



- 1 Comprehensive evaluation of typical planetary boundary
- 2 layer (PBL) parameterization schemes in China. Part I:
- 3 Understanding expressiveness of schemes for different

# 4 regions from the mechanism perspective

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34	$\label{eq:Abstract.} Abstract. \ The optimal choice of the planetary boundary layer (PBL) parameterization scheme is of$
35	particular interest and urgency to a wide range of scholars, especially for many works involving
36	models. At present, there have been many works to evaluate the PBL schemes. However, little
37	research has been conducted into a more comprehensive and systematic assessment of the
38	performance capability of schemes in key regions of China, especially when it comes to the
39	differences in the mechanisms of the schemes themselves, primarily because there's scarcely
<b>40</b>	sufficient observational data, computer resources, and storage support to complete the work. In
41	addition, there are many factors that influence the selection of schemes, such as underlying surface,
42	initial and boundary conditions, near-surface layer scheme, horizontal/vertical resolution, etc. In
43	this Part (i.e., Part I), four typical schemes (i.e., YSU, ACM2, BL and MYJ) are selected to
44	systematically analyze and evaluate near-surface meteorological parameters, PBL vertical structure,
45	PBL height (PBLH), and turbulent diffusion in five key regions (i.e., North China Plain, NCP;
<b>46</b>	Yangtze River Delta, YRD; Sichuan Basin, SB; Pearl River Delta, PRD and Northwest Semi-arid,
47	NS) of China in different seasons (i.e., January, April, July and October). The differences in the
<b>48</b>	simulated 2-m temperatures between the nonlocal closure schemes are mainly affected by the
49	down ward shortwave radiation, but to compare the nonlocal closure schemes with the local closure $% \mathcal{A}^{(n)}$
50	schemes, the effect of sensible heat flux needs to be further considered. In terms of temporal
51	variation, the simulated results for July are better than the other three months, and the simulated
52	results for nighttime are better than daytime. In terms of regional distribution, the temperature at
53	stations with higher elevation is easily underestimated, while overestimated with lower elevation.
54	The variation of relative humidity corresponds to temperature. The 10-m wind speed is under the
55	influence of factors like the momentum transfer coefficient and the integrated similarity functions
56	at night. The wind speeds are more significantly overestimated in the plains and basin, while less
57	overestimated or even underestimated in the mountains, as a result of the effect on topographic
58	smoothing in the model. Moreover, the overestimation of small wind speeds at night is attributable
59	to the inapplicability of the Monin-Obukhov similarity theory (MOST) at night. The model captures
60	the vertical structure of temperature well, while the wind speed is outstandingly overestimated
61	below 1000 m, largely because of the turbulent diffusion coefficient (TDC). The difference between
62	the MOST and the mixing length theory, PBLH and Prandtl number is cited as the reason for the
63	difference between the TDC of the YSU and ACM2 schemes. The TDCs of the BL and MYJ $$
64	schemes are affected by the mixing length scale, which of BL is calculated on the basis of the effect
65	of buoyancy, while MYJ calculates it with the consideration of the effect of the total turbulent
66	kinetic energy. The PBLH of the BL scheme is better than the other schemes because of the better
67	simulation results of temperature. The difference in the PBLH by the YSU and ACM2 scheme
68	mainly comes from the Richardson number and the jagged PBLH of the MYJ scheme is due to the
69	coarse vertical resolution and the threshold value.





70 In general, to select the optimal scheme, it is necessary to offer different options for different regions 71 with different focuses (heat or momentum). (1) Temperature field. The BL scheme is recommended 72 for January in the NCP region, especially for Beijing, and the MYJ scheme is better for the other 73 three months. The ACM2 scheme would be a good match for the YRD region, where the simulation 74 differences between the four schemes are small. The topography of the SB region is more complex, 75 but for most of the areas in the basin, the MYJ scheme is proposed, but if more stations outside the 76 basin are involved, the BL scheme is recommended. The MYJ scheme is applied to the PRD region 77 in January and April, and the BL scheme in July and October. The MYJ scheme is counselled for 78 the NS region. (2) Wind field. The YSU scheme is recommended if the main concern is the near-79 surface layer, and the BL scheme is suggested if focusing on the variation in the vertical direction. 80 The Part II will analyse and evaluate the factors that may influence the choice of the schemes and 81 the results of model. The final evaluation of the parameterization scheme and uncertainties will lay 82 the foundation for the improvement of the modules and forecasting of the GRAPES CUACE 83 regional model developed independently in China.

#### 84 1 Introduction

85 The planetary boundary layer (PBL) is the part of the troposphere that is directly influenced by the 86 force of the earth's surface with an hour or less timescale(R. B. Stull, 1988). Parameterization is the 87 determining factor in the predictive accuracy and skill as it determines key aspects of simulated 88 weather(Bauer et al., 2015; Williams, 2005). In numerical weather prediction, meagre 89 computational resources limit the resolution of the model. Following this reason, physical processes 90 cannot be resolved by the model in that the spatial scales are smaller than the model grid distance. 91 The physical module in the model that characterizes small scales relative to the model resolution is 92 called the sub-grid physical process parameterization scheme(Zhou et al., 2017). As a typical sub-93 grid parameterization scheme, the spatial scale of turbulence is limited by the PBL height (PBLH) 94 and cannot be resolved by mesoscale weather prediction models and macroscale global climate 95 models with horizontal grid distances of magnitude of  $\sim 10$  km and  $\sim 100$  km. Therefore, the physical 96 module in the model that describes the effect of sub-grid turbulence on resolvable atmospheric 97 motion is called the PBL parameterization scheme. Even in the high resolution large-eddy **98** simulation (LES), small-scale turbulence requires parametric closure to characterize the role of sub-99 grid turbulence(Deardorff, 1980). The PBL parameterization scheme controls the evolution of 100 momentum, heat, water vapor, and mass within the PBL, and the evolution of these parameters is 101 particularly affected by the turbulent diffusion coefficients (TDCs)(W. Jia and Zhang, 2021; 102 Nielsen-Gammon et al., 2010; Oke et al., 2017). Depending on the turbulence closure method, the 103 PBL parameterization schemes can be divided into three main categories: nonlocal closure schemes,





104 local closure schemes, and hybrid nonlocal-local closure schemes, and the above schemes have their 105 own advantages and disadvantages(Cohen et al., 2015; Hu et al., 2010; Wenxing Jia and Zhang, 106 2020; Xie et al., 2012). 107 Since the early 1980s, the vertical diffusion scheme based on local gradients of wind and potential 108 temperature (i.e., local K-theory) has been applied in the National Centers for Environmental 109 Prediction (NCEP). However, as pointed out by many scholars, this scheme has many deficiencies, 110 of which the most critical is that the mass and momentum transport within the PBL is mainly 111 accomplished by the large-scale eddies besides the local small-scale eddies(Roland B. Stull, 1984; 112 Wyngaard and Brost, 1984). Therefore, the new scheme developed later incorporates a counter-113 gradient flux term to characterize the turbulent transport processes in large-scale eddies, such as 114 Medium-Range Forecast (MRF) scheme (i.e., nonlocal closure) (Hong and Pan, 1996; Troen and 115 Mahrt, 1986). This scheme has also been commonly used in China's self-developed 116 Global/Regional Assimilation and PrEdiction System (GRAPES) model because of its 117 computational simplicity and its ability to produce plausible results under typical atmospheric 118 conditions (Ma et al., 2021). Nevertheless, the MRF scheme has gradually shown some 119 shortcomings, the most typical being that when the wind speed is strong, the resulting mixing is too 120 strong and thus the PBLH is too high to be realistic (Mass et al., 2002; Persson et al., 2001). To 121 overcome this critical problem, one of the most commonly used and popular PBL parameterization 122 scheme has been introduced, which is the Yonsei University (YSU) scheme(Hong et al., 2006). YSU 123 scheme adds an additional entrainment term to the MRF scheme for explicitly calculating the 124 entrainment process of heat and momentum fluxes(Noh et al., 2003). It is still unclear why this 125 scheme is popular among scholars, either because it gives the best simulation results or simply 126 because the code of this scheme is 1, which is more convenient for the model setting. To be contrast, 127 a newer scheme, as a nonlocal scheme of the same series, has been developed that further considers 128 the issue of gray-zone of sub-grid scale turbulence, but this scheme has been rarely used and 129 evaluated (Hong and Shin, 2013). 130 Repairing the defects of local K-theory is possible by developing nonlocal closure schemes on the 131 one hand, and higher-order local closure method on the other hand. The most representative is the 132 higher-order closure scheme of the M-Y series proposed by Mellor and Yamada, such as Mellor-133 Yamada-Janjic (MYJ) scheme and Mellor-Yamada Nakanishi and Niino Level 2.5/3 134 (MYNN2/MYNN3) scheme(Janjić, 1990, 1994; Mellor and Yamada, 1974, 1982; Nakanishi and 135 Niino, 2004). The higher-order closure schemes are capable of representing a well mixing PBL 136 structure, however, these schemes are computationally more expensive due to the addition of a 137 prognostic turbulent kinetic energy (TKE). In addition to the widely used local closure schemes of 138 the M-Y series, there is another local closure scheme that has been evaluated extensively. This

139 scheme is the Bougeault and Lacarrere (BL) scheme(Bougeault and Lacarrere, 1989), but there are





140	several differences between the BL and M-Y schemes. (1) In the parameterization of the turbulent
141	heat flux, an additional counter-gradient flux term is taken into account in the convective PBL, but
142	this counter-gradient term is different from that in the nonlocal closure scheme, which is a constant
143	$(= 0.7 \cdot 10^{-5} \text{ K cm}^{-1})$ in the BL scheme. (2) The turbulent diffusion coefficient in the BL scheme
144	is calculated similarly to the M-Y schemes, but the stability functions and mixing length are different
145	from M-Y schemes.
146	In addition to the typical nonlocal closure schemes and local closure schemes, there are also hybrid
147	nonlocal-local closure schemes, typically represented by the Asymmetric Convective Model version
148	2 (ACM2) scheme. ACM2 scheme operates based on the development of ACM1 that is modified
149	based upon the Blackadar convective model (Blackadar, 1962). The upward transport within the PBL
150	is mainly by buoyancy, which is transmitted upward from the lowest level to other levels, while
151	downward is transported level-by-level(Pleim, 2007). The deficiency of the ACM1 scheme is that
152	upward transport is not better represented when the vertical resolution of the model increases. In
153	response to compensating for the shortcomings of the ACM1 scheme, the ACM2 scheme adds level-
154	by-level transport to the upward level. The ACM2 scheme have the highest universality and was
155	most suitable for the study of meteorological elements in desert region(Meng Lu et al., 2018; Wang
156	et al., 2017).
157	At present, a total of 12 PBL parameterization schemes have been developed and evaluated in the
158	currently popular mesoscale Weather Research and Forecasting (WRF) model. The continuous
159	improvement of numerical simulation techniques brings opportunities for the update and
160	development of PBL parameterization schemes. Many scholars hope that by comparing the PBL
161	parameterization schemes, they can select one scheme that better reflects the changes in
162	meteorological parameters (e.g., temperature, relative humidity and wind speed/direction),
163	pollutants and the structures of the PBL. Recent review studies have shown that although many
164	studies on the evaluation and comparison of PBL parameterization schemes have been undertaken,
165	there is still no uniform conclusion on which PBL parameterization scheme performs best (Wenxing
166	Jia and Zhang, 2020). Moreover, most of the evaluation work on PBL parameterization schemes is
167	done for individual cases or a particular region(Avolio et al., 2017; Diaz et al., 2021; Falasca et al.,
168	2021; Ferrero et al., 2018; He et al., 2022; Shen et al., 2022). In spite of those, simulation results for
169	the PBL parameterization schemes are more uniform: $(1)$ the simulation of temperature is better
170	than that of relative humidity, and the simulation of wind speed and direction is worse. (2) The
171	simulation results of the nonlocal closure scheme are better under unstable conditions, while the
172	local closure scheme for stable conditions. However, these general conclusions are open to
173	speculation and debate(Wenxing Jia and Zhang, 2020). Many previous studies have been biased
174	towards the assessment of basic meteorological parameters, of course, which is the basic work. Due
175	to the indirect output of TDC by the model, there are fewer relevant studies to investigate the impacts





176	of turbulent diffusion on meteorological parameters. Moreover, turbulent diffusion is the key factor
177	to control the vertical mixing of momentum and scalars within the PBL. Even if there is not enough
178	turbulence observation data, it can be further analyzed and discussed according to the simulation
179	results. Aimed at remedying the current research deficiencies, this study first selects four typical
180	boundary layer parameterization schemes (nonlocal scheme: YSU, local scheme: MYJ and BL,
181	hybrid nonlocal-local scheme: ACM2) for five typical regions (NCP, YRD, SB, PRD and NS) in
182	China, and then assesses the performance capability of different PBL parameterization schemes in
183	different regions. The reasons for the differences in performance of meteorological parameters
184	between observation and simulation are illustrated in terms of temporal and regional variability.
185	Then, the mechanistic implications behind the differences are explored between schemes. In
186	addition, we further carry out the comparative analysis of the vertical structure of the PBL, turbulent
187	diffusion and the PBLH. The first part of this study (i.e., Part I) aims to be able to have a qualitative
188	and quantitative assessment of the PBL parameterization schemes in different regions for other
189	researchers to use as a reference when doing simulation studies. The second part (i.e., Part II)
190	focuses on the analysis of some uncertain factors that may affect the model simulation results,
191	chiefly including: the influence of meteorological initial and boundary conditions, underlying
192	surface (mainly considering the impact of urban and water bodies), near-surface layer (N-SL)
193	scheme (the PBL and N-SL schemes must match each other), the effect of model version update,
194	the influence of regional horizontal and vertical resolution, etc. We hope that we can dissect the
195	effect of uncertainties from some aspects that we are concerned about.

#### **196 2** Data and methods

#### 197 2.1 Data

198 Hourly meteorological observation data. The China Meteorological Administration (CMA) has 199 over 2400 automatic weather stations (AWSs), and the stations record variables such as temperature, 200 relative humidity, pressure, wind speed, wind direction and precipitation amount. In the NCP, YRD, 201 SB, PRD, NS regions, 576 stations, 455 stations, 341 stations, 128 stations, and 55 stations have 202 been selected, respectively (illustrated by gray cross in Fig. 1b-f). Observational data for four 203 months January, April, July and October 2016 have been selected and comparatively analyzed. 204 L-band radiosonde observation data. A total of 120 observation stations are equipped with L-band 205 radiosonde systems in China, which provide fine-resolution (1 Hz, and the rise rate is ~6 m s<sup>-1</sup>) 206 vertical profiles of temperature, relative humidity and wind speed and direction three times (08:00, 207 14:00 and 20:00 Beijing Time, BJT) a day (illustrated by red triangle in Fig. 1b-f). Four sounding 208 stations have been selected for each region, including different underlying surface conditions as 209 much as possible. In the NCP region, two plain stations (Beijing and Xingtai) and two mountain





210 stations (Zhangjiakou and Zhangjiu) have been picked. In the YRD region, one station closer to the 211 ocean (Shanghai), two stations with complex underlying surface (Anqing and Quzhou), and one 212 plain station (Nanjing) have been opted. In the SB region, three in-basin stations (Wenjiang, 213 Shapingba and Daxian) and one out-of-basin station (Hongyuan, with an altitude of 3491 m) have 214 been selected. In the PRD region, two plain stations (Qingyuan and Heyuan) and two stations 215 (Yangjiang and Shantou) closer to the ocean have been singled out. In the NS region, along the 216 Qilian mountains, four stations have been chosen, Mazongshan, with an altitude of 1770 m, Jiuquan, 217 with an altitude of 1477 m, Zhangye, with an altitude of 1460 m, and Minqin, with an altitude of 218 1367 m.

# 219 2.2 Model settings

220 In this study, we adopt the model WRF-ARW (Advanced Research Weather Research and 221 Forecasting) version 3.9.1 to evaluate the performance of PBL schemes. Long-term three-222 dimensional simulation experiments are conducted in 1 month of each season of 2016 (i.e., January, 223 April, July and October). Seven nested domains (D1, D2, D3, D4, D5, D6 and D7) are defined (Fig., 224 1a), with horizontal grid spacings of 75 km (74 × 74 grid cells, 9°N -59°N, 61°E -146°E), 15 km 225 (281 × 281 grid cells, 13°N - 51°N, 77°E - 136°E), 3 km (331 × 331 grid cells, 34°N - 44°N, 111°E 226 -124°E), 3 km (316 × 356 grid cells, 26°N - 34°N, 114°E - 126°E), 3 km (331 × 331 grid cells, 25° 227 N -33.5°N, 100°E-110°E), 3 km (236 × 301 grid cells, 18°N -25°N, 110°E -119°E), and 3 km (226 228 × 351 grid cells, 36°N -42.7°N, 94°E -107°E), respectively. Along the vertical direction, 48 vertical 229 layers are configured blow the top, and the model top is set to the 50 hPa. To resolve the PBL 230 structure finely, 21 vertical layers are set below 2 km (i.e., the specific setting of vertical levels is  $\sigma$ 231 = 1.000, 0.997, 0.994, 0.991, 0.988, 0.985, 0.980, 0.975, 0.970, 0.960, 0.950, 0.940, 0.930, 0.920,232 0.910, 0.895, 0.880, 0.865, 0.850, 0.825, 0.800). The initial and boundary conditions of 233 meteorological fields are set up by using the NCEP Global Forecast System (GFS) Final (FNL) 234 gridded analysis datasets, with a resolution of  $1^{\circ} \times 1^{\circ}$  (https://rda.ucar.edu/datasets/ds083.2/, last 235 access: 4 August, 2022). The Moderate Resolution Imaging Spectroradiometer (MODIS) dataset 236 includes 20 land-use categories(Broxton et al., 2014). The physical parameterization used in the 237 present model is listed in Table 1.







Figure 1. (a) Map of terrain height in the seven nested model domains. (b-f) Domain 3-7 correspond
to the North Plain China (NCP), the Yangtze River Delta (YRD), the Sichuan Basin (SB), the Pearl
River Delta (PRD) and the Northwest Semi-arid (NS), respectively. The locations of surface
meteorological stations and sounding stations are marked by the gray crosses, red triangles,
respectively. The purple dots indicate the major city sites that are our main focus in each region.
Table 1. A brief description of the parameterization scheme in the model.

Namelist option	Description	Input option	Reference
mp_physics	Morrison double- moment scheme	10	(Morrison et al., 2009)
ra_lw_physics	RRTMG scheme	4	(Iacono et al., 2008)
ra_sw_physics	RRTMG scheme	4	(Iacono et al., 2008)
cu_physics	Grell-3D scheme	5	(Grell and Dévényi, 2002)
sf_sfclay_physics	MM5 similarity scheme	1	(Jiménez and Dudhia, 2012)
sf_surface_physics	Noah land surface scheme	2	(Chen and Dudhia, 2001)
sf_urban_physics	Single-layer UCM	1	(Kusaka et al., 2001)





	scheme		
sf lake physics	CLM4.5 lake 1 scheme		(Gu et al., 2015)
bi_iane_physics			(========;====;===;
	YSU scheme	1	(Hong et al., 2006)
	MYJ scheme	2	(Mellor and Yamada, 1982)
bl_pbl_physics	ACM2 scheme	7	(Pleim, 2007)
	BL scheme	8	(Bougeault and Lacarrere, 1989)

245 All simulations embodied a total of 16 months. The 40 h simulation is conducted beginning from

246 00:00 UTC of 1d ago for each day (i.e., 492 simulation experiments), the first 16 h of each

simulation is considered as the spin-up period, and results obtained from the following 24 h

**248** simulations are analyzed for the present study.

### 249 2.3 Description of PBL parameterization schemes

#### 250 2.3.1 YSU scheme

251 The YSU is a first-order nonlocal scheme with an explicit treatment entrainment process at the top252 of the PBL:

253 
$$\frac{\partial c}{\partial t} = \frac{\partial}{\partial z} \left[ K_c \left( \frac{\partial c}{\partial z} - \gamma_c \right) - \overline{(w'c')_h} \left( \frac{z}{h} \right)^3 \right] (1)$$

254 where *c* denotes *u*, *v*,  $\theta$ , and the  $\gamma_c = b \frac{\overline{(w'c')_0}}{w_{s0}h}$  is the counter-gradient flux term, which increases 255 the nonlocal effect due to the large scale turbulence. *z* and *h* are the height of a level of the model

**256** and PBLH, respectively. The PBLH is defined by the bulk Richardson number method:

257 
$$h = Rib_{cr} \frac{\theta_{va}|U(h)|^2}{g(\theta_v(h) - \theta_s)}$$
(2)

where g is the gravity, and *Rib* is the critical bulk Richardson number, with a value of 0.25 under stable conditions and 0 under unstable conditions.  $\theta_{va}$  is the virtual potential temperature at the lowest model level,  $\theta_v(h)$  is the virtual potential temperature at *h*,  $\theta_s$  is the appropriate temperature near the surface ( $\theta_s = \theta_{va} + \theta_T$ ,  $\theta_T$  is the virtual temperature increment). Compared to the predecessor of the YSU scheme, the entrainment process is additionally treated explicitly (i.e., the last term on the right side of the Eq. (1))

264 Another key variable is  $K_c$ , which is the turbulent diffusion coefficient (TDC), and can be 265 expressed based on the Monin-Obukhov similarity theory (MOST) as:

$$266 K_c = \frac{\kappa u_* z}{\phi_c} \left(1 - \frac{z}{h}\right)^2 (3)$$

where  $u_*$  is the surface frictional velocity and  $\phi_c$  is dimensionless function, the expressions for different stability conditions are:

**269** i. Unstable and neutral conditions:





- **270**  $\phi_m = \left(1 16\frac{0.1h}{L}\right)^{-1/4}$  (4a)
- **271**  $\phi_h = \left(1 16\frac{0.1h}{L}\right)^{-1/2}$  (4b)
- 272 ii. Stable condition:
- 273  $\phi_m = \phi_h = \left(1 + 5\frac{0.1h}{L}\right)$  (4c)
- 274 The TDC of momentum (i.e.,  $K_m$ ) is first calculated in the model, and then the TDC of heat (i.e.,

275  $K_h$  is calculated, using the Prandtl number (i.e.,  $Pr = \frac{K_m}{K_k}$ ). The TDC controls the vertical mixing

process of momentum and scalars within the PBL, and it is crucial that it needs to be accuratelydescribed.

#### 278 2.3.2 MYJ scheme

- 279 The MYJ scheme is a one-and-a-half order local closure scheme with a prognostic equation for
- **280** turbulent kinetic energy (TKE,  $TKE = e = \frac{1}{2}(u'^2 + v'^2 + w'^2)$ ):
- $281 \qquad \frac{\partial \overline{e}}{\partial t} = -\frac{1}{\overline{\rho}} \frac{\partial}{\partial z} \overline{w' p'} \overline{w' u'} \frac{\partial \overline{u}}{\partial z} \overline{w' v'} \frac{\partial \overline{v}}{\partial z} \frac{\partial}{\partial z} \overline{w' e'} + \frac{g}{\theta_v} \overline{w' \theta_v'} \varepsilon$ (5)

The first term on the right side of Eq. (5) is a pressure correlation term which describes TKE is
redistributed by pressure perturbations, the second and third terms is a shear production/loss term,
the fourth term represents the turbulent transport of TKE, the fifth term describes the buoyant
production/consumption term, and the sixth term represents viscous dissipation of TKE. To close
the TKE equation, the turbulent fluxes must be parameterized. Based on the gradient transport
theory (i.e., K-theory), the turbulent fluxes can be indicated as:

 $\begin{array}{ll}
\mathbf{288} & \overline{w'u'} = -K_m \frac{\partial \overline{u}}{\partial z} & \text{(6a)} \\
\mathbf{289} & \overline{w'v'} = -K_m \frac{\partial \overline{v}}{\partial z} & \text{(6b)}
\end{array}$ 

**290** 
$$\overline{w'\theta'} = -K_h \frac{\partial\overline{\theta}}{\partial z}$$
 (6c)

**291** The TDC is proportional to the square root of TKE, and can be expressed as:

292 
$$K_m = S_m le^{1/2}$$
 (7a)

293  $K_h = S_h l e^{1/2}$  (7b)

294 where *l* is mixing length and can be described as  $l = \frac{l_0 \kappa z}{\kappa z + l_0}$ , where  $l_0 = \alpha \frac{\int_0^\infty z e^{1/2} dz}{\int_0^\infty e^{1/2} dz}$ ,  $\alpha$  is an

295 empirical constant (=0.1). When z converges to a very small value, l converges to κz. However, as

- **296** z converges to a very large value, l converges to  $l_0$ .
- **297** To obtain the  $S_m$  and  $S_h$  in Eq. (7),  $G_m$  and  $G_h$  are defined as:





**298**  $G_m = \frac{l^2}{2e} \left[ \left( \frac{\partial \overline{u}}{\partial z} \right)^2 + \left( \frac{\partial \overline{v}}{\partial z} \right)^2 \right]$ (8a)

**299** 
$$G_h = -\frac{l^2}{2e} \frac{g}{\theta_v} \frac{\partial \theta_v}{\partial z}$$
 (8b)

- **300**  $S_m$  and  $S_h$  are functions of  $G_m$  and  $G_h$ , and can be denoted as:
- **301**  $S_m(6A_1A_2G_m) + S_h(1 3A_2B_2G_h 12A_1A_2G_h) = A_2$  (9a)

**302** 
$$A_1(1+6A_1^2G_m-9A_1A_2G_h)-S_h(12A_1^2G_h+9A_1A_2G_h)=A_1(1-3C_1)$$
 (9b)

- **303** where  $[A_1, A_2, B_1, B_2, C_1] = [0.660, 0.657, 11.878, 7.227, 0.001].$
- **304** The PBLH in the MYJ scheme is defined as the height at which the TKE is reduced to a critical
- **305** value of  $0.1 \text{ m}^2 \text{ s}^{-2}$ .

## 306 2.3.3 ACM2 scheme

- 307 Unlike the YSU scheme, the ACM2 scheme applies the transilient matrix to deal with the
- **308** contribution of nonlocal fluxes. The governing equation can be expressed as:

$$309 \quad \frac{\partial C_i}{\partial t} = f_{conv} M u C_1 - f_{conv} M d_i C_i + f_{conv} M d_{i+1} C_{i+1} \frac{\Delta z_{i+1}}{\Delta z_i} + \frac{\partial}{\partial z} \left[ K_c (1 - f_{conv}) \frac{\partial C_i}{\partial z} \right]$$
(10)

310 The first three terms on the right side of Eq. (10) represent nonlocal mixing effect and the fourth

- 311 term represents local mixing effect. Where  $C_i$  is the variable at layer *i*, Mu is the nonlocal upward
- 312 convective mixing rate,  $Md_i$  is the downward mixing rate from layer *i* to layer *i*-1,  $\Delta z_i$  is the
- 313 thickness of layer *i*, and  $C_1$  represents the variable at the lowest layer in the model.  $f_{conv}$  is the

314 weighting factor for the nonlocal and local effects (i.e.,  $f_{conv} = \frac{K_h \gamma_h}{K_h \gamma_h - K_h \frac{\partial \theta}{\partial z}}$ , where the value of  $f_{conv}$ 

**315** ranges from 0 to 1, a larger *f<sub>conv</sub>* indicates stronger nonlocal mixing.

316 There are two methods to calculate the TDC, and the first method is the same as the YSU scheme,

- 317 i.e., Eq. (3), but there is also a very stable condition in the ACM2 scheme. In this case, the 318 dimensionless function can be expressed as  $\phi_m = \phi_h = \left(5 + \frac{0.1h}{L}\right)$ .
- **319** The second calculation principle is based on the mixing length theory, which uses mixing length
- **320** and stability function to calculate TDC:

321  $K_h = 0.01 + l^2 \sqrt{ss} f_h(Ri)$  (11)

322 where l is similar to the MYJ scheme, but  $l_0$  is a constant (=80), ss is the wind shear (ss =

**323**  $(\partial \overline{u}/\partial z)^2 + (\partial \overline{v}/\partial z)^2)$ , 0.01 denotes the minimum value of the TDC in the model, and  $f_h(Ri)$  is

- 324 the empirical stability functions of gradient Richardson number of heat.
- 325 i. when  $Ri \ge 0$ :

**326** 
$$f_h(Ri) = (1 - 25Ri)^{1/2}$$
 (12a)

327 ii. when 
$$Ri < 0$$

328 
$$f_h(Ri) = \frac{1}{1+10Ri+50Ri^2+5000Ri^4} + 0.0012$$
 (12b)





- **329** Similarly, the empirical stability functions of momentum can be indicated as:
- **330** i. when Ri < 0:
- **331**  $f_m(Ri) = Pr \cdot f_h(Ri) + 0.00104$  (13a)

332  $K_m = 0.01 + l^2 \sqrt{ss} f_m(Ri)$  (13b)

- 333 ii. when  $Ri \ge 0$ :
- $334 \qquad K_m = Pr \cdot K_h$ (13c)

 $\label{eq:335} {\ \ The ACM2 scheme has a range setting for the TDC in the model with a minimum value of 0.01 \ m^2}$ 

**336**  $s^{-2}$  and a maximum value that cannot exceed 1000 m<sup>2</sup> s<sup>-2</sup>.

337 The PBLH discrimination in the ACM2 scheme is similar to the YSU scheme, and is defined with338 the bulk Richardson number method. The difference is that the entrainment region at the top of the

**339** PBL is considered in the ACM2 scheme, and turbulence still exists due to the wind shear and thermal

- 340 penetration. Therefore, special processing of the PBLH is required under unstable and stable
- **341** conditions.

# 342 2.3.4 BL scheme

343 The BL scheme is also a one-and-a-half order local closure scheme, and the TDC is calculated in a

344 similar way to Eq. (7) of the MYJ scheme. Nevertheless, the function  $S_m$  and mixing length (i.e., l)

**345** are different from the MYJ scheme. In the BL scheme,  $S_m$  is a constant 0.4 and the *l* is divided into

346 upward and downward mixing length (i.e., *lup* and *ldown*), which are defined as:

347  $\int_{z}^{z+l_{up}} \beta[\theta(z) - \theta(z')] dz' = e(z) \quad (14a)$ 

348  $\int_{z-l_{\text{down}}}^{z} \beta[\theta(z') - \theta(z)] dz' = e(z) \quad (14b)$ 

 where,  $\beta$  is the buoyancy coefficient and the *l* is equal to the minimum of  $l_{up}$  and  $l_{down}$  (i.e., l =  $min(l_{up}, l_{down})$ ). It is worth noting that in the BL scheme the TDC of heat is equal to the TDC of momentum (i.e.,  $K_h = K_m$ ). In addition, the PBLH of the BL scheme is defined as the height at which the virtual potential temperature of a layer is greater than that of the first layer by 0.5 K.

To accommodate different methods of calculating PBLH for different schemes and to evaluate the
simulation performance of PBLH, two methods are employed to calculated PBLH using observed
data in this study, the first being the bulk Richardson number method, and the detailed calculation
principle is as follows (Miao et al., 2018):

357 
$$Ri(z) = \frac{(g/\theta_s)(\theta_z - \theta_s)(z - z_s)}{(u_z - u_s)^2 + (v_z - v_s)^2 + bu_*^2}$$
 (15)

**358** here z is the height, g is the gravity,  $\theta$  is the virtual potential temperature, u and v are the components **359** of the horizontal wind, b is a constant, and  $u_*$  is the friction velocity. The subscript "s" indicates **360** the near-surface. Since the friction velocity is much smaller in magnitude than the wind shear, the **361** b is set to 0, ignoring the effect of surface friction (Vogelezang and Holtslag, 1996; Seidel et al., **362** 2012). The PBLH is estimated as the lowest layer height when Ri reaches a critical value of 0.25.





- 363 The second method adopts the same calculation method as the BL scheme, i.e., the virtual potential
- **364** temperature method:
- **365**  $\Delta \theta_{\nu|PBLH} = \theta_{\nu 1} + 0.5$  (16)
- 366 The PBLH is the height when the virtual potential temperature exceeds the virtual potential
- **367** temperature of the first level by 0.5 K.

### 368 2.4 Evaluation of the model

- 369 To evaluate the PBL schemes and the performance of the model for estimating meteorological
- 370 variables, the statistical parameters used in this statistical analysis are defined as follows(Emery et
- **371** al., 2017):
- **372** Index of agreement (IOA):

$$373 \quad IOA = 1 - \frac{\left[\sum_{i=1}^{n} |X_{sim,i} - \overline{X}_{obs,i}|^2\right]}{\left[\sum_{i=1}^{n} (|X_{sim,i} - \overline{X}_{obs}| + |X_{obs,i} - \overline{X}_{obs}|)^2\right]}$$
(17)

374 Mean bias (MB):

375 
$$MB = \frac{1}{n} \sum_{i=1}^{n} (X_{sim,i} - X_{obs,i})$$
 (18)

**376** Root mean square error (RMSE):

377 
$$RMSE = \sqrt{\frac{1}{n} \sum_{i=1}^{n} (X_{sim,i} - X_{obs,i})^2}$$
(19)

**378** Normalized standard deviations (NSD):

**379** 
$$NSD = \frac{\sqrt{\frac{1}{n-1}\sum_{i=1}^{n} (X_{sim,i} - \overline{X_{sim}})^2}}{\sqrt{\frac{1}{n-1}\sum_{i=1}^{n} (X_{obs,i} - \overline{X_{obs}})^2}}$$
 (20)

**380** Relative bias (RB):

**381** 
$$RB = \frac{\overline{X_{sim}} - \overline{X_{obs}}}{\overline{X_{obs}}} \times 100\%$$
(21)

382 Where  $X_{sim,i}$  and  $X_{sim,i}$  represent the value of simulation and observation, respectively, *i* refers 383 to time and *n* is the total number of time series.  $\overline{X_{sim}}$  and  $\overline{X_{obs}}$  represent the average simulation 384 and observation.

The Taylor diagram is a compact tool that displays simultaneously the values of four statistical
parameters: IOA, NSD, RB, and RMSE. In particular, in these diagrams the perfect match of a
model with the observations would be the point with IOA=1, NSD=1, RB=0 and RMSE=0.

#### **388 3 Results and discussion**

In section 3.1, the mechanistic analysis of the PBL schemes for the simulation of near-surface
 meteorological parameters, including 2-m temperature, 2-m relative humidity, 10-m wind speed and
 direction. Section 3.2 gives an in-depth analysis of different schemes for PBL vertical structure. In
 section 3.3, the PBLH was evaluated for different schemes. In section 3.4, the reason for the





- **393** differences in turbulent diffusion are interrogated from the calculation principle of the schemes.
- 394 Section 3.5 summarizes the performance and expressiveness of different PBL schemes in different
- **395** regions, and recommends the optimal choice of PBL scheme.

# **396 3.1** surface meteorological variables

**397 3.1.1** 2-m temperature and relative humidity

398 To better analyze the variation of the time series, we selected representative stations in different 399 regions. Figure 2 shows the diurnal variation of 2-m temperature (i.e., T2) for four months (i.e., 400 January, April, July and October 2016) at representative sites (indicated in the orange dots in Fig. 1) 401 in the five regions. The model basically captures the daily variation characteristics of T2, but there 402 are significant differences between different regions and seasons. The simulated results for July are 403 closest to the observed values (Fig. 2 c1-c5), anywhere. Overall, the mean biases (MBs) of the 404 diurnal variation of T<sub>2</sub> predicted in July for the NCP, YRD, SB, PRD and NS regions are 0.61~1.19, 405 -0.02~-0.56, -0.32~-0.60, -0.38~-0.69, and 0.28~0.81 °C, respectively. However, a smaller value of 406 the mean bias does not mean that the simulated value of the model is closer to the observed value. 407 For example, if one overestimation and the other underestimation occur during the day and night, 408 the average results will cancel each other out, resulting in a small mean bias. Accordingly, more 409 statistical parameters are needed to further evaluate the optimal scheme. In the other three months 410 (January, April, and October), the simulated results of T2 are overestimated to varying degrees 411 during daytime in the YRD, SB and PRD regions, while in the NS region, T2 are underestimated to 412 varying degrees (Fig. 2 a2-b5, d2-d5 and Table 1). In the NCP regions, T2 presents underestimation 413 in January by the model with the YSU, ACM2, BL and MYJ schemes are -1.33, -1.21, -0.52 and -414 1.18 °C, respectively, while overestimation arises in the other three months (Table 1). In the five 415 regions, the simulation results of the nighttime T2 outperform those of the daytime T2 for almost 416 four months. At night, the MYJ scheme shows a significant underestimation of T2 for all months in 417 five regions compared to the other three schemes (Fig. 2 and Table 1). The simulation results of Hu 418 et al. (2010) and Xie et al. (2012) have also obtained the lowest temperature for the MYJ scheme 419 during the nighttime.









Figure 2. Time series of diurnal variation of observed and simulated 2-m temperature in five
regions for four seasons.

423 The 2-m temperature does not actually represent the air temperature at a height of 2 m, but it is a 424 diagnostic variable of the near-surface temperature. It is calculated from the surface temperature 425  $(T_s)$ , the sensible heat flux (HFX) and the heat transfer coefficient (*C<sub>h</sub>*). The T<sub>s</sub> is a prognostic 426 variable, which is obtained in the model through the energy balance equation:

427  $(1-\alpha)S\downarrow +L\downarrow -L\uparrow +G - HFX - LH = 0$  (22)

428 where  $\alpha$  is the albedo of the underlying surface,  $S \downarrow$  represents the downward of the shortwave 429 radiation,  $L \downarrow$  is the downward of the longwave radiation emitted by the cloud and atmosphere, 430  $L\uparrow$  is the upward of the longwave emitted by the ground surface, G is the ground heat flux, and it 431 is positive when heat transfers from the soil to the near surface, HFX is the sensible heat flux and 432 LH is the latent heat flux. 433 We compare the effects of the six variables mentioned above on  $T_s$  with the expectation that we can 434 further examine the reasons for the differences in T<sub>2</sub> variation between different schemes. The YSU 435 scheme is used as a control and analyzed in comparison with each of the schemes. 436 The nonlocal closure scheme (YSU) and the local closure scheme (MYJ) are compared first. 437 Theoretically, the greater the downward shortwave radiation ( $\mathbf{S} \downarrow$ ) becomes, the more energy reaches 438 the ground, and the higher the surface temperature  $(T_s)$  is. After comparing the YSU and MYJ 439 schemes, the surface temperature does not show a proportional change with the downward 440 shortwave radiation, and the  $\mathbf{S} \downarrow$  of the MYJ scheme is almost the same as that of the YSU scheme 441 (Fig. 3 a1-e1), but the T<sub>s</sub> of the MYJ scheme is the lowest (Fig. 4 a1-e1). Therefore, the  $\mathbf{S} \downarrow$  is not 442 the main factor that causes the difference in  $T_s$  between the two schemes. There is no significant 443 difference in the upward/downward longwave radiation between these two schemes (Fig. 3 a2-e3), 444 so the effect of longwave radiation on the  $T_s$  can also be excluded. During the daytime, the MYJ





445 scheme transfers less heat from the surface to the soil than the YSU scheme (Fig. 3 a4-e4), and the 446  $T_s$  of the MYJ scheme should be higher than that of the YSU scheme. But that's not how it has 447 turned out (Fig. 4 a1-e1). Thus, the ground heat flux (G) is also not a key factor that directly affects 448 the T<sub>s</sub>. The latent heat flux (LH) is mainly related to water vapor (or relative humidity), so further 449 attention is paid to the effect of sensible heat flux (HFX) on T<sub>s</sub> (Fig. 3 a5-e6). The HFX is determined 450 by the difference between the surface temperature and the 2-m temperature ( $T_s$ - $T_2$ ), and the heat 451 transfer coefficient (Ch). MYJ has the largest HFX, and transfers more heat from the surface to the 452 atmosphere, resulting in the largest energy loss at the surface, which should correspond to the 453 smallest Ts (Fig. 3 a6-e6, 4 a1-e1). The smallest difference between the two temperatures indicates 454 a smaller temperature gradient and more uniform mixing, symbolizing the largest  $C_h$ , which is also 455 true (Fig. 4 a2-e3). A larger  $C_h$  would lead to higher T<sub>2</sub> during the day. Although the T<sub>s</sub> of the MYJ 456 scheme is significantly lower than YSU scheme, it makes the  $T_2$  higher due to the large  $C_h$ . During 457 the daytime, the less heat is transferred from the surface to the soil in the MYJ scheme, which results 458 in lower soil temperature. During the nighttime, the difference in HFX and temperature gradient 459 between the two schemes decreases, and the lower soil temperature results in lower T<sub>s</sub> and T<sub>2</sub>.



460

461Figure 3. Time series of diurnal variation of (a1-e1) downward shortwave radiation  $(S \downarrow)$ , (a2-e2)462downward longwave radiation  $(L \downarrow)$ , (a3-e3) upward longwave radiation  $(L \uparrow)$ , (a4-e4) ground heat463flux (G), (a5-e5) latent heat flux (LH), and (a6-e6) sensible heat flux (HFX) by four PBL schemes464in five regions in January.465The VSS5The VSS

The differences between the YSU scheme and the ACM2 scheme are further explored. Except for
the NCP and NS regions, the S↓ of the ACM2 scheme is smaller than that of the YSU scheme in
the other three regions (i.e., YRD, SB and PRD) (Fig. 3 b1-d1). The HFX of the ACM2 scheme is
smaller than that of the YSU scheme (Fig. 3 b6-d6), the heat loss from the surface of the ACM2





469 scheme is less, and the Ts of the ACM2 scheme should be higher. However, the Ts corresponding to 470 the ACM2 scheme is lower than that of the YSU scheme (Fig. 4 b1-d1), reflecting that the  $S \downarrow$ 471 varies proportionally with the  $T_s$ , and it is the main factor controlling the  $T_s$  variation. In the ideal 472 case, assuming the same temperature gradient for the nonlocal schemes, the  $T_2$  of the YSU scheme 473 should also be higher than that of the ACM2 scheme when the Ts of the YSU scheme is higher than 474 that of the ACM2 scheme with the same  $C_h$ . But in fact, it can be seen that the  $C_h$  of the ACM2 475 scheme and YSU scheme are the same (Fig. 4 b2-d2), and the temperature gradient of the YSU 476 scheme is greater than that of ACM2 scheme (Fig. 4 b3-d3). The T<sub>2</sub> of the ACM2 scheme should be 477 slightly higher than the ideal case, closer to the  $T_2$  of the YSU scheme, and even may also exceed 478  $T_2$  of the YSU scheme. At night, the  $T_s$  of the YSU scheme is lower than that of the ACM2 scheme, 479 and the  $C_h$  of the YSU scheme is smaller than that of the ACM2 scheme (Fig. 4 a1-e2). Meanwhile, 480 the difference in HFX between the two schemes is not obvious at night, contributing to lower T<sub>2</sub> of 481 the YSU scheme. In both NCP and NS regions, there is no significant difference in downward 482 shortwave radiation between two schemes, and no noticeable difference between T<sub>2</sub> and T<sub>s</sub>.



483



487 Then, the reasons for the simulated temperature difference between the YSU scheme and the BL 488 scheme are demonstrated. During the daytime, the  $\mathbf{S} \downarrow$  of both schemes are the same (Fig. 3 a1-e1), 489 but the HFX of the BL scheme is smaller than that of the YSU scheme (Fig. 3 a6-e6), less heat is **490** loss at the surface, hence the T<sub>s</sub> should be higher than that of the YSU scheme (Fig. 4 a1-e1). The 491  $C_h$  of both schemes are the same, thus, the BL scheme has a higher T<sub>2</sub> (Fig. 2, 4a1-e2). At night, the 492 HFX of BL scheme is larger than that of YSU scheme, and more heat is transferred from atmosphere 493 to the surface, and the larger  $C_h$  resulting in higher T<sub>2</sub> (Fig. 3 a6-e6, 4 a1-e2). 494 Finally, we can also uncover the reasons for the difference between the local closure schemes (MYJ

495 and BL). The larger HFX of the MYJ scheme leads to a lower  $T_s$  in the daytime, while the

496 temperature gradient of MYJ scheme is smaller than that of the BL scheme, and C<sub>h</sub> is larger than

497 BL scheme (Fig. 3 a6-e6, 4). Therefore, the difference in T<sub>2</sub> between the two schemes is smaller

**498** than that in  $T_s$ . The  $T_2$  of the MYJ scheme is closer to that of the BL scheme.





499	In conclusion, the causes of temperature differences simulated by the nonlocal closure schemes
500	should first focus on the effect of the downward shortwave radiation (S $\downarrow$ ), and when it comes to the
501	local closure scheme, the effect of HFX should be further concerned. All of the above results have
502	been analyzed for January 2016, and the results for the other three months are similar (Figs. S1-S6).
503	In terms of regional distribution differences, $T_{2}\ in$ the northern and near mountainous regions of the
504	NCP region is significantly underestimated in the daytime for January, April and October, while $T_{\rm 2}$
505	in other regions shows an overestimation (Fig. 5, S7, S8 a1-e1). The range of overestimated areas
506	is smaller than the underestimated in January, only in a small part of the area south of Hebei Province
507	(Fig. 5 a1-e1). The relative bias (RB) of the underestimated (overestimated) $T_2$ with the YSU, ACM2,
508	BL and MYJ schemes are -0.60% (0.15%), -0.57% (0.17%), -0.43% (0.26%) and -0.60% (0.20%),
509	respectively in January. Also in these three months, temperature is overestimated at almost all
510	stations in the YRD region (RB=0.38%~0.50% in January, RB=0.49%~0.65% in April and
511	$RB{=}0.58\%{\sim}0.70\%$ in October) and underestimated at some stations along the coast (RB=-0.13%{\sim}-0.12\%)
512	0.24% in January, RB=-0.32%0.37% in April and RB=-0.23%0.28% in October) (Fig. 5, S7,
513	S8 a2-e2). The results show an overestimation of $T_2$ simulated in those stations in the basin for the
514	SB region as well as the simulation results of the stations in the plain for the NCP region, while for
515	the stations in the hilltop areas, the $T_2$ shows an underestimation (Fig. 5, S7, S8 a3-e3). In the PRD
516	region, the entire region exhibits an overestimation of $T_2$ , with the simulation results in October
517	$(RB=0.06\%\sim0.15\%)\ being\ significantly\ better\ than\ those\ in\ January\ (RB=0.59\%\sim0.81\%)\ and\ April$
518	(RB=0.56%~0.67%), with a lower degree of $T_2$ overestimation (Fig. 5, S7, S8 a4-e4). The BL
519	scheme simulates a higher $T_2$ and a large range of overestimated areas (about 167, 378, 252 and 100 $$
520	stations in NCP, YRD, SB and PRD regions). The NS region has a more complex topography and
521	higher elevation, and the $T_2$ is underestimated at almost all stations, with best simulation results in
522	October (RB=-0.04%~-0.21%) and worst in April (RB=-0.49%~-0.64%).
523	For July, when the temperature is higher, the simulation results are significantly different from the
524	other three months. The relative bias of $T_2$ simulated with the YSU, ACM2, BL and MYJ schemes $% T_2$
525	are 0.47%, 0.46%, 0.53% and 0.46%, respectively, in the NCP region. The overestimation results
526	are similar to daytime, with the most pronounced overestimation for the BL scheme. For the
527	southern region of the NCP, the $T_2$ is consistently overestimated regardless of the season (Fig. S9 $$
528	a1-e1). The $T_2$ at most stations are underestimated in the YRD region, which is different form the
529	other three months (Fig. S9 a2-e2). In summer, the temperature of the ocean, affected by the
530	subtropical high (prevailing southeasterly winds), is lower than that of the land, and the transport of
531	momentum is accompanied by the transport of heat from the sea to the land, causing the temperature
532	of the land to decrease. The $T_2$ of the basin area in the SB region is well reproduced, and no
533	significant overestimation occurs (Fig. S9 a3-e3). There is an underestimation of the T2 at most

534 stations in the PRD region, but to a lesser extent (Fig. S9 a4-e4). In contrast, for the NS region, the

537







- 535 temperature is overestimated for areas at lower elevations, while underestimated (or better
- **536** reproduced) for areas at higher elevations (Fig. S9 a5-e5).



542 The relative deviation of the nighttime T<sub>2</sub> simulations is less than that of the daytime, regardless of 543 the region and month (Fig. 6, S10-S12). The differences between the four schemes are more striking 544 at night compared to the daytime. The BL scheme simulates the highest T<sub>2</sub> and the MYJ scheme 545 simulates the lowest T<sub>2</sub> in the whole region (Fig. 6, S10-S12). Compared to the observed values, the 546 MYJ scheme is the best when all schemes overestimate the simulated temperature, but if there is an 547 underestimation, the MYJ scheme is no longer the best scheme. Later, a comprehensive statistical 548 evaluation of the schemes will be presented. For the NCP region, the overestimation and 549 underestimation in the whole region do not show a north-south divide (or a mountain-plain divide)





as in the daytime (Fig. 5, 6 a1-e1). In the YRD region, the temperature along the coastal area still shows a significant underestimation (Fig. 6, S10-S12 a2-e2). Similar to the daytime, the temperature at stations in the hill top areas of the SB region still presents an underestimation (Fig. 6, S10-S12 a3-e3). Most stations show the underestimation of  $T_2$  in July and October in the PRD region (Fig. S11-S12 a4-e4). In the NS region, the relative deviation of temperature simulations in October is greater than that in daytime (RB=0.17%~0.56%) (Fig. 5, S12 a5-e5).



556 557

#### Figure 6. Similar as Figure 5, but at night.

558 In summary, the simulation results of  $T_2$  have the following main characteristics. From the 559 perspective of differences between observations and simulations, (1) the simulation results for July 560 are better compared to the other three months. (2) The simulation results at night are better than 561 those at daytime, with less relative deviation. (3) The temperature is easily underestimated at higher 562 altitudes while overestimated in plains and basin areas. From the perspective of the differences 563 between the different schemes, (1) the differences in the performance of the four schemes are more 564 noticeable at night. (2) The difference in the simulation of temperature in the nonlocal closure





- 565 schemes is mainly attributed to the difference in downward shortwave radiation ( $S\downarrow$ ), and the 566 difference in the variation of sensible heat flux (HFX) needs to be further analyzed when the local 567 closure schemes are involved. (3) The BL scheme simulates the highest temperature and the MYJ 568 scheme for the lowest temperature. 569 The results for 2-m relative humidity (RH<sub>2</sub>) and T<sub>2</sub> correspond to each other, and the overestimation 570 of T<sub>2</sub> corresponds to the underestimation of RH<sub>2</sub>. The simulation of RH<sub>2</sub> still shows the best results 571 in July, with the highest simulated values for the MYJ scheme and the lowest for the BL scheme. 572 Except for the NS region, the simulated RH2 of the other four regions is almost underestimated.
- 573 This uniform trend in relative humidity may be due to errors in the initial field, which will be
- **574** discussed in Part II. Too much will not be repeated here (Figs. 2, 7).



576 Figure 7. Similar as Figure 2, but for 2-m relative humidity.

### 577 3.1.2 10-m wind speed and direction

575

Although the model simulates the diurnal cycle of wind speed, the wind speed shows different degrees of overestimation, and the mean bias is in the range of 0.86 m s<sup>-1</sup>~2.74 m s<sup>-1</sup> (Fig. 8, Table 2). This is also the conclusion reached in many previous studies and it is more widely accepted by the public. Except for the NCP region where the MYJ scheme has the largest mean bias (MB) value during the day (YSU: MB=1.6 m s<sup>-1</sup>, ACM2: MB=1.8 m s<sup>-1</sup>, BL: MB=1.7 m s<sup>-1</sup>, MYJ: MB=2.1 m s<sup>-1</sup>), while the BL scheme has the largest MB value at night regardless of the region (YSU: MB=1.3 m s<sup>-1</sup>, ACM2: MB=2.2 m s<sup>-1</sup>, MYJ: MB=1.7 m s<sup>-1</sup>) (Fig. 8, Table 2).







585

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586 Figure 8. Similar as Figure 2, but for 10-m wind speed.

587 Table 2. Mean bias of 2-m temperature, 2-m relative humidity and 10-m wind speed during day time

588 and nighttime by four PBL schemes in five regions and four seasons.

		Regions					
Variables			NCP	YRD	SB	PRD	NS
	Schemes/Se	asons					
		Jan	-1.33	0.91	0.20	1.84	-1.42
	VSU	Apr	0.27	1.35	0.77	1.74	-1.83
	130	Jul	1.41	-0.18	-0.20	-0.66	0.38
		Oct	0.91	1.69	0.68	0.25	-0.58
		Jan	-1.21	0.97	0.29	1.93	-1.29
	ACM2	Apr	0.34	1.28	0.64	1.79	-1.82
	ACM2	Jul	1.37	-0.21	-0.31	-0.49	0.27
T Day		Oct	0.98	1.54	0.61	0.21	-0.50
1 <sub>2</sub> -Day		Jan	-0.52	1.33	0.73	2.32	-0.76
	DI	Apr	0.89	1.71	1.07	2.00	-1.43
	BL	Jul	1.60	0.04	-0.16	-0.38	0.55
		Oct	1.49	1.88	1.09	0.46	-0.09
	MYJ	Jan	-1.18	0.92	0.32	1.70	-1.23
		Apr	0.45	1.19	0.87	1.66	-1.39
		Jul	1.38	-0.29	-0.28	-0.58	0.64
		Oct	0.83	1.56	0.68	0.18	-0.45
		Jan	-0.14	-0.27	-0.98	0.15	0.33
	Velt	Apr	0.04	-0.17	0.13	0.24	-0.23
	150	Jul	0.39	-0.53	-0.78	-0.54	0.68
		Oct	0.51	0.06	-0.52	-0.56	1.04
		Jan	0.15	0.03	-0.74	0.56	0.52
	ACM2	Apr	0.47	0.07	0.23	0.39	-0.07
	ACM2	Jul	0.57	-0.40	-0.72	-0.27	0.86
T <sub>2</sub> -Night		Oct	0.86	0.35	-0.29	-0.28	1.24
		Jan	0.88	0.34	-0.45	0.80	1.03
	ы	Apr	1.21	0.48	0.53	0.64	0.23
	BL	Jul	0.79	-0.08	-0.48	-0.21	1.07
		Oct	1.37	0.54	-0.06	-0.04	1.59
		Jan	-0.46	-0.59	-1.16	-0.28	0.21
	MYJ	Apr	-0.70	-0.62	-0.27	-0.19	-0.97
		Jul	-0.15	-0.83	-0.93	-0.80	-0.08





		Oct	0.01	-0.34	-0.74	-0.83	0.52
		Jan	-0.01	-9.95	-10.71	-6.62	7.19
	VCU	Apr	-7.09	-7.59	-5.75	-5.02	3.02
	150	Jul	-8.46	1.18	-0.16	3.18	-5.56
		Oct	-11.07	-5.62	-5.17	0.57	-1.02
		Jan	-0.92	-10.77	-11.63	-7.43	6.05
		Anr	-7.82	-8.32	-6.64	-7.12	2.03
	ACM2	Iul	-9.79	-0.65	-1 71	0.51	-6.18
		Oct	-11.83	-5.93	-6.23	-1.44	-2.09
RH <sub>2</sub> -Day		Ion	-11.05	-5.55	-0.25	-1.44	-2.09
		Jaii	-1.00	-10.55	-11.30	-7.54	2.40
	BL	Apr	-7.89	-8.32	-3.90	-0.31	2.30
		Jui	-8.02	0.68	0.92	1./1	-5.06
		Oct	-12.47	-6.10	-5.67	-0.38	-2.10
		Jan	4.47	-6.72	-7.69	-4.74	12.55
	MYJ	Apr	-4.39	-4.26	-3.08	-3.59	5.59
		Jul	-5.59	3.85	3.31	4.22	-3.01
		Oct	-7.44	-3.55	-2.61	2.47	2.72
		Jan	-5.11	-7.90	-10.62	-3.29	-0.03
	VSU	Apr	-9.65	-5.01	-6.63	0.47	-0.14
	150	Jul	-5.50	1.09	-1.33	2.48	-6.95
		Oct	-12.34	-1.44	-4.83	1.62	-7.79
		Jan	-6.38	-9.54	-11.03	-4.56	-0.48
		Apr	-11.35	-5.91	-7.05	-0.86	-1.15
	ACM2	Jul	-6.86	-0.64	-2.25	0.18	-7.96
		Oct	-14.11	-2.92	-5.69	-0.94	-8.85
RH <sub>2</sub> -Night		Jan	-8.40	-10.47	-11.89	-5.01	-2.36
		Anr	-13.51	-8.00	-8.25	-2.07	-2.16
	BL	Iul	-7.32	-1.75	-2 47	-0.23	-8.06
		Oct	-15.92	-3.51	-6.19	-0.25	-0.00
		Ion	-13.92	-3.51	-0.19	-0.91	-9.94
	MYJ	Anr	2.80	-5.80	-0.91	-0.09	5.84
		лрі Iul	-5.67	-1.05	-3.44	2.61	1.01
		Jui	-1.08	5.02	0.42	3.01	-1.91
		1	-/.18	1.04	-2.30	3.30	-1.43
		Jan	1.33	1.92	1.58	2.17	0.39
	YSU	Apr	1.86	1.97	2.04	2.25	0.93
		Jul	1.35	1.56	1.20	1.79	1.30
		Oct	1.68	2.11	1.54	1.70	0.93
		Jan	1.57	2.04	1.79	2.26	0.91
	ACM2	Apr	2.11	2.01	2.18	2.37	1.21
	1101112	Jul	1.43	1.62	1.30	2.02	1.50
WSDav		Oct	1.90	2.19	1.73	2.05	1.21
WD <sub>10</sub> -Day		Jan	1.50	2.02	1.63	2.40	1.01
	DI	Apr	1.85	2.04	1.93	2.44	0.86
	BL	Jul	1.21	1.54	1.12	1.72	1.10
		Oct	1.83	2.28	1.60	1.95	1.04
		Jan	1.63	2.26	2.14	2.67	1.10
	10/1	Apr	2.33	2.40	2.65	2.69	1.61
	MYJ	Jul	1.85	2.05	1.85	2.16	2.09
		Oct	2.01	2.58	2.17	2.12	1.55
		Jan	1.26	1.43	1.50	1.40	0.88
		Apr	1.51	1.49	1.67	1.22	1.07
	YSU	Jul	1.16	1.32	1.15	1.16	1.07
		Oct	1.42	1.45	1.40	1.14	1.04
WS10-Night		Ian	1.88	1 91	1.10	2 21	1.36
		Ann	2.15	1.91	2.07	1 70	1.50
	ACM2	Thi	2.15	1.61	2.07	1./7	1.55
		Oct	1.00	1.02	1.50	1./1	1.59
		Oct	1.98	1.92	1.80	1.98	1.39





	Jan	2.32	2.32	2.13	2.74	1.63
DI	Apr	2.72	2.28	2.38	2.33	1.73
DL	Jul	1.79	2.09	1.75	2.08	1.71
	Oct	2.38	2.44	2.06	2.38	1.78
	Jan	1.63	1.76	1.94	2.18	1.45
MAZI	Apr	1.79	1.70	2.12	1.71	1.53
MYJ	Jul	1.42	1.56	1.52	1.42	1.66
	Oct	1.74	1.81	1.83	1.68	1.65

589

590 Similar to T2, 10-m wind speed (i.e., WS10) is also a diagnostic variable of the near-surface wind

591 speed. For the YSU, ACM2 and BL schemes, in the revised MM5 surface layer scheme, WS10 is

592 calculated based on the Monin-Obukhov (M-O) similarity theory (Monin and Obukhov, 1954). The

593 dimensionless profile function of momentum is denoted as:

**594** 
$$\phi_m\left(\frac{z}{L}\right) = \frac{\kappa z}{u_*} \frac{\partial u}{\partial z}$$
 (23)

595 where  $\kappa$  is the von Karman constant,  $u_*$  is the friction velocity, z is the height, L is the Obukhov

596 length, integrating the Eq. (23) with respect to height z:

**597** 
$$du = \frac{u_*}{\kappa} \left[ \frac{dz}{z} - \frac{1 - \phi_m(\frac{z}{L})}{\frac{z}{L}} d\left(\frac{z}{L}\right) \right]$$
(24)

**598** integrate Eq. (24):

**599** 
$$\int_0^u du = \frac{u}{\kappa} \left\{ \int_{z_0}^z \frac{dz}{z} - \int_{\frac{z_0}{L}}^{\frac{z}{L}} \left[ 1 - \phi_m \left( \frac{z}{L} \right) \right] d\ln \left( \frac{z}{L} \right) \right\}$$
(25)

here, let  $\psi_m\left(\frac{z}{L}\right) = \int_0^{\frac{z}{L}} \left[1 - \phi_m\left(\frac{z}{L}\right)\right] d\ln\left(\frac{z}{L}\right)$ , where  $\psi_m\left(\frac{z}{L}\right)$  is the integrated similarity function 600

601 for momentum.

602 Therefore, Eq. (25) can be indicated as 
$$u = \frac{u_*}{\kappa} \left[ \ln\left(\frac{z}{z_0}\right) - \psi_m\left(\frac{z}{L}\right) + \psi_m\left(\frac{z_0}{L}\right) \right]$$
, where  $z_0$  is the

603 roughness length.

604 Based on the bulk transfer method, the momentum flux can be represented as  $\tau = \rho u_*^2 = \rho C_m u^2$ ,

605 where  $\tau$  is the momentum flux,  $C_m$  is the bulk transfer coefficient for momentum:

**606** 
$$C_m = \frac{u_*^2}{u^2} = \frac{\kappa^2}{\left[\ln\left(\frac{z}{z_0}\right) - \psi_m\left(\frac{z}{L}\right) + \psi_m\left(\frac{z_0}{L}\right)\right]^2}$$
 (26)

607 Thus, the wind speed at 10 m divided by the wind speed at a certain height can be written as:

$$608 \qquad u_{10} = \frac{u_*}{\kappa} \left[ \ln\left(\frac{z}{z_0}\right) - \psi_m\left(\frac{z}{L}\right) + \psi_m\left(\frac{z_0}{L}\right) \right] \cdot \frac{\left[ \ln\left(\frac{10}{z_0}\right) - \psi_m\left(\frac{z_0}{L}\right) + \psi_m\left(\frac{z_0}{L}\right) \right]}{\left[ \ln\left(\frac{z}{z_0}\right) - \psi_m\left(\frac{z}{L}\right) + \psi_m\left(\frac{z_0}{L}\right) \right]} = \frac{u_*}{\kappa} \left[ \ln\left(\frac{z}{z_0}\right) - \psi_m\left(\frac{z}{L}\right) + \psi_m\left(\frac{z_0}{L}\right) \right]$$

**609** 
$$\psi_m \left(\frac{z_0}{L}\right) \cdot \left(\frac{C_m}{C_{m10}}\right)^{1/2}$$
 (27)

610 where  $C_{m10}$  is the transfer coefficient for momentum at 10 m height:

(2)





$$611 \qquad C_{m10} = \frac{\kappa^2}{\left[\ln\left(\frac{10}{z_0}\right) - \psi_m\left(\frac{10}{L}\right) + \psi_m\left(\frac{z_0}{L}\right)\right]^2}$$
(28)

612 Comparing the  $C_m$  of the three schemes (i.e., YSU, ACM2 and BL schemes) at night,  $C_m$  is the 613 largest for the BL scheme, the second largest for the ACM2 scheme, and the smallest for the YSU 614 scheme (Fig. 9). Correspondingly, the BL scheme simulates the largest WS10, ACM2 the second 615 largest, and the YSU the smallest (Fig. 8). The larger  $C_m$  corresponds to the stronger mixing, which 616 transports more momentum from the upper to the lower layers, making WS<sub>10</sub> increase. Therefore, 617 the bulk transfer coefficient  $C_m$  controls the variation of WS<sub>10</sub> at night. During the daytime, the  $C_m$ 618 of the BL scheme is smaller than that of the other two schemes, and the corresponding WS<sub>10</sub> decrease 619 (Fig. 8, 9). However, the difference among the three schemes is smaller in daytime than that in 620 nighttime. The reason why the results of  $C_m$  and  $WS_{10}$  differ with the same calculation method is 621 because of the vertical variation of heat and momentum within the boundary layer involved in the 622 calculation. This will correlate to the vertical diffusion coefficients within the boundary layer that 623 will be discussed further in a later section.





625 Figure 9. Similar as Figure 2, but for momentum transfer coefficient  $(C_m)$ .

626 For the near surface scheme of the MYJ scheme, WS<sub>10</sub> is calculated according to the near surface627 flux profile relationship proposed by Liu et al. (1979):

$$628 \qquad u_0 - u_s = D_1 \left[ 1 - exp\left( -\frac{z_u u_s}{D_1 \nu} \right) \right] \left( \frac{F_u}{u_s} \right)$$
(29)

- 629 where 0 represents the value at height z above the surface where the molecular diffusivity still plays
- 630 a dominant role, s denotes the surface value,  $D_1$  denotes a near surface parameter,  $u_*$  is the friction
- 631 velocity,  $\nu$  is the molecular diffusivity for momentum (=1 × 10<sup>-5</sup>), and  $F_u$  is the momentum flux.

**632** Since 
$$1 - exp\left(-\frac{z_u u_*}{D_1 v}\right) \approx \frac{z_u u_*}{D_1 v}, \ u_0 - u_s = \left(\frac{z_u}{v}\right) F_u$$
.





- 633 The momentum flux in the surface layer above the viscous sublayer is represented by  $F_u =$
- 634  $\left(\frac{c_m}{\Delta z_e}\right)(u_{low} u_0)$ , here, the subscript *low* denotes the variables at the lowest model level,  $\Delta z_e$  is
- 635 either the equivalent height of the lowest model level that considers the presence of the "dynamical
- **636** turbulence layer" at the bottom of the surface layer (Janjić, 1990).  $C_m$  is the bulk transfer coefficient,
- 637 defined as:

$$638 \qquad C_m = \frac{\kappa u_*}{\ln\left(\frac{z_0 + z}{z_0}\right) + \psi_m\left(\frac{z_0 + z}{L}\right) - \psi_m\left(\frac{z_0}{L}\right)} (30)$$

- 639 In Eq. (29),  $z_u$  is still an unknown, such that  $\frac{z_u u_i}{D_1 v} = \xi$ , where  $\xi$  is a smaller constant (equal to 640 0.35 in the model). Here, the near surface parameter  $D_l$  is further defined as  $D_1 = C$ .
- 641  $\left(\frac{z_0 u_*}{v}\right)^{1/4}$ , where C is a constant (=30), the roughness length  $z_0$  as a function of  $u_*$  ( $z_0 = \frac{0.11v}{u_*} + \frac{1}{v}$ )
- 642  $\frac{0.018u_*^2}{g}$ , substituting  $D_1$  and  $z_u$  to Eq. (29):

643 
$$u_0 = \frac{\frac{\xi}{u_s} \left[ C \left( \frac{z_0 u_s}{v} \right)^{1/4} \right] \left( \frac{C_m}{\Delta x_e} \right) u_{low} + u_s}{1 + \frac{\xi}{u_s} \left[ C \left( \frac{z_0 u_s}{v} \right)^{1/4} \right] \left( \frac{C_m}{\Delta x_e} \right)}$$
(31)

644 The wind speed at a height of 10 m can be expressed as:

**645** 
$$u_{10} = \frac{F_u \Delta z_e}{C_{m10}} + u_0 = \frac{C_m (u_{low} - u_0)}{C_{m10}} + u_0$$
 (32)

**646** where  $C_{m10}$  is the transfer coefficient for momentum at 10 m height:

647 
$$C_{m10} = \frac{\kappa u_*}{\ln\left(\frac{z_0+10}{z_0}\right) + \psi_m\left(\frac{z_0+10}{L}\right) - \psi_m\left(\frac{z_0}{L}\right)}$$
 (33)

648 Therefore, the  $C_m$  of the MYJ scheme is significantly different from the other three schemes. Except 649 for the NCP region, although the  $C_m$  of the MYJ scheme is larger than the other three schemes at all 650 times of the day, the  $WS_{10}$  presents the maximum only during the daytime. This suggests that at 651 night, the wind speed simulated by the MYJ scheme also be influenced by other factors. For example, 652 the calculation method of the integrated similarity functions ( $\psi_m$ ) in the MYJ scheme is different 653 from the other three schemes. In the other three schemes, the  $\psi_m$  is calculated according to four 654 stability regimes defined in terms of the bulk Richardson number(Zhang and Anthes, 1982). In the 655 MYJ scheme, the  $\psi_m$  is calculated based on two stability regimes by the z/L (Paulson, 1970).









657 Figure 10. Similar as Figure 5, but for 10-m wind speed.

658 The reasons for the differences in  $WS_{10}$  simulation are further analyzed in terms of regional 659 distribution. During the daytime, wind speed is significantly overestimated at most sites throughout 660 the NCP region, which are centered in the plains and valleys, but is less overestimated and even 661 underestimated at some sites on the mountain tops (Fig. 10, S13-S15 a1-e1). Wind speed is 662 overestimated at almost all stations throughout the YRD and PRD regions (Fig. 10, S13-S15 a2-e2, 663 a4-e4). The WS<sub>10</sub> in the basin is importantly overestimated in the SB region, while less 664 overestimated at hilltop stations on the eastern side of the basin, with higher wind speed being more 665 pronounced in January (Fig. 10, S13-S15 a3-e3). In the NS region, wind speed is overestimated to 666 a lesser extent than in other regions, but for regions with lower wind speed, the relative bias (RB) **667** is larger, especially in July (Jan: RB=29.3%~49.1%, Apr: RB=32.0%~55.6%, Jul: 668 RB=44.2%~78.7%, Oct: RB=42.7%~65.7%). Comparing the simulation results of the four schemes, 669 the MYJ scheme simulates the most significantly overestimated wind speed and the least 670 overestimated for the YSU scheme (Fig. 10, S13-S15). In comparison with the four months, it is





671	found that the RB of the simulation is the largest for the month with slower wind speed (i.e., July).
672	At night, the wind speed is overestimated at almost all stations in the whole region of NCP, and the
673	overestimation is greater at the hilltop stations than during the day (Fig. 11, S16-S18). The other
674	four regions are more similar to the daytime (Fig. 11, S16-S18). However, by comparing the four
675	schemes, we find that the BL scheme has the most obvious overestimation, different from the
676	daytime, while the YSU scheme still has the lowest overestimation, the same as the daytime. In
677	general, wind speed is smaller at night, and the four schemes overestimate wind speed much more
678	than during the day. Averaging the RB of wind speed over the five regions and four months, the
679	daytime (nighttime) values for the YSU, ACM2, BL and MYJ schemes are 77.7% (92.4%), 85.6%
680	(123.6%), 80.2% (146.0%), and 100.8% (117.4%), respectively. This simulated misestimation of
681	low winds at night may mainly originate from the inapplicability of the M-O similarity theory. The
682	strong stable boundary layer usually occurs on nights with low winds(Monahan and Abraham, 2019;
683	Vignon et al., 2017). In this strong stable boundary layer, turbulence occurs weakly and
684	intermittently, the turbulence intensity is disproportionate to the mean gradient, and the M-O
685	similarity theory is no longer applicable(Acevedo et al., 2015; Sun et al., 2012). Ultimately, these
686	inapplicable functions affect the calculation of the bulk transfer coefficient and can further lead to
687	large deviations in the simulation of wind speed.









689 Figure 11. Similar as Figure 6, but for 10-m wind speed.

690 We further re-analyze the effect of topography on wind speed. The wind speed is overestimated for 691 plains and valleys and better reproduced/underestimated for mountain tops, mainly because of the 692 smoother topography in the model. This is rather because coastal stations in the plains, many of 693 which also have high wind speeds, are not well reproduced and still show significant overestimation 694 (Fig. 10), than the high wind speeds at the top of the mountains, which are better simulated. It is 695 assumed that the wind speed should be small in plain areas with complex underlying surface, but it 696 increases after the model has smoothed the terrain. The wind speed increases gradually with height, 697 and when the terrain at the top of the mountain is smoothed, the originally larger wind speed 698 decreases, and the wind speed will be closer to the observed value.







699

700 Figure 12. Similar as Figure 2, but for 10-m wind direction.

701 The model can basically simulate the changes of wind direction in the five regions, well capturing 702 the overall wind direction in each region (Fig. 12). In the NCP region, the simulation of wind 703 direction is poor in January compared to the other three months, with a high frequency of 704 northwesterly-northerly winds, overestimated by about 6.6% (Fig. 13 a1-d1). In addition, coupled 705 with larger wind speed, it causes the effect of advective transport to be amplified, thus affecting the 706 variation of pollutant concentrations(W. Jia and Zhang, 2021). The frequency of simulated 707 northeasterly winds in the YRD region is higher than that observed in January (~6.9%), April 708  $(\sim 6.9\%)$  and October  $(\sim 11.2\%)$ , while the frequency of southerly winds is higher in July  $(\sim 10.0\%)$ 709 (Fig. 13 a2-d2). The wind direction of SB region is poorly simulated since the topography is too 710 complicated in the SB region, with a basin in the middle and high topographic mountains all around 711 (Fig. 13 a3-d3). The low wind state in the middle of the basin is difficult to be captured. The 712 percentage of northeasterly winds simulated by the model in January, April, July and October are 713 22.9%~25.6%, 19.2%~20.6%, 14.9%~16.4%, 22.4%~24.2%, respectively, and the percentage of 714 observations are 15.1%, 12.0%, 10.9%, 16.1%, respectively. Similarly, the percentage of westerly 715 winds simulated by the model in January, April, July and October are 4.7%~5.1%, 4.7%~5.7%, 716 4.8%~5.5%, 3.8%~5.1%, respectively, and the percentages of observations are 9.9%, 12.4%, 12.4%, 717 12.3%, respectively. The model simulates a large proportion of northeasterly winds and a smaller 718 proportion of westerly winds (Fig. 13 a3-d3). In the PRD region, the frequency of northeasterly 719 wind occurrences in January and October is significantly overestimated by about 8.8% and 9.5%, 720 while the frequency of southerly winds is overestimated in April and July by about 14.5% and 17.3%, 721 and the frequency of southeasterly winds is underestimated (Fig. 13 a4-d4). The wind direction is 722 better simulated in the NS region, not significantly influenced by the complex terrain (Fig. 13 a5-









724 725

726

Figure 13. Wind-rose plots in five regions for four seasons are (a1-d1) NCP region, (a2-d2) YRD region, (a3-d3) SB region, (a4-d4) PRD region and (a5-d5) NS region, respectively.

727

# 728 3.2 Vertical structures

- 729 To better understand the performance of model in simulating PBL structure under different
- 730 underlying surface, four representative stations have been selected in each region, with stations in
- 731 plain areas, stations in mountains areas with high elevation, and stations near the sea.





# 732 3.2.1 temperature

733	Accurate simulation of the vertical structure of the PBL is very important for the evolution of
734	pollution, precipitation and typhoons. In the vertical direction, four typical sounding stations are
735	selected for each region at $08:00$ to better reflect the simulation of the vertical structure of the PBL
736	under different underlying surface conditions. Overall, the model captures the vertical structures of
737	the temperature. From the simulation results in January, the best reproduction of the temperature
738	simulation is found in the YRD region, in which the temperature is closer to the observed values
739	(Fig. 14 a2-d2). In addition to the NCP, SB and NS regions, a temperature inversion layer appears
740	in the lower layers at 08:00, and the NS region has the most significant temperature inversion (Fig.
741	14 a1-d1, a3-d3, a5-d5). The model does not simulate the temperature variation of the inversion
742	layer well, and shows significant differences from the observations. When there is a difference in
743	topography between the observed and simulated stations, the bias in the temperature is more
744	pronounced. These stations usually exist in complex topographic conditions, such as Zhangjiakou
745	and Zhangqiu stations in the NCP region, Shapingba in the SB region, and Mazongshan and Jiuquan
746	in the NS region (Fig. 14 b1, d1, c3, a5, b5). The topographic discrepancy caused by the lack of
747	high resolution may, on the one hand, account for it, resulting in more complex topography in the
748	grid points closest to the observation stations. On the other hand, there is also an urgent need for
749	finer underlying surface data to respond more closely to the observed real topography. The effect of
750	resolution and underlying surface will be discussed in detail in the Part II. Although the elevation
751	of the Hongyuan station in the SB region is higher, the difference in topographic height obtained
752	from observations and simulations are close to each other (Fig. 14 a3).









Figure 14. Average vertical profiles of observed and simulated temperature at 08:00 and 20:00 BJT
at four sounding stations for each region in January (Winter). The unobtrusive gray lines indicate
the simulated lines for all time periods, and the lines with shading indicate the average values and
shaded areas show the uncertainty range (the mean ±1 standard deviation).

758 There is an underestimation of temperature at stations with higher topography and overestimation 759 for lower topography, which is more consistent with the conclusions drawn from 2-m temperature 760 (Fig. 5, 14). However, the underestimation of temperature is not present throughout the vertical, but 761 is more pronounced in the lower layers, which are more influenced by the underlying surface. From 762 the differences of the four schemes, the MYJ scheme simulates the lowest temperature and largest 763 temperature gradient. Since the MYJ scheme simulates a weak turbulent diffusion of heat, a well 764 vertical exchange process cannot occur, bringing into a large temperature gradient (Fig. S19). The





765 differences of the four schemes gradually decrease with the increase of the height. The BL scheme, 766 which is also a local closure scheme, with a smaller vertical gradient in temperature, mainly because 767 this scheme adds a counter-gradient correction term to the heat flux, which is mainly applicable to 768 the convective PBL(Bougeault and Lacarrere, 1989). The presence of this term leads to an increase 769 in turbulent diffusion and a decrease in temperature gradient. However, it is worth noting that there 770 are still slightly stable stratifications at 08:00, and this term generates upward heat flux and reduces 771 the temperature gradient, which is closer to the results of the nonlocal closure schemes (YSU and 772 ACM2 schemes). The simulation results for the other three months are not as good as January, but 773 the simulation characteristics are similar to January (Figures not shown). The results at 20:00 are 774 similar to those at 08:00, and thus will not be repeated here (Figures not shown).

### 775 3.2.2 wind speed and direction

776 The simulation of wind speed vertical structure is much worse in comparison to temperature (Fig. 777 14, 15). The simulated results of wind speed in the vertical direction and 10-m wind speed are still 778 quite different. In the four months, the wind speed is almost overestimated at the lower altitude 779 stations below 1000 m in all the four regions except the NS region, and wind speed is less 780 overestimated in July than in the other three months (Fig. 15, S20-S22 a1-d4). However, for the NS 781 region, the wind speed is almost better simulated, or underestimated, and is significantly different 782 from the other four regions (Fig. 15, S20-S22 a5-d6). We can compare the Zhangjiakou station in 783 the NCP region with the Hongyuan station in the SB region, and find that the wind speeds at these 784 stations are almost not overestimated (Fig. 15 b1, a3). The effect of the model on terrain smoothing 785 contributes to it. Because the wind speed itself increases with the increase of height, and it decreases 786 when the model smooths over the terrain.

787 Unlike the 10-m wind speed, the simulation results of the 10-m wind speed have the smallest bias 788 for the YSU scheme, which is closer to the observed value (Table 2, Figs. 8, 10-11). Of course, this 789 phenomenon can also be found from the evolution of the wind speed in the vertical direction (Fig. 790 15, S20-S22). However, as the height increases, the bias of the YSU scheme gradually increases and 791 is greater than the other three schemes (Fig. 15, S20-S22). Such a large vertical gradient of wind 792 speed in the YSU scheme indicates a weak mixing in this scheme. From the turbulent diffusion 793 coefficients of the momentum at 08:00 in January, it is true that the YSU scheme simulates the 794 smallest turbulent diffusion coefficient below 1000 m (Fig. S23). While the BL scheme simulates a 795 smallest vertical gradient of wind speed, which corresponds to the largest turbulent diffusion 796 coefficient of momentum (Fig. S23). The time variation characteristics of the turbulent diffusion 797 coefficient will be deliberated and analyzed in detail later. 798 The simulation results of wind direction notes that the model can capture the characteristics of wind 799 direction well, and it also simulates well for the stations with more complicated topography and







**800** higher altitude (Figs. 15, S20-S22).



Figure 15. Similar as Figure 14, but for 10-m wind speed and direction.

803

# 804 3.3 PBLH

805 Based on the observed data, comparing the PBLH calculated by the two methods, it is found that
806 the results are mixed for two methods (Fig. 16). The results for January (IOA=0.70~0.90) are better
807 than the other three months (IOA=0.44~0.88 in April, IOA=0.51~0.86 in July, IOA=0.60~0.77 in
808 October), and the results in the NCP region are better than the other four regions (Fig. 16). The
809 PBLH in the NS region are more scattered, unlike the other regions where most of the PBLH are
810 concentrated below 500 m, especially in April and July (Fig. 16 b5-c5). On the whole, the difference





- Oct lan Apr Jul 150 IOA=0.88 MB=21.7 BB=14% **b1**) IOA=0.88 MB=-8.0 BB=-3% IOA=0.7 MB=-5.3 BB=-4% IOA=0.86 MB=66.2 BB=32% NCP(D03) PBLH<sub>OBS - R</sub> (m) Ē Ê (m) PBLH<sub>OBS</sub>. PBLHOBS PBLHOB 500 1000 PBLH<sub>OBS-PTv</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) IOA=0.9 MB=6.9 RB=2% IOA=0.79 MB=-10.3 RB=-5% IOA=0.68 MB=-30.2 RB=-11% IOA=0.77 MB=-32.3 RB=-11% 1×10<sup>-6</sup> Ē Ē Ē Ξ YRD(D04) PBLH<sub>OBS</sub> - RI (m PBLH<sub>OBS</sub>. PBLHOUS PBLHOS 50 50 50 PBLHOAS \_ PTv (m) PBLH \_*PTv* (m) PBLH<sub>OBS</sub> - PTv (m) PBLH \_ PTv (m) 3x10-IOA=0.68 MB=9.1 RB=7% IOA=0.80 MB=52.2 RB=62% b3 IOA=0.84 MB=49.7 RB=23% IOA=0.60 MB=43.2 RB=47% Ē Ê Ē Ē 100 100 Frequency 100 100 SB(D05) PPBLH<sub>085 -RI</sub> ( PBLHOBS PBLHo PBLHo 50 50 1000 1000 PBLH<sub>OBS-PTV</sub> (m) PBLH<sub>OBS - PTv</sub> (m) PBLH<sub>OBS - PTv</sub> (m) PBLHORS - PTV (m) IOA=0.70 MB=-4.5 RB=10% IOA=0.67 MB=26.6 RB=-2% IOA=0.51 MB=-11.1 RB=-4% IOA=0.65 MB=-19.8 BB=-6% PBLH<sub>OBS</sub> - RI (m) Ê í, **i** PBLHOBS PBLHOBS PBLHOR 10-6 500 1000 PBLH<sub>OBS</sub> - PT<sub>V</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) 500 1000 PBLH<sub>OBS - PTv</sub> (m) IOA=0.82 MB=5.8 RB=11% IOA=0.44 MB=228.7 RB=63% (h5 IOA=0.75 MB=-33.6 RB=-9% IOA=0.73 MB=-20.3 RB=-18% NS(D07) BLH<sub>085 - 8/</sub> (m) Ê 100 £ 1000 Ē 100 100 PBLH<sub>OBS</sub> -PBLH<sub>O85</sub> PBLH<sub>O8:</sub> 500 50 1000 PTV (**m**) 1000 - PTV (**m**) 1000 ETV (m) 1500 1500
- 811 in PBLH calculated by the two methods is more obvious in the NS region with more complex





PBLHOB

814 Figure 16. Density scatterplots of the PBLH by the two methods in five regions for four seasons. 815 Further, the mechanism of understanding the PBLH differences based on different stations from 816 different regions. The planetary boundary layer height (PBLH) for the YSU and ACM2 schemes is 817 calculated by the Richardson number method, and the PBLH for the BL scheme is obtained by the 818 virtual potential temperature method (see Section 2 for details). The same two methods are also used 819 to calculate the PBLH with sounding data for comparison (Eqs. 29 and 30). First, we compare the 820 PBLH calculated based on the observed data using Richardson number (Ri) with the PBLH 821 simulated by YSU and ACM2 schemes. The Ri is determined by both buoyancy term and the shear 822 term together (Eq. 29). The difference between simulated and observed temperature gradient is 823 smaller than the wind speed gradient within the PBL (Figs. 14-15). Therefore, the difference in Ri

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824	mainly comes from the variation in the shear term. The wind speed gradient simulated in both
825	schemes are greater than the observed values (Fig. 15), except for individual stations, which would
826	result in small values of Ri. Thus, the height of Ri up to $0.25$ would be high and the PBLH would
827	be high. Consequently, the PBLH simulated by the YSU and ACM2 schemes are higher than the
828	observed values at most stations. For example, in the case of the Quzhou station in the YRD region,
829	the simulated wind speed gradient at this station is much smaller than the observed value in January,
830	thus, the simulated PBLH is correspondingly smaller than PBLH calculated from observations (Figs.
831	15, S24 d2). Comparing the results of the other three months, we can also find similar conclusions
832	(Figures not shown). The wind speed gradient simulated by the YSU scheme is larger than that of
833	the ACM2 scheme, and therefore the PBLH is larger than that of the ACM2 scheme, except for the
834	Shanghai station in the YRD region, Shantou and Yangjiang stations in the PRD region (Fig. S24).
835	For the ocean, the PBLH simulated by the ACM2 scheme is higher than that of the YSU scheme,
836	while for most areas adjacent to the ocean, the PBLH simulated by the ACM2 scheme is on the high
837	side in the YRD and PRD regions (Fig. S25). The simulated PBLH of the BL scheme is in better
838	agreement with the PBLH calculated by the virtual potential temperature method, which is
839	substantially better than the other three schemes. The PBLH simulated by the MYJ scheme is mixed
840	(Fig. 17).











## 843 regions in January (Winter).

844 From the differences of regional distributions, the region with the best PBLH simulation results is 845 the YRD region in January, (IOA=0.64~0.74; MB=-6.4~40.5 m; RB=-2%~16%), followed by the 846 PRD region (IOA=0.53~0.72; MB=-35.4~1.4 m; RB=-12%~0.5%), and the worst simulation results 847 for the PBLH in the NS region (IOA=0.35~0.53; MB=24.9~101.5 m; RB=44%~200%) (Fig. 17). 848 Also, the PBLH simulation in the SB region is poorer and slightly better than that in the NS region, 849 noting that there is still much potential for the model to improve the reproduction of the PBLH in 850 complex terrain. From the simulation results of the four schemes, the PBLH simulated by the BL 851 scheme is the closest to the observed value, followed by the YSU and ACM2 schemes, and the 852 PBLH simulated by the MYJ scheme is the worst in that the simulation results of temperature are 853 better compared with other meteorological factors, so the method of judging the PBLH using the 854 virtual potential temperature will be more consistent with the observed values. While as the YSU





855 and ACM2 schemes using Richardson number method will involve the wind speed gradient, and 856 the vertical gradient of wind speed is poorly simulated below 1000 m. That's why it will affect the 857 judgment of the PBLH. If the simulation results of vertical gradient of wind speed can be improved 858 subsequently, then the simulation results of PBLH of these two schemes will be improved to some 859 extent. There are not enough observations to calculate the PBLH using TKE, so there will be some 860 differences with the PBLH simulated by the MYJ scheme. The mean bias of the simulation increased 861 in April and July when the PBLH is higher compared to January, with mean bias of -29.6~361.8 m 862 (6.5~603.9 m), -12.6~410.6 m (41.6~603.2 m), -34.1~301.1 m (3.2~683.9 m) and -14.5~96.3 m (-863 11.3~523.6 m) for the YSU, ACM2, BL and MYJ schemes in April (July), respectively. Similar to 864 January, the best simulation results have been obtained for the YRD (MB=7.8~72.4 m in April, 865 MB=28.5~66.5 m in July) and PRD (MB=-34.1~-12.6 m in April, MB=-11.3~54.8 m in July) 866 regions, and the worst for the NS region (MB=61.8~410.6 m in April, MB=523.6~683.9 m in July). 867 The results for October are more similar to those for January, with lower PBLH and better 868 simulations than those for April and July (Figures not shown).

#### 869 3.4 turbulent diffusion coefficient

870 Since the model itself does not directly output turbulent diffusion coefficients for all schemes, there 871 is relatively little direct comparison and analysis of this parameter. As seen in section 3.3.2 above, 872 the turbulent diffusion coefficient (TDC) plays a crucial role in the momentum vertical transport 873 within the PBL and also has an impact on the diffusion of other parameters, such as heat, water 874 vapor and pollutants (Ding et al., 2021; R. B. Stull, 1988). The accurate portrayal of the TDC directly 875 affects the evolution of the PBL structure. Based on the contents of section 2.3, the momentum TDC 876 is not equal to the heat TDC under unstable and neutral conditions for the YSU scheme, and the 877 momentum TDC is equal to the heat TDC under stable conditions. The ACM2 scheme uses the 878 MOST method to calculate the TDC as the YSU scheme, but also considers the TDC calculated by 879 the mixing length theory. The momentum TDC is not equal to the heat TDC in the MYJ scheme, 880 while the momentum TDC is equal to the heat TDC in the BL scheme. Because of the difference in 881 altitude of different stations, Bejing station in the NCP region, Nanjing station in the YRD region, 882 Daxian station in the SB region, Qingyuan station in the PRD region and Zhangye station in the NS 883 region were selected as representative stations to analyze the turbulent diffusion characteristics. 884 Here, the TDC of heat is taken as an example, the following basic characteristics have been found. 885 (1) the YSU and MYJ schemes have the largest TDC during the day, followed by the BL scheme, 886 and the ACM2 scheme has the smallest TDC (Fig. 18). (2) The TDC is largest in April and July, and 887 smallest in January and October (Fig. 18). (3) There are significant seasonal differences in the PBLH 888 for the NCP, SB and NS regions, while for the YRD and PRD regions. The difference in the PBLH 889 affects the variation of the turbulent diffusion, especially for the YSU and ACM2 schemes, where





890 the PBLH is used during the calculation of the turbulent diffusion. In the YSU scheme, the TDC of 891 momentum is calculated first, and then the TDC of heat is calculated with the Prandtl number (Pr). 892 Thus, the variation of the PBLH is proportional to the TDC (Fig. 18a). While in the ACM2 scheme, 893 the TDC of heat is calculated directly based on the dimensionless function of heat. Moreover, the 894 Pr in the YSU scheme varies with height, while the Pr in the ACM2 scheme is a constant (=0.8). It 895 is also worth noting that in the ACM2 scheme, another TDC is calculated using the mixing length 896 theory, and the change of the empirical stability function in the mixing length method changes the 897 TDC. Therefore, the YSU scheme calculates a large TDC of momentum, which also leads to a large 898 TDC of heat. The TDC of heat in the ACM2 scheme, on the other hand, will be affected by the 899 mixing length method, and differs from the calculation principle of the YSU scheme.





901 Figure 18. Time-height cross sections of heat turbulent diffusion coefficient simulated by (a) YSU
902 scheme, (b) ACM2 scheme, (c) BL scheme and (d) MYJ scheme for four seasons in the NCP region.
903 The gray line indicates the PBLH.

904 From the section 3.2.2 above, it is clear that the BL scheme has the strongest turbulent diffusion at
905 08:00, making the vertical gradient smaller, especially the wind speed is large. Similarly, this
906 phenomenon can be found in the daily variation of turbulent diffusion, and not only at 08:00, but
907 almost throughout the night (Fig. 18). The difference between MYJ and BL schemes is mainly
908 reflected in the calculation principle of mixing length, which is not directly related to the PBLH. In





909	the BL scheme, mixing length scale can be relative to the distance that a parcel originating from this
910	layer, can travel upward and downward before being stopped by buoyancy effects (Eq. 14).
911	Therefore, the vertical height below the temperature inversion layer at night, with the surface as the
912	lower boundary, is the length scale of turbulence, i.e., mixing length scale. While in the MYJ scheme,
913	the mixing length scale is equal to $z$ minus the integral depth scale, which is equal to the height of
914	the equal-area rectangle under the profile. It is worth noting that the mixing length scale in the BL
915	scheme mainly considers the effect of thermal and takes temperature gradient as the criterion, while
916	the turbulent length scale in the MYJ scheme is mainly determined based on the TKE. TKE is further
917	divided into horizontal TKE and vertical TKE. Horizontal TKE is mainly influenced by wind shear
918	and the turbulent eddy scale can reach 1.5~3 times the PBLH on the horizontal, and even reach 6
919	times the PBLH(Atkinson and Zhang, 1996). The vertical depth of an unstable layer capped by an
920	inversion is automatically selected as the length scale for turbulence in the BL scheme during the
921	daytime. Moreover, in the BL scheme, there is a counter-gradient correction term in the convective
922	PBL, which leads to a downward transport of dry and cool air, making the thermal reach a lower
923	height and a smaller length scale for turbulence. We also find that the PBLH of the MYJ scheme
924	exhibits a "sawtooth", and is more pronounced at night. This is mainly because the turbulence is
925	weaker at night, and presents intermittent characteristics, which, together with the judgment method
926	of PBLH and the coarse vertical resolution, can cause such variation of the PBLH. Although the
927	improvement of PBLH in the MYJ scheme cannot have a substantial effect on turbulent diffusion,
928	the threshold value of its determination method is open to question.

## 929 3.5 Discussion of optimal PBL schemes

930 To better understand the simulation performance of different PBL parameterization schemes for

931 different parameters in each region, this section will discuss the expressiveness of different PBL

932 schemes through the statistical approach. Figure 19 shows the Taylor statistics for the analysis of

933 near-surface meteorological parameters in four months in five regions.











943	For the NCP region, the 2-m temperatures are underestimated during the daytime in January (Fig.
944	2), while the BL scheme simulates the highest temperature, so the BL scheme performs optimally.
945	Although the temperatures are somewhat overestimated at night, the overestimated period is shorter.
946	The IOA of the four schemes is similar, with the ACM2 scheme having a slightly smaller bias (Fig.
947	19 a1). Combined with the regional distribution of all stations in the NCP region, the BL scheme is
948	recommended if the study area is mainly for Beijing, while the ACM2 and YSU schemes are
949	recommended for the south of the NCP in January, such as Shandong Peninsula and southern Hebei
950	province. For the other three months, temperatures are overestimated to varying degrees, both
951	during the day and at night (Fig. 2). The MYJ scheme performs best in all statistical parameters at
952	night (Fig. 19 b1-d1), while during the daytime, it slightly underperforms the YSU and ACM2
953	schemes in relative bias in January and April, but the difference is not very distinct. Therefore, for
954	the simulation of 2-m temperature in other three months, the MYJ scheme would be more
955	recommended. In the YRD region, the 2-m temperatures are overestimated during the daytime, the
956	BL schemes show overestimation at night, the MYJ scheme show underestimation, and the $YSU$
957	and ACM2 schemes perform optimally (Fig. 2 a 2-d2). According to the Taylor statistical parameters,
958	it can be seen that the ACM2 scheme performs better than the YSU scheme in the four months, and
959	based on that, the ACM2 scheme is recommended (Fig. 19 a2-d2). The 2-m temperature in the SB $$
960	region during the daytime is the same as in the YRD region, and the ACM2 scheme performs
961	$optimally (Fig. \ 19 \ a3-d3). \ However, the BL scheme performs optimally during the nighttime, except$
962	in April (Fig. 19 b3). The PRD region differs from the other regions in that the temperature
963	simulation is significantly higher in January and April, and the MYJ scheme performs best in both
964	daytime and nighttime (Fig. 19 a4-b4). In contrast, the temperature simulation bias less in July and
965	October, and the BL scheme performs best (Fig. 19 c4-d4). The 2-m temperature are almost
966	underestimated during the daytime and overestimated for the nighttime in the NS region, and the
967	MYJ scheme outperforms other schemes on account of its large diurnal temperature range (Fig. 19
968	a 5-d5). Of course, the BL scheme presents a slight advantage in the relative bias during the day time
969	in January (Fig. 19 a5).
970	The results of 2-m relative humidity are relatively uniform, and the MYJ scheme shows optimal
971	simulation performance in almost all months in all regions. Except for July in the YRD region, July
972	and October in the PRD region, and January and April in the NS region (Fig. 19 c2, c4-d4, a5-b5).
973	For the simulation of 10-m wind speed and direction, the YSU scheme shows a very clear advantage,
974	which is outstanding in all regions and all months (Fig. 19).
975	Several sounding stations with large differences between the observed and simulated altitudes are
976	removed. Then, the stations in each region are averaged to induce the variation characteristics from

**977** 100 m to 2000 m in vertical.







978

979 Figure 20. Statistics of temperature and wind speed in different layers at vertical height in January
980 (Winter), with circle size indicating the mean bias between simulations and observations, and
981 circles filled with color denoting the relative bias.

982 The BL scheme has the smallest simulation bias for the temperature in the vertical direction in 983 January for the NCP region, and performs optimally, which is associated with the discussion of the 984 optimal scheme for the 2-m temperature (Fig. 20 a1) to some degree. While in all the other three 985 months, the MYJ scheme has the smallest bias and is consistent with the conclusion of the 2-m 986 temperature (Fig. S26-S28 a1). In the YRD region, there is no clear difference between the four 987 schemes for the simulation of temperature in the vertical direction, and the deviation of the 988 simulation in January is less than 0.1 K (Fig. 20 b1). The optimal scheme for 2-m temperature can 989 be considered as a representative choice. The MYJ scheme has a better simulation of the vertical 990 profile of the temperature that is somewhat different from the most preferred scheme of 2-m 991 temperature in the SB region (Fig. 20, S26-S28 c1). This is mainly because the selected sounding 992 stations are basically located in the basin area with low elevation, and the temperatures are 993 overestimated, as is the 2-m temperature (Fig. 5). If the stations around the basin are not considered, 994 the simulation of 2-m temperature will also be overestimated and the MYJ scheme also perform 995 optimally. There is no complex topography in the PRD region, the results from the sounding stations 996 and surface layer can be well echoed. In the vertical direction, the MYJ simulates the vertical profile





997	of temperature better, particularly in January and April (Fig. 20, S30, d1). While in July and October,
998	the simulation results of the four schemes have little difference, especially in the lower level, which
999	can be represented by the optimal scheme of 2-m temperature (Fig. S27-S28, d1). For the NS region,
1000	there is an overestimation in the lower levels and an underestimation in the upper levels for each
1001	month. The height of overestimation is lower in January and April, around 200 m, while it can reach
1002	around 500 m in July and October (Fig. 20, S26-S28, e1). The positive deviation decreases as the
1003	height increases, but after reaching a certain height, the negative deviation increases again. In the
1004	process of decreasing the positive deviation with height, the MYJ scheme performs the best, which
1005	is consistent with the 2-m temperature. While the BL scheme performs slightly better when the
1006	negative deviation gradually increases with height.
1007	For wind speed, the YSU scheme is optimal for 10-m wind speed for all regions and months.
1008	However, the quite different in terms of the variation of vertical wind speed. Throughout all regions
1009	and months, the simulation bias of the BL scheme is the smallest and closer to the observation,
1010	which is also the results obtained from the section 3.2.2 (Fig. 20, S26-S28 a2-e2). The stronger
1011	turbulent diffusion of the BL scheme at 08:00 makes the wind speed more uniformly mixed in the
1012	vertical direction. And the vertical variation characteristics at 08:00 can be extended to the whole
1013	night. But during the daytime, the result may not be the same, after all, the vertical mixing of the
1014	YSU scheme is stronger, and does not produce such a large wind speed gradient.
1015	In general, in the selection process of the PBL scheme, if the focus is on temperature variation, such
1016	as temperature inversion, the optimal scheme for both 2-m temperature and vertical temperature can
1017	be considered, and there is basically no significant difference. However, if the focus is on the
1018	variation of wind speed, wind energy, then the vertical wind speed and in the surface layer need to
1010	

**1019** be evaluated and selected with comprehensive consideration.

# 1020 4 Conclusion

1021 The planetary boundary layer serves as a bridge between the ground and the free atmosphere, and 1022 its role cannot be ignored. Turbulence, as the primary motion within the PBL, controls the vertical 1023 mixing of heat, water vapor, momentum, and pollutants. Turbulence as a sub-grid-scale motion is 1024 usually parameterized in the model, i.e., PBL parameterization scheme. The most widely used 1025 mesoscale model (i.e., WRF), which has developed 12 schemes, includes nonlocal closure scheme, 1026 local closure scheme and hybrid nonlocal-local closure scheme. Across the world, there have been 1027 many evaluation studies for PBL parameterization schemes (reference Fig. 1 in Jia and Zhang, 2020). 1028 However, most of the studies have been conducted for individual stations in a small region with 1029 special individual cases for research and analysis, which are not well represented and applied. 1030 Meanwhile, there is a deficiency in understanding the mechanism of the scheme itself. In response, 1031 aiming at the current research deficiencies, four typical schemes (YSU, ACM2, BL and MYJ,





1032	covering each type scheme) are selected in this study to evaluate and analyze the near-surface
1033	meteorological parameters, vertical structure of the PBL, PBLH and turbulence diffusion in four
1034	months (i.e., January, April, July and October) in five typical regions of China (i.e., NCP, YRD, SB,
1035	PRD and NS regions).
1036	a. 2-m temperature. (1) In terms of time series and diurnal, the simulation results for July are better
1037	than the other three months, and better at night than daytime, with less deviation between simulation
1038	and observation. (2) in terms of regional distribution, temperatures at stations with higher elevations
1039	are easily underestimated (e.g., mountainous areas in the NCP region, areas around the SB basin,
1040	and the NS region), while overestimated at plains and basin (e.g., YRD, PRD and the SB basin
1041	regions), and the overestimation/underestimation is more significant during the daytime. (3) In
1042	terms of mechanism differences between schemes, the differences in the simulated temperatures of
1043	the four schemes are more pronounced at night. The differences in simulated temperatures between
1044	the nonlocal scheme mainly originate from downward shortwave radiation, while the effects of
1045	sensible heat flux (HFX) need to be further ruminated when comparing with the local closure
1046	scheme. when analyzing the HFX, the gradient of 2-m temperature and surface temperature, the
1047	variation of heat transfer coefficient need to be discussed in detail.
1048	b. 2-m relative humidity. The changes in relative humidity and temperature correspond to each other,
1049	and again the best simulation results are obtained in July. Except for the NS region, the relative
1050	humidity of the other regions is underestimated.
1051	c. 10-m wind speed. (1) The simulation bias is the largest for the MYJ scheme during the daytime
1052	(except for the NCP region), and the BL scheme presents the largest deviation at night in all regions,
1053	and the difference is not significant in the four months. The variation of 10-m wind speed is
1054	influenced by the momentum transfer coefficient, where a larger $C_m$ produces stronger mixing and
1055	transports more momentum from the upper layers to the lower layers. For the YSU, ACM2 and BL $$
1056	schemes, the $C_m$ and 10-m wind speed vary proportionally. In contrast, the MYJ scheme calculate
1057	principle of for MYJ scheme is different from the other schemes, and the $C_m$ is larger than other
1058	months almost all day. However, the wind speed simulated by the MYJ scheme is maximum only
1059	during the daytime, which indicates that it is influenced by integrated similarity functions. (2) In
1060	terms of regional distribution, the wind speed is more overestimated in plains and basins, and less
1061	overestimated or even underestimated in mountainous areas. This is chiefly due to the influence of
1062	the model on terrain smoothing. (3) The overestimation of smaller wind speed at night is more
1063	obvious in the four schemes, primarily owing to the non-application of the MOST. At night, the
1064	turbulence intensity is disproportionate to the mean gradient, and the M-O similarity theory is no
1065	longer applicable.
1066	<i>d. 10-m wind direction.</i> The simulation of wind direction in January for the NCP region worse than
1067	the other three months, and the frequency of simulated northwest-north winds is overestimated by





1068 about 6.6%. For the YRD region, the frequency of northeasterly winds is overestimated. The 1069 simulation of wind direction in the SB region is not as good as other regions due to the complex 1070 topography. The frequency of northeasterly winds is overestimated in January and October in the 1071 PRD region, and that of southerly winds is overestimated in April and July. The wind direction is 1072 better simulated for the NS region, and the difference is not very obvious. 1073 e. vertical distribution of PBL. The model can reproduce the vertical structure of temperature well, 1074 but the inversion temperature at the lower levels of many stations in complex terrain cannot be 1075 simulated well, mainly because there is a certain difference in the terrain height between observation 1076 and simulation. At 08:00, the MYJ scheme simulates the lowest temperature and the BL scheme for 1077 the highest temperature, and the difference is more conspicuous at the lower levels. The vertical 1078 structure of the wind speed is clearly not as good as the temperature. The wind speed is almost 1079 always overestimated below 1000 m, except for the NS region. Unlike the 10-m wind speed, YSU 1080 has the smallest deviation from the 10-m wind speed, while the BL scheme has the smallest bias in 1081 the vertical direction. The BL scheme has the largest turbulent diffusion and the strongest mixing at 1082 08:00. 1083 f. PBLH. The PBLH calculated based on the observed data using the two methods are better in

1084 January than in the other three months, and in the NCP region than in the other four regions. The 1085 wind speed gradient simulated by the YSU scheme is large, resulting in a small Richardson number 1086 (Ri), making the height higher when Ri reaches 0.25, and the PBLH is higher than that of the ACM2 1087 scheme. The PBLH simulated by the BL scheme is closer to the observation because the temperature 1088 gradient is best simulated. The MYJ scheme results in a jagged variation of the PBLH due to the 1089 determination of the threshold and the vertical resolution, and this phenomenon is especially 1090 obvious at night. In terms of regional distribution, the PBLH is best simulated in the YRD region, 1091 followed by the PRD region and worst in the NS region. The results are similar in January and 1092 October, when the PBLH is lower and the simulations are better than those in April and July.

1093 g. turbulent diffusion coefficient. (1) The TDC simulated by the YSU and MYJ schemes is the 1094 largest during the daytime, followed by the BL scheme, and the smallest by the ACM2 scheme. The 1095 TDC simulated by the BL scheme is the largest at night, and the other three schemes are about the 1096 same. (2) The TDC is maximum in April and July, and minimum in January and October. (3) The 1097 obvious difference in PBLH affects the turbulent diffusion of the YSU and ACM2 schemes. It is 1098 worth noting that the YSU scheme calculates the TDC of momentum first, and then uses Prandtl 1099 number (Pr) to calculate the TDC of heat, while the ACM2 scheme calculates the TDC of both 1100 momentum and heat. (4) The difference between the BL and MYJ schemes is mainly reflected in 1101 the calculation principle of mixing length. The buoyancy effect mainly affects the mixing length 1102 scale in the BL scheme, and the mixing length scale of MYJ scheme is influenced by the TKE. 1103 For the discussion of the optimal scheme, different schemes need to be proposed for different





1104 parameters. (1) Temperature. The BL scheme is recommended for January in the NCP region, 1105 especially for the Beijing, and the MYJ scheme is recommended for the other three months. The 1106 simulation difference between the four schemes is small in the YRD region, and the ACM2 scheme 1107 is recommended. The topography is more complex in the SB region, but the MYJ scheme is 1108 recommended for most areas within the basin, and the BL scheme is recommended for the SB region 1109 if more around basin is involved. The MYJ scheme is recommended for the PRD region in January 1110 and April, and the BL scheme is recommended for July and October. In the NS region, the MYJ 1111 scheme is recommended. (2) Relative humidity. The MYJ scheme is recommended for all regions 1112 in four months. (3) Wind speed. The YSU scheme is recommended if the main concern is the surface 1113 layer, and the BL scheme is recommended if the focus on the variation of wind speed in the vertical 1114 direction.

1115

# 1116 Code and data availability

1117 The source code of WRF version 3.9.1 can be found on the following website:
1118 https://www2.mmm.ucar.edu/wrf/users/download/, and the model settings file is named
1119 "3.9.1\_namelist.input", which can be found in the Supplement. In addition, the hourly
1120 meteorological observation data and L-band radiosonde observation data provided by the Chinese
1121 Academy of Meteorological Sciences, are available at https://doi.org/10.5281/zenodo.7792241 (Jia
1122 et al., 2023).

#### 1123 Author contributions

1124 Development of the ideas and concepts behind this work was performed by all the authors. Model
1125 execution, data analysis and paper preparation were performed by WJ. XZ and HW provide
1126 computing resources, and offer advice and feedback. YW, DW, and JZ support the data. WZ, LZ,
1127 LG, YL, JW, YY, and YL provides suggestions. All authors contributed to the manuscript.

#### **1128** Competing interests

1129 The authors declare that they have no conflict of interest.

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1135	
1136	Reference
1137	Acevedo, O. C., Mahrt, L., Puhales, F. S., Costa, F. D., Medeiros, L. E., and Degrazia, G. A.
1138	Contrasting structures between the decoupled and coupled states of the stable boundary layer.
1139	Quarterly Journal of the Royal Meteorological Society, 142(695), 693-702. doi:10.1002/qj.2693,
1140	2015.
1141	Atkinson, B. W., and Zhang, W. J. Mesoscale shallow convection in the atmosphere. Reviews of
1142	Geophysics, 34(4), 403-431. doi:10.1029/96rg02623, 1996.
1143	Avolio, E., Federico, S., Miglietta, M. M., Lo Feudo, T., Calidonna, C. R., and Sempreviva, A. M.
1144	Sensitivity analysis of WRF model PBL schemes in simulating boundary-layer variables in southern
1145	Italy: An experimental campaign. Atmospheric Research, 192, 58-71.
1146	doi:10.1016/j.atmosres.2017.04.003, 2017.
1147	Bauer, P., Thorpe, A., and Brunet, G. The quiet revolution of numerical weather prediction. Nature,
1148	525(7567), 47-55. doi:10.1038/nature14956, 2015.
1149	Blackadar, A. K. The vertical distribution of wind and turbulent exchange in a neutral atmosphere.
1150	Journal of Geophysical Research (1896-1977), 67(8), 3095-3102. doi:10.1029/JZ067i008p03095,
1151	1962.
1152	Bougeault, P., and Lacarrere, P. Parameterization of Orography-Induced Turbulence in a Mesobeta-
1153	-Scale Model. Monthly Weather Review, 117(8), 1872-1890. doi:10.1175/1520-
1154	0493(1989)117<1872:Pooiti>2.0.Co;2, 1989.
1155	Broxton, P. D., Zeng, X., Sulla-Menashe, D., and Troch, P. A. A Global Land Cover Climatology
1156	Using MODIS Data. Journal of Applied Meteorology and Climatology, 53(6), 1593-1605.
1157	doi:10.1175/jamc-d-13-0270.1, 2014.
1158	Chen, F., and Dudhia, J. Coupling an Advanced Land Surface-Hydrology Model with the Penn
1159	State-NCAR MM5 Modeling System. Part I: Model Implementation and Sensitivity. Monthly
1160	Weather Review, 129(4), 569-585. doi:10.1175/1520-0493(2001)129<0569:Caalsh>2.0.Co;2, 2001.
1161	Cohen, A. E., Cavallo, S. M., Coniglio, M. C., and Brooks, H. E. A Review of Planetary Boundary
1162	$Layer \ Parameterization \ Schemes \ and \ Their \ Sensitivity \ in \ Simulating \ Southeastern \ U.S. \ Cold \ Season$
1163	Severe Weather Environments. Weather and Forecasting, 30(3), 591-612. doi:10.1175/waf-d-14-
1164	00105.1, 2015.

- 1165 Deardorff, J. W. Stratocumulus-capped mixed layers derived from a three-dimensional model.
- **1166** Boundary-Layer Meteorology, 18(4), 495-527. doi:10.1007/BF00119502, 1980.





1167	Diaz, L. R., Santos, D. C., Käfer, P. S., Iglesias, M. L., da Rocha, N. S., da Costa, S. T. L., et al.
1168	Reanalysis profile downscaling with WRF model and sensitivity to PBL parameterization schemes
1169	over a subtropical station. Journal of Atmospheric and Solar-Terrestrial Physics, 222, 105724.
1170	doi:10.1016/j.jastp.2021.105724, 2021.
1171	Ding, H., Cao, L., Jiang, H., Jia, W., Chen, Y., and An, J. Influence on the temperature estimation of
1172	the planetary boundary layer scheme with different minimum eddy diffusivity in WRF v3.9.1.1.
1173	Geoscientific Model Development, 14(10), 6135-6153. doi:10.5194/gmd-14-6135-2021, 2021.
1174	Emery, C., Liu, Z., Russell, A. G., Odman, M. T., Yarwood, G., and Kumar, N. Recommendations
1175	on statistics and benchmarks to assess photochemical model performance. JAir Waste Manag Assoc,
1176	67(5), 582-598. doi:10.1080/10962247.2016.1265027, 2017.
1177	Falasca, S., Gandolfi, I., Argentini, S., Barnaba, F., Casasanta, G., Di Liberto, L., et al. Sensitivity
1178	of near-surface meteorology to PBL schemes in WRF simulations in a port-industrial area with
1179	complex terrain. Atmospheric Research, 264, 105824. doi:10.1016/j.atmosres.2021.105824, 2021.
1180	Ferrero, E., Alessandrini, S., and Vandenberghe, F. Assessment of Planetary-Boundary-Layer
1181	Schemes in the Weather Research and Forecasting Model Within and Above an Urban Canopy Layer.
1182	Boundary-Layer Meteorology, 168(2), 289-319. doi:10.1007/s10546-018-0349-3, 2018.
1183	Grell, G. A., and Dévényi, D. A generalized approach to parameterizing convection combining
1184	ensemble and data assimilation techniques. Geophysical Research Letters, 29(14), 38-31-38-34.
1185	doi:https://doi.org/10.1029/2002GL015311, 2002.
1186	Gu, H., Jin, J., Wu, Y., Ek, M. B., and Subin, Z. M. Calibration and validation of lake surface
1187	temperature simulations with the coupled WRF-lake model. Climatic Change, 129(3), 471-483.
1188	doi:10.1007/s10584-013-0978-y, 2015.
1189	He, J., Chen, D., Gu, Y., Jia, H., Zhong, K., and Kang, Y. Evaluation of planetary boundary layer
1190	schemes in WRF model for simulating sea-land breeze in Shanghai, China. Atmospheric Research,
1191	278, 106337. doi:10.1016/j.atmosres.2022.106337, 2022.
1192	Hong, SY., Noh, Y., and Dudhia, J. A New Vertical Diffusion Package with an Explicit Treatment
1193	of Entrainment Processes. Monthly Weather Review, 134(9), 2318-2341. doi:10.1175/mwr3199.1,
1194	2006.
1195	Hong, SY., and Pan, HL. Nonlocal Boundary Layer Vertical Diffusion in a Medium-Range
1196	Forecast Model. Monthly Weather Review, 124(10), 2322-2339. doi:10.1175/1520-
1197	0493(1996)124<2322:Nblvdi>2.0.Co;2, 1996.
1198	Hong, SY., and Shin, H. H. Analysis of Resolved and Parameterized Vertical Transports in
1199	Convective Boundary Layers at Gray-Zone Resolutions. Journal of the Atmospheric Sciences,
1200	70(10), 3248-3261. doi:10.1175/jas-d-12-0290.1, 2013.
1201	Hu, XM., Nielsen-Gammon, J. W., and Zhang, F. Evaluation of Three Planetary Boundary Layer
1202	Schemes in the WRF Model. Journal of Applied Meteorology and Climatology, 49(9), 1831-1844.





1204	Iacono, M. J., Delamere, J. S., Mlawer, E. J., Shephard, M. W., Clough, S. A., and Collins, W. D.
1205	Radiative forcing by long-lived greenhouse gases: Calculations with the AER radiative transfer
1206	models. Journal of Geophysical Research: Atmospheres, 113(D13).
1207	doi: <u>https://doi.org/10.1029/2008JD009944</u> , 2008.
1208	Janjić, Z. I. The Step-Mountain Coordinate: Physical Package. Monthly Weather Review, 118(7),
1209	1429-1443. doi:10.1175/1520-0493(1990)118<1429:Tsmcpp>2.0.Co;2, 1990.
1210	Janjić, Z. I. The Step-Mountain Eta Coordinate Model: Further Developments of the Convection,
1211	Viscous Sublayer, and Turbulence Closure Schemes. Monthly Weather Review, 122(5), 927-945.
1212	doi:10.1175/1520-0493(1994)122<0927:Tsmecm>2.0.Co;2, 1994.
1213	Jia, W., and Zhang, X. The role of the planetary boundary layer parameterization schemes on the
1214	meteorological and aerosol pollution simulations: A review. Atmospheric Research, 239.
1215	doi:10.1016/j.atmosres.2020.104890, 2020.
1216	Jia, W., and Zhang, X. Impact of modified turbulent diffusion of PM2.5 aerosol in WRF-Chem
1217	simulations in eastern China. Atmos. Chem. Phys., 21(22), 16827-16841. doi:10.5194/acp-21-
1218	16827-2021, 2021.
1219	Jia, W., Zhang, X., Wang, H., Wang, Y., Wang, D., et al. Comprehensive evaluation of typical
1220	planetary boundary layer (PBL) parameterization schemes in China. Part I: Understanding
1221	expressiveness of schemes for different regions from the mechanism perspective. Zendo [data set],
1222	https://doi.org/10.5281/zenodo.7792241, 2023.
1223	Jiménez, P. A., and Dudhia, J. Improving the Representation of Resolved and Unresolved
1224	Topographic Effects on Surface Wind in the WRF Model. Journal of Applied Meteorology and
1225	Climatology, 51(2), 300-316. doi:10.1175/jamc-d-11-084.1, 2012.
1226	Kusaka, H., Kondo, H., Kikegawa, Y., and Kimura, F. A Simple Single-Layer Urban Canopy Model
1227	For Atmospheric Models: Comparison With Multi-Layer And Slab Models. Boundary-Layer
1228	Meteorology, 101(3), 329-358. doi:10.1023/A:1019207923078, 2001.
1229	Ma, Z., Zhao, C., Gong, J., Zhang, J., Li, Z., Sun, J., et al. Spin-up characteristics with three types
1230	of initial fields and the restart effects on forecast accuracy in the GRAPES global forecast system.
1231	Geoscientific Model Development, 14(1), 205-221. doi:10.5194/gmd-14-205-2021, 2021.
1232	Mass, C. F., Ovens, D., Westrick, K., and Colle, B. A. DOES INCREASING HORIZONTAL
1233	RESOLUTION PRODUCE MORE SKILLFUL FORECASTS?: The Results of Two Years of Real-
1234	Time Numerical Weather Prediction over the Pacific Northwest. Bulletin of the American
1235	<i>Meteorological Society</i> , 83(3), 407-430. doi:10.1175/1520-0477(2002)083<0407:Dihrpm>2.3.Co;2,
1236	2002.
1237	Mellor, G. L., and Yamada, T. A Hierarchy of Turbulence Closure Models for Planetary Boundary
1238	Layers. Journal of Atmospheric Sciences, 31(7), 1791-1806. doi:10.1175/1520-





1239	0469(1974)031<1791:Ahotcm>2.0.Co;2, 1974.
1240	Mellor, G. L., and Yamada, T. Development of a turbulence closure model for geophysical fluid
1241	problems. Reviews of Geophysics, 20(4), 851-875. doi:10.1029/RG020i004p00851, 1982.
1242	Meng Lu, Zhao Tianliang, Yang Xinghua, Liu Chong, He Qing, and Jingxin, D. An assessment of
1243	atmospheric boundary layer schemes over the Taklimakan Desert hinterland. Journal of the
1244	Meteorological Sciences, 38(2), 157-166. 2018.
1245	Monahan, A.H., and Abraham, C.Climatological Features of the Weakly and Very Stably Stratified
1246	Nocturnal Boundary Layers. Part II: Regime Occupation and Transition Statistics and the Influence
1247	of External Drivers. Journal of the Atmospheric Sciences, 76(11), 3485-3504. doi:10.1175/jas-d-19-
1248	0078.1, 2019.
1249	Monin, A. S., and Obukhov, A. M. Basic laws of turbulent mixing in the surface layer of the
1250	atmosphere. Akad. Nauk SSSR Trud. Geofiz. Inst, 24, 163-187. 1954.
1251	Morrison, H., Thompson, G., and Tatarskii, V. Impact of Cloud Microphysics on the Development
1252	of Trailing Stratiform Precipitation in a Simulated Squall Line: Comparison of One- and Two-
1253	Moment Schemes. <i>Monthly Weather Review</i> , 137(3), 991-1007. doi:10.1175/2008mwr2556.1, 2009.
1254	Nakanishi, M., and Niino, H. An Improved Mellor-Yamada Level-3 Model with Condensation
1255	Physics: Its Design and Verification. Boundary-Layer Meteorology, 112(1), 1-31.
1256	doi:10.1023/B:BOUN.0000020164.04146.98, 2004.
1257	Nielsen-Gammon, J. W., Hu, XM., Zhang, F., and Pleim, J. E. Evaluation of Planetary Boundary
1258	Layer Scheme Sensitivities for the Purpose of Parameter Estimation. Monthly Weather Review,
1259	138(9), 3400-3417. doi:10.1175/2010mwr3292.1, 2010.
1260	Noh, Y., Cheon, W. G., Hong, S. Y., and Raasch, S. Improvement of the K-profile Model for the
1261	Planetary Boundary Layer based on Large Eddy Simulation Data. Boundary-Layer Meteorology,
1262	107(2), 401-427. doi:10.1023/A:1022146015946, 2003.
1263	Oke, T. R., Mills, G., Christen, A., and Voogt, J. A. (2017). Urban Climates. United States of
1264	America: Cambridge University
1265	Paulson, C. A. The Mathematical Representation of Wind Speed and Temperature Profiles in the
1266	Unstable Atmospheric Surface Layer. Journal of Applied Meteorology and Climatology, 9(6), 857-
1267	861. doi:10.1175/1520-0450(1970)009<0857:Tmrows>2.0.Co;2, 1970.
1268	Persson, P. O. G., Walter, B., Bao, J. W., and Michelson, S. A. (2001). 3 VALIDATION OF
1269	BOUNDARY-LAYER PARAMETERIZATIONS IN A MARITIME STORM USING AIRCRAFT DATA.
1270	Pleim, J. E. A Combined Local and Nonlocal Closure Model for the Atmospheric Boundary Layer.
1271	Part I: Model Description and Testing. Journal of Applied Meteorology and Climatology, 46(9),
1272	1383-1395. doi:10.1175/jam2539.1, 2007.
1273	Shen, W., Lu, Z., Ye, G., Zhang, Y., Chen, S., and Xu, J. (2022). Exploring the Impact of Planetary
1274	Boundary Layer Schemes on Rainfall Forecasts for Typhoon Mujigae, 2015. Atmosphere, 13(2).





1275	doi:10.3390/atmos13020220
1276	Stull, R. B. Transilient Turbulence Theory. Part I: The Concept of Eddy-Mixing across Finite
1277	Distances. Journal of Atmospheric Sciences, 41(23), 3351-3367. doi:10.1175/1520-
1278	0469(1984)041<3351:Tttpit>2.0.Co;2, 1984.
1279	Stull, R. B. (1988). Turbulence Closure Techniques. An introduction to boundary layer meteorology
1280	(Vol. 6). London: Atmospheric Sciences Library.
1281	Sun, J., Mahrt, L., Banta, R. M., and Pichugina, Y. L. Turbulence Regimes and Turbulence
1282	Intermittency in the Stable Boundary Layer during CASES-99. Journal of the Atmospheric Sciences,
1283	69(1), 338-351. doi:10.1175/jas-d-11-082.1, 2012.
1284	Troen, I. B., and Mahrt, L. A simple model of the atmospheric boundary layer; sensitivity to surface
1285	evaporation. Boundary-Layer Meteorology, 37(1), 129-148. doi:10.1007/BF00122760, 1986.
1286	Vignon, E., van de Wiel, B. J. H., van Hooijdonk, I. G. S., Genthon, C., van der Linden, S. J. A.,
1287	van Hooft, J. A., et al. Stable boundary-layer regimes at Dome C, Antarctica: observation and
1288	analysis. Quarterly Journal of the Royal Meteorological Society, 143(704), 1241-1253.
1289	doi:10.1002/qj.2998, 2017.
1290	Wang, C., Shen, Y., Luo, F., Cao, L., Yan, J., and Jiang, H. Comparison and analysis of several
1291	planetary boundary layer schemes in WRF model between clear and overcast days. Chinese J.
1292	Geophys. (in Chinese), 60(3), 924-934. 2017.
1293	Williams, P. D. Modelling climate change: the role of unresolved processes. Philos Trans A Math
1294	Phys Eng Sci, 363(1837), 2931-2946. doi:10.1098/rsta.2005.1676, 2005.
1295	Wyngaard, J. C., and Brost, R. A. Top-Down and Bottom-Up Diffusion of a Scalar in the Convective
1296	Boundary Layer. Journal of Atmospheric Sciences, 41(1), 102-112. doi:10.1175/1520-
1297	0469(1984)041<0102:Tdabud>2.0.Co;2, 1984.
1298	Xie, B., Fung, J. C. H., Chan, A., and Lau, A. Evaluation of nonlocal and local planetary boundary
1299	layer schemes in the WRF model. Journal of Geophysical Research: Atmospheres, 117(D12).
1300	doi:10.1029/2011JD017080, 2012.
1301	Zhang, D., and Anthes, R. A. A High-Resolution Model of the Planetary Boundary Layer-
1302	Sensitivity Tests and Comparisons with SESAME-79 Data. Journal of Applied Meteorology and
1303	<i>Climatology</i> , 21(11), 1594-1609. doi:10.1175/1520-0450(1982)021<1594:Ahrmot>2.0.Co;2, 1982.
1304	Zhou, B., Zhu, K., and Xue, M. A Physically Based Horizontal Subgrid-Scale Turbulent Mixing
1305	Parameterization for the Convective Boundary Layer. Journal of the Atmospheric Sciences, 74(8),
1306	2657-2674. doi:10.1175/jas-d-16-0324.1, 2017.
1307	