

Abstract

1 Introduction

and Northeast United States and southern Ontario, Canada. The regional climate is highly

sensitive to the Great Lakes due to the lakes' vast thermal inertia, potential source of moisture

to the overlying atmosphere, and contrasts in heat, moisture, roughness, and albedo compared

to surrounding land (Changnon and Jones, 1972; Scott and Huff, 1996; Chuang and Sousounis,

2003; Notaro et al., 2013a; Briley et al., 2021; Wang et al., 2022). During late autumn through

winter, when cold, dry continental air masses from Canada pass over the relatively mild Great

Lakes, the air masses are destabilized and moistened, leading to enhanced cloud cover and

precipitation downwind of the lakes (Niziol et al., 1995; Ballentine et al., 1998; Kristovich and

Laird, 1998; Notaro et al., 2013b; Shi and Xue, 2019). During the broader unstable lake season,

which spans from September to March and is characterized by amplified lake-effect cloud

cover and precipitation due to lake surface temperatures typically exceeding overlying air

temperatures, lake-effect snowfall typically peaks during December-January, and lake ice cover

is most extensive during February-March (Assel, 1990; Niziol et al., 1995; Kristovich and

Laird, 1998; Lam and Schertzer, 1999; Notaro et al., 2013b). The establishment of extensive

lake ice cover usually by mid-late winter dampens over-lake turbulent fluxes of heat and

- moisture, subsequently reducing resulting lake-effect precipitation (Brown and Duguay, 2010; Notaro et al., 2021). Specifically, increasing lake ice cover leads to a linear reduction in latent heat fluxes and nonlinear reduction in sensible heat fluxes (Gerbush et al., 2008). When relatively cool (warm) air masses pass over the Great Lakes during winter (summer), the relatively warm (cool) lake surface reduces (enhances) atmospheric stability and increases (decreases) deep convection, cloud cover, and precipitation (Scott and Huff, 1996; Holman et al., 2012; Bennington et al., 2014). The lakes' relatively low roughness compared to the surrounding land leads to strengthened over-lake wind speeds and potential shoreline convergence in support of enhanced lake-effect precipitation. Due to the lakes' large thermal inertia and resulting seasonal evolution in lake-air temperature contrast, the Great Lakes typically strengthen wintertime cyclones and summer anticyclones and weaken summertime cyclones and wintertime anticyclones (Notaro et al., 2013a). The basin is a preferred zone of
- wintertime cyclogenesis due to the relative warmth of the lake surfaces and consequential
- enhancement in low-level convergence (Petterssen and Calabrese, 1959; Colucci, 1976;
- Eichenlaub, 1978).

 Given the aforementioned substantial influence of the Great Lakes on regional climate, their representation and evaluation in both global and regional climate models have been the focus of several studies in the past decade. There is a wide spectrum among climate models regarding the treatment of large lakes. Due to their coarse spatial resolution, most global climate models (GCMs), including those from various phases of the Coupled Model Intercomparison Project (CMIP), either omit the Great Lakes entirely or offer a crude representation using wet soil, wetlands, ocean grid cells, or 1D lake models (Briley et al., 2021; Minallah and Steiner, 2021).

- Among regional climate models (RCMs) without lake models, many apply a rudimentary
- approach to estimate lake surface temperatures (LSTs) by extrapolating the closest ocean grid
- cell's sea-surface temperatures (SSTs), likely from Hudson Bay or the North Atlantic Ocean,
- from the initial and lateral boundary conditions datasets to the lake grid cell, potentially
- inducing vast biases and intra-lake discontinuities in LST and ice cover (Gao et al., 2012;
- Mallard et al., 2015; Spero et al., 2016; Xiao et al., 2016; Hanrahan et al., 2021). This approach
- is the default treatment of LSTs in the Weather Research and Forecasting (WRF) model
- (Hanrahan et al., 2021; Wang et al., 2022). Alternatively, the WRF Preprocessing System can

- designate time-averaged 2-m air temperatures to the underlying lake surfaces to provide
- estimated lower boundary conditions of LST based on the user-specified time window for
- temporal averaging and time lag for addressing thermal inertia (Wang et al., 2012; Mallard et
- al., 2015; Hanrahan et al., 2021; Wang et al., 2022). However, this approach still produces
- unrealistic LSTs and ice cover as the lakes cannot achieve equilibrium with the overlying
- atmosphere due to the lack of interactive lake-atmosphere feedbacks (Bullock et al., 2014;
- Spero et al., 2016).
- For those GCMs and RCMs that aim to incorporate coupled lake-atmosphere interactions, most
- apply 1D lake models (Perroud et al., 2009; Martynov et al., 2010; Stepanenko et al., 2010;
- Subin et al., 2012). Those include 2-layer bulk models founded in similarity theory such as the
- Freshwater Lake (FLake) model (Mironov et al., 2010), thermal diffusion models which
- parameterize eddy diffusivity such as the Minnesota Lake Water Quality Management Model
- (MINLAKE, Riley and Stefan, 1988) and the Hostetler model (Hostetler and Bartlein, 1990),
- Lagrangian turbulence models such as the Dynamics Reservoir Simulation Model (DYRSM,
- 122 Yeates and Imberger, 2003), and $k-\epsilon$ turbulence closure models with horizontally averaged
- velocity such as LAKE (Stepanenko and Lykossov, 2005; Stepanenko et al., 2011) and Simstrat
- (Goudsmit et al., 2002). Each of these different categories of 1D lake models has its own
- advantages and disadvantages (Perroud et al., 2009; Martynov et al., 2010; Stepanenko et al.,
- 2010; Subin et al., 2012). As demonstrated in these studies, the deficiencies include struggles
- with simulating seasonal stratification in FLake, insufficient mixing for deep lakes in the
- Hostetler model, and excessive mixing for shallow lakes in the computationally expensive
- turbulence models.
- Multiple modeling studies have assessed the performance of coupling RCMs to 1D lake models
- in the Great Lakes region. While this coupling permits the representation of key lake-
- atmosphere interactions and the heterogeneous spatiotemporal patterns of LSTs and lake ice
- cover, 1D lake models typically perform poorly at reproducing the lake thermal structure and
- seasonal ice evolution of large, deep lakes, such as Lake Superior, due to the overly simplified
- hydrodynamic processes. Common biases in 1D lake models include an anomalously early
- timing of both spring-summer stratification and autumn turnover, with positive biases in
- summer LST and negative biases in winter LST (Bennington et al., 2014; Mallard et al., 2014).

- The International Centre for Theoretical Physics (ICTP) Regional Climate Model version 4 (RegCM4), coupled to the 1D Hostetler lake model, yields a prolonged lake ice season with excessive ice cover due to the neglect of horizontal heat advection within the lakes (Notaro et al., 2013b). The coupling of a thermal diffusion lake model, the Lake, Ice, Snow and Sediment Simulator (LISSS, Subin et al., 2012) to the WRF model (available starting with version 3.6 of WRF) results in an early warm-up and overly rapid cool-down in the seasonal evolution of LSTs for deep lakes, along with an early onset of lake ice cover in support of its excessive abundance (Xiao et al., 2016). Mallard et al. (2014) found that WRF, coupled to FLake, produced the best performance for Lake Erie (the smallest and shallowest Great Lake) and the worst performance for Lake Superior (the largest and deepest Great Lake) among the Great Lakes in terms of simulated LST and ice cover biases. Often, modelers aim to reduce biases in the simulated vertical temperature profile of deep lakes in 1D models by artificially enhancing the vertical eddy diffusivity to crudely compensate for the absence of a dynamic circulation and vertical mixing processes (Subin et al., 2012; Bennington et al., 2014; Lofgren, 2014; Gu et al., 2015; Mallard et al., 2015), although such a non-physics based approach may only yield limited benefits to minimizing these biases (Xiao et al., 2016). The lack of fully resolved lake hydrodynamics in models, including dynamic 3D lake circulation, upwelling and downwelling, thermal bar formation, explicit horizontal mixing, and ice motion, along with overly simplified stratification processes, unrealistic treatment of eddy diffusivity, and the assumption of instantaneous mixing of thermal instabilities (Song et al., 2004; Martynov et al., 2010, 2012; Stepanenko et al., 2010; Bennington et al., 2014; Gu et al., 2015; Mallard et al., 2015; Sharma et al., 2018; Sun et al., 2020; Notaro et al., 2021; Hutson et al., 2024) has been the main obstacle in further improving climate simulations for the Great Lakes Basin.
- In recent years, a limited number of Great Lakes studies have aimed to enhance the
- representation of three-dimensional (3D) lake hydrodynamical processes and reduce the
- substantial biases in LST and ice cover associated with 1D lake models by coupling RCMs with
- 3D hydrodynamic models (Xue et al., 2017, 2022; Sun et al., 2020; Kayastha et al., 2023).
- These studies have responded to the urgent call for continued progress in coupling high-
- resolution RCMs with 3D lake models that address the complex processes and features of large,
- deep lakes, as highlighted in previous research (Martynov et al., 2010; Bennington et al., 2014;
- Briley et al., 2021; Notaro et al., 2021). Xue et al. (2017) developed a two-way coupled 3D

- lake-ice-climate modeling system, known as the Great Lakes-Atmosphere Regional Model (GLARM), by coupling RegCM4 with a 3D unstructured-grid hydrodynamic model, the Finite Volume Community Ocean Model (FVCOM, Chen et al., 2012). The resulting coupled 3D modeling system exhibited notable skill in reproducing the mean, variability, and trends in regional climate across the Great Lakes Basin and the physical characteristics of the Great Lakes, including their thermal structure and ice cover, significantly improving upon previous 175 RCM experiments coupled with 1D lake models. The updated version, GLARM-V2, has been utilized to generate future climatic and limnological projections for the Great Lakes region (Xue et al., 2022). Similarly, Sun et al. (2020) developed a lake-atmosphere-hydrology modeling system by coupling the Climate-WRF (CWRF) model with 3D FVCOM and compared its performance against CWRF coupled with the 1D LISSS. They found that the former configuration outperformed the latter in simulating LST, ice cover, and the vertical thermal structure in the Great Lakes. Kayastha et al. (2023) developed and validated the WRF- FVCOM Two-way Coupling (WF2C) model, showing WF2C improved upon past 1D lake model-based studies by significantly reducing the simulated summer LST bias, and revealing how coupled lake-atmosphere dynamics can influence summer LST by modifying surface heat fluxes through impacts on meteorological state variables. These studies underscore the advantages of coupling an RCM with a 3D lake hydrodynamic model for accurately depicting lake physical processes and lake-atmosphere feedbacks in the Great Lakes Basin. However, there is a notable absence of research dedicated to identifying the fundamental processes resolved in 3D lake models that contribute to these improvements, which is important to optimize effort allocation in future model development and improve our predictive understanding of the system. This knowledge gap is particularly significant for the Great Lakes during the winter seasons.
- This paper attempts to address this knowledge gap, by developing a new coupled lake-ice-
- atmosphere (CLIA version 1 or CLIAv1) modeling system for the Great Lakes by coupling the
- National Aeronautics and Space Administration (NASA)-Unified Weather Research and
- Forecasting (NU-WRF) regional climate model (RCM) with the three-dimensional (3D) Finite
- Volume Community Ocean Model (FVCOM). Note that CLIAv1 is hereinafter referred to as
- NU-WRF/FVCOM for the sake of particular attention given to comparing NU-WRF's
- performance during the cold season when two-way coupled with 3D FVCOM (NU-

- WRF/FVCOM) versus 1D LISSS (NU-WRF/LISSS). After a thorough validation of the
- coupled model, we conduct a series of process-oriented numerical experiments to identify the
- most important hydrodynamic processes that contribute to the superiority of the 3D lake model
- over the 1D lake model in enhancing lake-atmosphere coupling for the Great Lakes.

2 Model, Data, and Numerical Experiment Design

2.1 Atmosphere Model

NU-WRF is an observation-driven integrated regional modeling system, developed at NASA's

Goddard Space Flight Center (GSFC), that resolves chemistry, aerosol, cloud, precipitation and

land processes at satellite-resolvable scales (roughly 1–25 km) to improve the continuity

between microscale, mesoscale and synoptic processes. Developed as a superset of the

community WRF, NU-WRF unifies the NCAR - Advanced Research version of WRF model

(WRF-ARW) with the GSFC Land Information System (LIS, Kumar et al., 2006; Peters-Lidard

et al., 2007, 2015), the Goddard Chemistry Aerosol Radiation and Transport (GOCART) model

(Chin et al., 2000), the Goddard radiation and microphysics schemes (Shi et al., 2014), and the

Goddard Satellite Data Simulator Unit (G-SDU, Matsui et al., 2013, 2014). NU-WRF

simulations here utilize the Noah Land Surface Model, which simulates soil moisture and

temperature, skin temperature, snowpack depth and the energy flux and water flux terms of the

- surface energy balance and surface water balance (Mitchell, 2005). Currently, by default, the
- two-way lake-atmosphere interactions in NU-WRF are represented using the embedded 1D
- LISSS (Subin et al., 2012) from the Community Land Model version 4.5 (Oleson et al., 2013)
- 220 with modifications by Gu et al. (2015).

Notaro et al. (2021) conducted 20 simulations to identify the regionally optimal NU-WRF

configuration and schemes for the cold season period of November 2014-March 2015 in the

- Great Lakes region. The best model configuration was referred to as the "Morrison
- combination" and is used in this study. The "Morrison combination" includes Morrison

microphysics (Morrison et al., 2009), Rapid Radiative Transfer Model (RRTM, Mlawer et al.,

- 1997) longwave radiation physics, Community Atmosphere Model (CAM, Collins et al., 2004)
- shortwave radiation physics, Mellor-Yamada-Nakanishi-Niino Level 2.5 (MYNN2.5,
- Nakanishi and Niino, 2006, 2009) planetary boundary layer physics, and Mellor-Yamada-

- Nakanishi-Niino (MYNN, Nakanish, 2001) surface layer schemes. The improved simulations
- of air temperature and surface insolation using the Morrison combination primarily benefits
- from the Community Atmosphere Model's shortwave radiation scheme (Notaro et al., 2021).
- The Morrison combination is essentially the WRF configuration determined by Mooney et al.
- (2013) to produce the best simulated wintertime temperature simulation over Europe, who
- found that winter air temperatures are highly sensitive to the choice of radiation physics.
- The NU-WRF one-way nested configuration consists of an outer domain with 15-km grid
- spacing for the majority of North America and an inner domain with 3-km grid spacing for the
- Great Lakes region (Fig. 1), with the atmospheric vertical resolution assigned to 61 levels. The
- initial and lateral boundary conditions are provided by the Global Data Assimilation System 0-
- hour analysis. The cumulus parameterization option used for the outer domain is the Kain-
- Fritsch scheme (Kain and Fritsch, 1990; Kain, 2004) with resolved, unparameterized
- convection in the inner domain.

- **Figure 1.** NU-WRF nested domains (upper panel) and unstructured mesh used in FVCOM to represent
- 244 the Great Lakes in FVCOM (lower panel). The two dots denote the locations of Granite Island (87.4°W,
- 46.7°N) on Lake Superior and Spectacle Reef (84.1°W, 45.7°N) on Lake Huron.

2.2 Hydrodynamic Model

 The hydrodynamic model, FVCOM, is a free-surface, primitive equation hydrodynamic model that solves the momentum, continuity, temperature, salinity, and density equations and is closed physically and mathematically using turbulence closure submodels (Chen et al., 2012). Numerically, FVCOM employs the finite-volume method over an unstructured triangular grid and vertical sigma layers, optimizing flexibility and accuracy for complex terrains. The grid resolution adjusts from 1–2 km near coasts to resolve coastal geometry complexity, to 2-4 km offshore to improve computational efficiency (Fig. 1), with the model comprising 35,000 grid cells and 40 sigma layers. Vertical mixing processes are modeled using the Mellor–Yamada level-2.5 (MY25) turbulence closure model (Mellor and Yamada, 1982), while horizontal diffusivity is derived from velocity shear and grid resolution through the Smagorinsky (1963) formulation.

FVCOM also includes an unstructured-grid, finite-volume version of the Los Alamos

Community Ice Code (CICE), which describes ice thickness distribution in time and space.

CICE includes a thermodynamic model to compute local growth rates of snow and ice due to

vertical conductive, radiative, and turbulent fluxes. It also features an ice dynamics model to

simulate the ice pack velocity due to wind and ice-water stress, Coriolis effects, sea surface

slope, and internal stress, estimated with elastic–viscous–plastic rheology (Hunke and

Dukowicz, 1997). The transport model in CICE calculates the advective process of the areal

concentration, ice volumes, and other state variables. The ridging parameterization in CICE

 addresses mechanical redistribution, which transfers ice among thickness categories (Hunke et al., 2010).

In contrast, the default 1D lake model, LISSS, embedded in NU-WRF solves the 1D thermal

diffusion equation (i.e. lake thermal dynamics only) by segmenting the vertical stratification of

the lake into multiple distinct levels that include: snow (applicable when the snow's thickness

surpasses a specified minimum value); the combined section of lake water and ice, collectively

- identified as the "lake body"; and the bottom layers consisting of sediment, soil, and bedrock
- (collectively termed "sediment" unless specified differently). This structured division allows for
- simulating thermal dynamics within each segment, facilitating a prediction of temperature
- distribution and variations across the lake's depth (Subin et al., 2012).

3 Two-way Coupling of NU-WRF/FVCOM

 The development of interactively coupled model systems [see review by Giorgi and Gutowski Jr. (2015)] emerged quickly in the late 2000s driven by rapid technological advancement and the increase in computational capability. Model coupling is essential to multi-physics simulations representing various components of the Earth system. Over the past two decades, several coupling technologies for earth system modeling have been developed. Examples include the Earth System Modeling Framework (ESMF), the Model Coupling Toolkit (MCT), and the OASIS-MCT coupler, which is the latest version of the OASIS3 coupler interfaced with the Model Coupling Toolkit (MCT) that offers a fully parallel implementation of coupling field regridding and exchange (Valcke et al., 2012; Craig et al., 2017). Although coupling implementations can follow different approaches, their applications in geophysical simulations typically carry out several key functions, including interpolating and transferring the coupling fields between different model grids, managing data transfer between constitutive models at a desired coupling frequency, and coordinating the execution of the constituent models in a parallel computational environment (Valcke et al., 2012). In general, coupling data must be interpolated and transferred between the constituent models under several constraints, such as conservation of physical properties, numerical stability, consistency with physical processes, and computational efficiency.

- In the study, NU-WRF and FVCOM are run simultaneously, exchanging information
- bidirectionally at 1-hour intervals through the OASIS3-MCT coupler. FVCOM dynamically
- calculates the LST and ice cover, providing these as overlake surface boundary conditions to
- NU-WRF. Meanwhile, NU-WRF calculates and supplies the atmospheric forcings required by
- FVCOM, including surface air temperature, surface air pressure, relative and specific humidity,
- total cloud cover, surface winds, and downward shortwave and longwave radiation.

3.1 Data for Model Validation

3.2 Design of Numerical Experiments

We designed numerical experiments in two categories. In category 1, we evaluate the cold

season performance of the NU-WRF/FVCOM two-way coupling (case C1-1) against the NU-

WRF/LISSS 1D lake model (case C1-2). To ensure the objectivity of the comparison, both C1-

- 1 and C1-2 utilize an identical NU-WRF configuration (except for differences in lake
- treatment) as described in Section 2.1, following the optimal NU-WRF configuration for the
- study region as determined by Notaro et al. (2021). The comparison of C1-1 and C1-2 aims to

- examine the overall impact of using a 3D versus a 1D lake model configuration on simulating
- lake hydrodynamic conditions and the subsequent impact on the atmospheric state through
- lake-ice-atmosphere interactions from November 2014 to March 2015. The initial lake
- conditions of November 2014 were obtained from multiple years of FVCOM standalone
- simulations driven by Climate Forecast System Reanalysis (CFSR) forcing Xue et al. (2015).
- In category 2, a set of process-oriented numerical experiments is designed to identify the
- impact of various 3D hydrodynamical processes critical to the coupled Great Lakes system.
- These processes are either neglected or oversimplified by the NU-WRF/LISSS 1D lake model
- 338 while being resolved by the NU-WRF/FVCOM 3D lake model. Case C2-1 (NoIceTransp) is
- designed to examine the impact of ice transport associated with currents (Section 5.1). In this
- scenario, FVCOM is configured identically to C1-1, except that ice dynamics, ice velocity
- fields, and ice pack transport are disabled in FVCOM. Instead, only ice thermal dynamics are
- simulated to account for the spatio-temporal evolution of ice thickness distribution through
- thermodynamic growth and melting processes (Bitz and Lipscomb, 1999). Consequently, the
- ice model is simplified to function as an energy-conserving thermodynamic model, akin to that
- used in the 1D lake model.
- Case C2-2 (NoHeatAdv) analyzes the impact of 3D heat transport associated with lake
- circulation. FVCOM is configured identically to C1-1, except that the advective heat transport
- associated with current movement is disallowed in C2-2. This is realized by turning off the
- advection terms in the temperature equation in FVCOM, which is essentially an advection-
- diffusion equation that governs the distribution and evolution of temperature (Section 5.2).
- Therefore, the temperature calculation is simplified to imitate the 1D vertical diffusion equation
- used in the 1D lake model.
- Case C2-3 (NoShearProd) aims to assess the influence of 3D currents on calculation of
- turbulent mixing, a crucial factor in controlling the heat redistribution and thermal structure in
- the lakes. In this case, we exclude the turbulence shear production term that depends on
- currents in the turbulent kinetic equation (Section 5.3). In summary, the three cases in category
- 2 collectively reveal the significant impacts of currents in elements that are not accounted for in
- the LISSS 1D lake model, i.e. on ice transport, heat transport, and turbulent mixing intensity,

- 359 respectively. These experiments are summarized in Table 1.
- 360 **Table 1.** A summary of the numerical model experiments. The "3D currents" column shows if the
- 361 experiment resolves the 3D currents of the Great Lakes. The "Ice transport" column shows if the
- 362 experiment resolves the ice transport associated with currents in the Great Lakes. The "Heat advective
- 363 transport" column shows if the experiment resolves the 3D heat transport associated with Great Lakes
- 364 circulation. The "Shear production in turbulence" column shows if the experiment uses the turbulence
- 365 shear production term that depends on currents in the turbulent kinetic equation. The "Lake model"
- 366 column shows the lake model used in the experiment.

367

368 **4 Results**

369 **4.1 Lake Temperature and Ice Coverage**

370 The NU-WRF/FVCOM model (case C1-1) accurately captures the seasonal evolution of LSTs

371 across all of the lakes with lake-mean LST root-mean-square-error (RMSE) less than 0.4°C

372 (Fig. 2 upper panels). During November, the lakes are in the middle of their cooling period and

373 the LSTs decrease rapidly, yet at different paces, largely due to variations in the lakes' depth

 and latitude, which leads to strong spatial heterogeneity in LST (Fig. 3, left panels). The GLSEA data and the 3D lake model closely align in terms of the spatial LST patterns, with warmer waters of 10-12°C in the central and eastern basins of Lakes Erie and Ontario and 8- 10 $^{\circ}$ C in the southern basins of Lakes Michigan and Huron, while much cooler temperatures are found across Lake Superior, ranging between 4-6°C. The most notable underestimation of LST by the 3D lake simulation occurs in the southern basin of Lake Huron, while the model well captures the LSTs in the northern basin of Lake Huron. Transitioning to January 2015 (Fig. 3, right panels), at the onset of the ice season, NU-WRF/FVCOM accurately reflects the seasonal cooling of the lakes, showing a significant reduction in LSTs, while also well delineating the detailed temperature differences between the colder nearshore and relatively warmer offshore waters, in good agreement with the observational data. On the other hand, NU-WRF/LISSS (case C1-2) fails to capture the spatial heterogeneity in LSTs, but also generates a systematic cold bias of 2-3°C during January across nearly all of the lakes (Fig. 3, bottom panels). Such a cold bias was persistent in the NU-WRF/LISSS (Lake1D) simulation throughout the cold season, as detailed in Notaro et al. (2021).

 Figure 2. Time series of daily lake-averaged LST (°C, upper panels) and percent ice cover (lower panels) for the five lakes from GLSEA data (black lines) and NU-WRF/FVCOM 3D lake model simulations (red lines) during the simulation period of November 2014-March 2015. Both the temporal correlation and RMSE are reported in each panel.

Figure 3. Spatial patterns of monthly mean LSTs (°C) from GLSEA data (top panels), NU-

WRF/FVCOM 3D lake model simulations (middle panels), and NU-WRF/LISSSS 1D lake model

simulations (bottom panels) for November 2014 (left panels) and January 2015 (right panels).

 NU-WRF/FVCOM (Lake3D) also demonstrates its skill in capturing the evolution of the vertical thermal structure within the lake, which is particularly challenging in large and deep lakes (Bennington et al., 2014; Xue et al., 2017). As exemplified in Fig. 4, the in-situ thermistor measurement at Spectacle Reef on Lake Huron is located in a deep region with a water depth greater than 200 meters. The 3D model reproduces the conclusion of the summer stratification process until the end of November. The following turnover, a seasonal process where the surface water cools, becomes denser, and sinks—mixing with the warmer water from below— is also represented in the 3D lake model between December and January. Subsequently, the 406 winter inverse stratification, where colder water $(< 4^{\circ}C$) lies above warmer water due to the fact 407 that freshwater's density peaks at $4^{\circ}C$, is captured by the 3D model as it develops from February onward, although the model shows a stronger winter inverse stratification and earlier onset than observed. In contrast, NU-WRF/LISSS falls short of these detailed observations. Not only does it mispredict the occurrence of turnover and winter stratification much earlier than

- observed, but it also substantially underestimates the extent of mixing between the surface and
- deeper waters. This underestimation results in a flawed representation of excessive surface
- cooling and a substantial overestimation of the warming of the deep waters.

Figure 4. Seasonal evolution of daily vertical temperature (°C) profiles from the thermistor observations

(top panel), NU-WRF/FVCOM 3D lake model (middle panel), and NU-WRF/LISSS 1D lake model

(bottom panel) at Spectacle Reef in Lake Huron during November 2014-March 2015.

Correspondingly, NU-WRF/FVCOM resolves the spatiotemporal evolution of lake ice cover

very well across all of the lakes with RMSE of percent ice cover less than 8% for Lakes Huron,

Michigan, and Ontario and 11% and 18% for Lakes Superior and Erie, respectively (Fig. 2

- lower panels). The 3D lake model and GLSEA data exhibit similar seasonal trends both in
- timing and magnitude, with ice cover typically starting to rapidly increase in January, peaking
- in February and early March, and declining thereafter (Fig. 2). Lake Erie shows the earliest and
- sharpest increase in ice cover, peaking near 100% in early February and throughout mid-March,
- indicative of its shallower depth and weaker thermal inertia. Lakes Huron and Superior show a

- persistent increase in ice cover through February, with peak coverage of >90% occurring at the
- beginning of March. Lakes Michigan and Ontario exhibit more gradual increases and lower
- peaks in ice cover. The model appears to capture the general seasonal trends of the GLSEA
- data with high fidelity, although some discrepancies are evident, particularly over Lakes Erie
- and Superior (Fig. 2).

 NU-WRF/FVCOM performs reasonably well in mirroring the general spatial patterns of lake ice cover (Fig. 5, top and middle panels). For January, the GLSEA data shows a pronounced ice formation in the nearshore regions across the lakes, with the greatest ice concentration visible along the coastlines and very limited ice cover in offshore waters. The model captures this nearshore ice development quite well, although it suggests less ice cover in the offshore areas, particularly over Lake Erie. In February, the extent of ice cover varies dramatically across the lakes, including nearly full ice cover on Lake Erie and significant ice-free areas on Lake Ontario, as well as for Lakes Michigan and Huron, which have distinctly less ice cover in their southern and central basins, respectively. The model captures this variability very well, while slightly overestimating the ice cover in the central regions of Lake Superior. For March, the model successfully replicates the patterns of significant declines in ice cover in the western sections of the lakes, with much higher ice coverage in the eastern sections of the lakes.

- On the other hand, NU-WRF/LISSS (Lake1D) generates excessive ice cover during January,
- when both observations and NU-WRF/FVCOM suggested that the majority of the lakes were
- ice-free. In February, the excessive ice cover simulated by the NU-WRF/LISSS model persists,
- with near 100% ice coverage over all of the lakes, and the model fails to depict the large spatial
- variability across the lakes. Such a persistent overestimation of ice cover throughout the cold
- season by NU-WRF/LISSS was also reported in Notaro et al. (2021).

Figure 5. Spatial patterns of mean percent lake ice cover from GLSEA data (top panels), NU-

WRF/FVCOM 3D lake model simulations (middle panels), and NU-WRF/LISSSS 1D lake model

 simulations (bottom panels) for January 2015 (first column), February 2015 (second column), and March 2015 (third column).

4.2 Over-lake Latent and Sensible Heat Fluxes

 The improved LST and ice simulation by the 3D lake model translates to an improvement in the simulated over-lake latent and sensible heat fluxes, particularly for the ice-cover season (Fig. 6). The observations for upward latent and sensible heat fluxes from two eddy covariance flux towers at Granite Island on Lake Superior and Spectacle Reef on Lake Huron are compared against the simulated fluxes from NU-WRF/FVCOM (Lake3D) and NU-WRF/LISSS (Lake1D). The two lakes are selected for demonstration as they have the highest ice coverage during the simulation period. NU-WRF/LISSS reasonably simulates the magnitude and variability of the heat fluxes from November until mid-December, similar to the observations and NU-WRF/FVCOM, although with larger biases. However, it grossly underestimates the fluxes during the ice-cover season (January-March) by simulating a nearly constant near-zero flux. This is mainly due to the excessive ice cover simulated by the 1D lake model, which creates a physical barrier for air-lake energy fluxes. Since the 3D lake model more accurately

- simulates the LST and ice cover, it successfully captures the magnitude and variability of the
- heat fluxes, even during the ice-cover season, with RMSEs that are 50% lower than those from
- the 1D lake model (Fig. 6). Latent heat in Spectacle Reef is the only exception, where NU-
- WRF/FVCOM struggles to capture the magnitude of the upward latent heat flux due to the
- overestimated ice cover at the site. However, it still outperforms NU-WRF/LISSS in terms of
- capturing the seasonal trend in latent heat fluxes.

Figure 6. Time series of daily sensible (upper panels) and latent (lower panels) heat fluxes (W/m²) from GLEN observations (black lines), NU-WRF/FVCOM 3D lake model simulations (red lines), and NU- WRF/LISSS 1D lake model simulations (blue lines) at Granite Island on Lake Superior (left) and Spectacle Reef on Lake Huron (right). The RMSE and temporal correlations between the simulations and

GLEN observations are provided in each panel.

4.3 Over-lake Air Temperature and Wind

- Along with the improved simulation of the Great Lakes' physical characteristics and surface
- heat fluxes, NU-WRF/FVCOM improves the simulated over-lake atmospheric state across the
- Great Lakes, including air temperature and wind speed. The cold air temperature biases
- produced over the lakes by NU-WRF/LISSS are significantly reduced (Fig. 7) with better

 simulated, more intense upward heat fluxes in January. This improvement in the simulated air temperature at the two sites, Granite Island and Spectacle Reef, is clearly evident. Similar to the fluxes, NU-WRF/LISSS modeled air temperature diverges from the observations in January and February, with a noticeable cold bias. This cold bias is the result of significant suppression of the upward heat fluxes during those months in the 1D lake model due to excessive simulated ice cover. NU-WRF/FVCOM, on the other hand, produces a much warmer and more accurate over-lake air temperature for January and February due to its reasonable representation of upward heat fluxes. The simulated wind speed over the lakes is also improved, especially in January-February (Fig. 7). This advancement is attributed to the refined simulation of surface roughness (i.e., ice versus water), the water-air temperature gradient, and associated instability over the lakes due to decreased ice cover. Large wind spikes (16 m/s) in January-February are more accurately captured by NU-WRF/FVCOM.

 Figure 7. Time series of daily air temperature (°C, upper panels) at 2-m height (T2) and wind speed (m/s, lower panels) from GLEN observations (black lines), NU-WRF/FVCOM 3D lake model

- simulations (red lines), and NU-WRF/LISSSS 1D lake model simulations (blue lines) at Granite Island
- on Lake Superior and Spectacle Reef on Lake Huron during November 2014-March 2015. The RMSE
- and temporal correlations between the simulations and GLEN observations are provided in each panel.

5 Discussion

5.1 Impact of Ice Movement

 The 3D hydrodynamic model, FVCOM, includes an embedded unstructured-grid ice model capable of resolving several components for atmosphere-ice-water interactions (Gao et al., 2011). It includes a thermodynamic model that computes the local growth rates of snow and ice due to vertical conductive, radiative, and turbulent fluxes, aligning with features typically included in 1D lake models (Bitz and Lipscomb, 1999). More importantly, it features an ice dynamics model that predicts the ice pack's velocity field based on its material strength; a transport model that describes the advection of areal concentration, ice volumes, and other state variables; and a ridging parameterization that facilitates the transfer of ice among thickness categories (Hunke et al., 2010).

Case C2-1 (NoIceTransp) is designed to examine the impact of ice transport on LSTs and

overlying atmospheric conditions, compared to standard case C1-1 (Lake3D). In case C2-1, ice

dynamics, velocity fields, and ice pack transport are disabled in FVCOM. Instead, only ice

thermal dynamics are simulated, as in the 1D lake model. Figure 8 compares cases C1-1

(Lake3D) and C2-1 (NoIceTransp), illustrating their performance in simulating the observed

spatial pattern of ice coverage in March 2015, characterized by open water on the western side

of the Great Lakes and predominant ice cover on the eastern side (Fig. 8a). Utilizing a 3D lake

model that only accounts for ice thermal dynamics results in an overestimation of ice cover,

with near 100% lakewide ice cover in Lakes Superior, Huron, and Erie (Fig. 8b). However,

integrating ice dynamics, including transport influenced by wind and water-ice stress, results in

 excellent agreement with observations, highlighting the critical role of ice transport in accurate ice modeling (Fig. 8c). This pattern aligns with the modeled ice velocities, which attribute the eastward ice cover distribution to dominant eastward ice transport (Fig. 8d). Under cold winter conditions characterized by strong westerly winds, ice is driven eastward, maintaining open water in the lake's western part. This facilitates ongoing atmospheric interactions, allowing for heat release. Neglecting these dynamics leads to unrealistic ice accumulation by diminishing the influence of wind on surface water movement and mixing. This overaccumulation of ice cover hampers the efficiency of vertical turbulent mixing, which is essential for maintaining a warmer surface layer, thereby exacerbating ice formation and accumulation. The incorporation of ice dynamics into 3D lake models is thus essential for accurately simulating ice distribution, emphasizing the necessity of resolving ice transport to replicate observed patterns accurately.

 Figure 8. Spatial patterns of mean percent lake ice cover from GLSEA data (a), case 2-1 (NoIceTransp) simulations (b), and case 1-1 (Lake3D) standard simulations (c), along with simulated mean ice

velocities (m/s) during (d) March 2015.

546 **5.2 Impact of Heat Transport**

- 547 The 3D lake model also resolves the advective transport of heat associated with the simulated 548 circulation. The advective transport and turbulent mixing of temperature in the 3D lake model
- 549 are governed by following equation:

550
$$
\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} = \frac{\partial}{\partial z} \left(K_h \frac{\partial T}{\partial z} \right) + F_T
$$
 (1)

551 with the surface heat flux boundary condition:

552
$$
\frac{\partial T}{\partial t} = \frac{1}{\rho c_p K_h} \left[L W(x, y, t) - L H(x, y, t) - S H(x, y, t) \right]
$$
(2)

553 where *T* is the water temperature and *u, v*, and *w* are the *x, y,* and *z* components of the water

554 velocity, respectively. K_h is the vertical thermal diffusivity coefficient and F_T is the horizontal

555 diffusion term. ρ is water density, c_p is the specific heat capacity of water and

556 *LW* (x, y, t) , *LH* (x, y, t) , and *SH* (x, y, t) are net longwave radiation, upward latent heat and

557 sensible heat fluxes varying in space and time, respectively.

558 Case C2-2 (NoHeatAdv) analyzes the impact of 3D heat transport. In this case, the 3D 559 temperature advection terms $(u \frac{\partial T}{\partial x}, v \frac{\partial T}{\partial y}, w \frac{\partial T}{\partial z})$ are turned off.

 Comparing the standard simulation C1-1 (Lake3D) to case C2-2 (NoHeatAdv), Figure 9 demonstrates that, in the absence of advective heat transport by lake currents, the surface temperatures can remain consistent with the basic patterns observed in the standard 3D lake simulation throughout the entire simulation period. The differences in the time series of lake- wide average LSTs for the five lakes are small, with a maximum difference of 0.4°C between the two cases. The spatial patterns of LST biases, when compared with GLSEA, are generally 566 more noticeable, with the most significant positive biases (\sim 2 \degree C) concentrated around the coastal waters of the Great Lakes and eastern Lake Erie from January to March 2015 and larger 568 negative biases (\sim 3°C) in the central basin of Lake Huron in November 2014 in the

569 NoHeatAdv case.

- In addition, the vertical transport associated with upwelling, resolved by the 3D model, brings
- relatively warmer water from deep in the lake to the surface. This vertical transport mechanism
- cannot be represented in 1D lake models that only account for vertical diffusion. This can
- create significant local-scale differences along the coast, as shown on the western shore of Lake
- Superior in March 2015 [Fig. 9, bottom panels. Notice that the GLSEA is not able to well
- capture coastal upwelling (Ye et al., 2020)]. This underscores the importance of including
- advective heat transport to accurately resolve the redistribution of heat within the lake. The
- inclusion of advective dynamics, by facilitating both lateral and vertical redistribution, enables
- a more realistic simulation of the complex spatial heat patterns within large lake systems.

Figure 9. Spatial patterns of mean LSTs (°C) from GLSEA data (first column), case C1-1 (Lake3D)

- standard simulations (second column), and case C2-2 (NoHeatAdv) simulations (third column) from
- November 2014 (top row) to March 2015 (bottom row). Their monthly biases relative to GLSEA data are
- presented in the fourth and fifth columns, respectively.
- Capturing the evolution of the vertical thermal structure within the deep water is particularly
- challenging in lake models. As previously shown in Fig. 4, the in-situ thermistor measurement

- at Spectacle Reef on Lake Huron is located in a deep region with a water depth greater than 200
- meters. Case C2-2 (NoHeatAdv) generally reproduced the thermal patterns from case C1-1
- (Lake3D) in terms of both timing and intensity of summer stratification, fall turnover, and
- winter inverse stratification (Fig. 10a,b). While the comparison shows that the overall thermal
- structures are similar in both simulations, there is a noticeable difference within the subsurface
- layer, specifically between 50 to 100 meters in depth (Fig. 10c), suggesting that heat advection
- might have a more significant impact on temperature distribution in the subsurface layer of the
- water column in this case. Without accounting for heat advective transport, there appears to be
- artifacts of stepwise vertical thermal gradients in case C2-2.

- **Figure 10.** Mean vertical temperature (°C) profiles from a): case C1-1 (Lake3D) standard run and b):
- case C2-2 (NoHeatAdv) and c): their difference at Spectacle Reef in Lake Huron during November 2014-
- March 2015.

- To gain a deeper understanding of the results, we analyzed the heat balance to identify the
- contributions of different physical processes. This analysis involved examining each term in the
- temperature governing equation (Eq. 1) that is directly computed in FVCOM over the
- simulation period. The temperature change is driven by 3D advective heat transport, horizontal
- heat diffusion, and vertical diffusion due to turbulent mixing.

Figure 11. Monthly averaged vertical profile of each term of the temperature equation in the C1-1

(Lake3D) simulation from November 2014 to March 2015, output from location at Spectacle Reef in

607 Lake Huron. The temperature change rate $\left(\frac{\partial T}{\partial t}\right)$ is determined by 3D advection (blue), horizontal diffusion

(red), and vertical diffusion (purple).

 The analysis revealed the relative impact of physical processes on thermal changes during the winter months (Fig. 11). Significant cooling and decreasing temperatures were observed within the upper 100 meters of the water column, as indicated by the negative temperature change over 612 time $(\frac{\partial T}{\partial t})$ in this zone. In contrast, the water below 100 meters in depth remained largely unchanged and stable in temperature. In November and December, vertical turbulent mixing processes primarily controlled the cooling rate in the upper 100 meters, during which surface heat fluxes served as net losses from the lake along with vigorous turbulent mixing in the lake. Advection played a much less important role in temperature changes during this period. However, starting in January, 3D advection played an important role in redistributing heat in the 25-100 meter layer, offsetting some of the cooling induced by surface heat loss through mixing. In February and March 2015, advection proved to be significant at the lower boundary

- 620 of the surface mixed layer. These observations explain the larger temperature difference in the
- 621 subsurface layer between cases C1-1 (Lake3D) and C2-2 (NoHeatAdv) (Fig. 10), highlighting
- 622 the evolving balance between vertical diffusion and advection in controlling the epilimnetic
- 623 heat budget and temperature changes in large lakes during the cold season.

624 **5.3 Impact of Vertical Mixing**

 The analysis above (Fig. 11) highlights the dominant factor, vertical turbulent mixing, in determining seasonal lake temperature change. Note that we have already discussed the importance of ice transport associated with currents as well as the impact of advective heat transport. To understand the mechanism responsible for the differing performance between the 1D and 3D lake models in simulating vertical mixing, we examine how vertical turbulent mixing is calculated in these two types of models. The intensity of vertical mixing in both models is represented by vertical eddy diffusivity, which is determined by turbulent kinetic energy (q^2) . In the 3D hydrodynamic lake model, a sophisticated 3D turbulence closure model 633 is used, in which a prognostic equation predicts the change rate of q^2 based on its advection, and its turbulence production, including both shear-induced production (*Ps*) and buoyancy-635 induced production (P_b) , and its dissipation rate (ϵ) , as well as its diffusion. This equation is 636 complemented by either a separate prognostic equation for dissipation rate ($k-\epsilon$; Launder and Spalding, 1974) or a diagnostic equation for turbulent mixing length (Meller and Yamada, 638 1982).

639 The equation governing the evolution of turbulent kinetic energy (q^2) is

640
$$
\frac{\partial q^2}{\partial t} + u \frac{\partial q^2}{\partial x} + v \frac{\partial q^2}{\partial y} + w \frac{\partial q^2}{\partial z} = P_s + P_b - \epsilon + \frac{\partial}{\partial z} \left(K_q \frac{\partial q^2}{\partial z} \right) + F_q
$$
 (3)

641 where $q^2 = (\langle u'^2 \rangle + \langle v'^2 \rangle + \langle w'^2 \rangle)/2$, with u', v', w' represent the fluctuating components of 642 velocity in the x, y, z directions, respectively. The $\langle \rangle$ denotes averaging over time or space to 643 obtain the mean. Shear production is often approximated as $P_s = K_m \left(\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right)$, where 644 K_m is the vertical eddy viscosity coefficient. Buoyancy production is computed as $P_b =$ 645 $-\frac{g}{\rho_0}K_h\frac{\partial \rho}{\partial z}$, where g is acceleration due to gravity. ρ_0 is reference density of the fluid (e.g.,

- 646 ocean water or air). K_h is thermal diffusivity, $\frac{\partial \rho}{\partial z}$ is vertical gradient of density, indicating
- 647 stratification. The turbulent kinetic energy dissipation rate is represented as $\epsilon = q^3/Bl$, where
- 648 *i* is the turbulence length scale and B is an empirical constant. $K_{m,h,q} = q l S_{m,h,q}$, where $S_{m,h,q}$
- 649 are stability functions for $K_{m,h,q}$, respectively. K_q is the vertical diffusivity coefficient for
- 650 turbulent kinetic energy and F_q is horizontal diffusion of the turbulent kinetic energy.
- Figure 12 reveals that in the Great Lakes, shear production—induced by the vertical gradient of
- horizontal velocity in the water column—is the primary driver of subsurface turbulent mixing.
- Conversely, buoyancy production plays a secondary role, being at least one order of magnitude
- smaller than shear production in the first 50 meters of depth. This underscores the importance
- of including accurate current simulation when estimating the vertical turbulent mixing, which is
- crucial for accurately simulating heat exchange in the water column and ultimately determining
- the lake's thermal structure and ice formation.

Figure 12. Monthly averaged vertical profile of each term of the turbulence kinetic equation in the C1-1

(Lake3D) simulation from November 2014 to March 2015, at Spectacle Reef on Lake Huron. The change

- rate of turbulent kinetic energy (blue) is based on the 3D advection (red), and the turbulence production,
- including both shear-induced production (green) and buoyancy-induced production (cyan), the
- dissipation rate (black), and vertical diffusion (purple).

- Figure 13 compares the vertical temperature profiles between the standard simulation C1-1
- (Lake3D) and case C2-3 (NoShearProd). The NoShearProd case shows much stronger
- stratification, particularly from January to March. The absence of shear production leads to
- significantly reduced turbulent mixing and limiting heat exchange between surface and deeper
- waters, which results in a much colder surface layer (0-40 m) in January and much warmer
- deep waters (50-150 m) in February and March compared to the standard run. Consequently,
- the colder surface water temperature favors ice formation, leading to overestimated ice cover in
- the NoShearProd case compared to the standard simulation and observations, particularly in
- January and February (Fig. 14).

- **Figure 13.** Mean vertical temperature (°C) profiles from a): case C1-1 (Lake3D) standard run and b):
- case C2-3 (NoShearProd) and c): their difference at Spectacle Reef on Lake Huron during November
- 2014-March 2015.

 Figure 14. Spatial patterns of mean percent lake ice cover from GLSEA data (first column), case C1-1 (Lake3D; second column) and case 2-3 (NoShearProd; third column) for January 2015 (first row),

-
- February 2015 (second row), and March 2015 (third row).

LISSS, as true of other 1D lake models, was originally designed for small and shallow inland

- lakes and was not designed to resolve water currents (Subin et al., 2012; Notaro et al., 2021).
- Some other 1D lake models (Stepanenko and Lykossov, 2005; Stepanenko et al., 2011) employ
- a crude representation of average flow fields. Therefore, 1D lake models rely on empirical or
- semi-empirical relationships to estimate how wind stress affects the lake's turbulence and
- mixing without explicitly resolving 3D velocity fields. These thermal diffusion-based models
- often employ a latitude-dependent Ekman decay, accompanied by an empirical modification
- factor, to estimate a lumped eddy diffusivity coefficient as an approximation for surface wind-
- induced mixing (Xiao et al., 2016). Thus, the lack of accurate simulation of turbulent mixing

- processes makes the 1D model of limited capacity in accurately simulating the Great Lakes'
- thermal structure.

6 Summary and Conclusion

 In summary, a two-way coupled NU-WRF/FVCOM model (CLIAv1) has been developed toward the next generation of a regional climate model for the Great Lakes Basin for accurate representations of lake–ice–atmosphere interactions. NU-WRF/FVCOM significantly improved on the performance of NU-WRF coupled with an optimized 1D lake model, and accurately reproduced the physical characteristics of the Great Lakes (e.g., LST, ice cover, and thermal structure). This led to further improvements in simulated over-lake atmospheric conditions (e.g., air temperature, wind, latent and sensible heat) through two-way lake-atmosphere interactions.

 While 1D column lake models have been widely used in the simulations of inland lakes worldwide, small inland lakes and the Great Lakes exhibit fundamental differences in their physical characteristics, such as size and depth, which in turn influence their mixing behaviors, thermal structures, and circulation patterns. Inland lakes, generally much smaller (with a typical 705 average area of 1-10 kilometers) and much shallower (with a typical average depth of \sim 10m), respond more rapidly to atmospheric conditions. This leads to a fairly uniform horizontal pattern and a simpler mixing process in response to surface wind, due to their shallow depth and small thermal inertia. Therefore, 1D column lake models serve as an appropriate and efficient tool for simulating inland lake processes, particularly when the lake depth is shallower than 20 meters. In contrast, the vast size (e.g., Lake Superior alone covers about 82,100 square kilometers) and significant depth (e.g., the average depth of Lake Superior is 147 m, with a maximum depth of 400 m) of the Great Lakes result in complex hydrodynamic and thermal dynamics. This complexity causes the Great Lakes to exhibit many sea-like characteristics.

- This study has highlighted key physical processes that differentiate the large, deep Great Lakes
- from small, shallow inland lakes, and how these processes impact lake simulations.
- Specifically, we identified that ice dynamics, particularly ice transport, are vital in the Great
- Lakes, influencing ice cover formation and heat exchange between the lake and the atmosphere.

- Secondly, we show that advective heat transport, facilitates both lateral and vertical redistribution, enables a more realistic simulation of the complex spatial temperature patterns, particularly the predominance of advective heat transport in the subsurface layers. Thirdly, we identified the critical role of resolving shear production in turbulent mixing in the Great Lakes, which is the most influential factor that determines heat transfer and, subsequentially, lake thermal structure. Ice transport, heat transfer, and shear production in turbulence mixing are fundamentally linked to the 3D lake currents, which are missing or crudely represented in 1D lake models. Our findings underscore that circulation currents are pivotal in the physical limnology of the Great Lakes. Given the ongoing impact of climate change on these aquatic systems (Zhong et al., 2016; Woolway et al., 2021; Cannon et al., 2024), accurately incorporating 3D lake dynamics becomes crucial for projecting future thermal structures and ecosystem effects.
- Lastly, we acknowledge that there are multiple ways to tune the 1D lake column model or build an accurate empirical relationship between atmospheric conditions and the strength of mixing to improve 1D model simulations. However, the major challenge with this approach is that any empirical or simplified physical relationship carries significant risks of not holding in the future, especially in the context of climate change. While it may work well to calibrate the model based on a substantial amount of validation data, this approach has a much larger risk and lacks reliability if the model is used for climate projections where conditions change significantly. Therefore, we advocate for the complete integration of 3D hydrodynamic lake models in a two-way coupled fashion to project future changes in large freshwater systems. This method ensures that projections are based on physical processes, reducing the risk associated with empirical relationships and increasing the model's reliability for future climate scenarios.

Code and data availability

- The source codes of CLIAv1 with the two-way coupled FVCOM and NU-WRF used in this
- 744 study are available at https://doi.org/10.5281/zenodo.12746348 (Huang, 2024a) and
- https://doi.org/10.5281/zenodo.12746306 (Huang, 2024b) respectively. The GLSEA data were
- obtained from the NOAA Coastwatch website (https://coastwatch.glerl.noaa.gov/glsea/doc/)

- (GLSEA, 2023). The GLEN data were from the Lake Superior Watershed Partnership website
- (https://superiorwatersheds.org/GLEN/), with data compilation and publication provided by
- LimnoTech under Award/Contract 10042-400759 from the International Joint Commission
- (IJC) through a subcontract with the Great Lakes Observing System (GLOS).

Author contributions

- PX conceived the study. PX and CH developed the model code. PX designed the experiments.
- PX, MN, XZ, CH, MBK, and CZ conducted the analyses. PX, MN, MBK, and CZ wrote the
- original manuscript. All others contributed to revising the manuscript. All authors have read
- and agreed to the published version of the manuscript.

Competing interests

The authors declare that they have no conflict of interest.

Acknowledgments

- This is the Contribution No. 121 of the Great Lakes Research Center at Michigan
- Technological University. The study was funded by NASA's Modeling, Analysis, and
- Prediction Program (Grant 80NSSC17K0287 and Grant 80NSSC17K0291). Hydrodynamic
- modeling was also partly supported by the U.S. Department of Energy, Office of Science, under
- award number DE-SC0024446. The statements, findings, conclusions, and recommendations of
- authors expressed herein do not necessarily state or reflect those of the United States
- Government or any agency thereof.

References

