The Vlasiator 5.2 Ionosphere – Coupling a magnetospheric hybrid-Vlasov simulation with a height-integrated ionosphere model

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Abstract. Simulations of the coupled ionosphere-magnetosphere system are a key tool to understand geospace and its response to space weather. For the most part, they are based on fluid descriptions of plasma (magnetohydrodynamics, MHD) formalism, coupled an electrostatic ionosphere. Kinetic approaches to modeling the global magnetosphere with a coupled ionosphere system are still a rarity.

5   We present an ionospheric boundary model for the global near-Earth plasma simulation system Vlasiator. It complements the magnetospheric hybrid-Vlasov simulations with an inner boundary condition that solves the ionospheric potential based on field-aligned current and plasma quantities from the magnetospheric domain. This new ionospheric module solves the ionospheric potential in a height-integrated approach on an unstructured grid and couples back to the hybrid-kinetic simulation by mapping the resulting electric field to the magnetosphere's inner boundary.

10   The solver is benchmarked against a set of well-established analytic reference cases, and we discuss the benefits of a spherical Fibonacci mesh for use in ionospheric modeling. Preliminary results from coupled global magnetospheric-ionospheric simulations are presented, showing formation of both Region 1 and Region 2 current systems.

1 Introduction

15   The ionosphere is the upper part of Earth's atmosphere (in an altitude range of about 80 to 1000 km), in which a significant fraction of gas exists in an ionized state and hence electrodynamic effects are part of the atmospheric dynamics (Palmroth et al., 2021). At the interface between magnetospheric space plasma phenomena and atmospheric physics, the ionosphere plays a key role in space weather space weather effects (Pulkkinen, 2007), such as geomagnetically induced currents (Marshalko et al., 2023), Joule heating (Ahn et al., 1983; Billett et al., 2018; Palmroth et al., 2004) and auroral phenomena. Modeling ionospheric

20   physics and its interaction with the magnetosphere in a global context (including the solar wind, magnetosphere and ionosphere) is therefore a research focus of computer simulations. Multiple well-established simulation systems based on fluid modeling

of plasma flows (Magnetohydrodynamics, MHD) exist, such as SWMF (Gombosi et al., 2021), GAMERA (Lin et al., 2021) or GUMICS-4 (Janhunen et al., 2012).

The large-scale morphology of ionospheric current systems is characterized by the presence of two regions of field-aligned currents (Iijima and Potemra, 1976a, b): Region 1 forms a more poleward circumpolar oval structure, whereas Region 2 is located on the equatorward side of the auroral oval. MHD-based approaches have difficulties representing the Region 2 currents with high fidelity (Ridley et al., 2002), as they are unable to model ring current drift kinetics (Wolf et al., 2007) and overlapping multi-temperature plasmas and are hence underrepresenting the required pressure gradients (Zhang et al., 2011). Kinetic simulation treatment of magnetospheric plasma promises to improve upon this state of the art (Lin et al., 2014; Yu et al., 2022), but comes with massively increased computational requirements. Thanks to new simulation techniques and ever-growing computational capabilities of supercomputing facilities, going beyond the MHD approximation in global modeling has been an ongoing effort (Nishikawa et al., 2021; Palmroth et al., 2018) and opens novel ionospheric-magnetospheric coupling possibilities beyond the current methods.

Common to all coupled simulation approaches irrespective of their concrete plasma representation is the disconnect between the magnetospheric and ionospheric domains, which needs to be bridged by a coupling mechanism. This mechanism transports quantities from the magnetospheric domain into the ionosphere ("downmapping"), which usually encompasses information about field-aligned currents in the magnetosphere, precipitating particle fluxes, Poynting flux or other inputs to the ionosphere solver (Zhang et al., 2015), depending on the physics represented. Transport between the two domains can be modeled through first-principles or empiric formulas (Knight, 1973), but an adiabatic approach, in which further transport effects are neglected, is common (Paul et al., 2023). In the opposite direction, the coupling mechanism acts back onto the magnetospheric domain and affects plasma flows on the ionospheric boundary. Electromotive effects of ionospheric current closure get included by mapping ionospheric potential gradients to electric fields on the magnetosphere simulations' inner boundary ("upmapping"). The cross-polar cap potential (CPCP) is a commonly used diagnostic quantity, that subsums the strength of this feedback effect in a single scalar parameter for each hemisphere (Gordeev et al., 2011). Plasmasphere corotation (Vickers, 1976; Maus, 2017) can be introduced by the same upmapping mechanism through inclusion of a motional electric field contribution. Outflow of ionised atmospheric constituents (Strangeway et al., 2005) likewise find their way into the magnetospheric simulation domain in the upmapping process.

Vlasiator (Palmroth et al., 2018; von Alfthan et al., 2014) is a hybrid-Vlasov simulation system employed to model plasma processes in near-Earth space. Recent simulations (Palmroth et al., 2023; Juusola et al., 2018; Palmroth et al., 2017; Grandin et al., 2019) encompassing Earth's solar wind - magnetosphere system have shown that the employed hybrid kinetic approach has become computationally viable for high-fidelity studies of the magnetosphere. However, past Vlasiator simulations were constrained in their treatment of the ionospheric boundary, as the simulations' earthward boundary model was built to treat Earth as a magnetic obstacle sphere that absorbs all infalling plasma (Palmroth et al., 2018). This proved to be a sufficient model for the initial development goals (Palmroth, 2022) to construct a viable global ion-kinetic model, first in two spatial and three velocity dimensions (2D3V, Palmroth et al., 2018) and then in six-dimensional phase space (3D3V, Ganse et al., 2023) and for the initial science goals of foreshock (Turc et al., 2018), magnetosheath (Grandin et al., 2024) and magnetospheric (Palmroth

et al., 2023) dynamics. For all of these results, however, the studied effects could be well investigated without inclusion of ionospheric interaction. To fully represent global dynamics, a proper two-way magnetosphere-ionosphere coupling is required.

In this paper, we outline the newly implemented ionospheric boundary model in Vlasiator. While it shares no code with
60 earlier codebases, it builds on concepts and experiences from the GUMICS-4 (Janhunen et al., 2012) simulation system. The goal throughout the process was to provide a practical, physically motivated system for current closure of the inner magnetosphere, so that global effects of magnetospheric transients and their interplay with the ionosphere can be studied in a kinetic manner. The solver architecture was chosen, where possible, to avoid semi-empirical models and work from first principles, but deliberately simplifies many aspects of atmospheric physics in its current form, for which much more
65 comprehensive modeling approaches would need to be considered (such as Qian et al., 2014; Marchaudon and Blelly, 2015; Codrescu et al., 2012). The model can best be seen as a starting point for coupled hybrid-Vlasov and ionosphere modeling, in which proven mechanisms were combined to further investigate the possibilites that kinetic simulations offer in geospace modeling.

Section 2 describes the numerical setup, including mesh construction, coupling processes and ionosphere solvers. Specifi-
70 cally, in Subsection 2.3, three options for the height-integrated conductivity model that have been implemented in this framework are outlined. It concludes and explains our choice of the model in which longitudinal conductance is neglected. Section 3 presents verification test results of the ionospheric model and solver against known analytically solvable test cases. Finally, Section 3.3 presents an example Vlasiator magnetospheric run with the new ionosphere model enabled, highlighting the global-scale effects that a two-way coupled ionosphere model enables in the dynamics of a hybrid-Vlasov simulation of Earth's entire
75 magnetosphere.

2   The model

In a global magnetospheric simulation, an ionosphere model provides the inner boundary condition for the simulation domain, through which the ionospheric current systems are affecting the global magnetospheric plasma system. In general, inflowing current and plasma properties are provided as an input to the ionosphere and the solver supplies a predefined set of quantities
80 back to the encompassing magnetospheric model, based on which the outflowing plasma properties are affected by the ionospheric electric potential Φ. This outflow is fed back into the magnetospheric simulation model. Additional direct outputs of the model are ionospheric observable quantities. Some can be compared to ground-based observations, such as charge carrier concentrations and large scale magnetic field fluctuations. Others are of interest because they are not easily obtainable from measurements, such as conductivity maps (Laundal et al., 2022).
85 Coupling an ionosphere model to a hybrid-Vlasov simulation is unexplored territory, for which no thorough previous experience exists, due to the small amount of Vlasov simulation codes employed in space plasma simulations to date (Palmroth et al., 2018). The closest relatives to this approach, hybrid-PIC simulations (Lin et al., 2014), similarly to advanced coupled MHD simulations such as MAGE (Lin et al., 2021), handle their ionosphere inputs and outputs as macrophysical quantities, such as moments of the distribution function and electromagnetic field bulk quantities. This stems naturally from the formula-

tion of these models, since MHD equations themselves are constituted from macroscopic state variables. In a (hybrid-)Vlasov simulation, on the other hand, the ion kinetic distribution function $f_i(\boldsymbol{x}, \boldsymbol{v}, t)$ is the primary actor in the simulation domain. The ionosphere potential needs to directly influence its behaviour at the inner magnetospheric boundary.

Irrespective of the precise nature of the simulation's plasma model, electric currents relate to the magnetic field through Ampère's law,

$$\nabla \times \boldsymbol{B} = \mu_0 \boldsymbol{j}, \tag{1}$$

hence the field-aligned current (FAC) density $j_\parallel$ can be deduced from a simulation's magnetic field $\boldsymbol{B}$ state (μ_0 being the vacuum magnetic permeability). Note the formal difference between the three-dimensional current density $\boldsymbol{j}$ and the height integrated current density $\boldsymbol{J}$, used below. The implementation of the Vlasiator ionosphere is therefore able to take the same approach as the model employed in the GUMICS-4 MHD simulation system (Janhunen et al., 2012): it couples the field-aligned currents determined from the magnetospheric simulation's $\boldsymbol{B}$ onto a spherical shell ionosphere grid. The ionosphere is modeled as electrostatic and the third dimension (altitude) of the system is removed by treating all conductivity and current quantities as integrals from ground level up to an altitude of $200\,\mathrm{km}$. Formally, it thus constitutes a height-integrated model. In practice, this means that ionosphere electrodynamics are solved on a two-dimensional spherical mesh surface. The ionospheric potential corresponding to the tangential surface currents is solved by using a conductance tensor model that integrates the atmospheric density and ionisation columns. Any transient or travelling wave effects in the atmosphere are neglected in this process.

At the end of the ionosphere solver process, the electric potential Φ is obtained by solving Ohm's law in the ionospheric solver (see section 2.4). Φ is mapped upwards along the (equipotential) magnetic field lines to affect the source terms in the magnetospheric inner boundary, thus closing the loop with the magnetospheric simulation. In the case of Vlasiator, the Vlasov simulation's boundary cells dynamically affect their velocity distributions through the upmapped ionospheric potential (see section 2.5).

2.1 The ionosphere grid

The ionosphere is modelled on a spherical shell grid with a radius of $R_i = R_E + 100\,\mathrm{km}$ (where $R_E \approx 6371\,\mathrm{km}$ is Earth's radius), which acts as the effective ionospheric altitude in this model. The grid topology is a triangle mesh, structured as a spherical Fibonacci lattice (Keinert et al., 2015). The mesh resolution can be arbitrarily chosen from a minimum of 16 mesh points on the sphere up to 2^{24} mesh points. The paramount property of Fibonacci meshes is that they discretise the spherical surface into exactly equal-area triangle elements e_i, for any number N of mesh nodes. The resulting effective mesh resolution for an ionospheric Fibonacci mesh with N nodes is thus obtainable by equally dividing Earth's surface area and taking the square root: $l_{\mathrm{eff}} = \sqrt{4\pi R_E^2 / N}$.

As the ionosphere exhibits small-scale spatial structures especially at the auroral latitudes, one or multiple levels of mesh refinement can be performed by consecutively replacing one triangular mesh element by four smaller ones with double resolution, the process of which is outlined in Figure 1. At refinement boundaries, cell stitching (Schäfer et al., 2014) is employed

to maintain a watertight mesh geometry. Figure 2 exemplifies the refinement process with a low base mesh node count of $N = 256$, corresponding to an effective base grid resolution of $l_{\text{eff}} = 1368\,\text{km}$. This is subsequently refined in two stages.

125   The results in this paper were obtained from simulations based on a Fibonacci mesh (base $N = 2000$, $l_{\text{eff}} = 489\,\text{km}$) with three refinement stages around the auroral oval: All elements in the latitude range $\theta = 40° \ldots 90°$ were refined once ($l_{\text{eff}} = 244\,\text{km}$), elements with $\theta = 50° \ldots 90°$ were refined a second time ($l_{\text{eff}} = 122\,\text{km}$) and elements in $\theta = 60° \ldots 80°$ got a third refinement level ($l_{\text{eff}} = 62\,\text{km}$). Figure 3 presents the resulting mesh geometry around the northern polar cap.
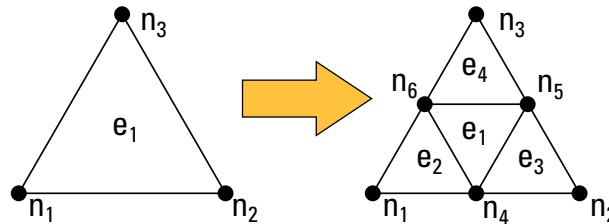
Figure 1. Mesh refinement step of the ionospheric triangle mesh. Each triangular element e_i chosen for refinement is replaced by 4 elements of double resolution, and three additional mesh nodes $n_4 \ldots n_6$ are introduced at the edges.
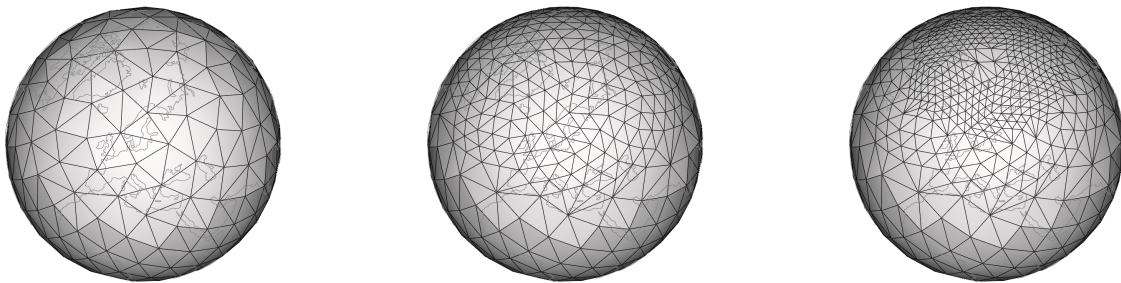
Figure 2. Ionosphere mesh geometry. **Left:** Spherical Fibonacci mesh with $N = 256$ nodes, without refinement of the polar region. **Centre:** One level of mesh refinement above $40°$ latitude. **Right:** A second level of mesh refinement between $60°$ and $80°$ latitude.

2.2   Downmapping of magnetospheric parameters

130   To obtain a bijective association between magnetospheric and ionospheric mesh cells, ionosphere mesh node coordinates are traced upwards along the fieldlines with an adaptive Euler tracing algorithm (Press et al., 1992) (compare Figure 4). Within the magnetospheric simulation domain, the magnetic field values are interpolated using the reconstruction method of Balsara (2017). In the gap region between the magnetospheric and ionospheric simulation domain, the fieldlines are traced along an ideal dipole field without a tilt. Each field line starts from the ionosphere grid radius and continues until the coupling radius r_{C}
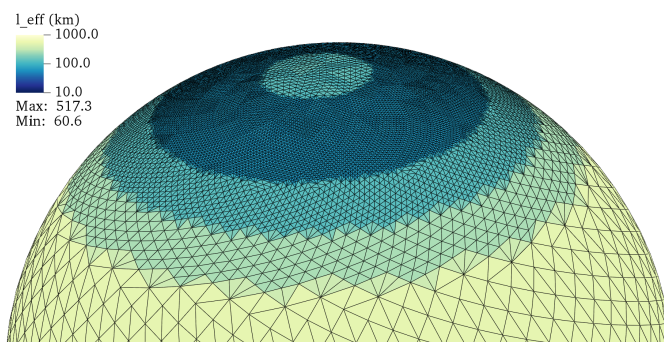
Figure 3. Resulting mesh structure near the northern polar cap after three levels of mesh refinement have been applied between $\theta \in [40°, 90°]$, $[50°, 90°]$ and $[60°, 80°]$, respectively. The colour shows the effective mesh resolution $l_{\mathrm{eff}} = \sqrt{2A_{\mathrm{cell}}}$ of each triangular grid cell.
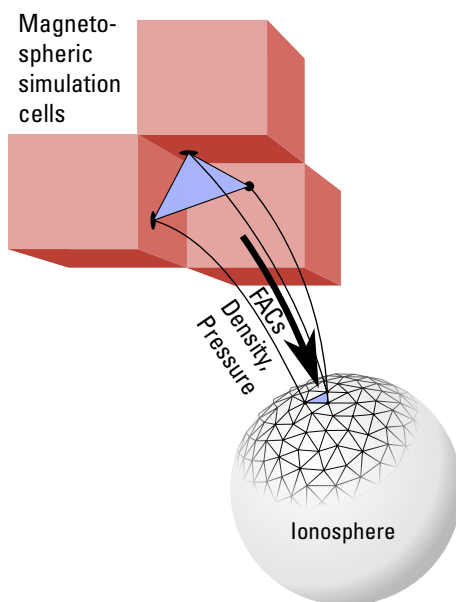
Figure 4. Downmapping of magnetospheric simulation parameters (electron density n_{e}, temperature T_{e} and FACs $J_{\parallel}$) is performed along magnetic field lines towards the ionospheric mesh node locations. The magnetospheric simulation cell data is linearly interpolated onto their cell surface locations.

is reached. r_C is a user-configurable parameter that is typically chosen to lie at least two full magnetospheric simulation cell sizes ($\Delta x = 1000\,\mathrm{km}$) outside of the Vlasov simulations' earthward boundary, to both prevent incorrect application of the finite difference scheme that calculates $\boldsymbol{j}_\parallel$ from eq. (1) and to smooth out any possible artefacts from the effectively non-spherical shape of Cartesian grid discretization ("staircasing" the spherical inner boundary). The traced cell coordinates do not in general hit the centres of magnetospheric cell locations. Hence the magnetospheric downmapping quantities are interpolated onto the upper endpoint of the fieldline using Balsara reconstruction (Balsara, 2017) before they are transported to the ionosphere mesh.

Field-aligned current density $j_\parallel$ produced in the magnetosphere, electron density n_e and temperature T_e are input quantities to the ionosphere model. To couple these magnetospheric quantities with the ionosphere model, they are transported from the magnetospheric grids (Ganse et al., 2023; Papadakis et al., 2022) onto the triangular ionosphere grid.

As Vlasiator simulations only carry full kinetic information about the particle distribution functions of ions and simplify electron dynamics to that of a massless fluid (thus making it a hybrid-Vlasov model), electron precipitation fluxes need to be inferred indirectly in a two-step process: First, a proxy for the precipitating electron distribution function is calculated from the magnetospheric boundary values that are downmapped from the Vlasov simulation cells, field-aligned current density $j_\parallel$, electron density n_e and temperature T_e. We employ a fixed temperature ratio of $\frac{T_i}{T_e} = 4$ in the same way as Janhunen et al. (2012). We assume the electron distribution to be a Maxwellian,

$$f_e(v) = n_e \left(\frac{2\pi k_B T_e}{m_e} \right)^{-3/2} \exp\left(-\frac{m_e v^2}{2 k_B T_e} \right), \tag{2}$$

where k_B is the Boltzmann constant, yielding a differential number flux DNF_e per energy E of

$$\mathrm{DNF}_e(E) = \frac{n_e E}{\sqrt{2 m_e}} (\pi k_B T_e)^{-3/2} \exp\left(-\frac{E}{k_B T_e} \right). \tag{3}$$

For lack of a model describing the precipitating electrons' pitch angle (ϑ) distribution in a hybrid simulation, we assume a cosine dependence in pitch angle (similar to Rees (1963), and further motivated by the results of Ergun et al. (2000)), thus getting to an angle-resolved differential number flux of

$$\mathrm{DNF}_e(E, \vartheta) = \mathrm{DNF}_e(E) \cos(\vartheta) \tag{4}$$

in the downwards-facing hemisphere, with $\mathrm{DNF}_e(E, \vartheta) = 0$ for $\vartheta > \pi/2$. Integration over the downwards-facing half sphere ($\vartheta = [0 \ldots \pi/2]$, $\varphi = [0 \ldots 2\pi]$) yields the omnidirectional differential energy flux:

$$\mathrm{DEF}_{\mathrm{omni}}(E) = \iint E \cdot \mathrm{DNF}_e(E, \vartheta)\, \mathrm{d}\varphi \, \sin\vartheta \, \mathrm{d}\vartheta \quad = \quad \frac{n_e E}{\sqrt{2 m_e \pi}} (k_B T_e)^{-3/2} \exp\left(-\frac{E}{k_B T_e} \right). \tag{5}$$

As the coupling radius r_C is typically situated outwards from the ionosphere by multiple Earth radii ($r_C \sim 5.6\,\mathrm{R_E}$ in recent simulation runs), transport delay of $j_\parallel$ and plasma quantities through this gap region to ionosphere needs to be modelled. Since the dynamic timescales of the inner magnetosphere (typical Vlasiator simulation timesteps are $\Delta t \approx 12\,\mathrm{ms}$) are much faster than ionospheric dynamics, the downmapped quantities are temporally smoothed by an exponential filter with a smoothing half-time of $t_{\mathrm{smooth}} = 4\,\mathrm{s}$. The choice of this value is motivated by the approximate Alfvén travel time from the magnetospheric inner boundary to the ionosphere and back. In the literature, similar effects have been achieved by the choice of solver time

interval between $1\,\mathrm{s}$ (Janhunen et al., 2012) and $15\,\mathrm{s}$ (Paul et al., 2023). In addition to smoothing of high-frequency signals in the downmapped quantities, this filter causes an effective transient propagation delay of $t_{\mathrm{smooth}}/2 \approx 2\,\mathrm{s}$. This process also numerically stabilizes the method, as otherwise instantaneous information transport across the entire magnetospheric inner boundary could occur, leading to unphysical feedback loops.

170 ## 2.3 Ionospheric conductance tensor

The field-aligned current density $\boldsymbol{j}_\parallel$ feeds a charge imbalance inside the ionospheric shell, so an in-plane electric field $\boldsymbol{E} = (E_{\perp 1}, E_{\perp 2}) = -\nabla\Phi$ forms to counter this charge imbalance. The Ohm's law

$$\boldsymbol{J} = \boldsymbol{\Sigma} \cdot \boldsymbol{E}, \tag{6}$$

where Σ is the conductance tensor and $\boldsymbol{J} = (J_{\perp 1}, J_{\perp 2})$ is the height-integrated current density on the sphere, can be used

175 to calculate this field (and with it, the ionospheric potential) if a suitable model for the anisotropic conductance tensor (in a coordinate system where $\hat{\mathbf{e}}_z \parallel \boldsymbol{B}$) is available:

$$\boldsymbol{\Sigma} = \begin{pmatrix} \Sigma_{\mathrm{P}} & -\Sigma_{\mathrm{H}} & 0 \\ \Sigma_{\mathrm{H}} & \Sigma_{\mathrm{P}} & 0 \\ 0 & 0 & \Sigma_{\parallel} \end{pmatrix} \tag{7}$$

A significant part of an ionosphere model's physical content lies in the modelling choices for Pedersen conductance Σ_{P}, Hall conductance Σ_{H} and field-aligned conductance $\Sigma_{\parallel}$. The values of these three conductance components are affected by multiple

180 physical processes in Vlasiator's ionosphere model:

On Earth's dayside, photoionisation from sunlight is dominating. Its magnitude can be obtained from radar observations, and Moen and Brekke (1993) modelled it as a function of solar zenith angle χ and $10.7\,\mathrm{cm}$ solar radio flux $F_{10.7}$ given in solar flux units ($1\,\mathrm{sfu} = 1 \times 10^{-22}\,\mathrm{W m^{-2}}$):

$$\Sigma_{\mathrm{P}}^{UV} = F_{10.7}^{0.49} \left(0.34 \cos\chi + 0.93\sqrt{\cos\chi} \right)\,\mathrm{S\,m^{-1}} \tag{8}$$

185 $$\Sigma_{\mathrm{H}}^{UV} = F_{10.7}^{0.53} \left(0.81 \cos\chi + 0.53\sqrt{\cos\chi} \right)\,\mathrm{S\,m^{-1}}. \tag{9}$$

The $F_{10.7}$ value is a user-definable input parameter to the simulation. For quiescent solar conditions, it is typically chosen to be $F_{10.7} = 100\,\mathrm{sfu}$.

UV photoionisation by starlight is assumed as an isotropic background, and is added as a constant value of $\Sigma_{\mathrm{H,P},\parallel}^{\mathrm{Star}} = 0.5\,\mathrm{S\,m^{-1}}$.

190 In high-latitude regions, conductivity caused by particle precipitation ionisation is a further dominating factor. In nature, this ionisation is predominantly caused by precipitating electrons colliding with neutral atoms and molecules.

Using the electron scattering profile model from Sergienko and Ivanov (1993), the differential energy flux is used to calculate an altitude-dependent ion production rate q in the altitude range of $h \in [60 \ldots 200]\,\mathrm{km}$, corresponding to the ionospheric region of high conductivity caused by precipitating particle ionization (Palmroth et al., 2021). Neutral density profiles going into this

calculation are obtained from the NRLMSIS00 model (Picone et al., 2002), from which a single representative atmosphere profile has been exported to $70°$ degrees northern latitude at midnight during spring equinox (Precise run parameters: Daily $AP = 25$, $F_{10.7} = 100$, $\mathrm{Lat} = 70°$, $\mathrm{Lon} = 0°$, Date is 2022-03-21, 12:00 PM Solar Local Time, height profile from 60 to 200 km with a step size of $1\,\mathrm{km}$). The user can choose to supply a different NRLMSIS00 output file to model specific events or situations.

The resulting production rate q is assumed to be in balance with neutral atom recombination rate $\alpha \approx 2.4 \times 10^{-13}\,\mathrm{m^3s^{-1}}$ (Schunk and Nagy, 2009, Table 8.5), leading to a stationary solution to the balance equation for the ionospheric E region,

$$\frac{\partial n_\mathrm{e}}{\partial \mathrm{t}} = q - \alpha n_\mathrm{e}^2 = 0. \tag{10}$$

We assume time-independence and obtain $n_\mathrm{e} = \sqrt{q/\alpha}$. The conductance values $\Sigma^{\mathrm{Precip}}_{\mathrm{H,P,\|}}$ are then calculated as a result of electron and ion contributions: (Schunk and Nagy, 2009):

$$\sigma_\mathrm{P} = \frac{n_\mathrm{i}\mathrm{e}^2}{m_\mathrm{i}\nu_\mathrm{i}}\frac{\nu_\mathrm{i}^2}{\nu_\mathrm{i}^2 + \Omega_\mathrm{ci}^2} + \frac{n_\mathrm{e}\mathrm{e}^2}{m_\mathrm{e}\nu_\mathrm{e}}\frac{\nu_\mathrm{e}^2}{\nu_\mathrm{e}^2 + \Omega_\mathrm{ce}^2} \tag{11}$$

$$\sigma_\mathrm{H} = -\frac{n_\mathrm{i}\mathrm{e}^2}{m_\mathrm{i}\nu_\mathrm{i}}\frac{\nu_\mathrm{i}\Omega_{ci}}{\nu_\mathrm{i}^2 + \Omega_\mathrm{ci}^2} + \frac{n_\mathrm{e}\mathrm{e}^2}{m_\mathrm{e}\nu_\mathrm{e}}\frac{\nu_\mathrm{e}\Omega_\mathrm{ce}}{\nu_\mathrm{e}^2 + \Omega_\mathrm{ce}^2} \tag{12}$$

$$\sigma_\| = \frac{n_\mathrm{e}\mathrm{e}^2}{m_\mathrm{e}\nu_\mathrm{e}} \tag{13}$$

where the ion and electron collision frequencies $\nu_\mathrm{i,e}$ are taken from Tables 4.5 and 4.6 in Schunk and Nagy (2009), respectively. Ω_{ci} and Ω_{ci} are the larmor frequencies of electrons and ions. Ion densities n_i are assumed to be in charge balance with the electrons. The ion contribution to $\sigma_\|$ has been neglected. Figure 5 shows an example of the resulting precipitation-based conductivity profiles. Height integration yields precipitation-based conductance contributions:

$$\Sigma^{\mathrm{Precip}}_{\mathrm{P,H,\|}} = \int_0^H \sigma_{\mathrm{P,H,\|}}(h)\mathrm{d}h, \tag{14}$$

numerically integrating through the NRLMSIS00 model output in an altitude range of $H = [60\ldots 200\,\mathrm{km}]$ at a step size of $1\,\mathrm{km}$.

Since the photoionization effects of sun- and starlight are mostly affecting the ionospheric F-layer, whereas precipitation effects most dominant in the E-Layer, overall height-integrated Hall and Pedersen conductivities are obtained by summing the individual components in quadrature, meaning

$$\Sigma_\mathrm{H} = \sqrt{\left(\Sigma^{\mathrm{Star}}_\mathrm{H}\right)^2 + \left(\Sigma^{\mathrm{UV}}_\mathrm{H}\right)^2 + \left(\Sigma^{\mathrm{Precip}}_\mathrm{H}\right)^2} \tag{15}$$

$$\Sigma_\mathrm{P} = \sqrt{\left(\Sigma^{\mathrm{Star}}_\mathrm{P}\right)^2 + \left(\Sigma^{\mathrm{UV}}_\mathrm{P}\right)^2 + \left(\Sigma^{\mathrm{Precip}}_\mathrm{P}\right)^2}. \tag{16}$$

This leaves the actual conductance tensor $\mathbf{\Sigma}$ to be assembled. This choice is affected by the orientation of the coordinate system for the solution of the ionospheric potential. Vlasiator's ionosphere model currently implements three different alternative formulations for this process:
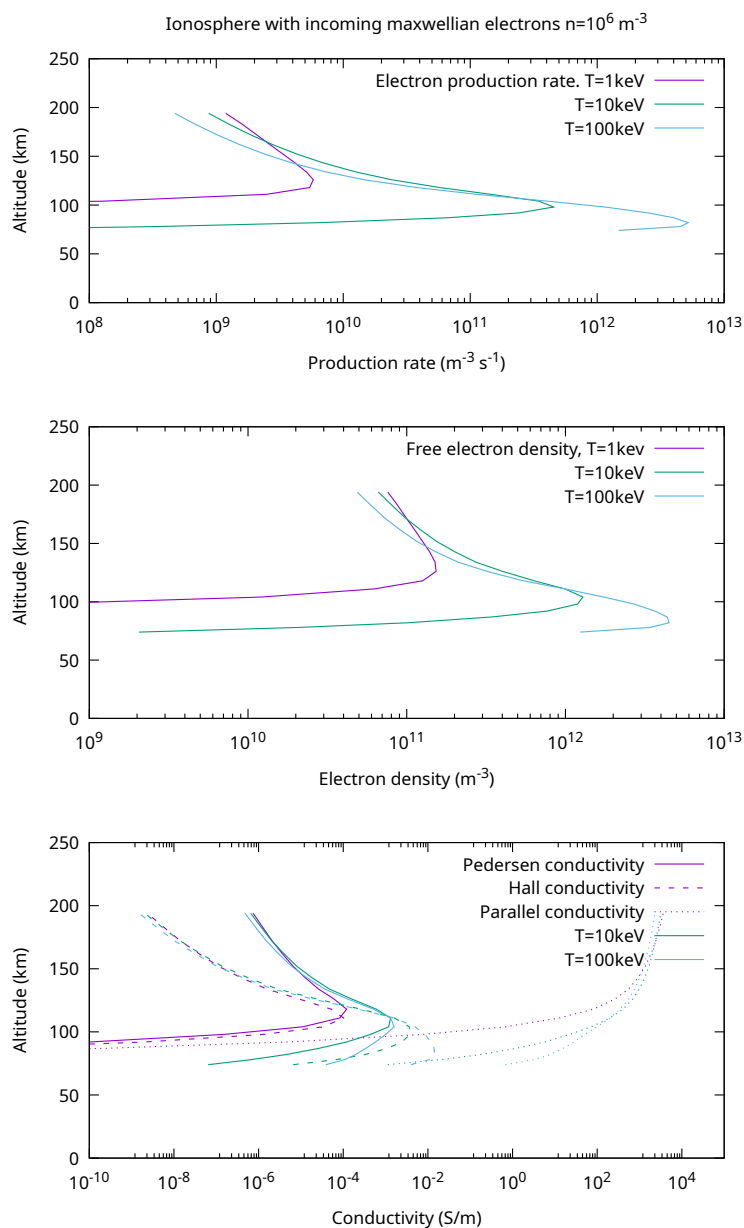
Figure 5. Ionization rate, charge carrier density and conductivity profiles for the three components of the ionospheric conductivity tensor obtained by equations (11) – (13). The shown profile is an example for a grid cell at latitude $70°$ north, with a precipitating electron density of $n_e = 10^6 \text{ m}^{-3}$ and temperature of $T_e = 1, 10,$ and 100 keV, respectively. Note that photoionisation contributions are not shown here, as they get added separately in the height-integrated conductivity terms of equations (15) and (16).

1. In the style of (Janhunen et al., 2012) or Paul et al. (2023), the local magnetic field direction at each grid node is assumed to be exactly radial to the sphere's surface, $\boldsymbol{B} \parallel \boldsymbol{r}$. The in-plane currents of the ionosphere are hence only affected by the two perpendicular conductance components, Σ_H and Σ_P, while $\Sigma_\parallel$ is effectively infinite, making $E_\parallel = 0$.

2. It would seem reasonable to employ the actual dipole magnetic field in the calculation of $\boldsymbol{\Sigma}$. In this case, $\Sigma_\parallel$ obtains a finite value. Analogous to Ridley et al. (2004), its value can be simply chosen to be an arbitrary high conductance value, such as $\Sigma_\parallel = 1000\,\mathrm{Sm}^{-1}$, which is motivated by the assumption that parallel charge separations neutralize near-instantly.

3. The full conductance tensor, using all components as outlined above, and rotated according to the local magnetic field vector.

The fact that the conductance tensor gets employed only in a height-integrated fashion makes inclusion of the parallel conductance problematic: Height integration along a fieldline would effectively mix parallel and perpendicular conductivities in the solution plane. The parallel conductance contribution would vanish at the poles, and become increasingly dominant at more equatorial latitudes. Our experiments have thus shown that both options 2. and 3. result in unphysically large conductivities at low latitudes, favouring current closure over the equator and hence greatly reduced polar potential values. Other models, formulated in polar coordinates, solve this issue by introducing a magnetic field dip factor (Goodman, 1995; Merkin and Lyon, 2010; Paul et al., 2023), but this approach does not readily translate to the spherical Fibonacci grid employed here. In the following, results will only be shown using the conductance model 1.

2.4 Potential solver

The purpose of the ionosphere solver is to find an electric field $\boldsymbol{E}$ that solves equation (6) two-dimensionally on the ionospheric sphere, given a field-aligned current distribution $j_\parallel$. The two additional horizontal current components are forming the tangential current vector $\boldsymbol{J} = (J_{\perp 1}, J_{\perp 2})$. The electric field can be expressed as the gradient of the ionospheric potential, $\boldsymbol{E} = -\nabla\Phi$ and thus the equation to solve becomes:

$$\boldsymbol{J} = \boldsymbol{\Sigma} \cdot (-\nabla\Phi). \tag{17}$$

On the surface of the height-integrated ionosphere model sphere, the field-aligned currents can be substituted in the ionospheric continuity equation $\nabla \cdot \boldsymbol{J} = -j_\parallel$. By taking the divergence of equation (17), it can be rewritten as

$$\nabla \cdot \boldsymbol{J} = \nabla \cdot [\boldsymbol{\Sigma} \cdot (-\nabla\Phi)] = -j_\parallel. \tag{18}$$

In our implementation, the solution to this equation is obtained via a finite element approach using Galerkin's method (Ern and Guermond, 2004), in which the individual triangular mesh elements of the ionosphere grid are used to build trapezoidal test functions for $\nabla\Phi$. The resulting sparse matrix equation is solved with a modified conjugate gradient solver (Press et al., 1992). Since the spherical ionosphere mesh forms a compact manifold with no boundary, the potential Φ on the sphere has a gauge degree of freedom, which causes the finite element solver matrix to be positive semidefinite (with an eigenvalue $\lambda \approx 0$

corresponding to the gauge freedom). This makes naïve implementation of a conjugate gradient solver numerically unstable. To alleviate this instability, an gauge constraint to the potential is introduced such that the potential of mesh nodes near the
255 equator, at a configurable shielding latitude of $\theta_{\text{Shield}} \in [0° \dots 70^{c}irc]$ are pinned to zero potential.

2.5 Upmapping

The potential Φ produced by the ionosphere solver is upmapped to the magnetospheric simulation grid along the same magnetic field trajectories as the downmapping in section 2.2. At this point, the effect of Earth's rotation is taken into account, as the potential was solved without regard for any motional electric field caused by corotation. At the magnetospheric inner boundary
260 cells, an effective electric field is hence calculated as a sum of ionospheric potential and motional electric field from ionospheric corotation,

$$\boldsymbol{E} = -\nabla\Phi - c(L)\left(\boldsymbol{\Omega}_E \times \boldsymbol{r}\right) \times \boldsymbol{B} \tag{19}$$

with Earth's rotation vector $\boldsymbol{\Omega}_E = \hat{\boldsymbol{e}}_z \, 2\pi/24\,\text{h}$ and the corotation factor $c(L) \in [0\dots 1]$ as a function of L-shell parameter:

$$c(L) = \begin{cases} 1 & \text{for } L \leq 5 \\ 0 & \text{for } L > 5. \end{cases} \tag{20}$$

265 The threshold value of $L = 5$ has been chosen to roughly match the plasmapause location, which forms a discontinuous boundary between corotating and convecting plasma (Maus, 2017).

The electric field thus obtained affects the ion distribution function $f_i(\boldsymbol{x},\boldsymbol{v},t)$ in each inner boundary cell of the Vlasov simulation. As a first simple approach, the inner boundary cells are set up to contain a shifted Maxwellian distribution with a bulk velocity given by the $\boldsymbol{E} \times \boldsymbol{B}$-drift $\frac{\boldsymbol{E}\times\boldsymbol{B}}{B^2}$:

$$270 \quad f_i(\boldsymbol{x},\boldsymbol{v},t) = n_i \left(\frac{m_i}{2\pi k_B T_i}\right)^{3/2} \exp\left(-\frac{m_i\left(\boldsymbol{v} - \frac{\boldsymbol{E}\times\boldsymbol{B}}{B^2}\right)^2}{2k_B T_i}\right). \tag{21}$$

We further impose the same $\boldsymbol{E} \times \boldsymbol{B}$ drift velocity on the first layer of cells adjacent to the ionospheric boundary, where the existing (potentially non-thermal) velocity distribution functions are shifted so that the first moment ($\boldsymbol{v}_{\text{bulk}}$) of their distribution coincides with the drift speed. These distribution functions then participate in Vlasiator's global magnetosphere dynamics and influence near-Earth plasma flows, pressure balance and wave dynamics. As Vlasiator is a kinetic simulation model, the
275 choice of f_i in the boundaries is in no way limited to Maxwellians though. Future investigations will study what other outflow distribution functions are sensible here, and investigate their effects on the overall magnetospheric system.

3 Verification and results

Verification and validation of the ionosphere model were performed using a number of test cases: In the following, we check correctness of the numerical implementation of the potential solver by benchmarking the solver against spherical harmonics-
280 shaped $j_\parallel$ patterns and an literature-established analytically-solvable distribution of $j_\parallel$ that is close to a real convective pattern

(section 3.1). Finally, we perform an integration test of the whole coupled magnetosphere-ionosphere model by performing a Vlasiator global run with the included ionosphere model, and validate the resulting ionospheric phenomenology in section 3.3.

3.1 Solver verification and convergence

Verification of the potential solver was performed in a test setup, in which the ionospheric conductance tensor Σ was set to
285 identity $\mathbf{I}$. With this choice, equation (18) reduces to the Laplace equation,

$$\Delta \Phi = j_\parallel. \tag{22}$$

On a spherical surface, the eigensolutions of this equation are the spherical harmonic functions $Y_m^l(\theta, \phi)$ with degree l and order m. Given an input $j_\parallel$ distribution that is composed of a single spherical harmonic, the solver should thus produce a potential Φ with the same shape (up to a proportionality constant). Figure 6 shows a selection of spherical harmonic (up to
290 $l = 6$) $j_\parallel$ distributions. The colour scale of the potential plots was chosen quite narrow to particularly highlight how finite grid resolution leads to small discrepancies in regions of polarity change, while the overall potential morphology matches the expected spherical harmonic shape.
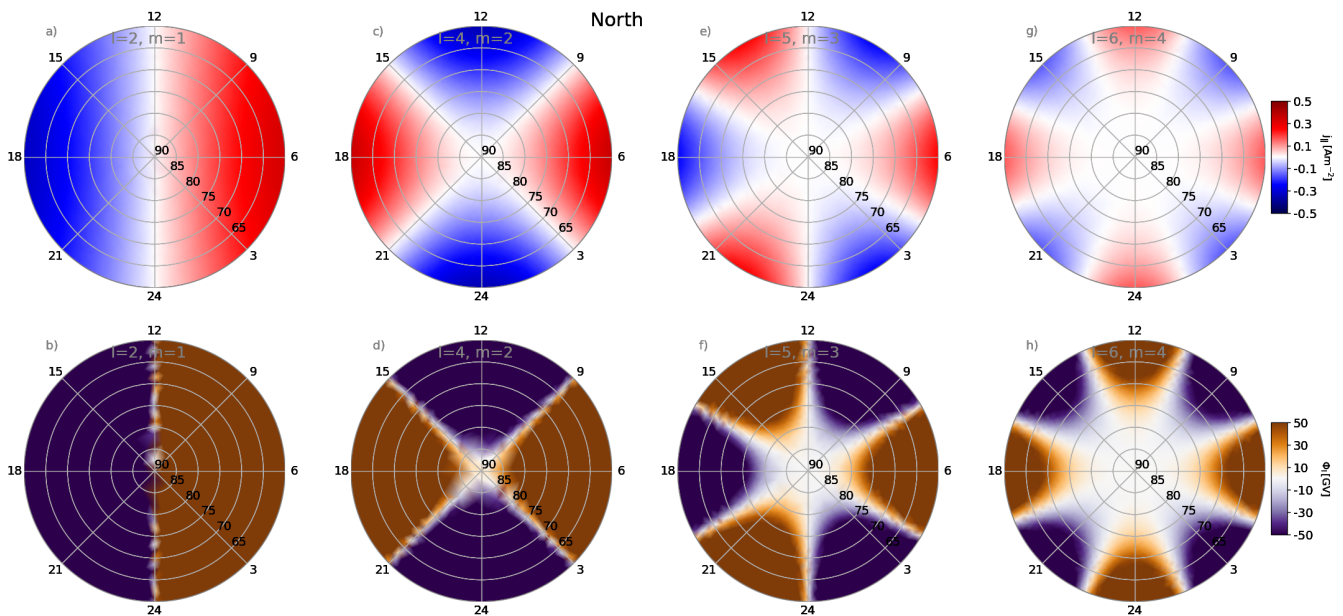
Figure 6. Top row: Spherical harmonic $j_\parallel$ distributions with $(l = 2, m = 1)$, $(l = 4, m = 2)$, $(l = 5, m = 3)$ and $(l = 6, m = 4)$, respectively.
Bottom row: The corresponding potential patterns, when solved with $\Sigma = \mathbf{I}$, showing that the ionosphere solver correctly reproduces the eigensolutions of the system.

To systematically test how well a given Fibonacci mesh with node count N resolves a given spherical harmonic $j_\parallel$ distribution, a parameter study was conducted. Defining a correlation product of two functions X and Y over the full sphere S

295 (evaluated on the Fibonacci point set sites with N points) as

$$\langle X \cdot Y \rangle = \frac{1}{\|X\| \, \|Y\|} \int_S X \cdot Y \, \mathrm{d}S, \tag{23}$$

where $\|X\| = \int_S X^2 \mathrm{d}S$ is the norm of X on the sphere, we expect a value of $\langle \Phi \cdot Y_m^l \rangle = 1$ for a numerically correct solution of the test case. Any deviation from this values indicates numerical error, in this case due to insufficient resolution of the mesh. Figure 7 presents this for a set of spherical harmonic numbers up to $l = 13$ and mesh node counts in the range $N = 20 \ldots 500$. It

300 can be seen that low grid point numbers only resolve low orders l of spherical harmonics, but the solver converges by increasing the node count N. Note that the result for high-order spherical harmonics are entirely dominated by aliasing for low values of N, and hence the curves for $l = 8$ and $l = 13$ are only shown for $N > 189$ and $N > 198$, respectively.
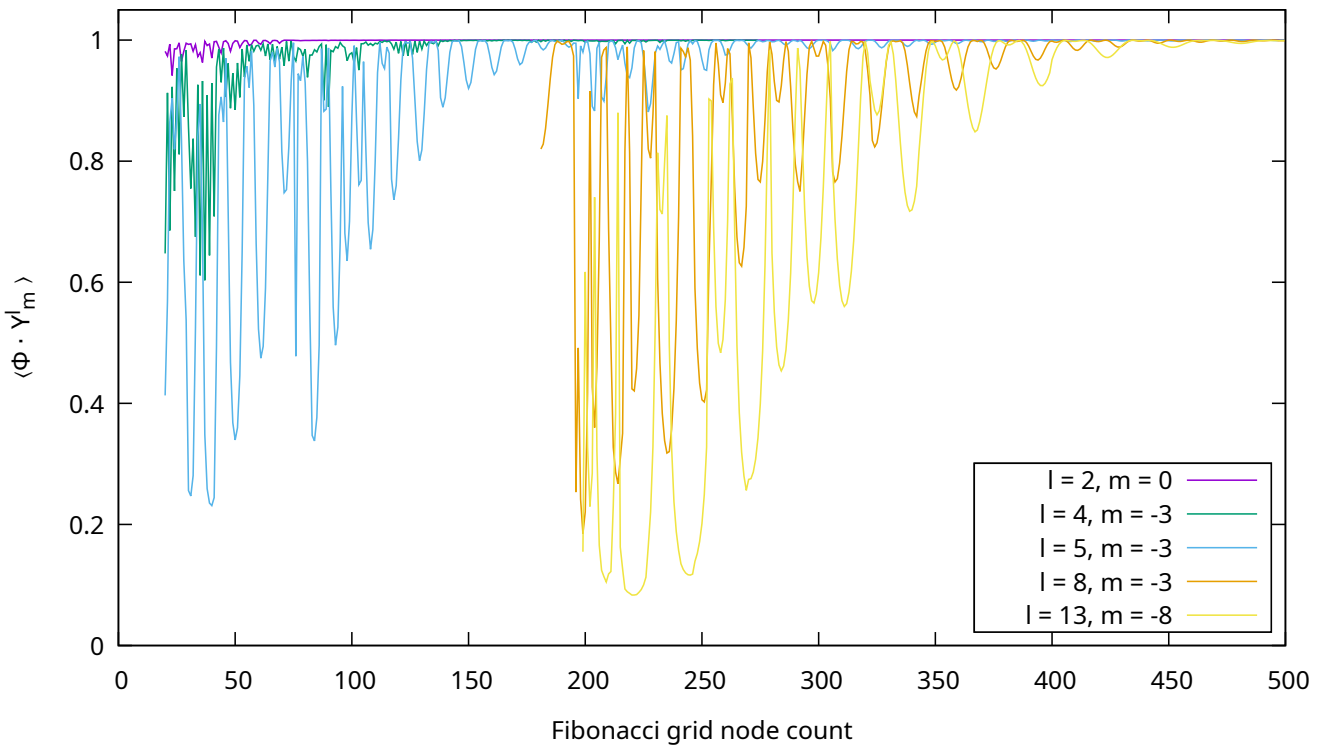
Figure 7. Spherical harmonic solver convergence test for different mesh resolutions N: An FAC distribution $j_\| = Y_m^l(\theta, \phi)$, with given l and $m \in [-l \ldots l]$ is solved to give the potential $\Phi(\phi, \theta)$. The normalized spherical function correlation $\langle \Phi \cdot Y_m^l \rangle = 1$ ((23)) acts as a measure how well Y_m^l is an eigenfunction of the solver at resolution N. Low harmonics l are easily resolved, even with low N. For higher orders l, the Fibonacci mesh point number N needs to be increased.

3.2 Physically-motivated analytic test case

For a second test case that is more closely representative of the physical situation in Earth's ionosphere, we replicate the veri-
fication test proposed in Merkin and Lyon (2010). In this test case, a longitude-dependent field-aligned current ring, specified
by

$$
j_\parallel = J_0 \begin{cases} \sin\theta\sin\phi, & \text{if} \quad \theta_0 \leq \theta < \theta_0 + \Delta\theta \\ 0, & \text{otherwise,} \end{cases}
\tag{24}
$$

is used as the solver input, where $j_0 = 1\,\mu\text{A/m}^2$, $\theta_0 = 56°$, $\Delta\theta = 12°$. Note that Merkin and Lyon (2010) give these angles
as colatitudes measured from the pole instead. Solving with a constant ionospheric conductance of $\Sigma_H = 0$, $\Sigma_P = 10$ S and
an equatorial shieldling latitude $\theta_{\text{Shield}} = 45°$, the resulting potential is presented in Figure 8. As eq. (24) specifies $j_\parallel$ to be
completely symmetric wrt. longitude ϕ, and since $\Sigma_H = 0$ in this test, the resulting map of Φ shows the same symmetry. In
fact, $\Phi(\theta, \phi)$ factorizes into a purely θ-dependent function $\hat\Phi(\theta)$ and $\sin\phi$,

$$
\Phi(\theta, \phi) = \hat\Phi(\theta)\sin\phi.
\tag{25}
$$

The right panel of Figure 8 plots the values of $\hat\Phi$ a scatterplot. Merkin and Lyon (2010) provide an analytic solution for
the potential in this test, which is shown as a solid line. There is excellent agreement between our solution and the analytic
prediction, with the only significant deviation of the two close to the shielding latitude, in the area of coarse mesh resolution
(compare Figure 3).
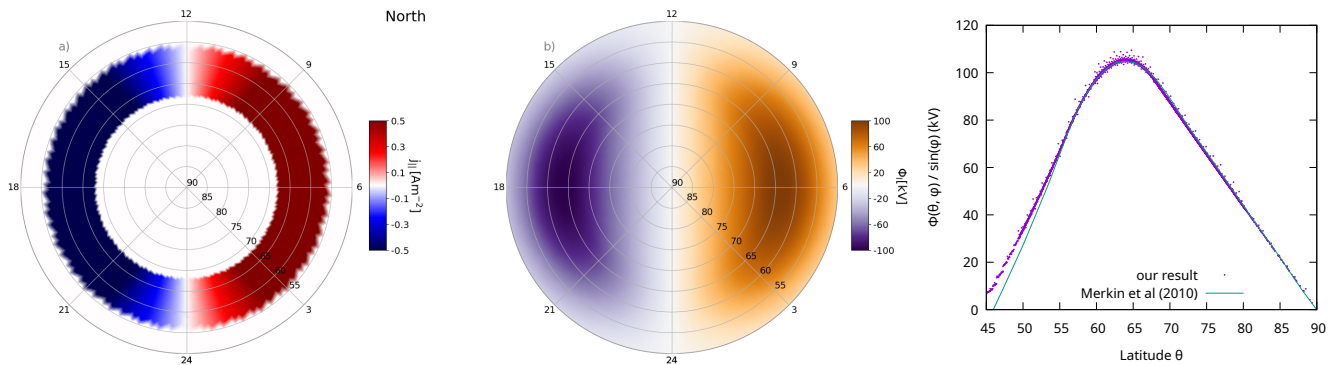
Figure 8. Ionosphere solver test case proposed by Merkin and Lyon (2010). **Left:** Input is an longitude-dependent current ring between
$\theta = 56°$ and $68°$. **Centre:** Solved with a constant ionospheric conductance of $\Sigma_H = 0$ and $\Sigma_P = 10$ S yields the presented ionospheric
potential. **Right:** Comparison of the potential with the analytic solution from Merkin and Lyon (2010).

3.3 Coupled global magnetosphere - ionosphere simulations

To verify and validate the new magnetosphere-ionosphere coupling in Vlasiator, a large-scale global simulation run was per-
formed, in which the ionosphere model presented above was coupled to the magnetospheric simulation. In this manuscript,

we present only the initial states of the simulation up to $t = 500\,\mathrm{s}$, to verify the formation of a physically reasonable coupled magnetosphere-ionosphere system and check the run's stability. Proper scientific analyses of the resulting phenomena with a fully-formed magnetosphere will be subject of upcoming publications.

3.3.1 Simulation setup

325 The simulation domain is spatially 3-dimensional, with simulation box extents of $x = [-110 \cdots + 50]\,R_{\mathrm{E}}$, $y, z = [-57.8 \cdots + 57.8]\,R_{\mathrm{E}}$ in the GSM (Geocentric Solar Magnetic) coordinate frame. Earth's dipole was initialized with its nominal strength of $8\mathrm{e}15\ \mathrm{Tm}^3$ and no tilt. The inflowing solar wind conditions were chosen with a proton density of $10^6\,\mathrm{m}^{-3}$, a purely antisunward velocity of $v_x = -750\,\mathrm{km\,s}^{-1}$ and a Maxwellian velocity distribution with $T = 5 \times 10^5\,\mathrm{K}$. The interplanetary magnetic field was chosen to lie purely southward with a strength of $B_z = -5\mathrm{nT}$. Figure 9 shows an overview plot of the simulation at

330 $t = 500\,\mathrm{s}$. A more thorough discussion of the same setup of the magnetospheric domain is available in Palmroth et al. (2023); the critical difference to the run presented here is the inclusion of the ionosphere model as inner boundary of the magnetospheric simulation, situated at $r_M = 3 \times 10^4\,\mathrm{km}\,(\approx 4.7\,R_{\mathrm{E}})$. The coupling radius is at $r_C = 35.7\,\mathrm{km}\,(\approx 5.6\,R_{\mathrm{E}})$, which is linked to the ionosphere at at $R_{\mathrm{i}} = 6471\,\mathrm{km}\,(= 1\,R_{\mathrm{E}} + 100\,\mathrm{km})$. The spherical Fibonacci mesh is refined with the strategy outlined in section 2.1, resulting in $N_n = 4958$ nodes and $N_{\mathrm{e}} = 9912$ elements.

335 The downmapping process described in section 2.2 results in ionospheric solver input quantities shown in Figure 10, namely $j_\parallel$ (panel a), precipitating electron population density n_{e} (panel b), temperature T_{e} (panel c) and precipitating energy flux $W_{\mathrm{precipitation}}$ (panel d).

3.3.2 Field-aligned current patterns

In the beginning of the simulation, where the simulation box is filled with homogeneous plasma flowing at solar wind velocities,
340 the ionospheric field-aligned current pattern is a smooth hemispheric convection shape. Figure 11 shows the evolution of field-aligned currents in the northern polar region for multiple snapshots of the global simulation at $t = 50\,\mathrm{s}$, $100\,\mathrm{s}$, $250\,\mathrm{s}$ and $500\,\mathrm{s}$. The initial convective pattern gradually transforms into a ring of multiple overlapping current regions. Region 1 and 2 current patterns (Iijima and Potemra, 1976a, b) establish themselves in the later simulation stages (for example at $t = 250\,\mathrm{s}$). The fully formed current system state at $t = 500\,\mathrm{s}$ shows a clear qualitative improvement over earlier Vlasiator investigations of $j_\parallel$
345 without a feedback mechanism to the magnetosphere, such as Horaites et al. (2023). Inclusion of the ionosphere in the Vlasiator simulation apparently strengthens the Region 2 system by ensuring current balance around the magnetosphere-ionosphere boundary.

3.3.3 Ionospheric conductivity

The solar (Eqs. 8 and 9) and starlight UV ionisation contributions are causing dayside conductivity enhancement right from
350 the start of the global simulation run. Precipitating particles form structures of enhanced conductivity in the auroral regions as
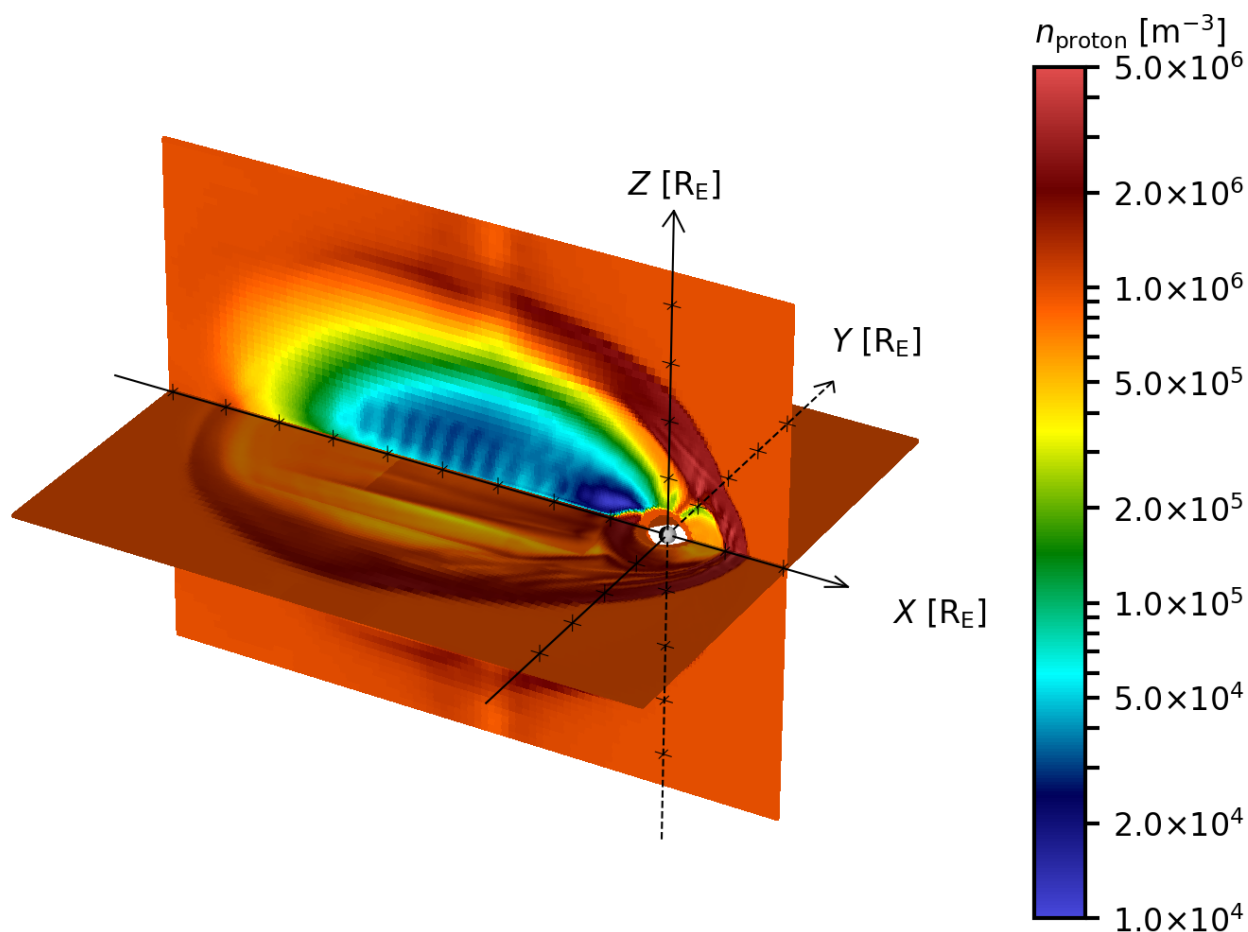
Figure 9. Large-scale overview of the magnetospheric simulation setup. Earth's magnetosphere selfconsistently forms in a simulation box with extents of GSM $x = [-110 \cdots +50] \, R_{\mathrm{E}}$, $y, z = [-57.8 \cdots +57.8] \, R_{\mathrm{E}}$. Solar wind inflow from the boundary in $+x$ direction forms the bow shock and affects ionospheric dynamics. State of the simulation at $t = 500 \, s$, showing the magnetosphere and magnetotail.
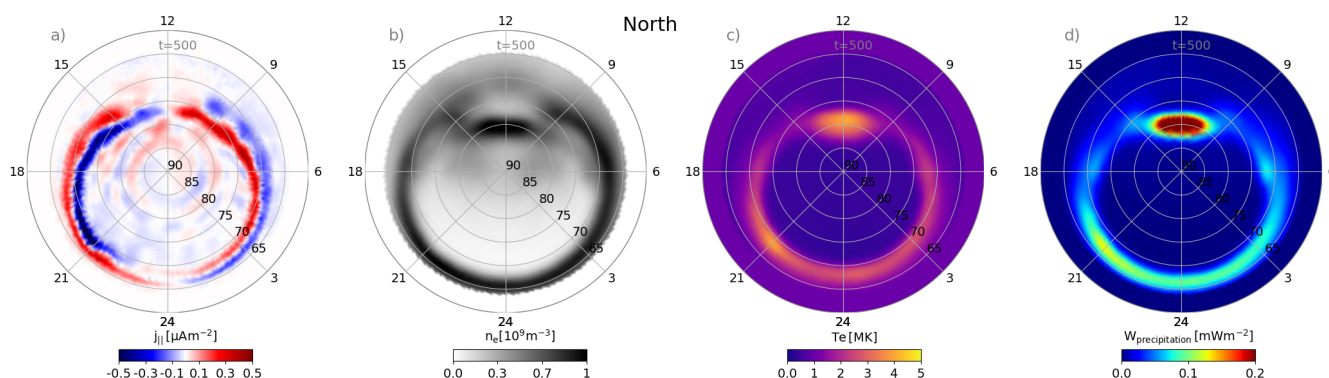
Figure 10. Magnetospheric simulation quantities from the global magnetosphere simulation presented in Figure 9, downmapped onto the ionospheric mesh. **a)**: Field-aligned current density $j_\parallel$ **b)**: effective precipitating electron number density n_e **c)**: temperature of the precipitating electron population T_e. **d)**: electron precipitation flux $W_{precipitation}$. At $t = 500\,s$, these show patterns consistent with a fully formed magnetosphere, including cusp precipitation and circumpolar precipitating electrons at auroral latitudes.
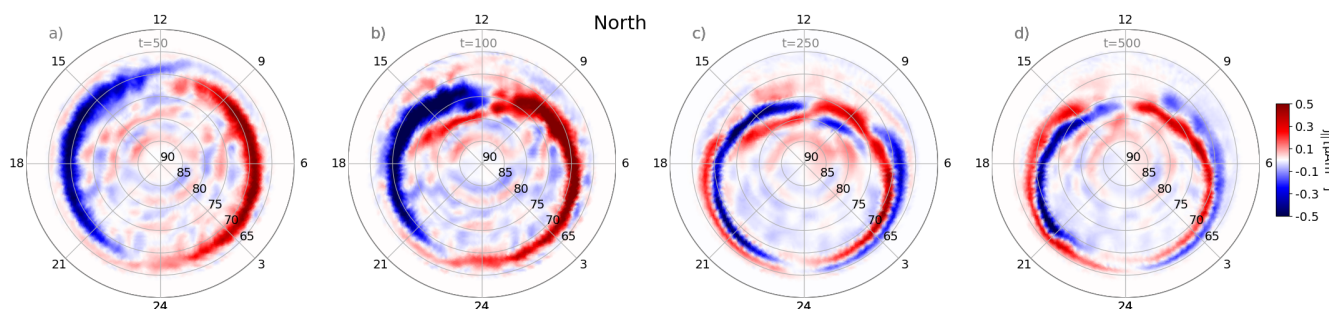
Figure 11. Evolution of field-aligned current density $j_\parallel$ in the northern polar region in a global hybrid-Vlasov simulation coupled to the ionosphere model. The states at $t = 50\,s$ (**a**), $100\,s$ (**b**), $250\,s$ (**c**) and $500\,s$ (**d**) are shown, demonstrating how the simulation initializes from a purely convective flow pattern to a properly-formed ionosphere with Region 1 and 2 current systems.

the magnetosphere forms, and have established a steady state by $t = 500\,\mathrm{s}$. Figure 12 shows conductance maps in Σ_P and Σ_H at $t = 500\,\mathrm{s}$ in the global simulation.

The sunward halves of both Σ_P and Σ_H are dominated by UV-ionization, plus a clearly visible peak feature at the approximate location of the polar cusp. At auroral latitudes between $\sim 63°$ to $70°$, conductivity enhancement due to particle precipitation is apparent. As neither the dayside nor nightside reconnection region had time to develop transient outflow features at this point of the simulation, the auroral oval remains still mostly smooth. Some inhomogeniety of conductance is visible on the night side, with a slight increase both the dusk sector (around MLT =21 h) and dawn (around MLT = 3 h). The auroral region's oval has a sharp equatorward cutoff, as to the magnetospheric inner boundary radius of $R_B = 4.7\,\mathrm{R_E}$ does not allow any downmapping from lower latitudes. This phenomenon is expected to improve, as the magnetosphere boundary radius gets decreased in future Vlasiator simulations.
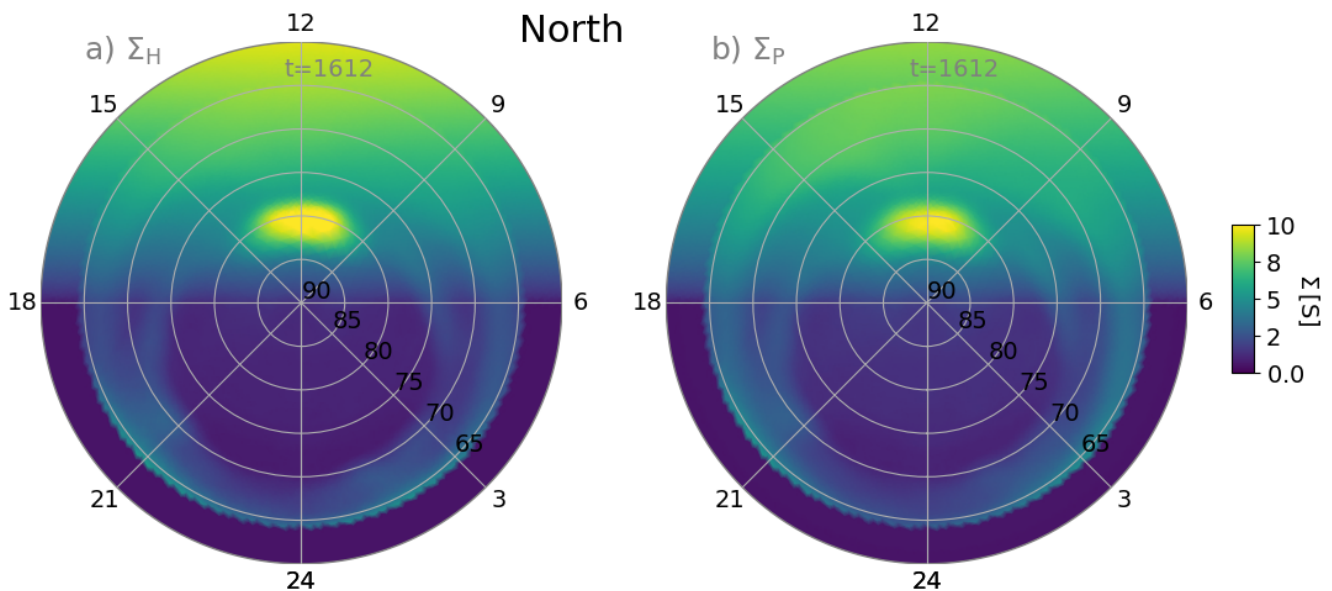
Figure 12. Ionospheric Hall Σ_H and and Pedersen Σ_P conductances in a global magnetosphere-ionosphere simulation ($t = 500\,\mathrm{s}$). Only the northern polar region is shown. The sunward direction (MLT = 12h) is on top, where conductance enhancement from UV radiation dominates. Particle precipitation further leads to enhanced Hall conductance in the auroral oval. High amounts of cusp precipitation lead to strongly enhanced conductance on the dayside at around $\theta \approx 80°$.

3.3.4 Evolution of ionospheric potentials

The ionospheric potential Φ, at the very start of the simulation, is likewise dominated by a pure convection pattern phenomenology, and shows a fully symmetric, hemispheric cross-polar cap potential distribution. The panel a) of Figure 13 presents this state at $t = 50\,\mathrm{s}$ in the simulation. We calculate the CPCP by taking the maximum and minimum value of Φ in each hemisphere,
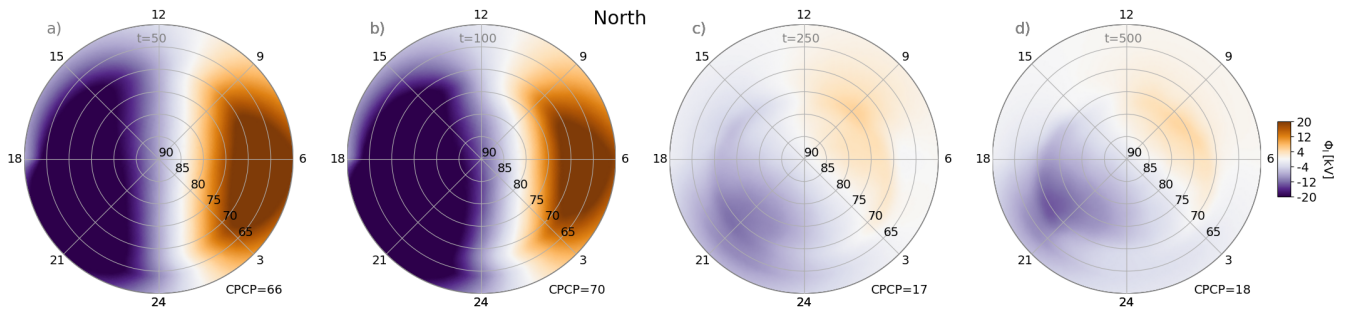
Figure 13. Ionospheric potential at $t = 50\,\mathrm{s}$ (**a**), $t = 100\,\mathrm{s}$ (**b**), $t = 250\,\mathrm{s}$ (**c**) and $t = 500\,\mathrm{s}$ (**d**) in the global magnetospheric simulation. The initial, hemispherical potential distribution caused purely by convective effects slowly develops into a more complex pattern, as precipitation and magnetospheric structure effects form.

and taking their difference. An initial peak of CPCP of $\sim 70\,\mathrm{kV}$ quickly dissipates, and settles into a latent stable state with CPCP$\approx 18\,\mathrm{kV}$ by $t = 500\,\mathrm{s}$ (Figure 13 d).

4 Discussion

We have implemented a new ionospheric conductivity and current systems solver, for the first time coupling a global hybrid-Vlasov simulation to a height-integrated ionosphere model. The solver implementation was verified by benchmarking against a set of analytic test cases (section 3) and checking physical validity in a global, magnetosphere-coupled simulation setup (section 3.3). The multipole tests in section 3.1 served to verify the mesh geometry and solver convergence behaviour under different resolution constraints. We have demonstrated that a spherical Fibonacci mesh forms a suitable and versatile base grid for Ionospheric simulations, allowing fine control over the desired mesh resolutions, especially when combined with a mesh refinement mechanism. This test can also be utilized in the converse manner, to choose a suitable mesh resolution N in order to resolve physical phenomena at the scale lengths of a specific spherical harmonics function $Y_m^l(\theta, \phi)$, as presented in Figure 7.

The current ring test (section 3.2) then specifically addressed the solver fidelity when calculating the potential Φ resulting from a semi-realistic distribution of $j_\parallel$. The result in Figure 8 shows correctness of our solver implementation, but also highlights the importance of grid resolution, as a mismatch between our solution and the analytic curve is visible in the low latitude regions, where a lower mesh resolution was chosen. The low latitude shielding boundary ($\theta_\mathrm{Shield} = 45°$ in this test) is implemented as a Dirichlet boundary condition in the ionospheric potential finite element solver (Ern and Guermond, 2004), which makes the solution sensitive to even small variations of the boundary location due to mesh element placement resolution. In actual physics runs, the shielding latitude is chosen far enough away from the auroral regions, so discretisation errors are confined to regions with little or no contribution to the actual ionospheric current dynamics. Note that both the multipole and ionospheric current ring tests are run as specified without a hall conductivity $\Sigma_\mathrm{H} = 0$, in order to be analytically tractable. A

suitable, well-established analytic test case for verification of ionospheric solvers that includes a nonzero Σ_H is still missing, and would complement physics-based validation studies such as Chartier et al. (2023).

The resulting first output from global simulation run data shows satisfactory fidelity in representation of the ionospheric current structures. The global run results show that the CPCP magnitude is in line with those of the GUMICS-4 MHD sim-
390   ulation (Gordeev et al., 2013). In our results, it seems that Region 1 currents appear as soon as the magnetosphere has fully formed. After initialization transients of the simulation have passed and pressure gradients in the inner magnetosphere establish themselves, Region 2 currents also become apparent. Hence we conclude that after about $t \approx 500\,\mathrm{s}$, the ionospheric model's reproduction of current systems becomes sufficiently realistic to study their effects on global magnetospheric phenomena. The observed simulation behaviour of the downmapped $j_\parallel$ values has shown to be very sensitive to the choice of coupling radius
395   R_C from the magnetospheric domain. If chosen too close to (within ~ 2 simulation cells of) the inner simulation boundary, the resulting $j_\parallel$ patterns are strongly affected by simulation edge artefacts and the strength of current patterns in the ionospheric domain is decreased. Similar safety distance between inner boundary and coupling radius has likewise been reported by Ridley et al. (2004).

To facilitate verification against MHD simulations (e.g., Palmroth et al., 2006), as the first step the Vlasiator ionospheric pre-
400   cipitation was decided to mimic the choices for electron precipitation in MHD simulations (Janhunen et al., 2012). The choice of precipitation only makes use of magnetospheric simulation data on the macroscopic level, that is, through the moments of the ion distribution function, even though kinetic information is available. While the magnitude of the conductivities are aligned with those in MHD simulations (Palmroth et al., 2006), reliance on ion population data to infer electron precipitation leads to misplacement of ionospheric conductivity structures in longitude. This is because the drift motions of ions and electrons,
405   which should be oppositely directed, cannot currently be taken into account separately. In some locations, such as the polar cusps, electron precipitation fluxes may be overestimated, as they are directly tied to proton fluxes in the current implementation. Work is ongoing to develop more accurate models that take the plethora of kinetic simulation data from Vlasiator more effectively into account and to provide for a more sophisticated precipitating distribution function model (e.g., Zhang et al., 2015). Further research will be required to overcome this limitation. One option would be the implementation of a conductance
410   model that bypasses modeling of electron precipitation and semi-empirically constructs the values of Σ_H and Σ_P as functions of MLT and $j_\parallel$ (such as Robinson et al., 2020; Wang and Zou, 2022).

The magnetic field dipole model that Vlasiator employs is a simple, untilted dipole with the magnitude matching Earth's magnetic field. As such, it neglects many intricacies that proper empirical magnetic field models such as Tsyganenko and Andreeva (2015) would provide. This is partially by design, as the Vlasiator philosophy is to start investigation of phenomena
415   on a clean background, and to increase complexity in a second step. Analysis of a more complete magnetic field model in a hybrid-Vlasov simulation is an interesting avenue of research in itself and will be part of future investigations. The example global simulation run presented here, with its steady solar wind speed and fluctuation-free southward IMF, is likewise an idealisation that served to verify the nominal behaviour of all simulation components. Performing runs with a time-varying inflow condition, as in Zhou et al. (2022), will allow the study of resulting magnetospheric and ionospheric transients in kinetic
420   physics. In the meantime, studies of proton precipitation from Vlasiator (Grandin et al., 2019, 2020), for which full kinetic data

is available in the simulation, have been carried out, and showed good agreement with satellite observations (Grandin et al., 2023). Including ionospheric conductivity contributions from proton precipitation, through a model such as Fang et al. (2013), promises interesting future avenues of research for studying kinetic interactions of ion-scale phenomena and their ionospheric correlates, such as dayside reconnection and FTE processes.

5 Conclusions

We have implemented a new ionosphere solver for Vlasiator, a hybrid-Vlasov plasma simulation code targeting global magnetospheric dynamics. The coupling of a hybrid-Vlasov magnetospheric simulation with an ionospheric current model employs similar methods established through global MHD modeling, but requires careful consideration in the coupling process.

Section 2 presented our chosen spherical Fibonacci mesh structure, motivated and described our downmapping and precipitation models and outlined the solver mechanism. The numerical implementation was verified through a set of test cases in Section 3. Mesh and solver behave as expected, and pass the standard test cases well. Preliminary results from a global magnetosphere-ionosphere simulation were shown in Section 3.3, in which the overall ionospheric response to the magnetosphere simulation was confirmed to be consistent with results from fluid-based modelling efforts. Some preliminary evidence of new kinetic-physics features is present, thorough analysis of which will be topic of future publications.

The model presented here provides a solid foundation for further studies of kinetic plasma simulations coupling to ionospheric modeling.

Code availability. The Vlasiator simulation code is distributed under the GPL-2 open source license (Palmroth and the Vlasiator Team, 2024). The ionosphere model discussed here has been included since the Vlasiator release version 5.2 (Palmroth and the Vlasiator Team, 2022).

Data visualisation was performed using the Analysator toolkit (Battarbee et al., 2021), which is likewise available under an open source licence.

Data availability. Full simulation data for the presented analysis is stored in the University of Helsinki Datacloud. Data presented in this paper can be accessed by following the data policy on the Vlasiator website https://helsinki.fi/vlasiator.

Author contributions. UG conceptualized the study. UG, YK and AL implemented and verified the ionosphere model. UG wrote the original manuscript UG and AW implemented and performed the visualisations. UG, YK, KP, MG, MA, HZ, MB and LK contributed to Vlasiator code development. MG, KH, LJ, AW and FK contributed to the derivations of solver mechanisms and analysis of ionospheric results in reference to physical observations. UG, YK and JS performed the simulation runs presented here. MP is the Vlasiator PI, and she participated in conceptualisation of the study, and supervised the work. All authors reviewed and commented the paper.

Competing interests. No competing interests are present.

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