

# Skin Sea Surface Temperature schemes in coupled ocean-atmosphere modeling: the impact of chlorophyll-interactive e-folding depth.

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**Abstract.** In this paper, we explore different prognostic methods to account for skin sea surface temperature diurnal variations in a coupled ocean-atmosphere regional model of the Mediterranean Sea. Our aim is to characterize the sensitivity of the considered methods with respect to the underlying assumption of how the solar radiation shapes the warm layer of the ocean. All existing methods truncate solar transmission coefficient at a warm layer reference depth which is constant in space and time; instead, we develop a new scheme where this latter is estimated from a chlorophyll dataset as the e-folding depth of solar transmission, therefore allowing it to vary in space and time depending on seawater's transparency conditions. Comparison against satellite data shows that our new scheme, compared to the one already implemented within the ocean model, improves the spatially averaged diurnal signal, especially during winter, and the seasonally averaged one in spring, and autumn, while showing a monthly, basin-wide averaged bias smaller than 0.1 K year-round. In April, when most of the drifters' measurements are available, the new scheme mitigates the bias during nighttime, keeping it positive but smaller than 0.12 K during the rest of the monthly-averaged day. The new scheme implemented within the ocean model improves the old one by about 0.1 K, particularly during June. All the methods considered here showed differences with respect to objectively analyzed profiles confined between 0.5 K during winter and 1 K in summer for both the eastern and the western Mediterranean regions, especially over the uppermost 60 m. More in detail, the new scheme reduces the RMSE on the top 15 m in the central Mediterranean for summertime months, compared to the one already implemented one within the ocean model. Overall, the surface net total heat flux shows that the use of a skin SST parametrization brings the budget about 1.5 W/m<sup>2</sup> closer to zero on an annual basis, despite all simulations showing an annual net heat loss from the ocean to the atmosphere. Our "chlorophyll-interactive" method proved to be an effective enhancement of existing methods, its strength relying on an improved physical consistency with the solar extinction implemented in the ocean component.

## 38 1 Introduction

39 Air-sea fluxes govern the energy exchange at the ocean-atmosphere interface. A reliable representation of  
40 the Sea Surface Temperature (SST) diurnal cycle, i.e. the typical SST oscillation/excursion between night and  
41 day mainly due to solar heating, is crucial to accurately estimate air-sea heat fluxes (Kawai and Wada, 2007,  
42 Soloviev and Lukas, 2013), whose direct measurement is very difficult. Indeed, diurnal warming events can  
43 often exceed 5 K depending on weather conditions (Soloviev and Lukas, 1997) and geographical location,  
44 typically at tropical and mid-latitudes but also occasionally at high latitudes (Karagali and Høyer, 2013). Large  
45 diurnal warming events can lead to changes in air-sea heat flux locally reaching up to  $60W/m^2$  (Fairall et al.,  
46 1996, Ward, 2006, Kawai and Wada, 2007, Marullo et al., 2010, Marullo et al., 2016) on a variety of scales,  
47 ranging from the short regional ocean weather ones to large seasonal or long-term ones.

48  
49 Therefore, there is a wide interest in the development of models to accurately reconstruct SST diurnal variations  
50 in order to improve the representation of air-sea energy exchanges, especially, but not solely, within the coupled  
51 ocean-atmosphere modeling framework (Penny et al., 2019).

52 The net energy flux across the air-sea interface results from four contributions: the net solar radiation; latent  
53 and sensible heat fluxes, and the net thermal radiation. The last three contributions depend on SST and have a  
54 direct impact in determining ocean heat uptake or dynamical processes such as deep-water formation (Chen  
55 and Houze Jr., 1997). Ideally, the most accurate flux estimate would imply the knowledge of the temperature  
56 right at the atmosphere-ocean separation interface. From an observational point of view, the skin SST is the  
57 temperature immediately adjacent to the ocean surface (~10-20 microns depth) that is measurable, typically  
58 from infrared radiometers, and thus a key parameter to understand heat flux exchange (Minnet et al., 2019).  
59 Indeed, following what is measurable by current sensors, the GHRSSST-PP (i.e. the Global ocean data  
60 assimilation experiment High Resolution SST Pilot Project) introduced the distinction between skin, sub-skin,  
61 depth, and foundation SST (Donlon et al., 2007), which can be respectively regarded as successive, better-to-  
62 worse approximations to the ideal target, i.e. SST right at the interface, which is actually impossible to measure.  
63 However, in most of the widely used ocean models and configurations, the too-coarse vertical resolution does  
64 not allow to direct modeling skin SST (the first model layer being only around 0.5 - 1 meter thick, e. g. the  
65 ocean model NEMO – see the sketch in figure 1). Therefore, one must use schemes to reconstruct skin SST  
66 variations. Sadly, the only thing one can be sure about is that in general no model will be able to perfectly  
67 reproduce skin SST diurnal variations, and there are different ways to approach this challenging problem, each  
68 one still with its own limitations (see Kawai and Wada, 2007 and references therein). Simplified approaches  
69 widely employed in ocean and atmosphere state-of-the-art models parameterize the skin SST dynamics via the  
70 distinction of two main effects: the cool skin and the warm layer. Due to its interactions with the atmosphere,  
71 the temperature right at the interface of separation is supposed to be almost anywhere and anytime lower than  
72 the temperature of the waters infinitesimally close to it, resulting in the ocean being covered with a thin cool  
73 skin layer. One of the very first and simpler models assumes this cool skin temperature difference as  
74 proportional to the ratio between heat fluxes and kinematic stress (Saunders, 1967), via the Saunders' constant.

75 The cool skin effect is very important in obtaining accurate estimates of the latent and sensible heat flux,  
76 especially because its consideration modifies specific humidity at the ocean surface, which is one of the factors  
77 in the bulk formula. Indeed, latent and sensible heat fluxes are defined as the heat transfer across the  
78 ocean/atmosphere interface due to turbulent air motions (the former including the one resulting from  
79 condensation or evaporation). For example, a recent study in the South China Sea showed that during nighttime  
80 the cool skin temperature difference is around 1 K, and there's currently a large uncertainty in the Saunders'  
81 constant (Zhang et al., 2021). A warm layer (in which diurnal warming effectively takes place) develops below  
82 this cool skin, and its extent reaches a depth at which the penetration of solar radiation can be neglected (usually  
83 fixed to 3m by most existing parameterizations – see section 3.3 for more details). Diurnal warm layer  
84 anomalies (which can sometimes exceed 3K) can potentially impact both the atmosphere and ocean mean state  
85 on a variety of spatial (ranging from regional, basin-wide to global ones) and temporal scales (relevant for  
86 weather or seasonal forecast to long-term climatic trends) (Donlon et al., 2007). The skin SST diurnal warming  
87 amplitude increases under low surface winds (smaller than 2 m/s) and intense solar radiation (higher than  
88 typical daily peaks, around 900 W/m<sup>2</sup>) conditions, smaller in winter and at the poles than in summer and in  
89 the tropics. The accuracy of skin SST models, and therefore their ability to reconstruct skin SST diurnal  
90 variations is crucial especially in heat budget closure problems, which are still a subject of active debate  
91 especially in climate change hot spot regions such as the Mediterranean domain (see Marullo et al., 2021 and  
92 references therein). Skin SST schemes are also crucial for assimilating daytime SST data from satellite sensors  
93 (Penny et al., 2019; Storto and Oddo, 2019, Jansen et al., 2019), with obvious impact on the accuracy of  
94 numerical weather and ocean predictions; a correct account of skin SST diurnal variations in turn is crucial for  
95 flux calculations, which is already a very delicate problem also from an instrumental point of view.

96 Within these prognostic schemes, seawater's transparency conditions (e.g., estimated using chlorophyll  
97 concentration) have great implications in the way solar radiation is absorbed within the ocean's uppermost layer  
98 (Morel and Antoine 1994). Ohlmann et al. 2000 quantified with the help of radiative transfer calculations effects  
99 of physical and biological processes on solar radiation transmission, classifying as main factors chlorophyll  
100 concentration, cloud cover and solar zenith angle. Ohlmann and Siegel 2000 and Lee et al. 2005 are further  
101 examples of how radiative transfer models are used to develop solar transmission parameterization which is fit  
102 to the sum of exponentials (the number of terms in the sum depending on the variable which has been  
103 considered). To the best of our knowledge, these ideas have not been implemented nor tested within the  
104 prognostic scheme for skin SST present in the ocean model NEMO, which just relies on chlorophyll-calibrated  
105 coefficients though (Gentemann et al 2009).

106 Our main aim here is therefore to improve existing skin SST prognostic schemes, investigating the impact of  
107 variable seawater's transparency conditions in modeling solar radiation extinction in the upper ocean. The use  
108 of chlorophyll concentration as a proxy for seawater's transparency is not new. In fact, given its covariance  
109 with Secchi disk depth (estimated from reflectance at various wavelength), it has been often applied by the  
110 ocean color community to study the dynamics of oligotrophic gyres (Leonelli et al., 2022 and references  
111 therein). The paper is structured as follows: after this introduction, we describe the data and coupled modeling

112 system in section 2. The mathematical context in which we developed our new method, whose novelty stands  
113 in allowing the warm layer's extent to vary in space and time according to a chlorophyll-concentration  
114 climatology follows in section 3. In section 4 we present results, discussing them and drawing conclusions in  
115 section 5.

## 116 **2 Data and Modeling System**

117 We describe here the data and the coupled regional modeling system used in this study. Our description  
118 here is functional to the scope of this paper, and far from a complete depiction of each dataset. We refer readers  
119 to the documentation and relevant literature for detailed information on each dataset and model.

### 120 121 **2.1 Operational MED DOISST within CMEMS**

122 The MEDiterranean Diurnal Optimally Interpolated Sea Surface Temperature (MED DOISST) product,  
123 operationally distributed and freely available within the Copernicus Marine Environmental Service (CMEMS)  
124 provides gap-free (L4) hourly mean maps of sub-skin SST at  $1/16^\circ$  horizontal resolution over the Mediterranean  
125 domain, covering from 2019 to present. Sub-skin SST is defined as the temperature at the base of the cool skin  
126 layer, typically sensed by microwave radiometers, and representative of a depth of few millimeters from the  
127 ocean's surface (Minnet et al., 2019).

128  
129 This product combines satellite data acquired from the Spinning Enhanced Visible and InfraRed Imager  
130 (SEVIRI) and model data from the Mediterranean Forecasting System (MedFS), respectively used as  
131 observations and first guess for an optimal interpolation, giving a L4 field representative of subskin SST (see  
132 Pisano et al., 2022 and references therein). In all diagnostics involving these data (and presented in the following  
133 sections), regions where the percentage of valid SEVIRI measurements is lower than 50% have been masked  
134 out both in CMEMS MED DOISST and our experiments.

### 135 **2.2 iQuam in-situ data**

136 SST from drifter data were used for validation purposes and acquired from the iQuam (In situ SST Quality  
137 Monitor) archive (Xu and Ignatov, 2014). The iQuam provides high-quality and quality controlled (QC) in-situ  
138 SST data collected from various platforms, such as drifters, Argo Floats, ships, tropical and coastal moored  
139 buoys. iQuam SST data are also provided along with quality level flags ranging from 0 to 5, with 5  
140 corresponding to the highest quality level (Xu and Ignatov, 2014). For this study, SST with quality level equal  
141 five were selected from drifters only, since they provide the temperature measurement closest to the surface  
142 (compared to the other available instruments), ranging between 20-30cm (depending on the drifter type).

143  
144 Additionally, we interpolated model outputs on drifters' location in time and space. Table S1 outlines the  
145 number of available measurements for each given month and hour of the day. A total number of 555919 records

146 were available after the quality flag and platform selection, with the month of April being the most populated  
147 one, with 222996 measurements, and 10361 measurements at 9:00 am.

148

### 149 **2.3 EN4 objective analysis**

150 EN4, the quality controlled subsurface ocean temperature and salinity profiles and objective analyses, were  
151 used to assess the impact on the temperature vertical profiles. To facilitate the comparison, we made use of the  
152 objective analyses after bias corrections of Expendable Bathythermograph (XBT) calibrations (Gouretski and  
153 Reseghetti, 2010, Gouretski and Cheng, 2020), which give a gridded version of the dataset on a 1-degree regular  
154 grid. In the comparison, model outputs were interpolated on this grid.

### 155 **2.4 Mediterranean Chlorophyll concentration**

156 Chlorophyll data were used to estimate e-folding depths' seasonality (see Methods, Section 3). These data  
157 are a daily interpolation at 0.3 km horizontal resolution over the Mediterranean domain, and result from a  
158 merging between multiple sensors (MERIS - MEdium Resolution Imaging Spectrometer from ESA, SeaWiFS  
159 - Sea-viewing Wide Field-of-view Sensor and MODIS - Moderate Resolution Imaging Spectroradiometer from  
160 NASA, VIIRS - Visible Infrared Imager Radiometer Suite from NOAA, and most recently the Copernicus  
161 Sentinel 3A OLCI - Ocean and Land Colour Instrument), as detailed in the product description (see Volpe et  
162 al., 2019 and references therein for further details).

163

### 164 **2.5 ECMWF Atmospheric Reanalysis - ERA5**

165 We used heat fluxes (net solar radiation, latent and sensible heat fluxes, net thermal radiation) from ERA5  
166 at 0.25° horizontal and hourly temporal resolution (Hersbach et al., 2020) as reference for comparing  
167 performances across simulations with different skin SST schemes. Despite their possible biases in air-sea fluxes,  
168 atmospheric reanalyses today are still widely thought to provide the best gap-free and dynamically consistent  
169 reconstructions of the atmosphere system (Valdivieso et al., 2017, Storto et al., 2019).

### 170 **2.6 Mixed Layer Depth 1969-2013 Climatology**

171

172 Data from a mixed layer depth (MLD) climatology was used to test to what extent our modified scheme  
173 correctly represents the seasonality of the mixed layer.

174 This monthly gridded climatology was produced using MBT, XBT, Profiling floats, Gliders, and ship-based  
175 CTD (Conductivity, Temperature, Depth) data from different databases and carried out in the Mediterranean  
176 Sea between 1969 and 2013. As for the model outputs, MLD is calculated with a  $\Delta T = 0.1^\circ\text{C}$  criterion relative  
177 to 10m reference level on individual profiles (Houpert et al., 2015a, Houpert et al., 2015b).

178

## 179 2.7 ISMAR Mediterranean Earth System Model (MESMAR)

180 MESMAR is a newly developed coupled regional modeling framework for the Mediterranean region (Storto  
181 et al., 2023). MESMAR includes the following components:

- 182 • the ocean model: NEMO v4.0.7, with horizontal resolution of about 7 km, 72 unevenly spaced vertical  
183 levels (the first and the last levels being respectively about 0.5m and 200m thick) and a timestep of 7.5  
184 minutes (NEMO System Team, 2019);
- 185 • the atmosphere model: WRF v4.3.3, with 41 vertical hybrid levels and horizontal resolution of about 15  
186 km, covering the European branch of the international Coordinated Downscaling Experiment (EURO-  
187 CORDEX) domain, and a timestep of 1 minute (Skamarock et al., 2019);
- 188 • an interactive runoff model: HD v5.0.1, with a timestep of 30 minutes and 1/12° degree horizontal  
189 resolution over Europe (Hagemann et al., 2020);
- 190 • the coupler: OASIS3-MCT, coupling the three models with a coupling frequency of 30 minutes, and  
191 using the SCRIP library to interpolate fields between different model grids (Craig et al., 2017) ;

192 We report in figure 2 a graphical summary of different grids. Further details of its implementation, tuning, and  
193 performances are described in (Storto et al., 2023).

194

## 195 3 Methods

196 Many schemes to reconstruct the skin SST diurnal variations rely on the existence of a cool skin and a warm  
197 layer, respectively in the upper micrometers and few meters of the ocean, whose dynamics strongly depends on  
198 wind conditions and solar radiation extinction within the upper ocean. To explain the rationale behind the  
199 developments in our new method, we need to recap here some elements of this theory, which is mostly based  
200 on Zeng and Beljaars, 2005 (named ZB05 hereafter) work.

201 We start from the one-dimensional heat transfer equation in the ocean:

$$202 \quad \frac{\partial T}{\partial t} = \frac{\partial}{\partial z} (K_w + k_w) \frac{\partial T}{\partial z} + \frac{1}{\rho_w c_w} \frac{\partial R}{\partial z} \quad (1)$$

203 in which the subscript  $w$  refers to water properties,  $T$  is seawater temperature ( $K$ ),  $K_w$  ( $m^2 s^{-1}$ ) is the turbulent  
204 diffusion coefficient,  $k_w$  ( $m^2 s^{-1}$ ) is the molecular thermal conductivity,  $\rho_w$  ( $Kg m^{-3}$ ),  $c_w$  ( $J Kg^{-1} K^{-1}$ ) are respectively  
205 seawater density and heat capacity per unit volume,  $R$  ( $W m^{-2}$ ) is the net solar radiation flux, defined as positive  
206 downward.

207 **3.1 Cool Skin**

208 We assume that there exists an oceanic molecular sublayer of depth  $\delta$ , where  $K_w$  is negligible, and  
 209 temperature can be assumed constant in time, since it is always cooler than temperature of the underlying  
 210 seawater (Donlon et al., 2007, Zeng and Beljaars, 2005). Then integration of eqn. (1) gives,  $\forall z \in [0, -\delta]$

$$211 \quad k_w \frac{\partial T}{\partial z} + \frac{1}{\rho_w c_w} [R(z) - R_s] - k_w \frac{\partial^2 T}{\partial z^2} \overset{\mathcal{O}(z^2)}{\phantom{= const,}} = const, \quad (2)$$

212 where  $R_s$  is the net solar radiation at the surface (constant, open-ocean albedo, since the Mediterranean Sea is  
 213 an ice-free basin), assuming this constant to be the top boundary condition at  $z = 0$ :

$$214 \quad \rho_w c_w k_w \left. \frac{\partial T}{\partial z} \right|_{z=0} = Q = LH + SH + LW, \quad (3)$$

215 in which  $LH$ ,  $SH$ ,  $LW$  are respectively the surface fluxes of latent, sensible heat and net  
 216 long wave radiation.

217 Thus, eqn. (2) can be rewritten as

$$218 \quad \rho_w c_w k_w \frac{\partial T}{\partial z} = Q + R_s - R(z). \quad (4)$$

219 Making a further integration we get the cool skin temperature difference:

$$220 \quad T_s - T_{-\delta} = \frac{\delta}{\rho_w c_w k_w} (Q + f_s R_s), \quad (5)$$

221 where  $T_s$  and  $T_{-\delta}$  are respectively the temperature at the upper (air-sea interface) and lower limits of the cool  
 222 skin layer, while  $f_s$  is the fraction of solar radiation absorbed in this layer:  
 223

$$224 \quad f_s = \frac{1}{\delta} \int_{-\delta}^0 \left( 1 - \frac{R(z)}{R_s} \right) dz,$$

225 which depends on the way radiation gets absorbed within the cool skin. Being time-independent, the cool  
 226 skin temperature difference is a diagnostic variable in the scheme.

227 Eq. (5) is analogous to Saunders' model. Indeed, Saunders, 1967 was one of the first to construct a theory for  
 228 the ocean "cool skin" effect, i.e. the observed temperature at the air-sea interface is generally cooler than the  
 229 temperature of the water at about 10 cm depth, especially during nighttime. This effect takes place mainly  
 230 because of the transfer of energy between the ocean and the atmosphere, realized via heat loss and momentum  
 231 transfers (wind stress). In a nutshell, at the end of its derivation (Saunders, 1967), he obtains the following  
 232 expression for the temperature difference across the cool skin,  $\Delta T_c$  (K):

$$233 \quad \Delta T_c = \lambda \frac{Q_{v_w}}{k_w(\tau/\rho_w)^{1/2}}, \quad (6)$$

234 where  $\lambda$  is the Saunders' proportionality constant,  $Q$  ( $W m^{-2}$ ) has already been defined above,  $\tau/\rho_w$  ( $m^2 s^{-2}$ ) is the  
 235 kinematic stress (ratio between wind stress module and seawater density), and  $\nu_w$ , ( $m^2 s^{-1}$ )  $k_w$  ( $m^2 s^{-1}$ ) are  
 236 respectively the kinematic viscosity and thermal conductivity of seawater. Saunders' formulation was originally  
 237 conceived for low, nonzero wind conditions and neglecting the effect of solar radiation (which however  
 238 recognized its role and added a discussion on how to account for it in the model only at the end of his paper).  
 239 As noticed by Fairall et al. 1996, Artale et al., 2002 (named A02 hereafter), with a constant  $\lambda$ , eqn. (6) becomes  
 240 problematic in limiting cases of low and very high wind speeds (greater than  $7 m/s$ ), because the wind stress  
 241 in the denominator limits its validity. Thus, A02 proposed to include a wind dependence in Saunders' constant,  
 242 in order to still have a finite, nonzero cool skin to bulk temperature difference even when the wind speed goes  
 243 to zero or becomes very high. This scheme has proven to have good performances compared to other schemes  
 244 also on a mooring site in the Pacific Ocean (Tu and Tsuang, 2005).

245

### 246 3.2 Warm Layer

247 Below the skin layer, turbulent transfer is much more effective, and  $k_w$  can be neglected in favor of  $K_w$ .  
 248 Integrating eqn. (1) within the  $[-d, -\delta]$  layer, we get:

$$249 \quad \frac{\partial}{\partial t} \int_{-d}^{-\delta} T dz = \frac{Q + R_s - R(-d)}{\rho_w c_w} - K_w \left. \frac{\partial T}{\partial z} \right|_{z=-d}, \quad (7)$$

250 where  $d$  is a reference depth which can be assumed as the depth at which the diurnal  
 251 cycle can be omitted.

252 The turbulent diffusion coefficient can be expressed as (Large et al., 1994):

$$253 \quad K_w = k u_{*w} (-z) / \phi_t \left( \frac{-z}{L} \right), \quad (8)$$

254 in which  $k = 0.4$  is the Von Karman constant,  $z$  is negative in the ocean,  $u_{*w}$  is the friction velocity in the water  
 255 (this being the air friction velocity multiplied by the square root of air to sea density ratio), and the stability  
 256 function  $\phi_t$  discriminates between a stable and an unstable regime, depending on the sign of its argument, which  
 257 is the ratio of the vertical coordinate to the Monin Obukhov length  $L$ : positive for the stable and negative for  
 258 the unstable one. Assuming  $z$  to be negative in the ocean, the change of sign entirely depends on the Monin  
 259 Obukhov length, which is a length characterizing the prevalence of buoyancy variations induced turbulence  
 260 over the one generated by wind shear effects. This in turn is strongly dependent on the sign of the net heat flux  
 261  $Q$ . If  $Q > 0$ , i.e. the ocean gains heat from the atmosphere, and we have the stable regime: the diffusion  
 262 coefficients decrease with increasing depth, favoring the downward heat transfer within the water column. The  
 263 opposite case, which favors transfer of heat from the ocean to the atmosphere, can be modeled in different ways  
 264 (see While et al., 2017 and references therein).

265 Assuming a temperature dependence, for  $d \gg \delta$  of the form

$$266 \quad T = T_{-\delta} - \left[ \frac{z + \delta}{-d + \delta} \right]^\nu (T_{-\delta} - T_{-d}), \quad \nu \text{ empirical parameter} \quad (9)$$



267 eqn. (7) simplifies to

$$\frac{\partial}{\partial t} (T_{-\delta} - T_{-d}) = \frac{Q + R_s - R(-d)}{d\rho_w c_w} \frac{\nu + 1}{\nu} - \frac{(\nu + 1)k u_* w}{d\phi_t (d/L)} (T_{-\delta} - T_{-d}) \quad (10)$$

269 In ZB05 scheme (Zeng and Beljaars, 2005), eqs. (5, 10) are the coupled equations for the cool skin (diagnostic  
270 part) and warm layer (prognostic part) respectively. Being time dependent, the determination of the warm layer  
271 temperature difference at time  $t$  requires the knowledge of the one at the previous time step, and thus is the  
272 prognostic variable in the scheme. Assumptions on the fraction of solar radiation within the warm layer and the  
273 cool skin depth usually follow Fairall et al., 1996 parameterization, whose detail are given in the Supplementary  
274 Material section.

275

### 276 3.3 Solar transmission expression

277 The expression of the solar transmission in Zeng and Beljaars, 2005 is

$$\frac{R(-d)}{R_s} = \sum_{i=1}^3 a_i e^{-db_i}, \quad (a_1, a_2, a_3) = (0.28, 0.27, 0.45),$$

$$(b_1, b_2, b_3) = (71.5, 2.8, 0.07)m^{-1},$$

279 following Soloviev formulation (Soloviev, 1982) (S82 in the following), which is very widely used in  
280 atmosphere models (such as WRF, Skamarock et al., 2019).

281

282 So far this is not the only possibility: a formulation with 61 coefficients has been developed by Jerlov, 1968,  
283 which is based on different water types classified based on chlorophyll concentration and particulates, for light  
284 in the visible spectrum.

285

286 Formulations with 9 coefficients (reported in Table 2) have been proposed to include such effects: for example  
287 Soloviev and Schlüssel, 1996 use a different coefficient for the first term depending on Jerlov's optical water  
288 type, while Gentemann et al., 2009 include solar angle in the parameterization, keeping the value of the of the  
289 first coefficient as in the case of pure-water. Without knowing what the Jerlov water type is, what is currently  
290 implemented in NEMO is to take  $b_1$  as the average between coefficients for I, IA, IB, II and III Jerlov's optical  
291 water types. This formulation is widely employed in ocean models (such as in the optional skin SST routine of  
292 NEMO, see While et al., 2017), with the reference depth  $d$  fixed to 3 m. So, the solar transmission follows as:

$$\frac{R(-d)}{R_s} = \sum_{i=1}^9 a_i e^{-db_i} \quad (13)$$

294 Ideally, one would like to have a reference depth representative of the one at which the transmission of solar  
295 radiation is negligible, and if we take it as the depth at which transmission drops to  $1/e$  from its surface value,  
296 we get a value which can be different from  $d = 3$  m, as we can see from figure 3a. Allowing for a realistic time  
297 and space varying value of  $d$  represents the main novelty of our work.

298

299 From this viewpoint, choosing a value of  $d = 3 \text{ m}$  while using the solar extinction formulation as in Soloviev,  
300 1982 or Soloviev and Schlusel, 1996 would lead to underestimating the penetration of solar radiation into the  
301 warm layer. Another possibility, which constitutes our modification to the scheme already implemented in  
302 NEMO, is to reconstruct a chlorophyll profile from its surface values following what is already implemented  
303 in the NEMO module for radiation calculations (Jerlov, 1968, Morel et al., 1989, Lengaigne et al., 2007), and  
304 employ an R-G-B+Chl-a scheme to calculate radiation as a function of depth. Then, from eqn. (13) with only  
305 4 terms (one for chlorophyll, and three for R-G-B, expressed in lookup tables), one can numerically derive the  
306 warm layer reference depth as the e-folding depth of the light extinction profile (see Fortran source files in the  
307 [Zenodo](#) repository, de Toma (2024)).

308

309 This would give a constant transmission throughout the basin, but with a spatially and temporally varying e-  
310 folding depth and defines our new prognostic scheme for skin SST warm layer calculation, thus embedding in  
311 it the ocean color information coming from Chl-a. Everything else is left unchanged, both the refinements of  
312 Takaya et al., 2010, which include the effect of Langmuir circulation and a modification of the Monin-Obukhov  
313 similarity function under stable conditions (T10 hereafter), and the A02 model for cool skin, which has been  
314 demonstrated to improve the scheme respectively under wavy and windy conditions.

### 315 **3.3.1 E-folding depth estimates**

316 Mediterranean Chlorophyll climatology data (see section 2.4) were re-gridded onto a  $0.25^\circ$  regular  
317 longitude/latitude grid, and tabulated coefficients within NEMO were used to retrieve the transmission,  
318 accounting for chlorophyll variations. E-folding depths then can be estimated as the depth at which transmission  
319 drops to  $1/e$  from its surface value. It can be noticed from figure 3b that also the e-folding depth varies with  
320 seasonality, with typical values ranging from about 3 to 4.5 meters. This is the central point of our modification  
321 to the prognostic scheme. In our setup we extracted pixelwise and at each time step of the NEMO model the e-  
322 folding depth used within the prognostic scheme.

## 323 **3.4 Overview of the simulations performed.**

324 With the coupled ocean-atmosphere regional system we performed a set of four simulations, forced by  
325 ERA5 in the atmosphere and ORAS5 (Zuo et al., 2018) in the ocean and covering three years (from 2019 to  
326 2021), with hourly outputs (a synthesis is provided in Table 1). In cases where a skin SST scheme is active, we  
327 substitute the SST, i.e. temperature on the first NEMO level, with the skin SST coming out from the scheme:

- 328 1. a control run, in which no skin SST prognostic scheme is activated, therefore the diurnal SST variations  
329 in the uppermost ocean layer (0.5 m thick) only come from the variability represented by the ocean model  
330 at about 0.5 m of depth, considering also the 0.5 hours frequency of the coupling. We will refer to this  
331 experiment in the following as *ctrlnoskin*;

- 332 2. a run in which the ZB05 scheme in WRF (Zeng and Beljaars, 2005) is active – we shall refer to this case  
333 in the following as *wrfskin*;
- 334 3. a run in which the existing scheme within NEMO, which employ the 9-coefficient parameterization for  
335 light extinction coefficients (Gentemann et al., 2009 – G09 hereafter), the scheme for the cool skin as  
336 modified in A02, and refinements of the stability function, in the warm layer formulation as in T10 – we  
337 shall refer to this as the *nemoskwrite* case;
- 338 4. a fourth simulation in which we modified the reference depth for the basis of the warm layer from  $z = 3$   
339  $m$ , to an e-folding depth (i.e. the depth at which radiation gets diminished by  $1/e$  from its surface value),  
340 which is allowed to vary temporally and spatially because it is estimated from R-G-B light extinction  
341 coefficients and chlorophyll concentration (see section 3). We will refer to it as *modradnemo*, being the  
342 experiment where our modification to the skin SST scheme is implemented and tested.

343 The reason behind the choice of the above mentioned period of three years 2019-2021 is twofold: firstly, it  
344 allows a validation against all the measurements from different data sources (satellite, drifters and objectively  
345 analyzed profiles), and secondly, it is a good trade-off between the needs of keeping a reasonable computational  
346 load, data volume for the analysis, and guarantees a minimal robustness of our finding, compared to a simulation  
347 which covers just one year. However, we do not discard the possibility to extend the time coverage in our plans  
348 for future works.

## 349 4 Results

350 In this section, we present skill scores against satellite, drifters and temperature profiles data (see section 2)  
351 from the set of the simulations performed, aimed at characterizing the impacts of our modified skin SST scheme.  
352 Since we are mainly acting to improve skin SST diurnal variations reconstruction in the ocean component, the  
353 main focus is on the difference between the *nemoskwrite* and *modradnemo*, while the *ctrlnoskin* and *wrfskin*  
354 ones are included as further reference elements (the latter being not directly comparable because the  
355 atmospheric model sees the ocean foundation SST and employ the scheme just to diagnose the skinSST).  
356

### 357 4.1 Comparison with CMEMS MED DOISST

358 We calculated the mean diurnal warming amplitude in each season as the seasonally averaged diurnal  
359 warming amplitudes (diurnal warming amplitude being defined for each day as the difference between daytime  
360 maximum and nighttime minimum of SST), which can be cast into the following equation:

$$361 \langle \text{DWA} \rangle_{\text{seas}} = \frac{1}{N_{\text{seas}}} \sum_{i=1}^{N_{\text{seas}}} \left\{ \max_{h_i \in [10:00, 18:00]} \text{SST}(h_i) - \min_{h_i \in [00:00, 06:00]} \text{SST}(h_i) \right\}, \quad (14)$$

362 where seas = DJF, JJA, MAM, SON is the given season,  $N_{seas}$  is the number of days in that particular season  
363 and  $h_i$  is the local time in hours for any given day.

364 Seasonally averaged diurnal warming amplitudes are shown in figure 4. On average, the maximum amplitude  
365 is reached in summer, with the wrfskin simulation peaking at about 3 K, thus overestimating the mean diurnal  
366 cycle compared to CMEMS MED DOISST (the monthly biases with respect to CMEMS foundation SST both  
367 in the western and the eastern part of the Mediterranean Sea stay below 1 K year-round for all of the simulations  
368 performed – see figure S1 in Supplementary Materials). The nemoskwrite simulation yields a pattern very  
369 similar to CMEMS MED DOISST in summer, but underestimates the signal in the remaining seasons. Outside  
370 the Summer season, our modifications yield a slight improvement (see modradnemo, last row of figure 4). As  
371 expected, the control run in which no skin SST method is active, generally underestimates the diurnal signal  
372 everywhere. Compared to nemoskwrite, the modradnemo simulation improves JJA mean diurnal warming  
373 amplitude, especially over the Southern Mediterranean Sea, while in central and Northern part of the basin  
374 tends to overestimate the signal by about 0.5 K with respect to CMEMS-DOI data. Furthermore, a general  
375 underestimation is present also in DJF, with the modradnemo simulation showing the smallest differences with  
376 respect to CMEMS-DOI data.

377 The spatial average over the whole Mediterranean domain is shown in figure 5, confirming the general  
378 underestimation of the control run and the overestimation of the wrfskin and modradnemo in all seasons except  
379 winter.

380 Computing spatial averages highlights that modradnemo slightly improves the mean diurnal warming  
381 amplitude signal during wintertime, while in all the other seasons the best agreement is gained by using the  
382 nemoskwrite setup (ZB05 with T10 and A02 modifications), at least according to the verification against  
383 CMEMS MED DOISST.

384 On a monthly timescale, figure 6 confirms that the control simulation tends generally to have a negative bias  
385 of the diurnal amplitude, for the whole simulated period. The wrfskin (ZB05 scheme) shows a warm bias during  
386 summertime months, shown just as a reference. The comparison between nemoskwrite (ZB05+A02+T10) and  
387 modradnemo (chl e-folding depth) shows improvement of our scheme (modradnemo) over the old one  
388 (nemoskwrite) especially in May, but not in April, June, July, August and September, despite in the rest of the  
389 period the amplitude of the bias is slightly reduced.

## 390 **4.2 Comparison with iQuam Star HR-Drifters**

391 The bias with respect to drifter measurements averaged over drifters positions as a function of the month  
392 and time of the day is shown in figure 7. All the schemes present a systematic cool bias in autumn (SON) for  
393 most of the hours of the day. During April and June, the modradnemo simulation significantly reduced the

394 warm bias with respect to observations, compared to the nemoskwrite case, keeping it however generally  
395 positive. This is quite reasonable, since drifters measurement can be thought representative of a depth which  
396 can also be below the subskin level (typically of the order of some centimeters). Consistently with figure 6, the  
397 wrfskin has a larger positive bias than modradnemo in June.

398  
399 Further, as shown by figure 8, the bias between CMEMS MED DOISST and drifters is generally positive  
400 anytime except in late spring/summer and autumn during nighttime. This pattern arises because of the  
401 composite effect of having a temperature representative of the subskin level where and when there are data  
402 from radiometers, and a temperature of about 1 *m* depth from the MEDFS system as first guess of the optimal  
403 interpolation over cloudy regions (Pisano et al., 2022). However, the modified scheme significantly reduces the  
404 difference, yielding a bias closer to the one of CMEMS MED DOISST with respect to drifters, especially during  
405 April, which is the month in which the number of observations from drifters is definitely larger.

### 406 **4.3 Comparison with EN4 objective analysis**

407 Bias corrected vertical profiles gathered in an objective analysis were used to assess differences across  
408 schemes along the water column. To summarize we report here only a macro subdivision into the eastern and  
409 the western Mediterranean Sea, respectively in figures 9, 10. Model outputs were remapped on the same vertical  
410 and horizontal grid. Looking at the mean profile averaged over all grid points in the given area, the agreement  
411 is better for all simulations during summertime months, both for the eastern and the western region (see figs.  
412 9c, 10c), showing in particular that the modradnemo simulation outperforms the nemoskwrite one. This is also  
413 true for the wintertime season in the eastern Mediterranean (see fig.9b). On the other hand, in the western  
414 Mediterranean all simulations tend to overestimate the signal, with our modified scheme doing a better job with  
415 respect to the nemoskwrite case, with an average profile which is about 0.4°C closer to the EN4 profile.  
416 However, below about 80 *m* depth differences across schemes vanish.

417  
418 Looking in more detail at the RMSE on the top 15 *m* depth between each simulation and EN4 as a function of  
419 the month and more detailed region subdomains shown in figure 11a, we can see how in general all simulations  
420 present the same pattern for the region outside of Gibraltar Strait, which can be thought an effect related to the  
421 presence of the relaxation to horizontal boundary conditions, while for all the remaining regions and months  
422 the control run, the wrfskin and the modradnemo present a similar pattern, with the modradnemo reducing the  
423 RMSE in most of the regions and for most of the months, especially with respect to nemoskwrite, and this is  
424 particularly true over the central Mediterranean Sea, in regions like Thyrrenian and Adriatic Seas.

### 425 **4.4 Heat fluxes and vertical propagation**

426 In this section we aim to characterize the differences of each scheme with respect to the control simulation.  
427 We do this by specifically looking at the seasonality of Mixed Layer Depth (MLD), vertical profiles of

428 temperature in specific months and regions, and via the comparison of the net surface heat fluxes over the whole  
429 Mediterranean Sea.

430  
431 Compared to the Mixed Layer Depth climatology from 1969 to 2013 (Houpert et al., 2015a, Houpert et al.,  
432 2015b, section 2.7), all of the tested schemes seem to have a similar impact on Mixed Layer Depth's seasonality,  
433 with larger differences with respect climatological values being mostly located in the Eastern Mediterranean  
434 Sea and during wintertime/spring (Figure 12). It may seem from this picture that there's not such a huge change  
435 to prefer one method over the other considered in this paper, and this may also be because of the short period  
436 simulated (2019-2021). Figure 13 show how our modified scheme allows more (less) vertical propagation of  
437 the diurnal signal during summer (winter) with respect to schemes with constant e-folding depth in all central  
438 regions of the Mediterranean domain (regions 2, 3, 4 as defined in figure 11a), when all of them are referenced  
439 to the control simulation temperature daily minimum.

440 Indeed, from figure 13b, we can see that when all the temperature profiles for each simulation are referenced  
441 to the ctrlnoskin daily minimum, there is a much wider diurnal warming signal for most of all the considered  
442 depths level, with modradnemo representing an intermediate situation between the wrfskin and the nemoskwrite  
443 simulation. This is probably due to the inclusion of chlorophyll-interactive variations, which allow for a better  
444 representation of the variability of the mixed layer dynamics.

445 Estimates of the mean Mediterranean heat exchange between ocean and atmosphere based on previous studies  
446 range from -11 to +22  $W/m^2$ , with an evident dominance of negative estimates, i.e., heat loss from the ocean to  
447 the atmosphere (Jordà et al., 2017, Pettenuzzo et al., 2010). Some other studies suggest that the Mediterranean  
448 heat budget is close to a neutral value, -1  $W/m^2$  (Ruiz et al., 2008) or +1  $W/m^2$  (Criado et. al., 2012). Many  
449 factors can contribute to such wide variability among different estimates, such as differences in the  
450 parameterizations employed, initial and boundary conditions, and the way the physical processes, especially  
451 through the Strait of Gibraltar are modeled (Macdonalds et al., 1994, Gonzales, 2023).

452  
453 As shown by table 3, all simulations on an annual basis give a negative, non-closed balance for the net surface  
454 heat flux, and modifications to include skinSST, performing very similarly one to another, bring the budget by  
455 1.5  $W/m^2$  closer to zero, while ERA5 data show a positive net surface heat flux close to 5  $W/m^2$ . However, all  
456 estimates fall into the (large uncertain) literature-based estimates. On seasonal timescales, the inclusion of  
457 skinSST diurnal variations has the following effects:

- 458 • less net heat loss to the atmosphere during wintertime with respect to the control run (wrfskin differing  
459 from the ctrlnoskin by about 6 $W/m^2$ , while nemoskwrite and modradnemo having a similar impact, with  
460 a difference of about 4 $W/m^2$  with respect to the control run);
- 461
- 462 • in springtime, all simulations show a positive imbalance, with the highest difference with respect to the  
463 control run of about 1  $W/m^2$  in the modradnemo simulation;

- during summer, our modified scheme brings on average about  $3 \text{ W/m}^2$  more than the control simulation into the basin, yielding an estimate which is closer to ERA5;
- in autumn, our scheme cools down more than the control (about  $2 \text{ W/m}^2$ ), being the farthest simulation from ERA5 estimate, while traditional schemes tend to have a less negative net heat input.

All seasons except spring show larger difference with respect to ERA5 fluxes, with underestimation in summer, and overestimation during winter and autumn, resulting in a bias of about  $10 \text{ W/m}^2$  with respect to the net heat flux annual budget in ERA5.

## 5 Summary and Conclusions

In this paper we studied the sensitivity of a regional coupled ocean-atmosphere-hydrological discharge regional model on the Mediterranean Sea to prognostic schemes for skin sea surface temperature. Specifically, we developed a new scheme which allows for spatial and temporal variations of the warm layer's extent according to seawater's transparency conditions. This is possible by using tabulated solar extinction coefficients already used in the ocean model, and inverting the functional form which determines how the solar radiation varies along the vertical direction to find the depth at which this latter drops to  $1/e$  from its surface value.

We simulated the period 2019-2021, analyzing hourly model outputs, and comparing aggregated results with satellite, objectively analyzed and drifters data. Overall, the comparison with data shows that the new scheme improves what is already implemented in NEMO, e.g. mean diurnal warming amplitudes are closer to satellite observations in winter, spring and autumn, not being much worse than other existing schemes in summer, at least looking at maps of mean diurnal warming amplitude grouped by seasons. Looking to the typical temperature profile in both the eastern and the western Mediterranean Sea, non-negligible differences across schemes stay confined in the topmost  $20\text{m}$  ( $100\text{m}$ ) of depth during summertime (wintertime). Regionally, typical profiles are warmer than EN4 observation year-round for western regions (regions -1,1,2) especially in winter, while regions in the east show a smaller RMSE in the topmost meters for basically all the regions and months when comparing modradnemo to nemoskwrite. The Adriatic Sea has a systematically higher RMSE with respect to EN4 in all the tested methods, for the whole period considered. In the central regions, the new scheme penetrates temperature anomalies more (less) during summer (winter) months, having a less intense mean diurnal warming amplitude signal in summer, especially over the upper few meters (the converse holds for wintertime values). Therefore, with respect to the ctrlnoskin simulation, nemoskwrite shows the coldest signal, the wrskin the hottest, and our modification modradnemo constitutes the middle situation, with milder summer and winter than the control run. Our interpretation is that within modradnemo, the chl-interactive e-folding depths allow, where and when necessary, the warm layer to become a little deeper than in the already existing scheme (nemoskwrite), depending on chl-variations. For these space-time points, solar penetration is increased and so it tends to make warmer the warm layer. Therefore, future research efforts should be devoted

498 to the better characterization of this aspect, especially to understand if the modified vertical penetration of heat  
499 has some particular effect on the dynamics of the mixed layer (see Song and Yu, 2017 and references therein).  
500 Furthermore, testing the implementation within NEMO of more sophisticated radiative transfer models (such  
501 as the one of Ohlmann and Siegel 2000), or the development of deep learning based parameterizations are  
502 underway as future research efforts. On a long-term perspective, the method needs to be tested also in other  
503 areas and for longer periods, which can increase the results' certainty and allow for usage in investigating  
504 impacts on relevant climate large-scale phenomena, where the role of an improved diurnal warming signal  
505 could be more relevant (Bernie et al., 2007, Bernie et al., 2008). These includes phenomena and physical  
506 processes such as propagation of Marine Heat Waves (MHW) or deep water formation and deep convective  
507 events.

508

509



510 *Code and data availability*

511

512 The NEMO ocean model code (v4.0.7) is available at <https://forge.ipsl.jussieu.fr/nemo/wiki>.

513

514 The WRF atmospheric model code (v4.3.3) is available at <https://github.com/wrf-model/WRF>.

515

516 The HD hydrological discharge model (v5.1) is available at <https://zenodo.org/record/5707587#.Y-0VQ3bMKUk>.

518

519 The frozen version of the MESMARv1 code used in this manuscript is available at:  
520 <https://doi.org/10.5281/zenodo.7898938>.

521

522 CMEMS MED DOISST Data downloaded from [CMEMS portal](#).

523

524 Chlorophyll data are freely available from [CMEMS portal](#).

525

526 The iQuam data version of this study used is V2.1, downloaded from the National Environmental Satellite,  
527 Data, and Information Service Satellite Applications and Research [NOAA NESDIS STAR portal](#).

528

529 Gridded analyses of EN4 profiles are distributed from the [MetOffice Hadley Centre Observations](#) (we used  
530 version 4.2.1).

531

532 ERA5 data are freely available after registration on the [Climate Data Store \(CDS\) by Copernicus Climate](#)  
533 [Change Service \(C3S\)](#).

534

535 MLD data are distributed on a 0.25 degree regular grid, and freely available from the [Sea Open Scientific Data](#)  
536 [Publication SEANOE portal](#).

537

538

539 Minimal data and scripts used within the manuscript to reproduce the figures in the manuscript are available at  
540 this link:

541 <https://zenodo.org/records/10451206>

542

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549

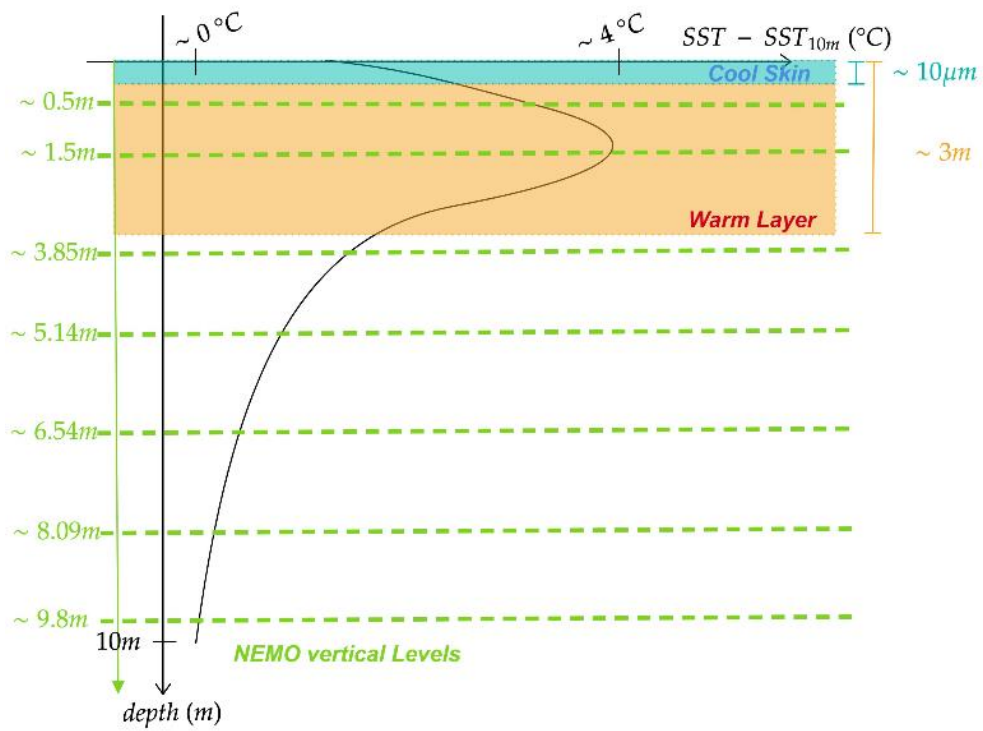
550 *Author Contributions.* VdT and AS conceived the study and designed the experiments to conduct, VdT  
551 Performed the simulation and data analysis, data downloading and writing of the first draft, VdT, DC, YH, CY,  
552 VA, AP, DC, RS and AS equally contributed to discuss and interpret the results, finalizing the draft.

553

554 *Competing Interests.* All authors declare they have no competing interests.

555 **Figures**

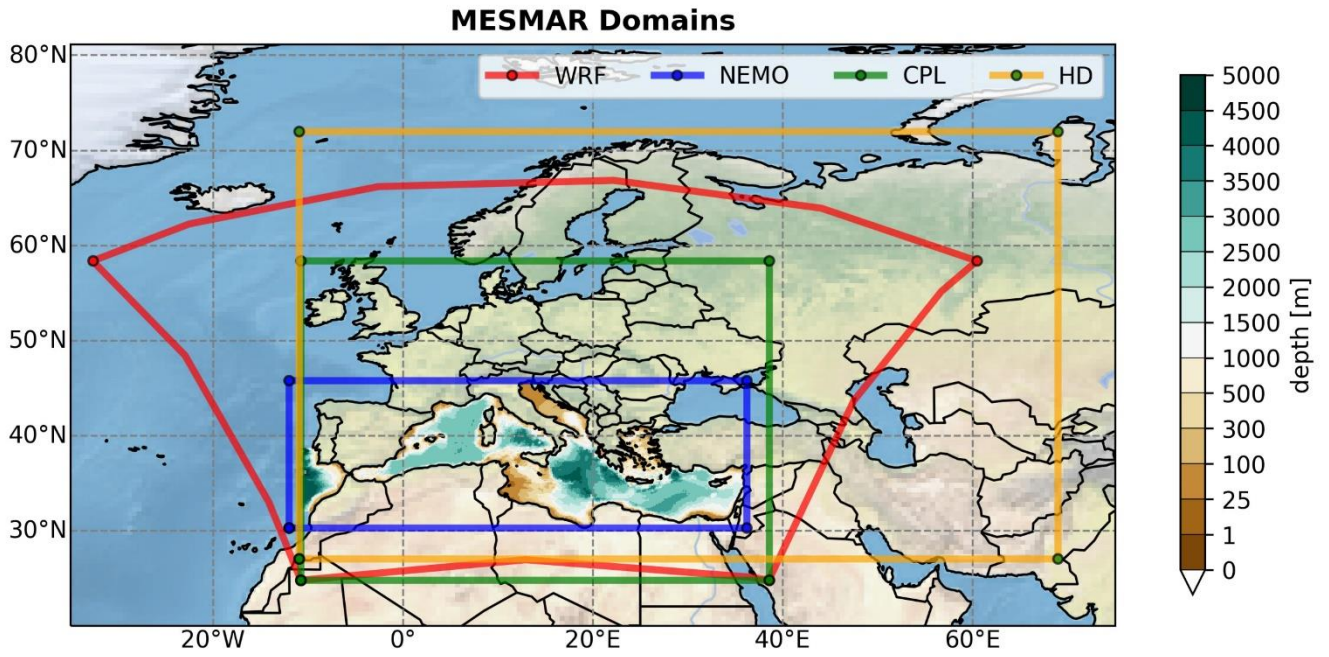
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558 **Figure 1.** Sketch of the cool skin and warm layer adapted from Donlon et al., 2007. Vertical discretization  
559 of NEMO levels is shown in green (not perfectly in scale with the underlying y-axis).  
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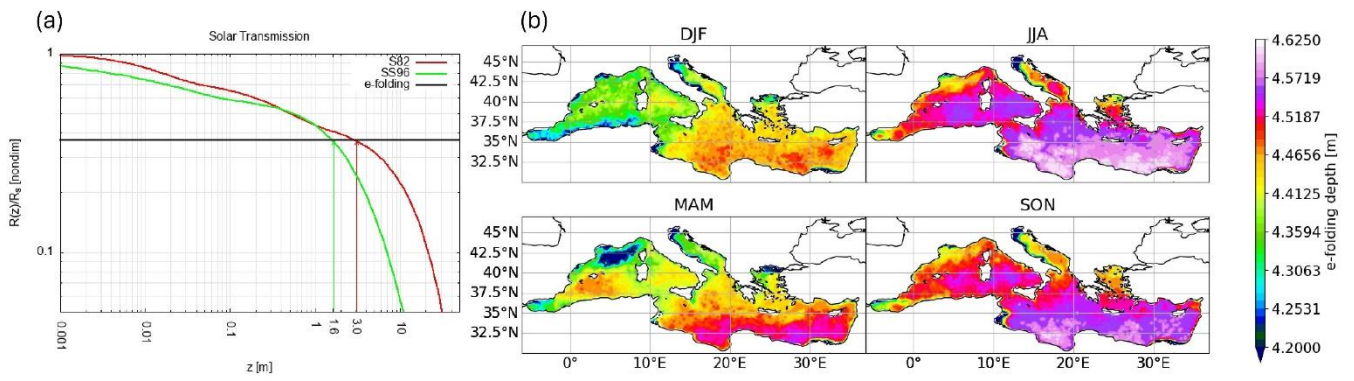
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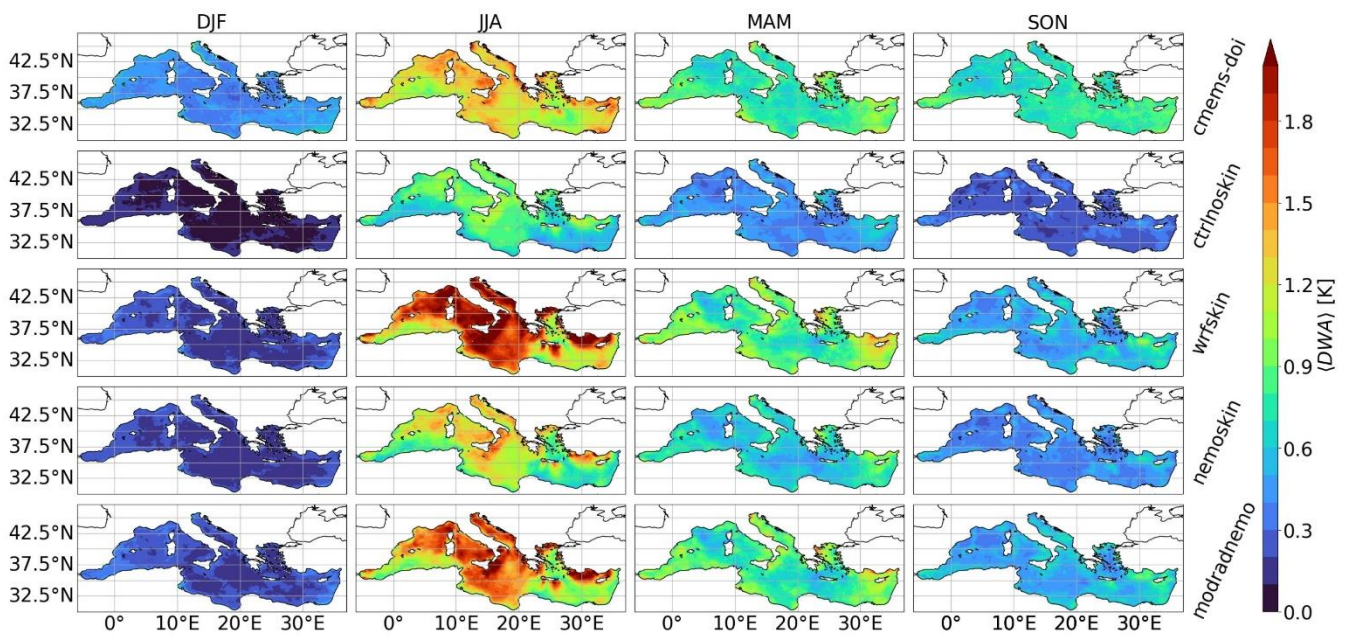
564 **Figure 2:** The modeling system domain: WRF, NEMO, HD and boundaries for the coupling mask are  
565 respectively in red, blue, orange, and green. Contour filled plot shows the ocean model bathymetry.

566



569 **Figure 3:** Panel 2a shows two different formulations frequently used for the transmission coefficient expression:  
 570 the red curve shows the formulation of Soloviev, 1982, while the green curve the one defined in Soloviev and  
 571 Schlussel, 1996. Panel 2b shows e-folding depth estimates from Mediterranean Chlorophyll climatology of  
 572 Volpe et al., 2019: lowest values touch the 2.5 meters.

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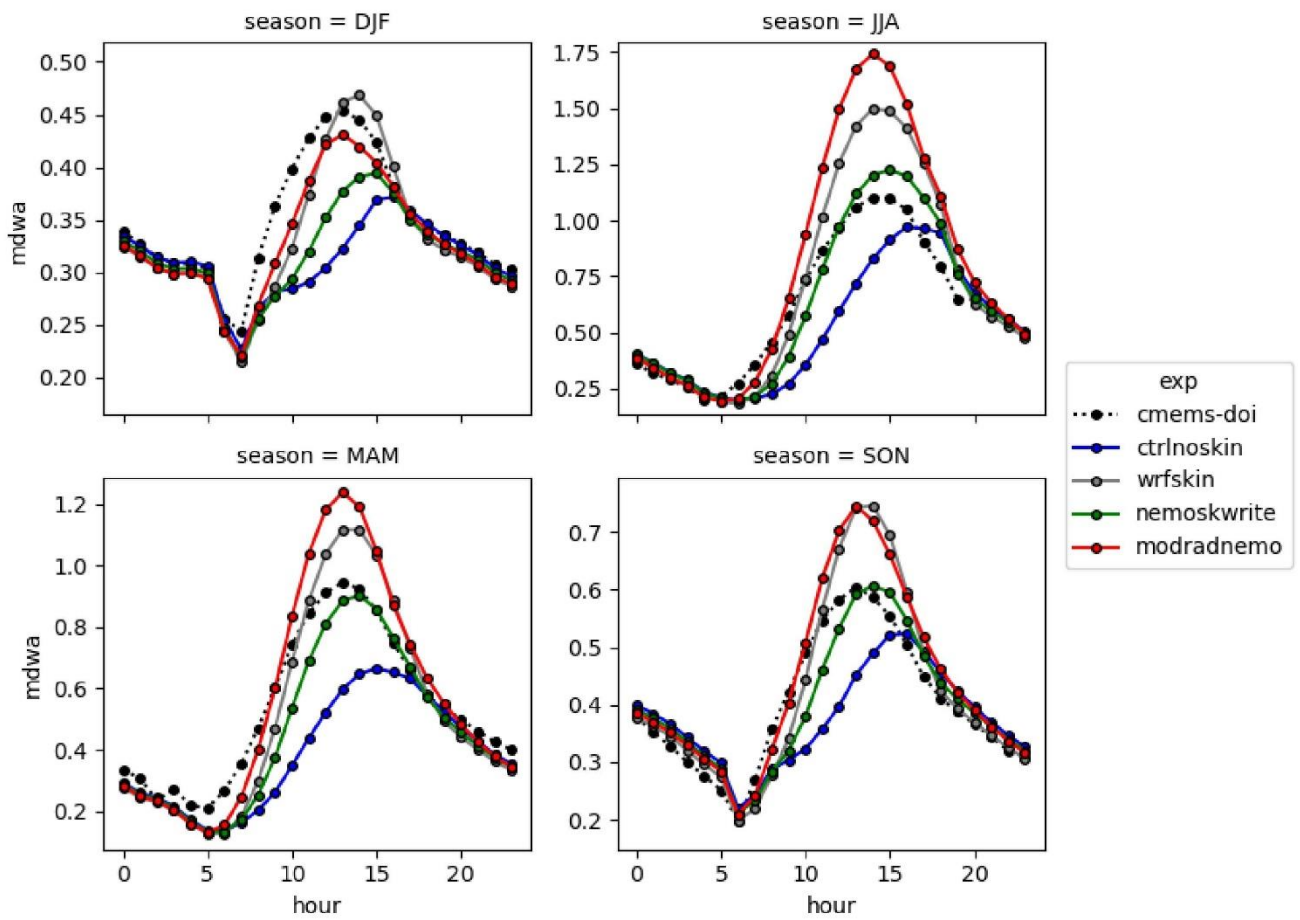
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576 **Figure 4:** Mean diurnal warming amplitude averaged over seasons (on columns), for each case (row): the first  
577 row is the CMEMS MED DOISST data, followed in order by the control simulation, wrfskin, nemoskwrite and  
578 modradnemo.

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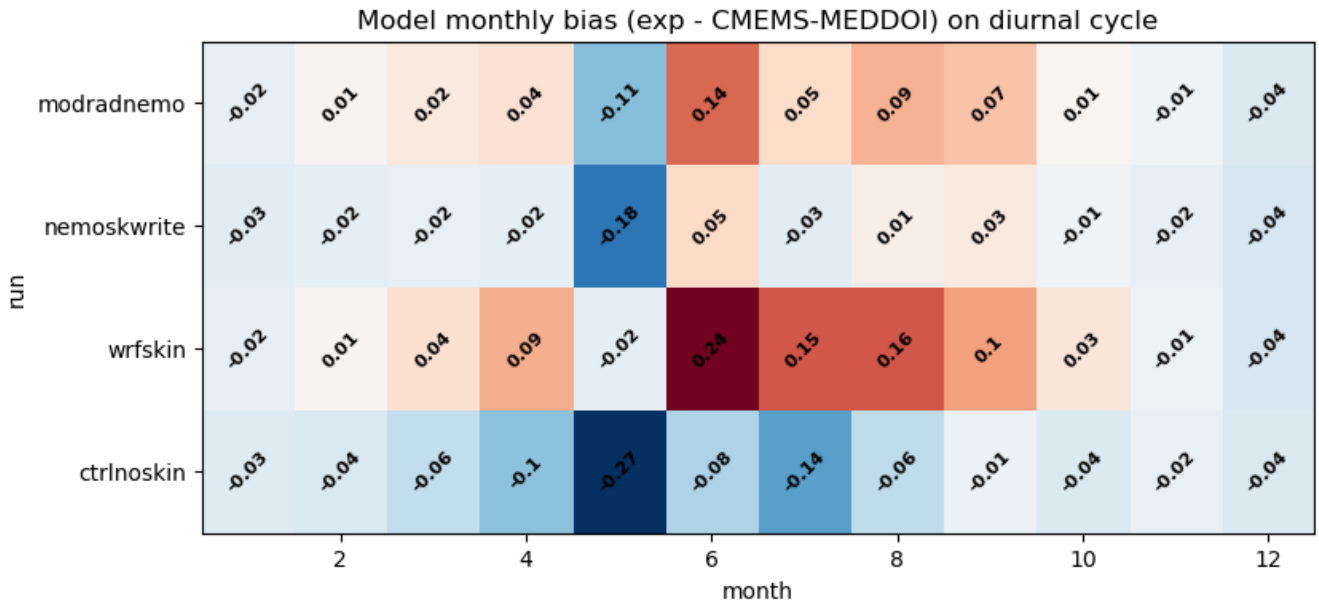
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583 **Figure 5:** Seasonality of the diurnal cycle averaged over the whole Mediterranean Sea, masking out regions in  
584 time and space where the percentage of model data in CMEMSDOI is greater than 50%.  
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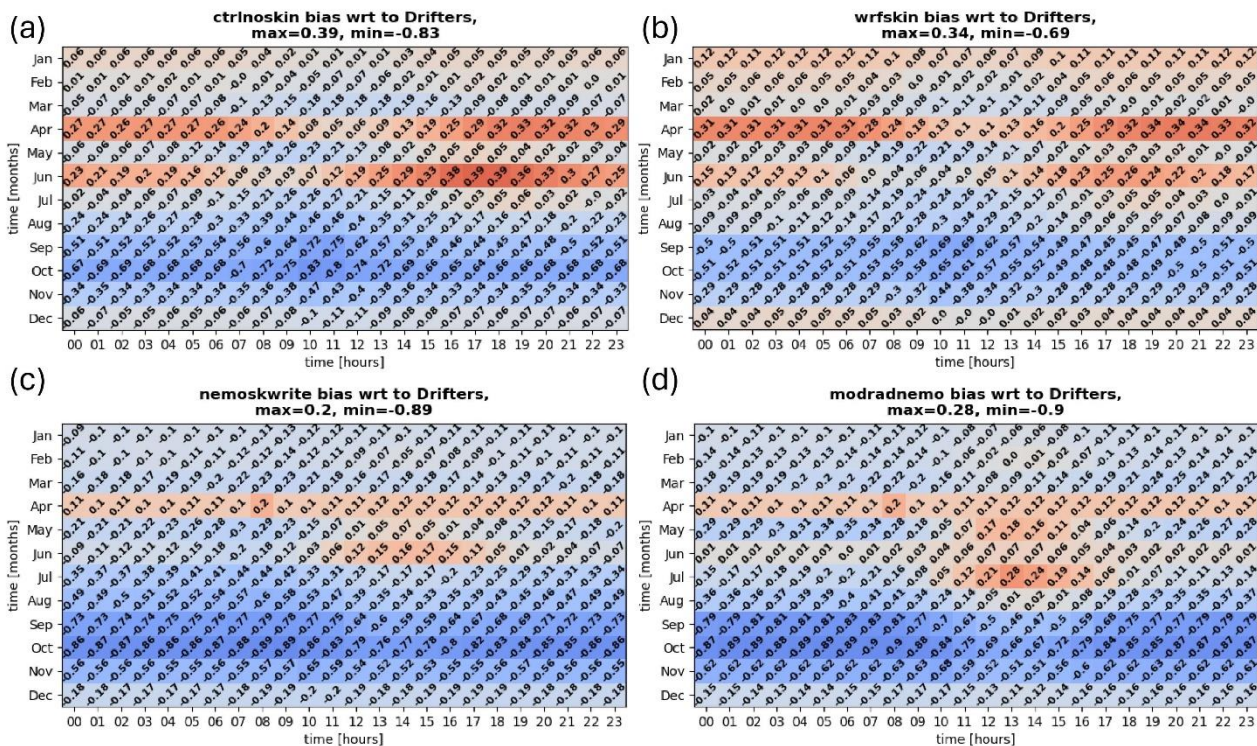
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589 **Figure 6:** Monthly averaged values for the time series of spatial mean diurnal cycle over the Mediterranean Sea  
590 (bias with respect to CMEMS MED DOISST)





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593

**Figure 7:** Bias with respect to measurements averaged over drifters' locations as a function of the month and the time of the day. Panels 6a, 6b, 6c, 6d show respectively the results for all the simulations carried out in the present study. Confidence on these numbers can be supported by the numbers of measurements reported in table S1.

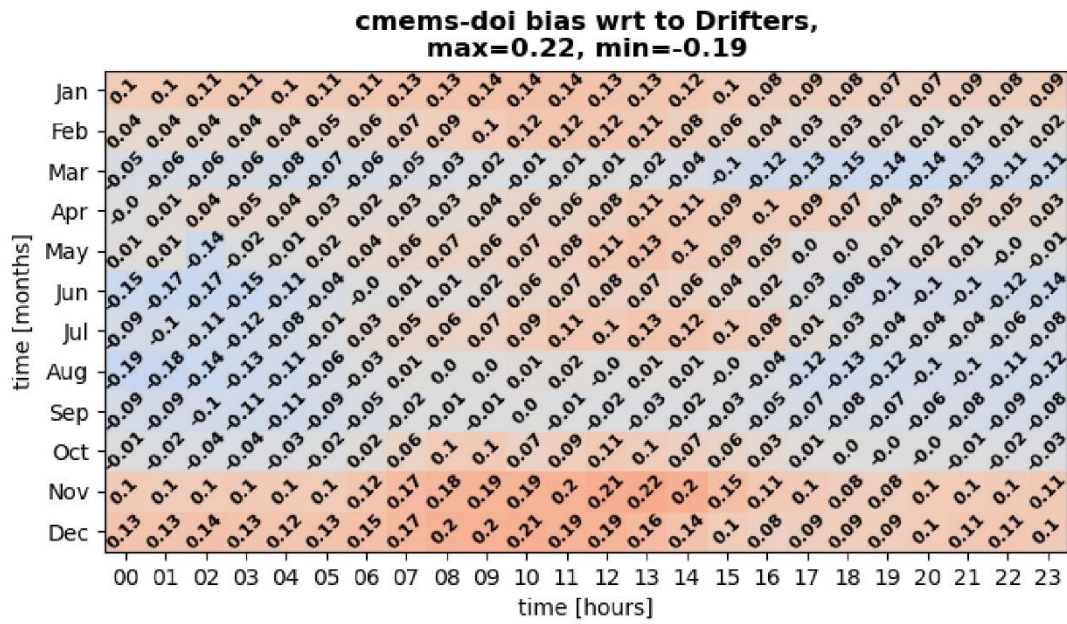
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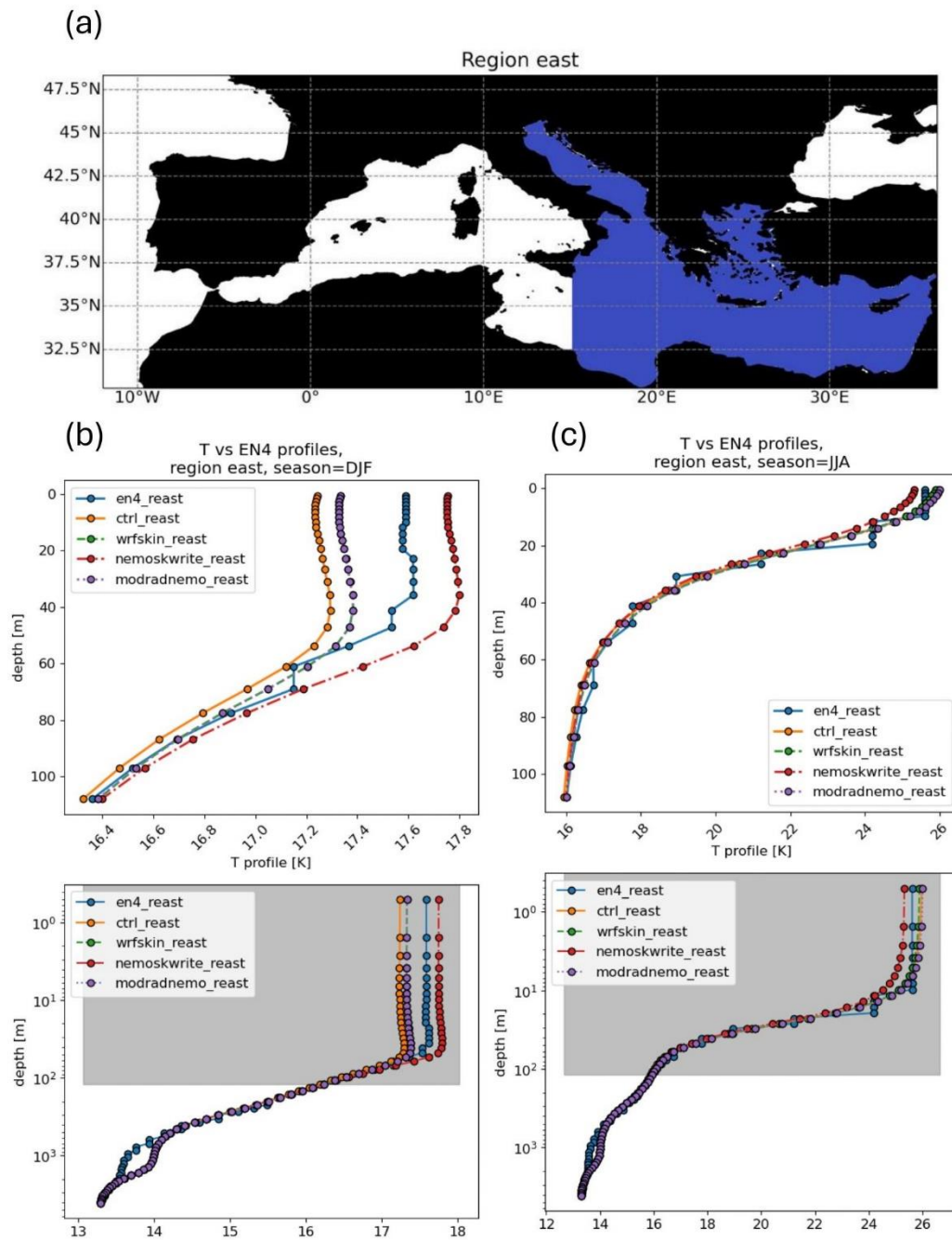
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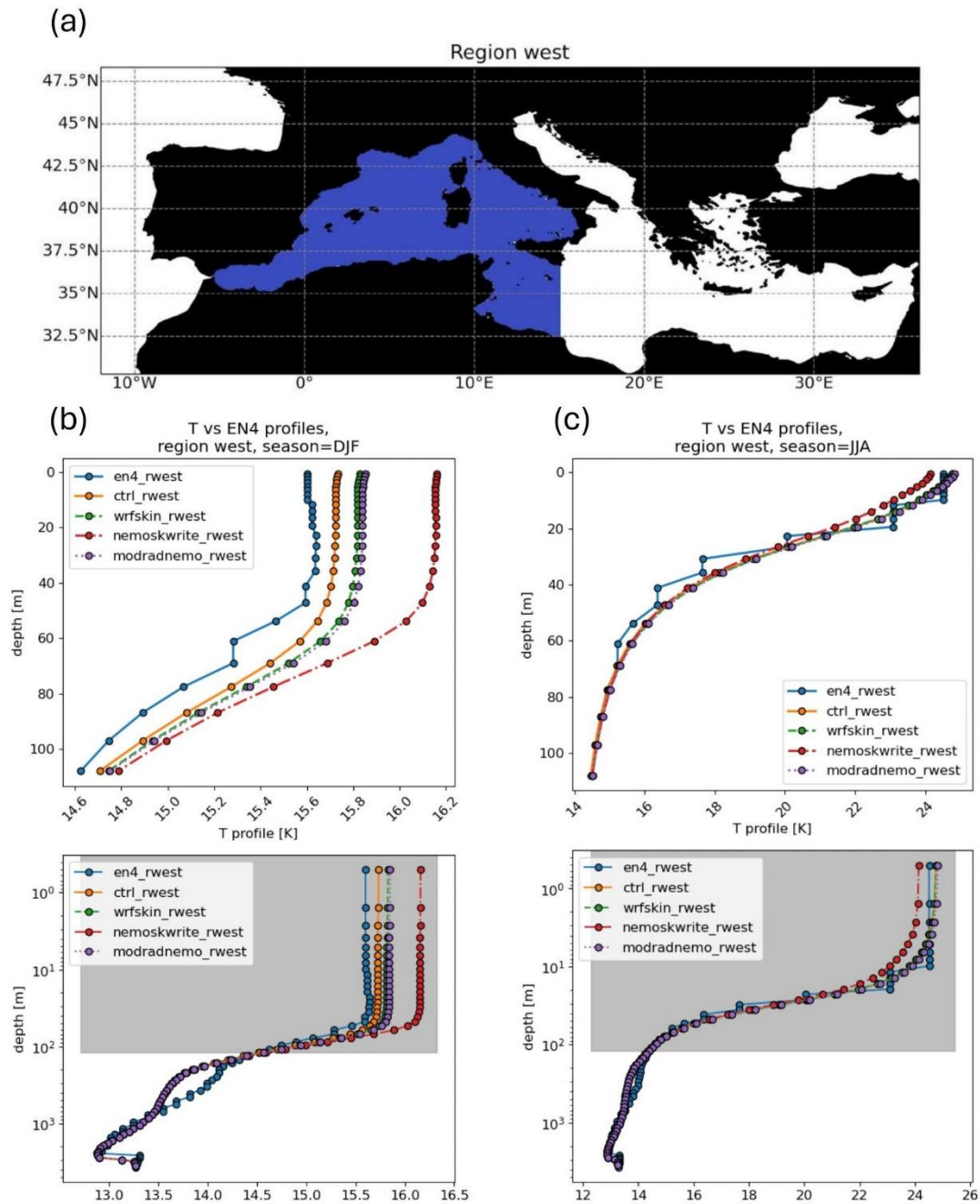
601 **Figure 8:** Bias with respect to measurements averaged over drifters' locations as a function of the month and  
 602 time of the day, for CMEMS MED DOISST data.

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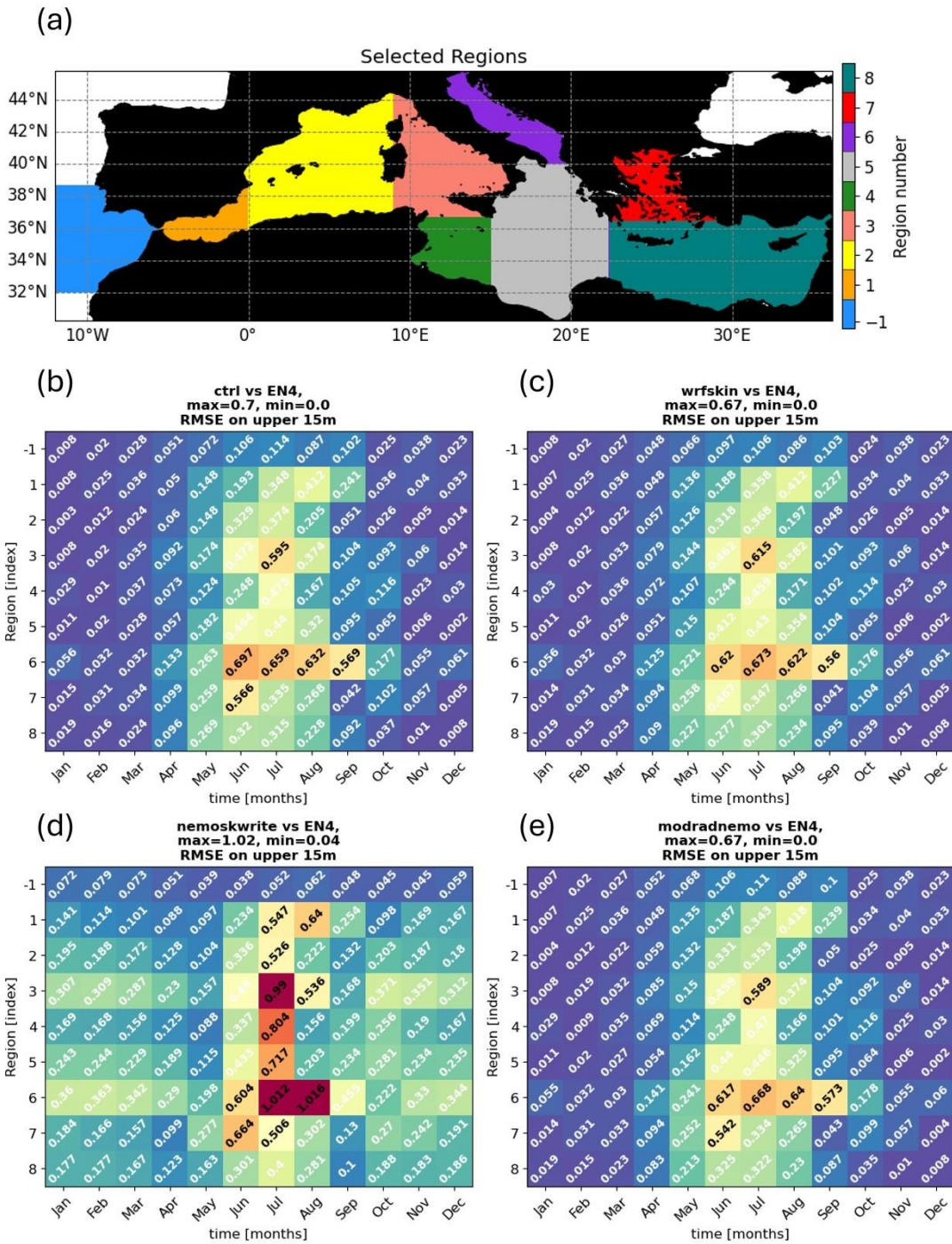
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**Figure 9:** Spatial average of profiles within the eastern Mediterranean Sea, during winter and summer. Panel 9a shows the eastern region, while 9b, 9c show respectively wintertime and summertime spatially averaged profiles within the top 100 m in the upper part, on the bottom the whole depth range on a logarithmic scale.

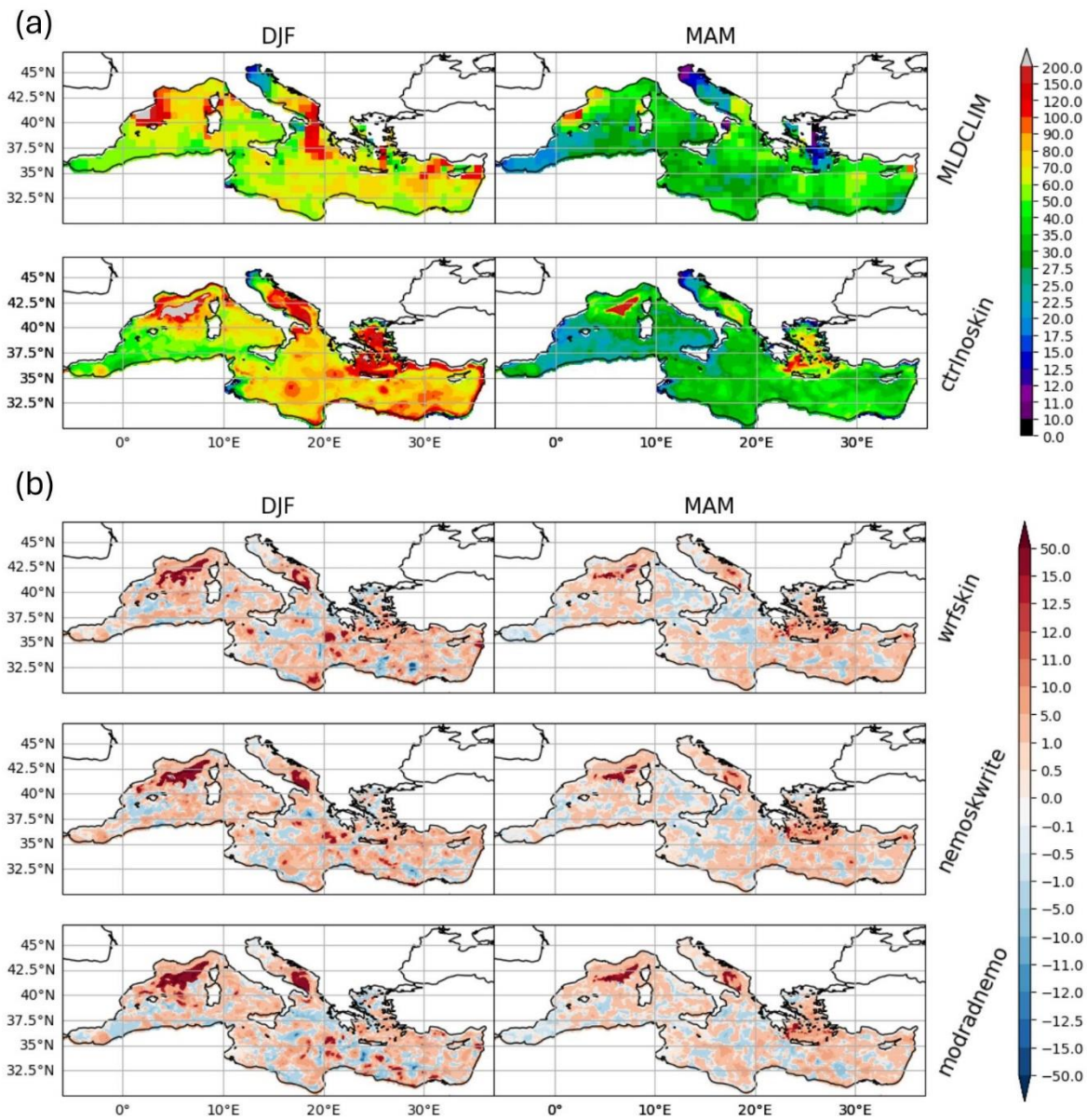


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612 **Figure 10:** Spatial average of profiles within the eastern Mediterranean Sea, during winter and summer. Panel  
 613 10a shows the eastern region, while 10b, 10c show respectively wintertime and summertime spatially-averaged  
 614 profiles within the top 100 m in the upper part, on the bottom the whole depth range on a logarithmic scale.  
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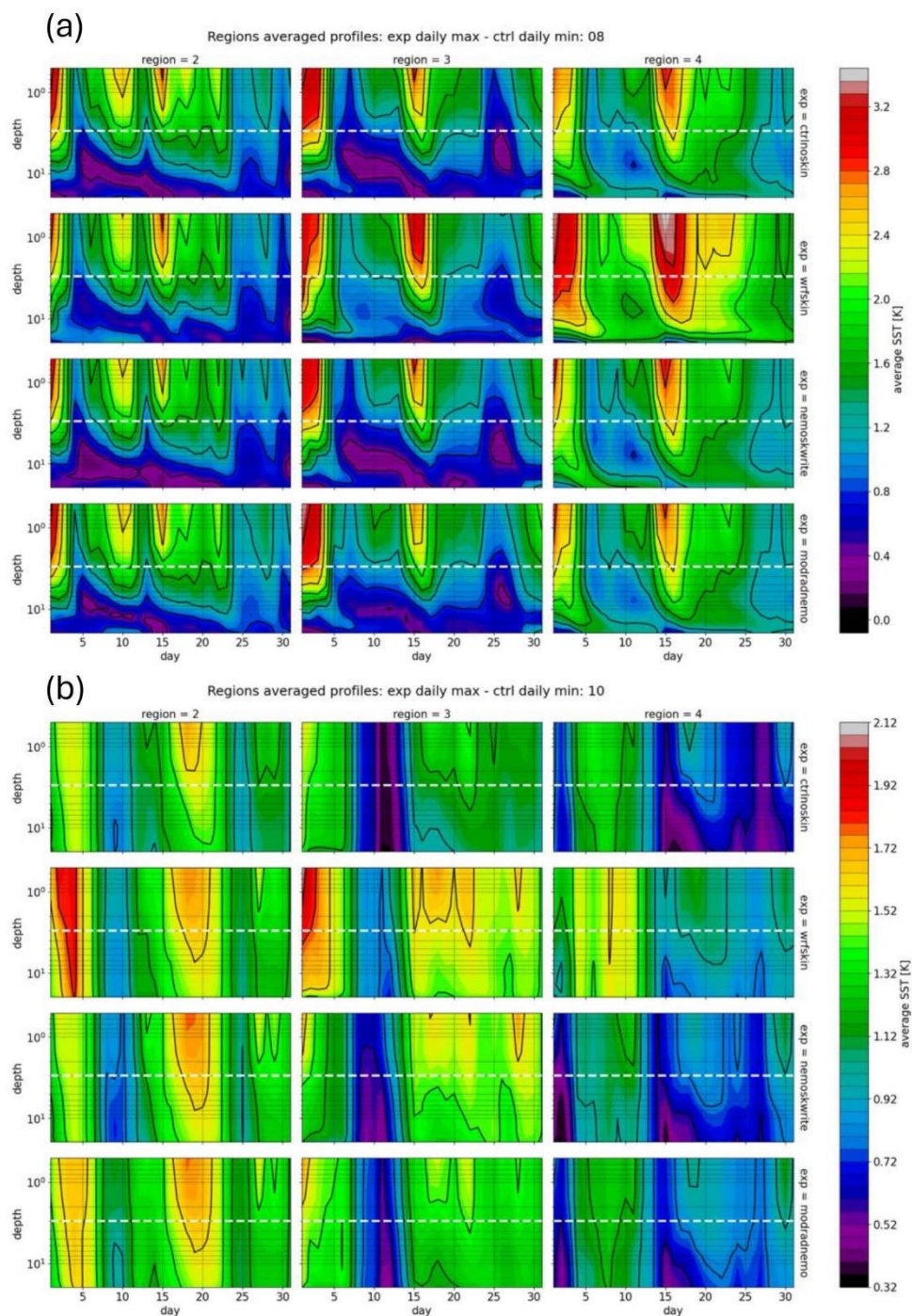


616  
 617 **Figure 11:** RMSE on the top 15m of the difference between regionally averaged profiles between each  
 618 simulation and EN4, displayed as a function of the region and the particular month. Division in regions is  
 619 reported in panel 11a, while 11b, 11c, 11d, 10e show respectively the results for all the simulations carried out  
 620 in the present study.  
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**Figure 12:** Maps of DJF, MAM of mixed layer depth for the climatology and for the control simulation in panel (a). Panel (b) shows the difference of the control with respect to each simulation. Units are meters.

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644 **Figure 13:** Hovmoller plots for spatial average of model outputs temperature profiles in the regions 2,3,4 as  
 645 defined by figure 11a. Each row shows the difference between daily maxima for the given experiment minus  
 646 the daily minima for the control simulation. The white dashed line traces the  $z = 3m$  line of the depth used as  
 647 reference for the base of the warm layer as in ZB05 scheme Zeng and Beljaars, 2005. Panel 13a shows August,  
 648 panel 13b shows October.

649

650

651 Tables

652

Simulation	Scheme active	Extinction coefficients in Warm Layer
ctrlnoskin	None	None
wrfskin	ZB05	SS82
nemoskwrite	ZB05+A02+T10	G09
modradnemo	ZB05+A02+T10	R-G-B + chl e-folding

653

**Table 1:** Overview of the simulations performed

654

Wavelength [ $\mu m$ ]	i	$a_i$	$b_i$ [ $m^{-1}$ ]
0.3-0.6	1	0.2370	$1.488 \times 10^{-1}$
0.6-0.9	2	0.3600	$4.405 \times 10^{-1}$
0.9-1.2	3	0.1790	$3.175 \times 10^1$
1.2-1.5	4	0.0870	$1.825 \times 10^2$
1.5-1.8	5	0.0800	$1.201 \times 10^3$
1.8-2.1	6	0.0246	$7.937 \times 10^3$
2.1-2.4	7	0.0250	$3.195 \times 10^3$
2.4-2.7	8	0.0070	$1.279 \times 10^4$
2.7-3.0	9	0.0004	$6.944 \times 10^4$

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**Table 2:** Parameters for the Transmission coefficient following Soloviev and Schlu"ssel, 1996, in which the first coefficients is the average between the one corresponding to I, IA, IB, II, and III Jerlov optical water types. This is currently implemented in NEMO.



simulation	DJF	MAM	JJA	SON	Annual
ctrlnoskin	-173.31	133.92	75.56	-66.40	-7.55
wrfskin	-168.83	134.19	76.51	-65.87	-5.97
nemoskwrite	-169.28	133.79	76.77	-65.72	-6.10
modradnemo	-169.06	134.87	78.16	-68.13	-6.04
ERA5	-140.36	133.24	81.96	-53.46	5.35

658 **Table 3:** Averaged surface net heat flux over the Mediterranean Sea ( $W/m^2$ ): seasonal and annual spatial  
659 averaged mean values.

660

661

662

## 663 **References**

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