grey).**

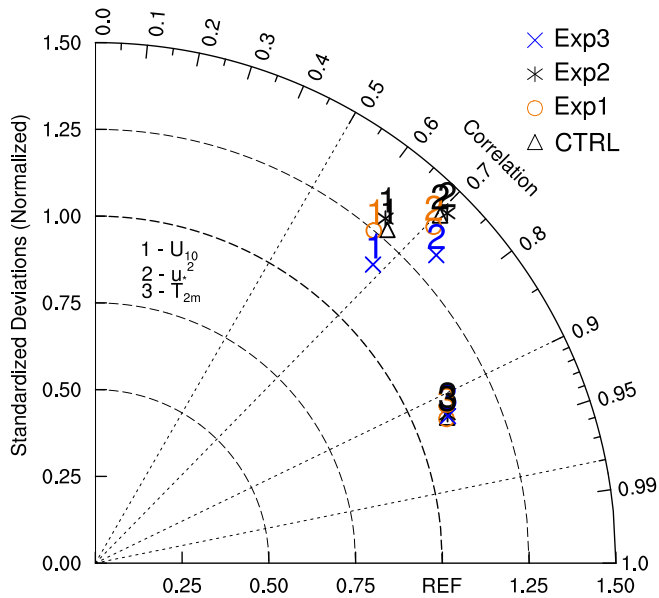
744



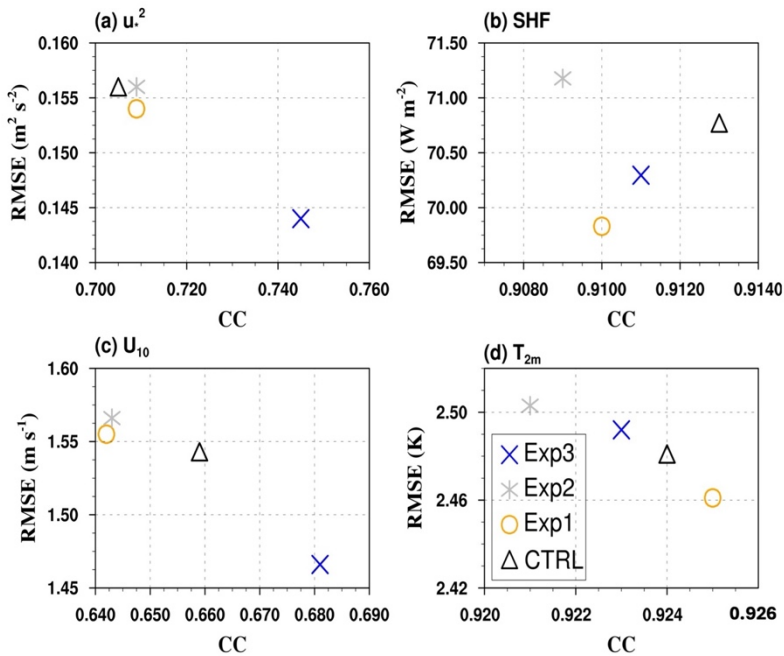
748 Figure 5: Variation of model simulated (a) ζ with Ri_B , (b) C_D and (c) C_H with ζ from different experiments using different ψ_m and
 749 ψ_h corresponding to F96 (CTRL), BD71 (Exp1), CL73 (Exp2), and KY90 (Exp3) under convective conditions. The red circles in (b)
 750 denote the observed C_D with ζ at the location of flux tower. The mean values of observed C_D in each sublayer are shown with green
 751 solid circles along with standard deviations in the form of error bars. Depending upon the data availability, two or three bins of
 752 equal width are chosen in each sublayer. The background colour corresponds to different sublayers in convective conditions (Kader
 753 and Yaglom 1990), from the dynamic sublayer ($0 \geq \zeta > -0.04$; light grey) to the free convective sublayer ($\zeta < -2$; dark grey).



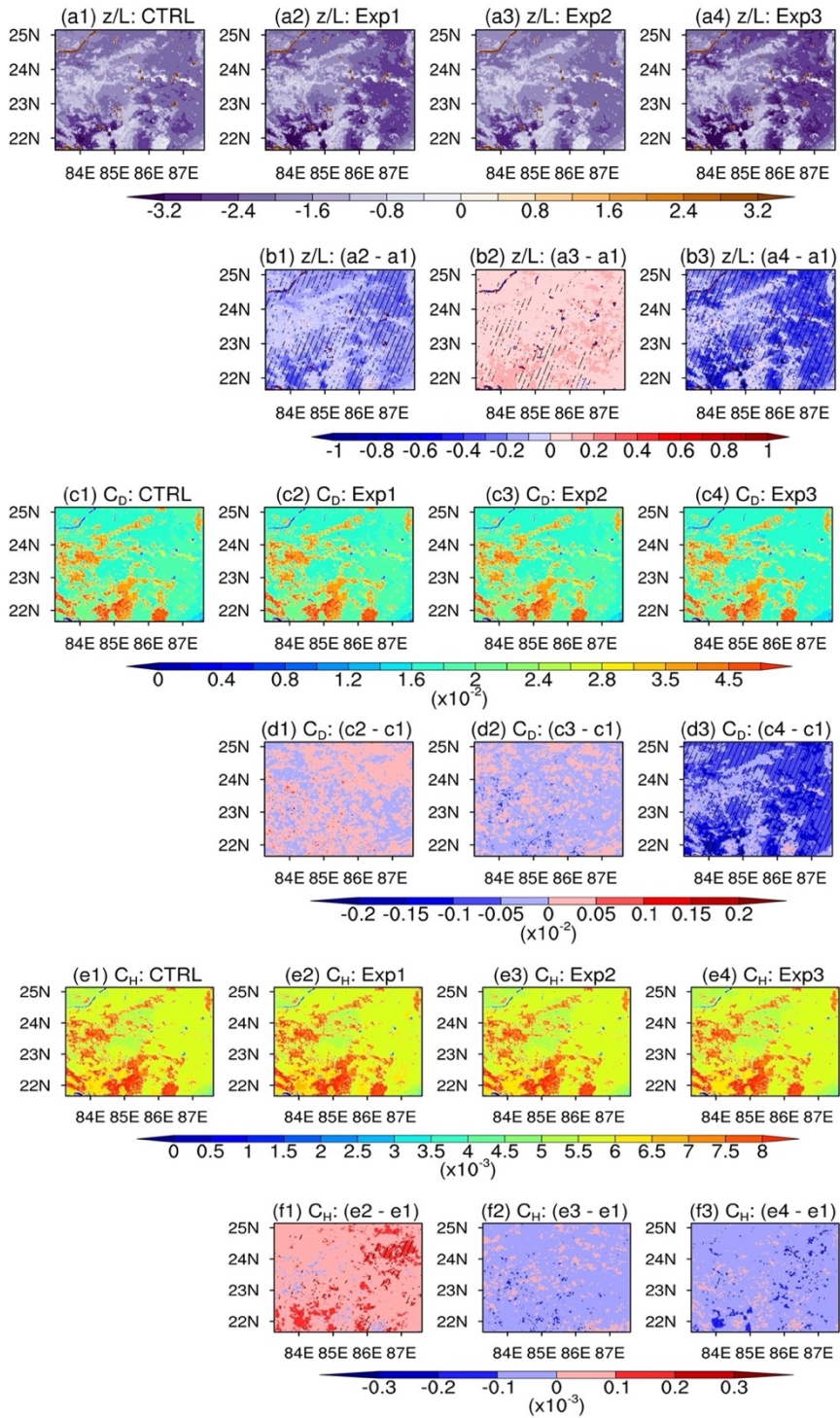
755 Figure 6: Scatter plot of model simulated (a) u_*^2 ($\text{m}^2 \text{s}^{-2}$) (representative of momentum flux), (b) SHF (W m^{-2}) (sensible heat flux), (c)
 756 U_{10} (m s^{-1}) (wind speed at 10 m height), and (d) T_{2m} (K) (temperature at 2 m height) vs observed values at the location of flux tower
 757 at Ranchi (23.412oN, 85.440oE), India (centre point of the domain) during pre-monsoon season (MAM).



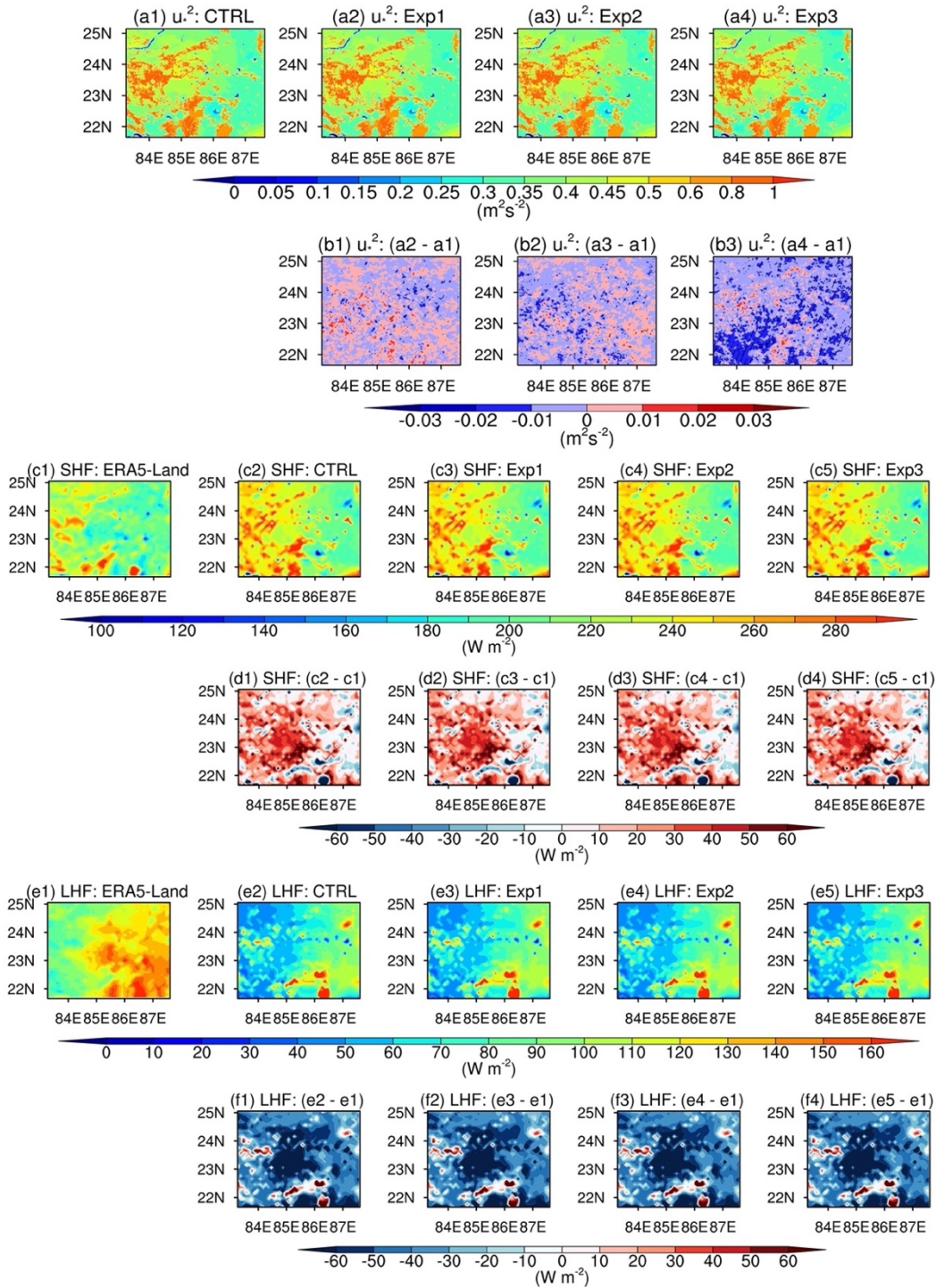
760 Figure 7: Taylor diagram showing the correlation coefficient, normalized standard deviations for U_{10} , u_*^2 , and T_{2m} from different
 761 experiments together with CTRL simulation with respect to observations derived from flux tower installed at Ranchi (23.412°N,
 762 85.440°E), India.



763 Figure 8: Scatter plot between correlation coefficient (CC) and root mean square error (RMSE) for (a) u_*^2 , (b) SHF, (c) U_{10} , and (d)
 764 T_{2m} simulated by various experiments (Exp1-3) together with CTRL simulation for pre-monsoon season (MAM; 2009) at the location
 765 of the flux tower (23.412°N, 85.440°E).

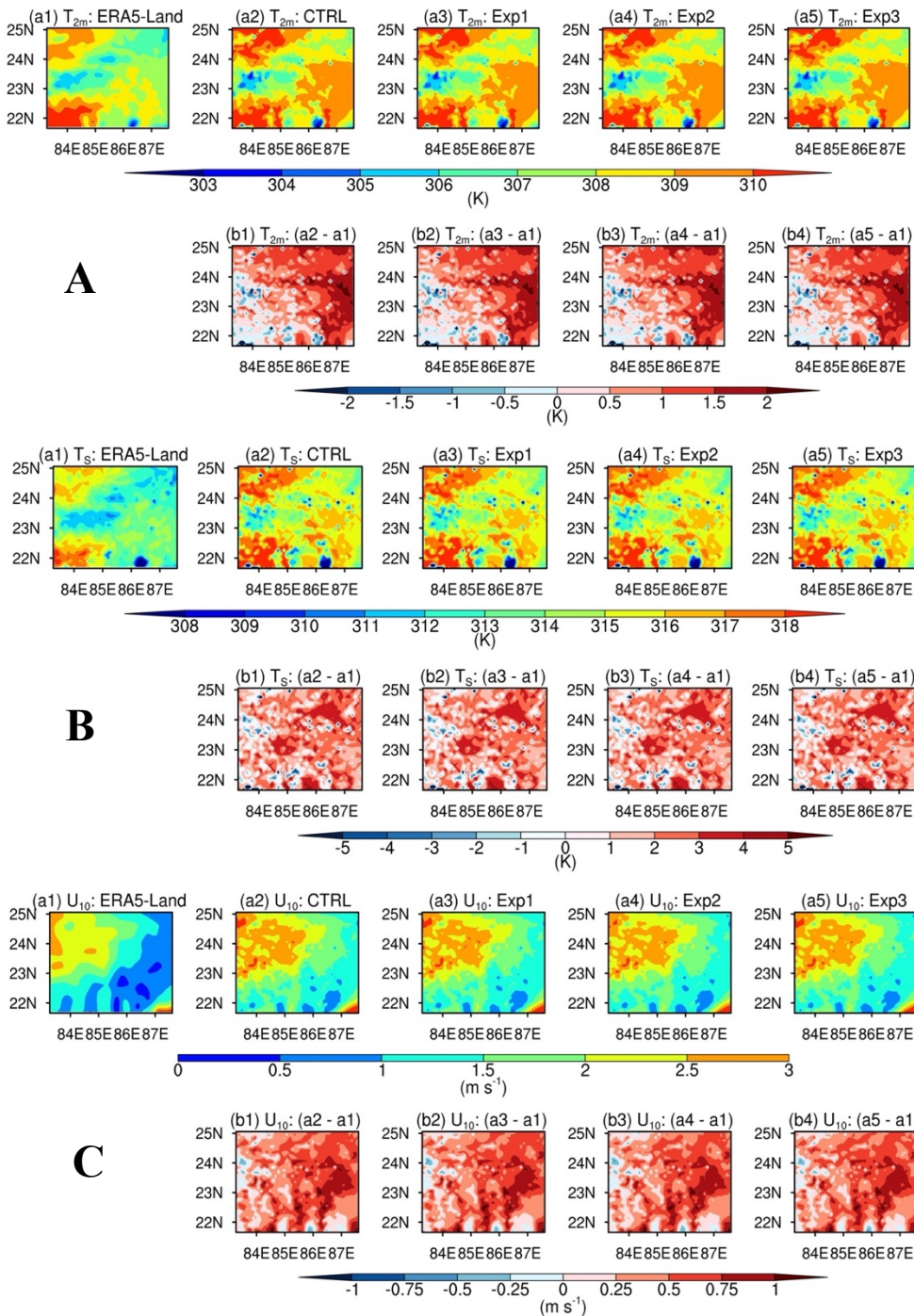


766 **Figure 9: Mean spatial distribution of model simulated ζ (1st row), C_D (3rd row) and C_H (5th row) from different experiments and**
767 **their differences with respect to CTRL simulation averaged during daytime for whole simulation period. Hatched regions show**
768 **significant differences at 95% confidence level in experiments with respect to CTRL simulation.**

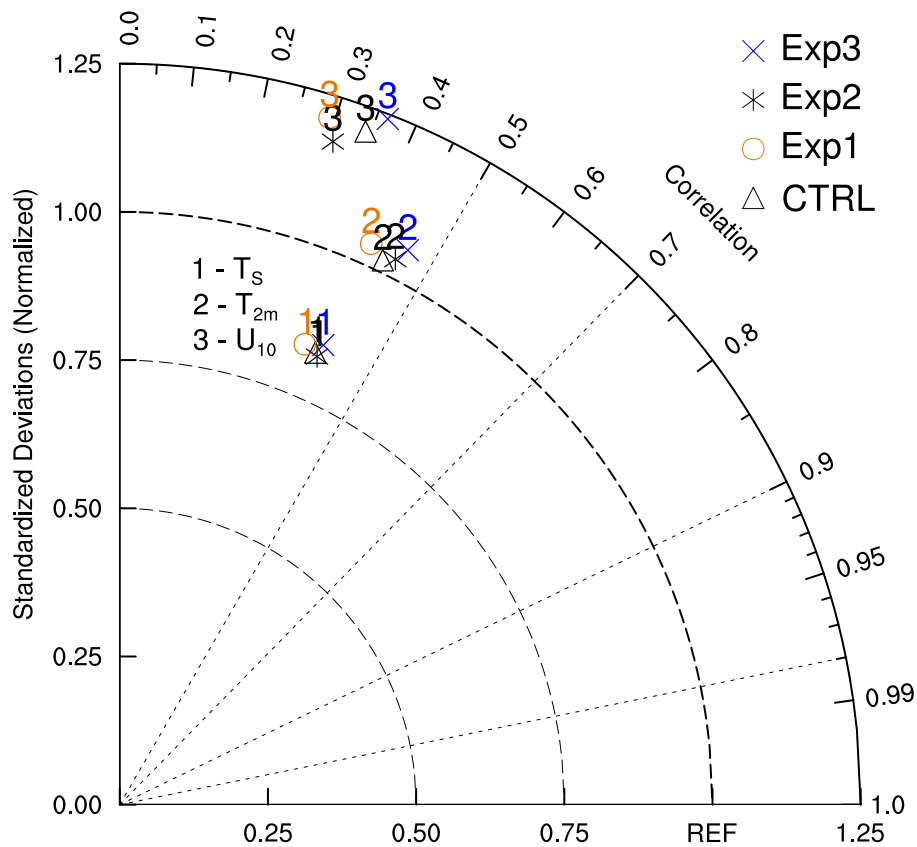


769 Figure 10: Mean spatial distribution of simulated u_*^2 (1st row) from different experiments and their differences (2nd row) with respect
770 to CTRL simulation. SHF and LHF from ERA5-Land reanalysis and simulated using various experiments and their differences
771 with respect to ERA5-Land data averaged during daytime for the whole simulation period are shown. Hatched regions show
772 significant differences at 95% confidence level in experiments with respect to CTRL simulation.

773
774
775
776
777
778
779
780
781
782
783
784
785
786
787
788
789
790
791
792
793
794
795
796
797
798
799
800



801 **Figure 11: In upper panel (A), mean spatial distribution of T_{2m} from ERA5-Land reanalysis (a1) and simulated using different**
802 **experiments (a2-a5) and their differences with respect to ERA5-Land reanalysis (b1-b4) averaged during daytime for the whole**
803 **simulation period. Middle (Lower) panel is same as the upper panel but for T_S (U_{10}).**



805 **Figure 12: Taylor diagram showing the correlation coefficient, normalized standard deviations for T_s (K), T_{2m} (K), and U_{10} ($m s^{-1}$)**
 806 **from different experiments together with CTRL simulation with respect to ERA5-Land reanalysis dataset averaged during strong**
 807 **convective conditions (hours during daytime in which ζ is smaller than -10) for whole simulation period.**

808

809

810

811

812

813

814

815

816

817

818

819 **Table 1. Description of various simulations conducted in this study.**

Experiments	Description
CTRL	Simulation using default surface layer scheme with F96 functions
Exp1	Simulation using surface layer scheme with BD71 functions
Exp2	Simulation using surface layer scheme with CL73 functions
Exp3	Simulation using surface layer scheme with newly installed KY90 functions

820
821
822
823
824
825
826
827
828
829
830
831
832
833
834
835
836
837
838
839
840

841 Table 2: Comparison statistics for u_*^2 ($\text{m}^2 \text{s}^{-2}$), SHF (W m^{-2}), U_{10} (m s^{-1}), and T_{2m} (K) simulated using different experiments together
842 with CTRL simulation with respect to observations derived from flux tower at Ranchi (India) for MAM season. The mean absolute
843 error (MAE), root mean square error (RMSE), mean bias (MB), index of agreement (IOA), and correlation coefficient (CC) are
844 shown.

MAM		u_*^2 ($\text{m}^2 \text{s}^{-2}$)	SHF (W m^{-2})	U_{10} (m s^{-1})	T_{2m} (K)
CTRL	MAE	0.09	43.46	1.20	1.82
	RMSE	0.16	70.77	1.54	2.48
	MB	0.03	34.88	0.83	0.93
	IOA	0.82	0.89	0.73	0.95
	CC	0.71	0.91	0.66	0.92
Exp1	MAE	0.09	42.72	1.20	1.81
	RMSE	0.15	69.83	1.56	2.46
	MB	0.03	33.06	0.81	0.90
	IOA	0.82	0.89	0.72	0.96
	CC	0.71	0.91	0.64	0.93
Exp2	MAE	0.09	43.55	1.20	1.84
	RMSE	0.16	71.18	1.57	2.50
	MB	0.03	34.49	0.81	0.87
	IOA	0.82	0.89	0.72	0.95
	CC	0.71	0.91	0.64	0.92
Exp3	MAE	0.08	42.96	1.16	1.83
	RMSE	0.14	70.30	1.47	2.49
	MB	0.03	33.47	0.78	0.91
	IOA	0.84	0.89	0.75	0.95
	CC	0.74	0.91	0.68	0.92

845

846

847

848

849

850

851

852

853

854

855 Table 3: Comparison statistics for T_{2m} (K), T_s (K), and U_{10} ($m s^{-1}$) simulated using different experiments together with CTRL
856 simulation with respect to ERA5-Land reanalysis data averaged during strong unstable stratification (hours during daytime in
857 which ζ is smaller than -10) for whole simulation period. The percent mean bias (Bias %), pattern correlation coefficient (PCC),
858 and root mean square error (RMSE) are shown.

MAM	T_s (K)			T_{2m} (K)			U_{10} ($m s^{-1}$)		
	Bias (%)	RMSE	PCC	Bias (%)	RMSE	PCC	Bias (%)	RMSE	PCC
CTRL	1.26	4.01	0.40	0.64	2.13	0.43	-4.96	0.44	0.34
Exp1	1.26	4.03	0.37	0.64	2.16	0.40	-4.43	0.45	0.29
Exp2	1.25	3.99	0.40	0.63	2.10	0.45	-5.39	0.44	0.31
Exp3	1.24	3.97	0.41	0.62	2.10	0.46	-0.28	0.47	0.36

859

860

861