Modeling Commercial-Scale CO_2 Storage in the Gas Hydrate Stability Zone with PFLOTRAN v6.0

Michael Nole[1], Jonah Bartrand[1], Fawz Naim[2], and Glenn Hammond[1]

[1]Pacific Northwest National Laboratory, Richland, 99354, USA
5   [2]Ohio State University, Columbus, 43210, USA

Correspondence to: Michael Nole (michael.nole@pnnl.gov)

Abstract

Safe and secure carbon dioxide (CO_2) storage is likely to be critical for mitigating some of the most dangerous effects of climate change. In the last decade, there has been a significant increase in activity associated with reservoir characterization
10   and site selection for large-scale CO_2 storage projects across the globe. These prospective storage sites tend to be selected for their optimal structural, petrophysical, and geochemical trapping potential. However, it has also been suggested that storing CO_2 in reservoirs within the CO_2 hydrate stability zone (GHSZ), characterized by high pressures and low temperatures (e.g., Arctic or marine environments), could provide a natural thermodynamic barrier to gas leakage. Evaluating the prospect of commercial-scale, long-term storage of CO_2 in the GHSZ requires reservoir-scale modelling capabilities designed to account
15   for the unique physics and thermodynamics associated with these systems. We have developed the HYDRATE flow mode and accompanying fully implicit parallel well model in the massively parallel subsurface flow and reactive transport simulator PFLOTRAN to model CO_2 injection into the marine GHSZ. We have applied these capabilities to a set of CO_2 injection scenarios designed to reveal the challenges and opportunities for commercial-scale CO_2 storage in the GHSZ.

1 Introduction

20   Large-scale deployment of carbon capture and storage (CCS) projects is likely to be critical for constraining future global temperature increase due to climate change, yet major uncertainties exist regarding potential injectivity of CO_2 in subsurface reservoirs (Lane et al., 2021). Evaluating CO_2 storage sites for long-term sequestration requires synthesizing sophisticated laboratory, field, and modelling tools to assess the CO_2 trapping potential of a prospective subsurface reservoir during a large-scale injection over a significant post-injection performance period. CO_2 trapping in sequestration applications involves
25   potentially both physical and chemical trapping mechanisms which include structural, solubility, capillary, and mineralization trapping (Al Hameli et al., 2022). Currently, large-scale carbon capture and storage (CCS) projects around the globe are each storing over 400,000 metric tons (Mt) of CO_2 annually (Snæbjörnsdóttir et al., 2020); the U.S. Department of Energy's CarbonSAFE initiative aims to develop CO_2 storage complexes across the U.S. that would be capable of storing total volumes exceeding 50 million metric tons (MMt) of CO_2 each (Sullivan et al., 2020). To achieve this vision, a diverse set of potential

30 reservoir host rocks and environments is being considered. This includes speculation about the feasibility of offshore CO_2 sequestration and mineralization, such as in Cascadia Basin basalts offshore the U.S. Pacific Northwest (Goldberg et al., 2018). Shallow sub-sea environments are not only isolated from the atmosphere by a large water body, but they can also exist in a unique pressure and temperature regime conducive for forming gas hydrate.

35 Gas hydrate is a solid-phase, non-stoichiometric mixture of low molecular weight gas molecules occupying free spaces in a solid water lattice. Hydrates of several different gases occur abundantly in nature, but since the hydrate phase is only stable at high pressures and low temperatures it is only found naturally on Earth in the pore space of soils in either permafrost or sub-sea environments. CH_4 hydrate is of interest for its potential as a natural gas energy resource (Collett, 2000; Oyama et al., 2017; Singh et al., 2022), for its potential role in global carbon cycling as the climate changes (Ruppel and Kessler, 2017), and
40 for its role as a geohazard (Zander et al., 2018; Kaminski et al., 2020). On the other hand, CO_2 hydrate, which can form at similar pressures and temperatures as CH_4 hydrate, is increasingly being explored as a potential means to permanently sequester CO_2 as a climate change mitigation solution that comes with additional safety factors beyond those typically encountered in more conventional CO_2 sequestration scenarios. These include the fact that CO_2 hydrate is immobile in sediment pore space, which adds a thermodynamic barrier to gas escape, and the ability of the ocean to dissolve leaked CO_2
45 in marine environments which isolates CO_2 from release into the atmosphere (Tohidi et al., 2010). An added benefit of CO_2 injection into the CH_4 hydrate stability zone is that CO_2 hydrate can be more thermodynamically favourable than CH_4 hydrate, meaning it could theoretically be possible to use CO_2 to replace CH_4 in the hydrate phase, thus sequestering CO_2 while producing natural gas from a CH_4 hydrate deposit (Koh et al., 2016).

50 CO_2 sequestration in gas hydrate form can only occur in a finite bounded temperature and pressure range. In terms of a soil column, there exists a depth-bounded gas hydrate stability zone (GHSZ) in the subsurface in which CO_2 hydrate can form. In a marine environment, the CO_2 GHSZ typically begins several meters above the seafloor, but hydrate does not form freely in the water column (except for, e.g., as a gas bubble crust [Fu et al., 2021]) because the guest molecule gas (e.g., CO_2) typically cannot become concentrated enough in the water to do so. Therefore, the seafloor is typically the shallowest extent of hydrate
55 formation in marine systems. Working downward through the GHSZ, pressure increases roughly hydrostatically and temperature increases along a geothermal gradient. Increases in pressure stabilize hydrate, while increases in temperature destabilize hydrate. The geothermal temperature change effect on hydrate stability outweighs the hydrostatic pressure change effect, so there exists a depth below the seafloor where the temperature is too high to form hydrate, known as the base of the gas hydrate stability zone (BHSZ). Overall, the specific thickness of the bulk GHSZ is dependent on pressure, temperature,
60 porewater salinity and gas composition (Sloan and Koh, 2007).

For a potential host reservoir within the CO_2 GHSZ, evaluating the long-term CO_2 storage potential of the reservoir would require considering the thermodynamic trapping mechanism of solid gas hydrate formation in addition to traditional trapping

mechanisms. Several experimental studies have demonstrated the process of CO_2 trapping and hydrate conversion in the CO_2 GHSZ at the lab-scale, showing how conversion of CO_2 into a solid phase adds an additional safety factor (Gauteplass et al., 2020; Rehman et al., 2021). An experimental study of layered sediments using different injection strategies demonstrated the need to consider thermal management when designing a CO_2 injection in the GHSZ and suggested multilateral perforated horizontal wells may achieve the most optimal CO_2 conversion efficiency (Pang et al., 2024). However, reservoir-scale modelling studies of the transport and thermodynamic phenomena associated with injection of CO_2 in commercial volumes into the GHSZ are lacking.

We present several new capabilities developed in the open source, massively parallel multiphase flow and reactive transport simulator PFLOTRAN (Hammond et al., 2014) to model reservoir-scale injection of CO_2 in the CO_2 GHSZ. We have extended PFLOTRAN's HYDRATE mode capabilities to model free-phase CO_2 flow properties and CO_2 hydrate phase behaviour. Additionally, we introduce a fully coupled parallel well model that can be used to model CO_2 injection into heterogeneous media and can adapt to changes in flow properties associated with hydrate formation in the vicinity of the wellbore. Finally, we add a new fully coupled salt mass balance to consider salinity and salt precipitation effects in the GHSZ. We demonstrate these capabilities on a series of test problems designed to elucidate the challenges and opportunities associated with commercial-scale injection of CO_2 into the GHSZ.

2 Methods

PFLOTRAN's HYDRATE mode was originally developed to model CH_4 generation, transport, and structure 1 (SI) gas hydrate formation in deep marine and Arctic terrestrial reservoirs. PFLOTRAN's HYDRATE mode has been benchmarked against other reservoir simulators for modelling CH_4 gas production from hydrate reservoirs (White et al., 2020). It has been used to predict shallow gas generation and gas hydrate formation offshore the eastern U.S. (Eymold et al., 2021), to study relationships between gas generation and slope stability along the U.S. Atlantic margin (Carty and Daigle, 2022), and to model gas hydrate accumulation offshore Norway (Frederick et al., 2021). An extension of HYDRATE mode to include salinity coupling was developed to investigate viscous fingering and convective mixing in layered marine sediments during CH_4 hydrate formation over geologic time (Fukuyama et al., 2023). Here, we have redeveloped PFLOTRAN's HYDRATE mode to optionally consider CO_2 as the working gas; to couple fully implicitly with a new parallel well model; to include a new fully coupled salt mass balance; and to consider variable salinity effects on H_2O-CO_2-$NaCl$ mixtures and the CO_2 hydrate phase boundary.

2.1 Governing Equations

A system of three mass balance equations, one energy balance, and one well equation is now solved fully implicitly in PFLOTRAN's HYDRATE mode. The mass conservation equations take the following form:

$$95 \quad \frac{\partial}{\partial t} \phi \sum_{\alpha=\mathrm{l,g,h,i,s}} \left(s_\alpha \rho_\alpha x_j^\alpha\right) + \nabla \cdot \left(\boldsymbol{q}_\mathrm{l} \rho_\mathrm{l} x_j^\mathrm{l} + \boldsymbol{q}_\mathrm{g} \rho_\mathrm{g} x_j^\mathrm{g} - \phi s_\mathrm{l} D_\mathrm{l} \rho_\mathrm{l} \nabla x_j^\mathrm{l} - \phi s_\mathrm{g} D_\mathrm{g} \rho_\mathrm{g} \nabla x_j^\mathrm{g}\right) = Q_j + Q_{\mathrm{w},j} , \tag{1}$$

where phase α can be liquid (l), gas (g), hydrate (h), ice (i), or salt precipitate (s); component j includes water, gas (CO_2, CH_4, or air), and salt (NaCl); s_α is the saturation of phase α; ρ_α is the density of phase α; x_j^α is the mole fraction of component j in phase α; q_l is the liquid Darcy flux vector; q_g is the gas Darcy flux vector; D_l is the liquid phase diffusivity; D_g is the gas phase diffusivity; ϕ is the porosity; and Q_j includes any non-well sources/sinks of component j; and $Q_{\mathrm{w},j}$ is a source/sink of component j from a well. Solid phases are considered immobile and include the hydrate, ice, and salt precipitate phases. Mole fractions of components in the solid phases are fixed: by the hydration number in the hydrate phase, as pure water in the ice phase, and as pure salt in the salt precipitate phase. Formation of gas hydrate and ice therefore results in salt exclusion and aqueous dissolved salinity enhancement, which affects the hydrate phase boundary and gas solubility in the brine.

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The energy conservation equation takes the form:

$$\sum_{\alpha=\mathrm{l,g,h,i,s}} \left(\frac{\partial}{\partial t}\left(\phi s_\alpha \rho_\alpha U_\alpha\right) + \nabla \cdot \left(\boldsymbol{q}_\alpha \rho_\alpha H_\alpha\right)\right) + \frac{\partial}{\partial t}\left((1-\phi)\rho_\mathrm{r} C_p T\right) - \nabla \cdot (\kappa \nabla \mathrm{T}) = Q_e + Q_{\mathrm{w},e} , \tag{2}$$

where U_α is the internal energy of phase α, H_α is the enthalpy of phase α, ρ_r is the rock density, C_p is the heat capacity of the rock, κ is the composite thermal conductivity of the medium, T is the temperature, Q_e includes any non-well heat sources/sinks, and $Q_{\mathrm{w},e}$ is a heat source/sink imposed by the well (e.g., a heater in addition to a fluid injection). Exothermic hydrate formation (and vice versa, i.e., endothermic hydrate dissociation) is captured here by a decrease in internal energy of the hydrate phase during formation; this typically results in either an increase in system temperature or a change in phase saturations in three-phase systems. As we will show later, this phenomenon is important during CO_2 injection in the short term and it can continue to buffer conversion between phases for hundreds of years; similar effects have been shown for natural CH_4 hydrate systems where the base of the gas hydrate stability zone is shifted due to climactic changes (Owulunmi et al., 2022).

2.2 Well Model

120 A fully implicit, parallel well model has also been incorporated into HYDRATE mode. A well model can more accurately represent the insertion of a (comparably) small cylindrical wellbore into a reservoir grid cell than a standard source/sink term. Given a prescribed surface injection rate of CO_2 into the well, the well model solves for pressure variation along a wellbore and dynamically adjusts flow rates into the reservoir in response to changes in reservoir physical properties like permeability. This phenomenon can be critical to capture in a horizontal well or injection into a heterogeneous reservoir in the gas hydrate stability zone, where near-wellbore formation (or dissociation) of gas hydrate can significantly alter reservoir permeability and

thus injection behaviour. The well model developed for HYDRATE mode is a hydrostatic well model based off the design of White et al. (2013) but with key modifications including full parallelization to run flexibly on very large, unstructured grids and the addition of a thermal component; as we show here, injection temperature could be one of the most important design considerations for CO_2 storage in the gas hydrate stability zone. The well model developed here accounts for the enthalpy of the injected CO_2 at the prescribed temperature and wellbore pressure using the same equation of state (EOS) as the reservoir.

Solving a hydrostatic well model involves solving one extra conservation equation per well in addition to the reservoir mass and energy conservation equations. PFLOTRAN uses a fully implicit Newton-Raphson nonlinear solution search method; for each distinct well, only one extra row and one extra column are added to the Jacobian. For each reservoir cell intersected by a well, well pressure is computed at the centroid of the well section crossing through the reservoir. All well pressures are determined from the bottomhole pressure, a primary variable (see Section 2.4). The well model conservation equation is compact and reads as follows:

$$\sum_i Q_{w,j}^i = q_{w,j} , \tag{3}$$

where i is the discrete reservoir cell index through which the wellbore passes, $Q_{w,j}^i$ are the reservoir source/sink terms of phase j associated with a well in reservoir grid cell i, and $q_{w,j}$ is the prescribed surface injection rate of phase j.

Well flux at each reservoir grid cell is computed as a function of the pressure difference between the well and the reservoir cell (free phase CO_2 pressure for gas injection) scaled by the well index as follows:

$$Q_{w,j} = -\frac{WI\rho_j}{\mu_j}\left(P_w - (P_r + \rho_j\boldsymbol{g}\Delta z_{w-r})\right) \tag{4}$$

where P_w is the well node pressure, P_r is the reservoir pressure of phase j in the grid cell associated with a given well node, and Δz_{w-r} is the vertical distance between well node centre and reservoir cell centre. The well can be oriented in any direction in 3D. To account for permeability anisotropy and arbitrary well orientations, the well index, WI, is calculated using a 3D extension of the Peaceman equation (White et al., 2013):

$$WI = \sqrt{WI_x^2 + WI_y^2 + WI_z^2} \tag{5}$$

$$WI_x = \frac{2\pi\sqrt{k_y k_z}L_x}{\ln\left(\frac{r_{0,x}}{r_w}\right)+s} \; ; \; WI_y = \frac{2\pi\sqrt{k_x k_z}L_y}{\ln\left(\frac{r_{0,y}}{r_w}\right)+s} \; ; \; WI_z = \frac{2\pi\sqrt{k_x k_y}L_z}{\ln\left(\frac{r_{0,z}}{r_w}\right)+s} \tag{6}$$

$$r_{0,x} = 0.28 \frac{\left(\left(\frac{k_y}{k_z}\right)^{0.5}\Delta z^2 + \left(\frac{k_z}{k_y}\right)^{0.5}\Delta y^2\right)^{0.5}}{\left(\frac{k_y}{k_z}\right)^{0.25} + \left(\frac{k_z}{k_y}\right)^{0.25}}; \; r_{0,y} = 0.28 \frac{\left(\left(\frac{k_x}{k_z}\right)^{0.5}\Delta z^2 + \left(\frac{k_z}{k_x}\right)^{0.5}\Delta x^2\right)^{0.5}}{\left(\frac{k_x}{k_z}\right)^{0.25} + \left(\frac{k_z}{k_x}\right)^{0.25}}; \; r_{0,z} = 0.28 \frac{\left(\left(\frac{k_x}{k_y}\right)^{0.5}\Delta y^2 + \left(\frac{k_y}{k_x}\right)^{0.5}\Delta x^2\right)^{0.5}}{\left(\frac{k_x}{k_y}\right)^{0.25} + \left(\frac{k_y}{k_x}\right)^{0.25}} \tag{7}$$

where s is the well skin, r_w is the wellbore radius, WI_α are well indices, k_α are the absolute permeabilities of the medium, L_α are the projections of each well leg, and $r_{0,\alpha}$ are the Peaceman radii in each principal direction $\alpha \in \{x, y, z\}$.

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The bottomhole pressure of the well is solved as a primary variable as part of a fully coupled system of equations (Section 2.4). Pressures of each well segment are then directly calculated as functions of the bottomhole pressure by a hydrostatic adjustment to the depth of each well segment centre, where the density of the gas phase is updated as a function of pressure and injection temperature.

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2.3 Constitutive Relationships

Diffusive flux is modelled using Fick's Law with diffusivities computed as functions of temperature and salinity for CO_2 (Cadogan et al., 2014) and NaCl (Reid et al., 1987). Advective fluxes of mobile phases are computed by employing a two-phase Darcy's Equation:

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$$q_\alpha = -\frac{k k_\alpha^r}{\mu_\alpha} \nabla (P_\alpha - \rho_\alpha \boldsymbol{g} z) \tag{8}$$

where k is the intrinsic permeability of the medium, k_α^r is the relative permeability of phase α, μ_α is the viscosity of phase α, P_α is the pressure of phase α, $\boldsymbol{g}$ is the gravity vector, and z is depth. Relative permeability is computed as a function of phase saturations according to one of a suite of standard relative permeability relationships available in PFLOTRAN. Phase densities and viscosities are computed as functions of temperature, pressure, and salinity according to several options in PFLOTRAN (e.g., the Span-Wagner equation of state for CO_2, IF97 equation of state for water). Salt (NaCl) mass can be distributed in either the aqueous or solid salt precipitate phases, and aqueous dissolved salt mass affects brine flow properties like density (Haas, 1976), viscosity (Phillips et a., 1981), enthalpy, and diffusivity (Cadogan et al., 2014; Belgodere et al., 2015).

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Gas phase pressure and liquid phase pressure are related as a function of gas phase saturation through a choice of capillary pressure functions available in PFLOTRAN. When the gas hydrate phase is present, a capillary pressure associated with the hydrate phase is computed using the same capillary pressure function as the gas phase. This capillary pressure is used in the Gibbs-Thomson equation vis-à-vis the Young-Laplace equation to determine the hydrate 3-phase equilibrium temperature depression required to precipitate hydrate in pores as follows:

$$\Delta T_{\mathrm{m}} = -\frac{T_{\mathrm{mb}} P_{\mathrm{c}}}{\Delta H_{\mathrm{m}} \rho_{\mathrm{h}}} \tag{9}$$

where ΔT_{m} is the change in the hydrate melting temperature, P_{c} is the hydrate phase capillary pressure, T_{mb} is the bulk melting temperature, ΔH_{m} is the specific enthalpy of the phase transition, and ρ_{h} is the density of solid hydrate. A similar method is often used to compute ice freezing temperature depression vis-à-vis the Clausius-Clapeyron equation. This effect is typically only significant in fine-grained sediments and/or at very high effective hydrate phase saturations (Anderson et al., 2003).

When both gas hydrate and free gas occupy significant fractions of the pore space, as would be likely during CO_2 injection, their combined presence in the pore system should be accounted for through an effective saturation that is passed to the capillary pressure function. At three-phase (aqueous, free gas/CO_2 phase, gas hydrate) equilibrium, the chemical potential of CO_2 in the gas hydrate phase at a given hydrate capillary pressure must equal that of CO_2 in the free gas phase at a different free gas capillary pressure and dissolved CO_2. At bulk thermodynamic equilibrium, free gas and gas hydrate are stable together at a single pressure and temperature. In porous media, capillary effects on both the hydrate phase and gas phase lead to a window of possible pressures and temperatures over which three-phase equilibrium can be maintained (Clennell et al., 1999). To incorporate this effect and maintain thermodynamic reversibility, we adopt the approach of Liu and Flemings (2011) and require free gas and gas hydrate to partition the large pore space equally when both are present (Nole et al., 2018). This partitioning scheme results in the following effective saturations of free gas and gas hydrate:

$$s_{\alpha,\mathrm{eff}} = \begin{array}{ll} 2s_\alpha\,, & s_\alpha < s_\beta \\ s_\alpha + s_\beta\,, & \mathrm{otherwise} \end{array} \tag{10}$$

where $s_{\alpha,\mathrm{eff}}$ is the effective saturation of nonwetting phase α and β is the other nonwetting phase in a 3-phase system where liquid water is the wetting phase.

The presence of gas hydrate in the pore space of a reservoir decreases the reservoir's permeability below its intrinsic (water-saturated) permeability. We model permeability reduction as a function of hydrate saturation as follows (Dai and Seol, 2014):

$$k_{\mathrm{eff}} = \frac{(1-s_{\mathrm{h}})^3}{(1+2s_{\mathrm{h}})^2} \tag{11}$$

where k_{eff} is the effective permeability coefficient and is multiplied by intrinsic permeability to compute the effective absolute permeability.

Heat transfer occurs through mobile fluid phase flow, phase transitions, thermal conduction, and injection/production. Fluid phase enthalpies are computed using corresponding equations of state. For the CO_2 phase, the Span-Wagner equation of state is recommended (Span and Wagner, 1996), and for water the IF97 equation of state with salinity extensions are available in PFLOTRAN. Enthalpies of the solid gas hydrate (Handa, 1998) and salt (Lide and Kehiaian, 2020) phases are computed as functions of temperature. Several options for composite thermal conductivity can be used; the default thermal conductivity function is a linear scaling function of phase saturations:

$$\kappa = \kappa_{dry} + \phi \sum_{\alpha=l,g,h,i,s} s_\alpha \kappa_\alpha \tag{12}$$

where κ_{dry} is the dry rock thermal conductivity and κ_α is the thermal conductivity of phase α.

The presence of salt has several impacts on system behaviour. Solid salt precipitation can occur due to near-wellbore dry-out (water evaporating into the free CO_2 phase) during the injection period or due to rapid hydrate formation, whereby water and CO_2 forming a hydrate excludes salt and thus increases dissolved salt concentrations. Salt precipitation reduces permeability, which affects CO_2 injectivity, gas flow, and liquid imbibition. This occurs either at the injection site if enough CO_2 is injected to dry out the water or far into the future when free phase CO_2 has undergone conversion to very high hydrate saturations (though at this point, permeability reduction due to salt precipitation is dwarfed by the presence of solid hydrate at high saturations). Aqueous dissolved salt concentration affects the density of the aqueous phase; the presence of gradients in salt concentration drives convective mixing. Salt exclusion during hydrate formation locally increases salt concentrations, which can produce this phenomenon (Fukuyama et al., 2023). Dissolved salt also affects gas solubility and shifts the three-phase equilibrium pressure of gas hydrate.

CO_2 equilibrium phase partitioning is computed using the method of Spycher and Pruess (2010); the CO_2 hydrate – free phase CO_2 – water three-phase equilibrium curve is determined from a polynomial fit of data from Men et al. (2022) up to 283.2K and then a linear function at higher temperatures (Figure 1). This 3-phase equilibrium boundary is consistent with the CO_2 hydrate phase boundary used in STOMP-HYD up to 100 MPa (McGrail et al., 2007) and is similar to the phase boundary proposed by Garapati et al. (2011) with the exception that here the phase boundary is monotonically increasing.
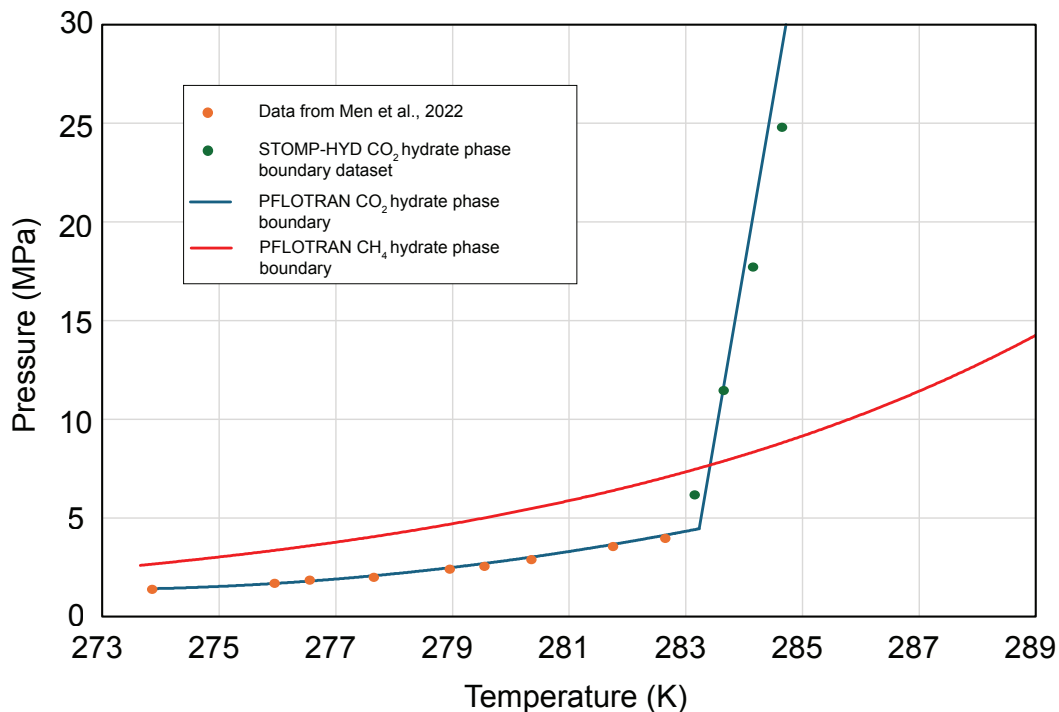
Figure 1: Three-phase equilibrium pressure of CO₂ hydrate. At temperatures below 283.2 K, a curve fit of experimental data is used. Above 283.2 K, a line is used that is consistent with STOMP-HYD. The CH₄ hydrate phase boundary is plotted for reference (Moridis, 2003)

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2.4 Phase States and Primary Variables

PFLOTRAN's HYDRATE mode solves mass conservation, energy conservation, and well flux conservation equations for a set of three components (CH_4/CO_2/air, H_2O, NaCl) over five phases (aqueous, gas component-rich/gas, gas hydrate, ice, salt

255   precipitate). This results in solving a set of four partial differential equations for all cells in the domain plus one coupled well equation per cell containing the bottom segment of a well. Therefore, PFLOTRAN's fully implicit solution solves for four primary variables everywhere plus one extra primary variable per well in the domain.

The reservoir (non-well) equations use primary variable switching depending on the thermodynamic state of a grid cell.

260   HYDRATE mode contains 13 phase states with four primary variables per phase state (Table 1). For example, cells in the fully liquid (aqueous) saturated state solve for liquid pressure, dissolved gas mass fraction, temperature, and total salt mass per unit liquid mass as primary variables. Secondary variables like phase densities, viscosities, and enthalpies are computed at equilibrium from the primary variables through use of various equations of state. Precipitated salt saturation is computed by

determining whether bulk salt concentration (total salt mass per mass of liquid phase) exceeds dissolved salt solubility and converting the excess salt mass into a solid phase (permeability updates according to Verma and Pruess, 1988). If dissolved gas mass fraction exceeds solubility and aqueous pressure, temperature, and dissolved salt mass fraction lie within the GHSZ, the cell will transition into the hydrate-aqueous state and primary variables will update accordingly. Upon entering the hydrate-aqueous state, PFLOTRAN then switches primary variables and solves for gas pressure, hydrate saturation, temperature, and salt concentration.

Table 1: Phase states and primary variable combinations in PFLOTRAN's HYDRATE mode

Phase State	Primary Variables	Phase State	Primary Variables
L (aqueous)	P_l, x_l^g, T, m_l^s	AI (aqueous-ice)	P_l, x_l^g, S_l, m_l^s
G (gas component-rich)	P_g, P_a, T, m_b^s	HGA (hydrate-gas-aqueous)	S_l, S_h, T, m_l^s
GA (two-phase gas-aqueous)	P_g, S_g, T, m_l^s	HAI (hydrate-aqueous-ice)	P_g, S_l, S_i, m_l^s
HG (hydrate-gas)	P_g, P_a, T, m_b^s	HGI (hydrate-gas-ice)	S_i, S_h, T, m_b^s
HA (hydrate-aqueous)	P_g, S_h, T, m_l^s	GAI (gas-aqueous-ice)	P_g, S_g, T, m_l^s
HI (hydrate-ice)	P_g, S_h, T, m_b^s	HGAI (hydrate-gas-aqueous-ice)	S_l, S_g, S_i, m_l^s
GI (gas-ice)	P_g, S_i, T, m_b^s		

P_l = liquid pressure, P_g = gas pressure, P_a = gas-rich gas (air) component partial pressure, x_l^g = aqueous dissolved gas mass fraction, T = temperature, m_l^s = salt mass fraction per unit aqueous mass, m_b^s = total salt mass per unit bulk volume, S_g = gas saturation, S_h = hydrate saturation, S_l = liquid saturation, S_i = ice saturation

For the well equation, the bottomhole pressure of the well is solved as a primary variable. Given a user-defined well flow rate, each well's bottomhole pressure is solved fully implicitly as part of the full reservoir flow solution. At a given bottomhole pressure, the well model solves for all other pressures in the well by working upward and iteratively solving hydrostatic pressure of the injection fluid where variable fluid density in the well is updated as a function of pressure and temperature. Once the hydrostatic pressures are determined, mass and energy fluxes between well and reservoir at the well segment centroids are computed by adding a hydrostatic adjustment to reservoir pressures to align reservoir pressures with well segment centroids and then employing Equation 4. Fluxes between coupled wells and reservoir cells are therefore functions of the well primary variable (bottomhole pressure) and reservoir cell primary variables (e.g., gas/liquid pressure).

Using fully implicit coupling and a Newton-Raphson solution search method, insertion of a coupled well into the domain adds additional fill to the Jacobian matrix used to compute solution updates. Critically, the well model adds extra connectivity beyond the typical stencil for two-point flux calculations. Thus, the structure of the Jacobian matrix is altered by introducing wells. This alteration is typically minimal, but it would likely become more computationally demanding as the number of wells in the domain is increased or as the number of screened segments per well is increased. The number of wells, the extent to

290   which each well increases the fill of the Jacobian, and the strength of the coupling between well and reservoir are all likely to affect the overall performance of simulations using the coupled well model. PFLOTRAN uses neighbour cell ghosting to parallelize computations: for a given subset of grid cells in the model domain that are owned by a particular process, say, Process "N", all of the off-process grid cells that border these cells are "ghosted" onto Process N, meaning that copies of state variables for those cells are kept up to date for computing updated flux terms at processor boundaries. Cell ghosting is

295   determined by the numerical stencil and the distribution of cells on processors in parallel. Since adding a well introduces connectivity beyond the original stencil, PFLOTRAN updates the ghosting stencil to include all off-process reservoir cells connected by a given well, allowing for full incorporation of well terms in the Jacobian. When a well is turned off (or well flow rates are set to 0), that well equation is not solved.

3 Results

300   We demonstrate our developments by applying the software to a set of two hypothetical CO_2 injection scenarios in marine environments within the GHSZ. In the first example, liquid CO_2 is injected slowly into a simple 1D homogeneous sediment column through a partially screened well beneath the GHSZ. The second example simulates commercial-scale injection of supercritical CO_2 into a 2D radial domain with heterogeneous layering, where a well is screened within and beneath the GHSZ. The CO_2 plume in this model is tracked for 10,000 years as it transitions from a supercritical phase to a dense liquid phase and

305   then into the gas hydrate phase.

3.1 1D Liquid CO_2 Injection into a Homogeneous Reservoir

In this scenario, a relatively slow injection is designed to illustrate the multiphase and thermodynamic processes associated with injecting CO_2 into the GHSZ. A 1D, 1,000-m homogeneous domain is initialized to hydrostatic conditions where the top

310   of the domain is held at seafloor pressure of 10 MPa, bottom water temperature of 5°C, and geothermal gradient of 20° C/km. The top Dirichlet boundary condition is set to the initial bottom water temperature, pressure, zero gas mass fraction, and constant salinity (0.035 kg/kg). The bottom boundary at 1,000 mbsf is a Neumann zero flux boundary. The domain is discretized into 200 grid cells in the vertical dimension; individual grid cells measure 25 m in the horizontal x-dimension by 1 m in the horizontal y-dimension by 5 m in the vertical z-dimension. A well extends to 500 meters below seafloor (mbsf) and

315   is screened for 25 m from 475 mbsf to 500 mbsf. Dense liquid phase CO_2 is injected at 20° C at an injection rate of 15,000 kg/yr for 50 years. Pressure of the injection varies along the well depending on the bottomhole pressure, but CO_2 remains in the liquid phase for the entirety of the injection. Use of the well model will lead to variable CO_2 injection rates in each of the screened well segments and correspondingly variable injection enthalpy as a function of well segment pressure; this effect is less noticeable in this homogeneous case than in the heterogeneous case. This simulation is run for 200 years. While it is likely

320   that bottom water temperature would fluctuate over 200 years, bottom water temperature variation is site specific; PFLOTRAN

can apply time-dependent bottom water temperature as a boundary condition, but for this study a constant temperature was used for simplicity. A constant reservoir porosity is set to 0.35, and constant isotropic permeability is set to 1×10^{-13} m^2. A Van Genuchten capillary pressure function is used, where Van Genuchten m is set to 0.5, $\alpha = 1 \times 10^{-4}$ Pa^{-1}, and $S_{lr} = 0.05$. Corey relative permeability functions are used, where $m = 0.5$, $S_{lr} = 0.3$ and $S_{gr} = 0.05$.

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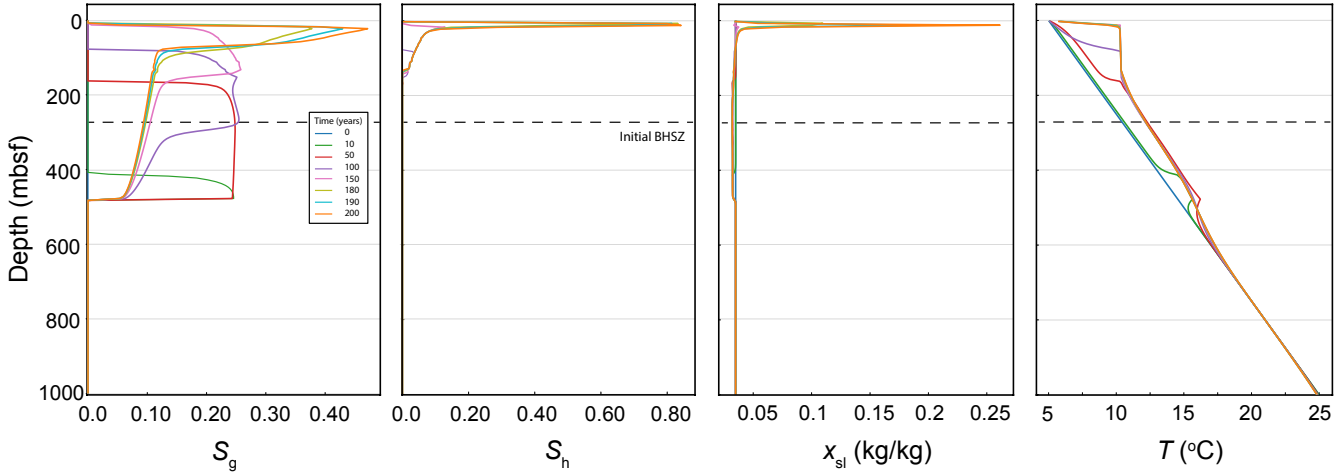
Figure 2: Gas saturation (S_g), hydrate saturation (S_h), dissolved salt mass fraction (x_{sl}), and temperature (T) during and after a CO$_2$ injection beneath the GHSZ. Dashed line indicates the initial (pre-injection) BHSZ, though this is perturbed by upward migration
330   **of warm gas.**

During the 50 years of injection, the injected CO$_2$ remains beneath the base of the CO$_2$ hydrate stability zone (Figure 2). Therefore, it can only exist as free-phase CO$_2$. CO$_2$ migrates upward in characteristic fashion: buoyancy and pressure forces drive gas (free-phase CO$_2$) saturations exceeding the residual gas saturation to migrate upward in the sediment column. By
335   100 years, the free-phase CO$_2$ front has reached the base of the GHSZ. At this point, some free-phase CO$_2$ converts to CO$_2$ hydrate. Exothermic hydrate formation keeps the reservoir temperature at the three-phase equilibrium temperature while free-phase CO$_2$ and hydrate coexist. As the gas plume migrates upward over time, more gas converts into gas hydrate. While the gas supply is strong and hydrate is still forming, the temperature of the reservoir is pushed well above the background (initial) geothermal temperature. The CO$_2$ hydrate 3-phase equilibrium temperature is approximately 10° C at this depth, which is why
340   the temperature throughout the three-phase zone remains at roughly 10° C.

After the injection period ends, hydrate accumulates in high saturations near the top of the domain due to the bottom water temperature and pressure being fixed; the resulting permeability reduction causes gas to pool and salt to concentrate where hydrate is forming, leading to salt precipitation. This kind of scenario is unlikely to occur in a more realistic reservoir in 2D

345 and 3D where permeability reduction would cause gas to migrate laterally and therefore not cause such significant pooling effects. But this model illustrates how exothermic hydrate formation can lead to a thick (albeit transient) three-phase stability zone throughout which free-phase CO_2 can remain mobile. Therefore, the fact that CO_2 hydrate is stable in the upper hundred meters of sediments at a particular hydrostatic pressure and geothermal temperature is an insufficient condition for ensuring CO_2 trapping by hydrate formation. Permeability reduction associated with gas hydrate formation can act to slow free-phase 350 CO_2 migration, but, at least at early time, a combination of thermodynamic and other structural trapping mechanisms is likely necessary to ensure the long-term sequestration of most of the injected CO_2 in the subsurface GHSZ.

3.2 Supercritical CO_2 Injection into a 2D Heterogeneous Reservoir

In this scenario, a commercial-scale CO_2 injection is modelled under more realistic reservoir and injection conditions (Figure 3). A 2D, heterogeneous cylindrical domain extends from the seafloor down to 600 mbsf with a radius of 3.765 km. The model 355 domain consists of 40 grid cells in the horizontal dimension increasing in thickness from 7.38 m at the model centre to 364.36 m at the far edge. The model contains 55 cells in the vertical dimension with varying thickness, each corresponding to a different layer in the model. The model consists of interbedded sand and mud units as might be found within the marine GHSZ. High and low permeability layers alternate with synthetic heterogeneity; similarly, the model contains heterogeneous porosity and capillary entry pressure. All other physical properties are kept constant between layers. A Brooks-Corey capillary pressure 360 function along with Burdine relative permeability functions for liquid and gas phases are used for all layers. For all layers, Brooks-Corey λ =0.8311 and $S_{rl} = S_{rg} = 0.0597$. The capillary entry pressure (the inverse of which is expressed by the Brooks-Corey α parameter) varies between reservoir layers (Figure 4). Rock density is set to 2,650 kg/m^3, dry rock thermal conductivity is set at 2.0 W/m-C, and soil compressibility is modelled with a linear compressibility function using a soil compressibility of 1.0E-8 Pa^{-1} and a reference pressure of 10 MPa. Seafloor pressure is set to 10 MPa, bottom water temperature 365 is 5° C, and seafloor salinity is 0.035 kg/kg. The model is initialized at hydrostatic pressure, constant salinity, and a geothermal gradient of 20° C/km. The top and outer edges of the domain are kept at the initial conditions. The bottom boundary condition is set to no liquid or gas flux, constant salinity, and a constant heat flux to preserve the geothermal gradient.

The well used in this scenario extends from the seafloor to 300 mbsf and is cased for the first 100 m. The rest of the well is 370 screened; flow from well to reservoir is possible only in the screened interval. Given the conditions outlined above, the initial BHSZ for CO_2 hydrate is at approximately 250 mbsf. Therefore, the well in this scenario extends through the GHSZ and 50 m below the initial BHSZ. Care was taken to ensure that the required well pressures to achieve the specified CO_2 injection rate were realistic; the 100 m depth of well casing was chosen so that the well pressures resulting from our prescribed injection rate did not exceed the lithostatic pressure. This calculation is approximate and does not consider the fracture gradient; the fact 375 that well pressures can easily approach the lithostatic pressure in these settings means that reservoir integrity should be rigorously evaluated when performing site-specific evaluations of CO_2 injectivity in the shallow subsurface.
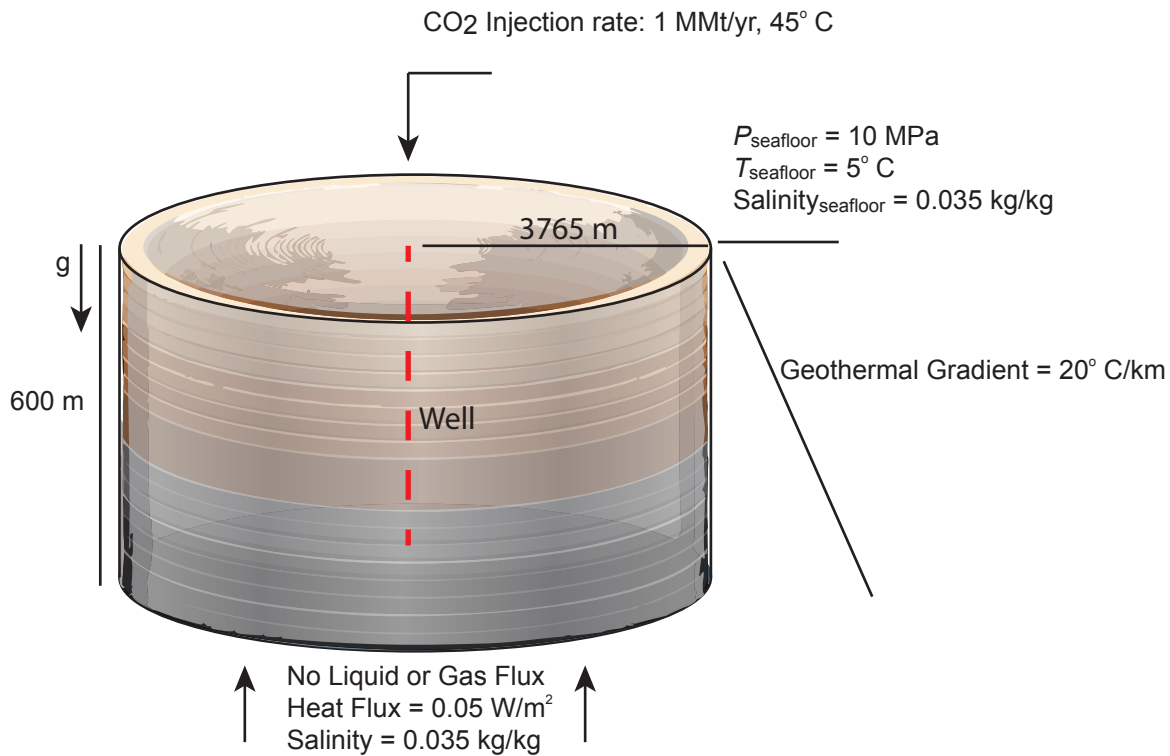
Figure 3: Illustration of a heterogeneous layered 2D cylindrical model, with specific model parameters labeled

380 The model is run for 10,000 years; CO_2 is injected continuously at a rate of 1 million metric tons (MMT) per year for 50 years to meet a storage target of 50 MMT of CO_2. The CO_2 is injected at a constant temperature of 45° C; injection pressure will vary along the wellbore depending on the hydrostatic pressure of the well, and this in turn will affect the enthalpy of the injected gas. After 50 years, the well is shut off; over time, warm supercritical CO_2 will cool into a dense liquid CO_2 phase and then eventually a gas hydrate phase.

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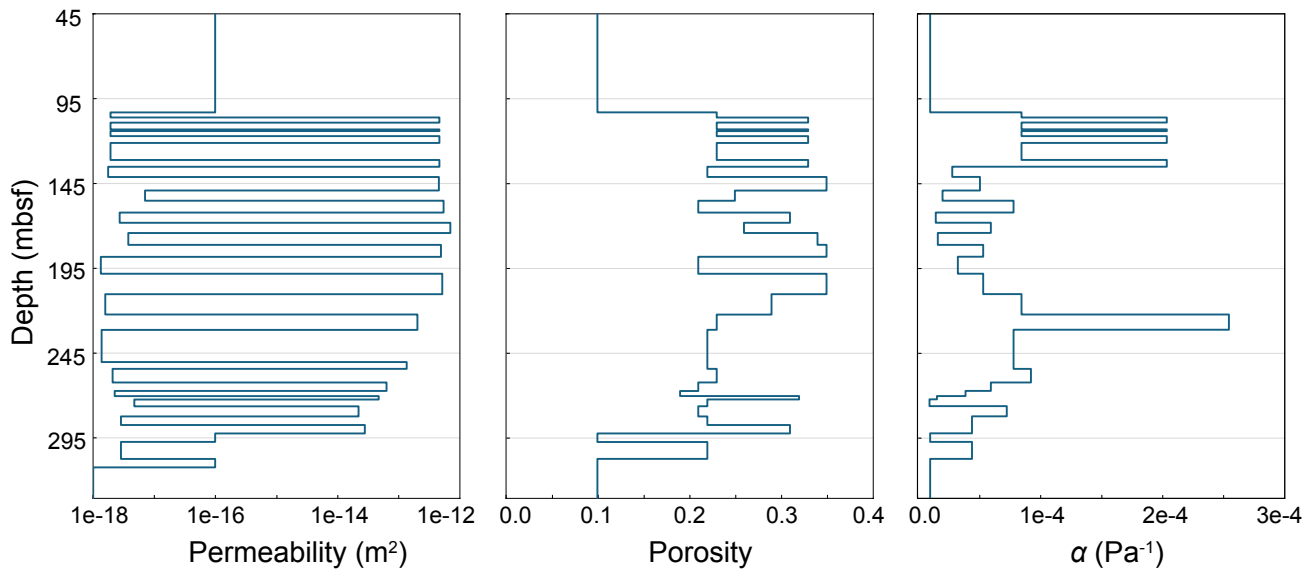
Figure 4: Depth-varying physical properties (permeability, porosity, and inverse capillary entry pressure [α]) of the heterogeneous 2D injection scenario near the injection interval. Physical properties held constant above and below these snapshots.

During the 50-year injection interval, gas flows predominantly into the high-permeability reservoir intervals (Figure 5, Figure 6). This is because the well model adjusts how the mass of injected CO_2 is distributed to each interval in the injection region (Figure 7) given the total gas injection rate, hydrostatic well pressure, reservoir pressure, and well index, where well index is a function of reservoir permeability. Since early hydrate formation in the reservoir units elicits a permeability and pressure response, the well flow rate into individual units evolves over time during the injection. In some units, well flow rate drops, and these drops are then compensated by increases in flow rates in other units. Likewise, the pressure in the well evolves over time in response to hydrate formation and relative permeability of the mobile fluids.

By the end of the injection period, gas has flowed preferentially in the radial direction along high permeability flow paths. On the outer edges of the gas plume, free phase CO_2 combines with water to form a gas hydrate phase where pressures and temperatures are within the gas hydrate stability zone. Since pure CO_2 is being injected through the well and since water is miscible in the CO_2 phase, high gas saturations in the near-wellbore cells cause salt concentrations in those cells to increase above salt solubility. This salting out effect results in small amounts of salt precipitate saturation in the pore space at the end of the injection.
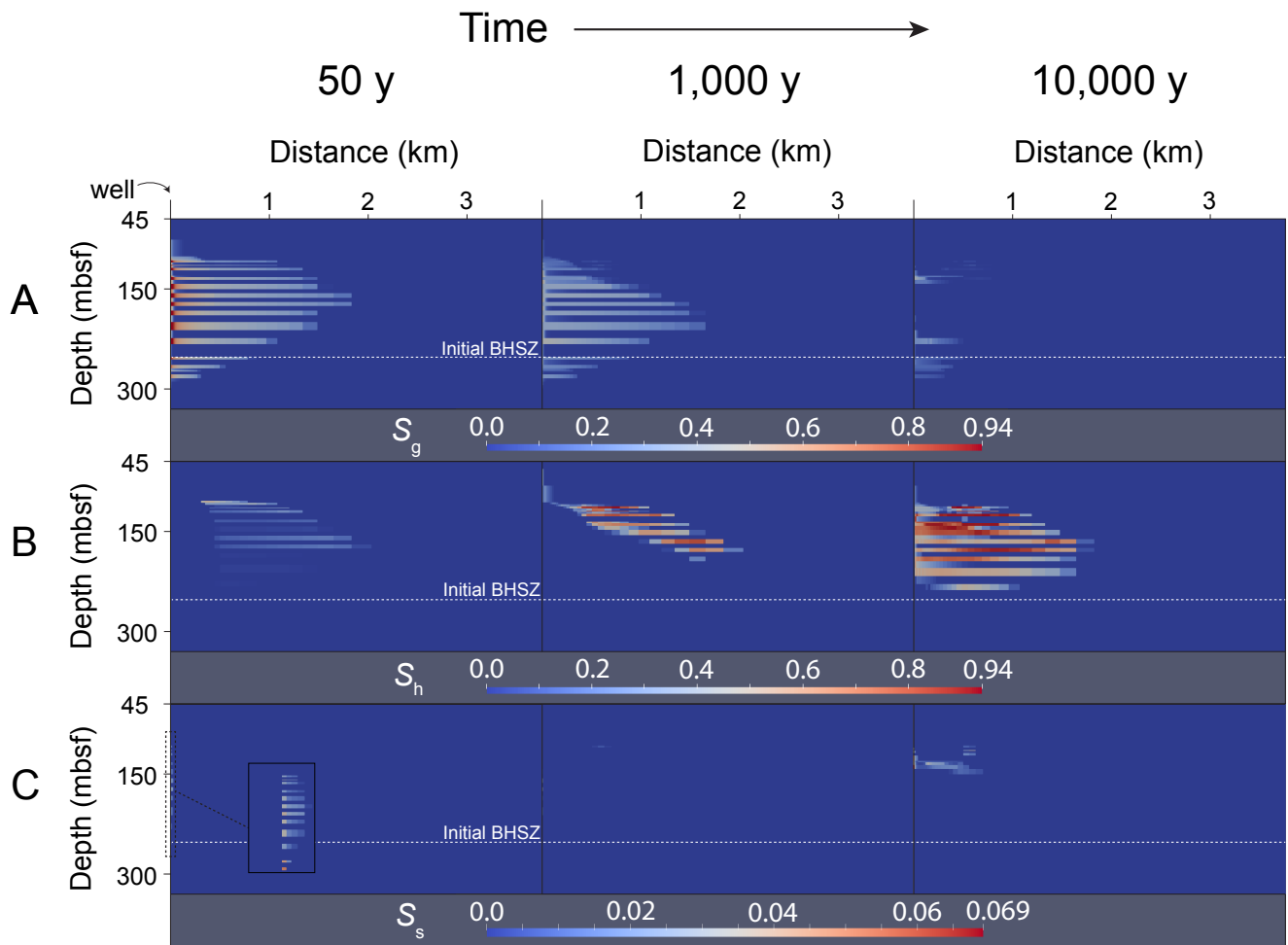
Figure 5: Snapshots of saturations over time in the vicinity of the wellbore. Depth is represented by the vertical axis, and radial distance from the well is represented by the horizontal axis. A) Gas (free-phase CO_2) saturation (S_g), B) hydrate saturation (S_h), and C) salt precipitate saturation (S_s) distribution at 50 years, 1,000 years, and 10,000 years of simulation time. A zoomed-in cut-out shows near-wellbore salt precipitate saturations at 50 years.

As water imbibes back into the near-wellbore cells between 50 and 1,000 years, gas saturations in those cells drop and salt re-dissolves (salt precipitate saturations near the wellbore drop toward 0). During this time, the temperature of the injected fluid is dropping toward the initial geothermal temperature. As this happens, free phase CO_2 combines with available water to form gas hydrate. Exothermic hydrate formation props up temperatures during hydrate formation and slows the process of CO_2 conversion into gas hydrate. In some areas at the upper edges of the CO_2 plume, where the system is furthest into the GHSZ, very high conversion of CO_2 to gas hydrate is achieved in a relatively short amount of time. Since hydrate formation only involves water and CO_2 components, salt exclusion during rapid hydrate formation results in local buildup of salt

concentrations. Some cells in the model associated with rapid hydrate formation therefore exhibit some solid salt precipitation by 1,000 years.

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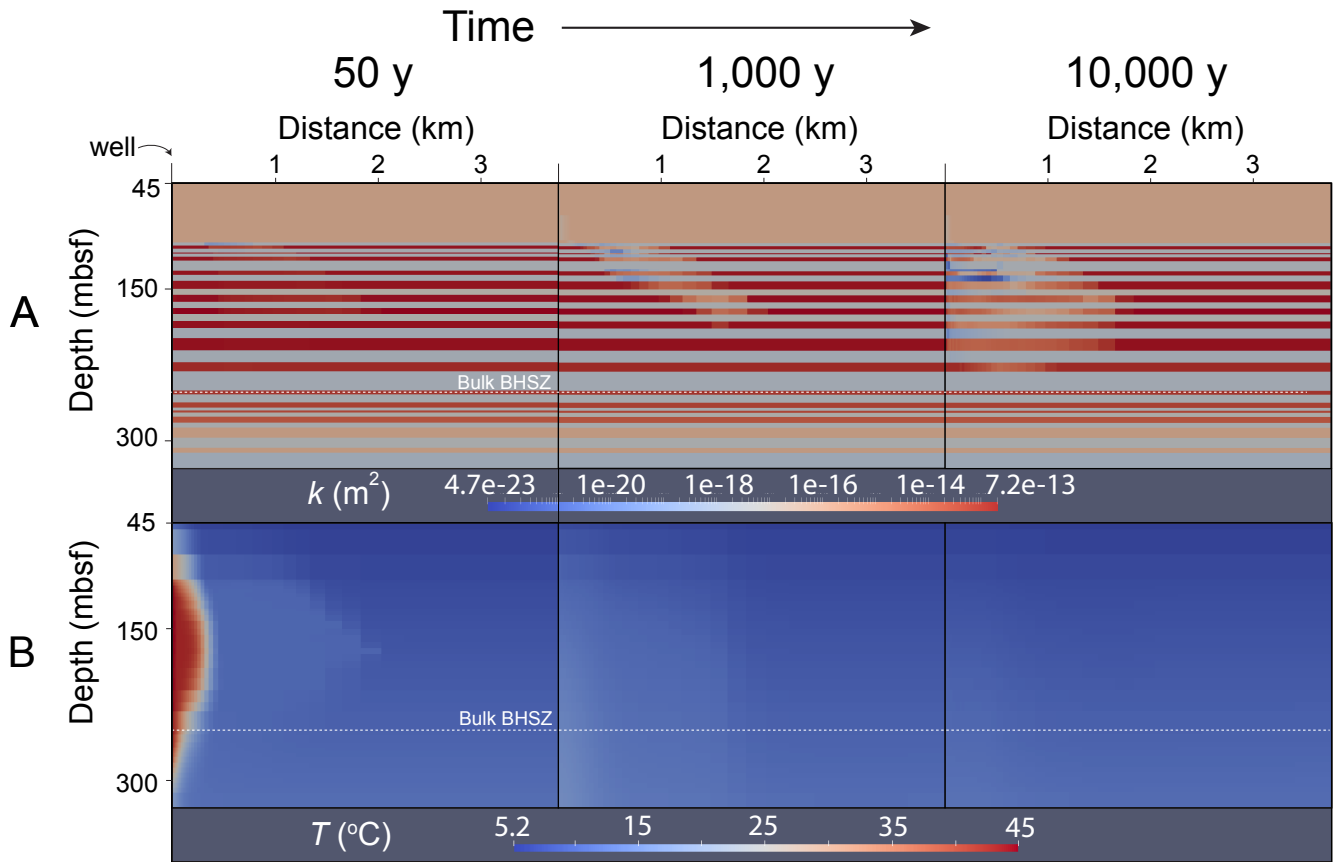
Figure 6: A) Permeability (k) and B) temperature (T) distribution at 50 years, 1,000 years, and 10,000 years of simulation time. Depth is represented by the vertical axis, and radial distance from the well is represented by the horizontal axis.

425   After 10,000 years, most of the injected CO_2 has converted into gas hydrate. High gas hydrate saturations have built up in the near-wellbore area since the initial temperature of the injection has decayed away toward the steady-state geothermal temperature profile. Hydrate formation has significantly decreased the permeability of the host reservoir, and gas has migrated into the other layers to form hydrate. A region of three-phase coexistence (liquid water, free phase CO_2, and gas hydrate) is still present after 10,000 years because of a combination of exothermic hydrate formation and slow imbibition rates of cool

430   liquid water due to significant permeability reduction.
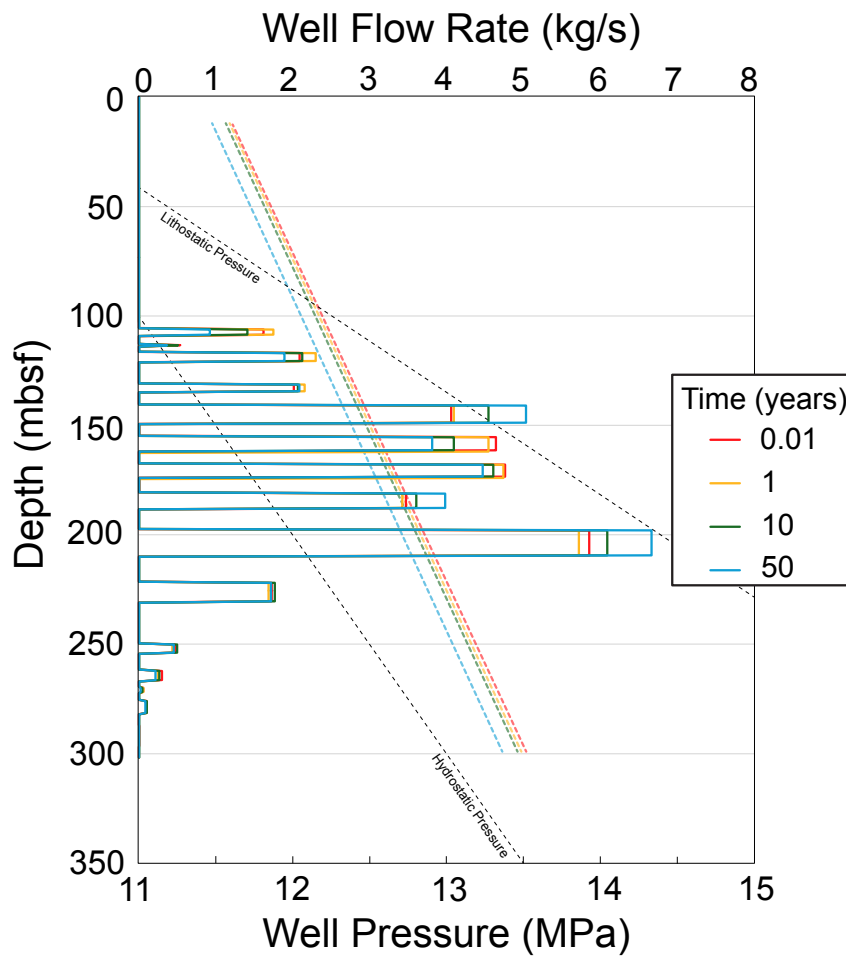
Figure 7: CO₂ mass flow rate in each well segment (solid lines) and supercritical CO₂ phase pressure in the well (dashed lines) during the 50-year injection period.

Discussion

The two models presented here were selected to demonstrate some of the key dynamic coupled processes associated with CO_2 injection into the gas hydrate stability zone. In the 1D homogeneous model, CO_2 is injected beneath the GHSZ and forms a free phase which migrates upward due to buoyancy and pressure forces. Once it enters the GHSZ, conversion of CO_2 into an immobile hydrate phase is limited by the rate at which heat can diffuse away and availability of water. When thermal conduction and liquid water flow are limited, the system can maintain 3-phase equilibrium temperature for decades (or thousands of years as is shown in the 2D model). This thermal buffering phenomenon has been observed in models of natural CH_4 hydrate formation and dissociation in marine sediments and can occur on geologic timescales depending on free gas phase flow rate or rate of environmental change (You and Flemings, 2018; Oluwunmi et al., 2022). Clearly, such a scenario would

445   not be ideal for permanent CO_2 sequestration as much of the CO_2 remains in a free phase and accumulates very close to the seafloor. Permeability reduction due to hydrate formation adds a physical trapping mechanism analogous to a low permeability sealing facies. The fact that this permeability reduction is the primary mechanism for preventing CO_2 flow to the surface in the 1D model suggests that physical/structural trapping should be considered just as important or more important than thermodynamic trapping when evaluating a reservoir within the GHSZ for long-term CO_2 storage.

450

The 2D cylindrical model was designed to incorporate more realistic reservoir physical properties and include an injection rate more viable for commercial-scale CO_2 storage in the GHSZ. In this scenario, CO_2 was injected into a layered reservoir that is bounded by low permeability facies that inhibit direct flow of CO_2 to the seafloor. Instead of injecting beneath the GHSZ, a high-temperature supercritical CO_2 phase is injected directly into and beneath the GHSZ. Near-wellbore gas hydrate formation

455   is prevented by the high temperature of the injection during the injection period. Hydrate formation does occur during the injection period at the edges of the CO_2 plume; the associated changes in fluid mobility and permeability alter the pressure in the well and cause well flow rates to fluctuate layer-by-layer. Therefore, even if the CO_2 injection temperature is designed to prevent near-wellbore hydrate formation, hydrate formation in the far-field should be considered when designing a CO_2 injection insofar as it could appreciably affect wellbore pressure. Salt precipitation can occur near the wellbore during injection

460   due to "salting out" effects of dry CO_2 injection. Salt can also precipitate later in time as CO_2 converts to hydrate faster than the pore water can freshen through either aqueous imbibition or salt diffusion. In either case, salt precipitate saturations appear to be minimal for the scenario modelled here, but salt precipitation could appreciably decrease permeability under a configuration with either more rapid CO_2 injection or more rapid conversion of CO_2 to hydrate. In some regions of this model, hydrate saturations become very high at late times and lower the permeability of host reservoir units by several orders of

465   magnitude. This makes for effective sealing of CO_2 by conversion to an immobile phase and by impeding flow of the free CO_2 phase. This phenomenon also has the effect of pushing gas into less intrinsically permeable layers and ultimately smoothing the distribution of gas hydrate throughout the model domain.

Potential leakage of CO_2 into the ocean at the seafloor is clearly a risk when injecting CO_2 into shallow marine reservoirs. In

470   the 2D model, advective leakage of CO_2 is prevented by the low permeability of the top facies, but CO_2 still migrates upward diffusively and thus there is some diffusive leakage of CO_2 at the seafloor. Beyond just considering the bulk permeability of the top facies, however, care should also be taken to inject CO_2 at pressures that remain below the fracture pressure of the overburden. While we do not model permeability evolution due to fracturing in this model, we designed our model scenarios to avoid fracturing (well model pressures in the injection interval were designed to remain below the lithostatic pressure). In a

475   site-specific scenario, the fracture gradient should be well characterized and incorporated into CO_2 injection designs.

Although the models presented here only consider injection of CO_2 into reservoirs without naturally occurring gas hydrate deposits, the presence of other hydrates (e.g., CH_4 hydrate) in a reservoir could affect the injectivity of CO_2 by altering the

physical properties (e.g., permeability) of the reservoir. The stability of all potential hydrates could be altered by injection of
480 CO_2, so site characterization should consider whether and what kinds of hydrates might exist in a potential subsea host reservoir. This information could be used to parameterize heterogeneous initial physical properties of a model, modify phase behaviour as a function of gas mixtures, or to develop a more sophisticated model of the interactions between multiple gases forming/dissolving multiple hydrates, but such models are beyond the scope of this work.

Conclusions

485 We present several new developments in PFLOTRAN's HYDRATE mode including a new option to model CO_2 as the working gas, a new salt mass balance for considering effects of salinity gradients and salt precipitation, and a new fully coupled hydrostatic well model. We demonstrate these new capabilities on two test problems designed to explore the coupled processes relevant to CO_2 injection into the marine gas hydrate stability zone for the purpose of permanently sequestering CO_2. CO_2 sequestration in the gas hydrate stability zone is a potentially promising technique for secure storage of CO_2 because of the
490 associated favourable conditions for converting injected CO_2 into solid gas hydrate form, which is immobile in the pore space. However, no reservoir modelling studies to date have demonstrated what commercial-scale CO_2 injection into the gas hydrate stability zone might look like. We show through a 1D homogeneous model that it is critical to consider multiple trapping mechanisms in addition to the thermodynamic trapping accompanied by conversion of CO_2 into hydrate form. We then expand to a 2D heterogeneous cylindrical model with a commercial-scale 1 MMT/yr CO_2 injection rate to underscore the interplay
495 between structural trapping, thermodynamics, and permeability alteration on the migration and conversion of CO_2. We demonstrate how our fully implicit well model adapts to changes in flow properties during CO_2 injection, and how injection of a warm supercritical CO_2 phase can facilitate near-wellbore injectivity but can lead to pressure change in the well. In the future, this capability could be used to more rigorously evaluate the potential for secure CO_2 storage in greater volumes, at larger (3D) scales, with more site-specific inputs, and with more exotic well designs including multiple wells or horizontal
500 wells.

Code and Data Availability

The software developments described here were released on August 23, 2024 with PFLOTRAN version 6.0 (www.pflotran.org). PFLOTRAN is open source and freely available under a GNU LGPL Version 3 license at https://bitbucket.org/pflotran/pflotran. Software inputs and a snapshot of the PFLOTRAN v6.0 Bitbucket repository are
505 available on Zenodo at https://zenodo.org/records/13619874. The files on Zenodo include PFLOTRAN input decks for both model scenarios and associated Span-Wagner EOS database files referenced by those input decks.

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Author Contribution:

MN: software development, model conceptualization, formal analysis, methodology, writing; JB: model conceptualization, writing; FN: model conceptualization, writing; GH: software development, review and editing.

Competing Interests: The authors declare that they have no conflict of interest.

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