



NEWTS1.0: Numerical model of coastal Erosion by 1

Waves and Transgressive Scarps 2

- Rose V. Palermo^{1,2}, J. Taylor Perron³, Jason M. Soderblom³, Samuel P. D. Birch³, Alexander G. 3
- 4 Hayes⁴, Andrew D. Ashton⁵

¹ U. S. Geological Survey, St. Petersburg Coastal and Marine Science Center, St. Petersburg, Florida 33701, USA

- ² MIT-WHOI Joint Program in Oceanography/Applied Ocean Science & Engineering, Cambridge and Woods Hole, MA, USA
- 56789 ³Department of Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of Technology, Cambridge, MA. USA
- 10 ⁴Department of Earth, Atmospheric and Planetary Sciences, Cornell University, Cambridge, MA, USA
- 11 ⁵Department of Geology and Geophysics, Woods Hole Oceanographic Institution, Woods Hole, MA, USA
- 12 Correspondence to: Rose V. Palermo (rpalermo@usgs.gov)

13 Abstract: Models of rocky coast erosion help us understand the physical phenomena that control

14 coastal morphology and evolution, infer the processes shaping coasts in remote environments,

15 and evaluate risk from natural hazards and future climate change. Existing models, however, are

16 highly complex, computationally expensive, and depend on many input parameters; this limits

17 our ability to explore planform erosion of rocky coasts over long timescales (100s to 100,000s

- 18 years) and a range of conditions. In this paper, we present a simplified cellular model of coastline
- 19 evolution through uniform erosion and wave-driven erosion. Uniform erosion is modeled as a
- 20 constant rate of retreat. Wave erosion is modeled as a function of fetch, the distance over which
- 21 the wind blows to generate waves, and the angle between the incident wave and the shoreline.

22 This reduced complexity model can be used to evaluate how a detachment-limited coastal

23 landscape reflects climate, sea level history, material properties, and the relative influence of

24 different erosional processes.

25 1 Introduction

26 Rocky coastlines are erosional coastal landforms resulting from the landward 27 transgression of a shoreline through bedrock. They make up approximately 80% of global coasts 28 (Emery and Kuhn, 1980) and often erode slowly through the impact of waves (Adams et al., 29 2002, 2005), abrasion by sediment (Sunamura, 1976; Robinson, 1977; Walkden & Hall, 2005; 30 Bramante et al., 2020), and chemical weathering (Sunamura, 1992; Trenhaile, 2001). Rocky 31 coastlines protect coastal communities from erosion and flooding, provide sediment for estuaries, 32 marshes, and beaches, serve as important habitats (such as kelp forests), and support tourism 33 economies.

34 The imprint that each erosional mechanism leaves on the shoreline may be further 35 complicated by sea-level changes, accumulation and redistribution of sediment, heterogeneities 36 in the bedrock, or climate forcings. Wave-driven erosion occurs at a rate proportional to the 37 wave power (Huppert et al., 2020). Therefore, over long time scales, waves tend to erode more 38 exposed parts of coastlines preferentially, blunting headlands while preserving the shapes of 39 sheltered embayments. Uniform erosional processes, like dissolution or mass backwasting, erode 40 at a nearly uniform rate everywhere along a coastline and result in smooth, rounded coastal 41 features punctuated by skewed, pointy promontories or headlands (Howard, 1995).





Although the relative influence of uniform erosion processes, such as dissolution, and
wave-driven erosion are still being quantified (Trenhaile, 2015), the shape of coastlines may
offer a means to infer dominant processes in remote environments where in situ measurements
are impractical. One such example are arctic coasts, where local field data are sparse. A reduced
complexity model of long-term, planform evolution of erosion-dominated coasts can provide

- 47 insights about the importance of wave erosion relative to uniform erosion, such as backwasting
- 48 of permafrost (Günther et al., 2013). Here, we present a reduced-complexity model of
- 49 detachment-limited coastal erosion by uniform erosion and wave erosion. We test the model by
- 50 comparing our numerical solution of erosion with an analytical solution and test for model result
- 51 sensitivity to grid resolution and input parameters.

52 2 Background

53 2.1 Previous Models of Coastal Erosion

54 2.1.1 Models of uniform erosion

Shorelines formed by dissolution in karst landscapes have received some attention,
mostly in the context of cave collapse features or sinkholes (Johnson, 1997; Martinez et al.,
1998, Yechieli et al., 2006). However, most research has focused on the initial formation of these
features; studies of the long-term retreat of coastlines due to dissolution are focused on the
meter-scale erosion of coastal notches through mechanical and biochemical erosion and by
dissolution (Trenhaile 2013; Trenhaile, 2015) and to our knowledge have not been evaluated
over a larger spatial scale.

62 Howard (1995) modeled the retreat of a closed basin scarp as a uniform erosion process. 63 Howard's approach identifies gridded domain points as either interior or exterior to the 64 escarpment and erodes the escarpment edge at a constant rate in all directions originating from 65 adjacent points (Howard, 1995). In his model experiments, the escarpment retreats uniformly toward the interior of the domain from the exterior. This uniform scarp retreat is analogous to 66 67 coastline retreat in response to dissolution of a uniform substrate. Although Howard's model was 68 designed for a different, subaerial system, uniform erosion of a lake shoreline can be described 69 with the same process law, as we assume the planform lake shoreline also erodes at the same rate 70 in all directions.

71 2.1.2 Models of wave-driven erosion

72 Models of rocky-coastline geomorphology have historically focused on the erosion of the 73 cross-shore profile through sea-level rise (Walkden and Hall, 2005; Young et al., 2014), wave 74 impacts (Adams et al., 2002, 2005; Huppert et al., 2020), and the competing effects of sediment abrasion and sediment cover (Kline et al., 2014; Young et al., 2014; Sunamura 2018; Trenhaile, 75 76 2019). But recent work has explored the alongshore variability (Walkden and Hall, 2005) and 77 planform evolution of these features (Limber & Murray, 2011; Limber et al., 2014; Sunamura, 78 2015; Palermo et al., 2021), with particular focus on either the relationship between planform 79 morphology and retreat rates following storms (Palermo et al., 2021) or the persistence of an 80 equilibrium coastline shape consisting of headlands interspersed with pocket beaches due to





variable lithology, grain size, or sediment tools and cover (Trenhaile, 2016; Limber & Murray,
2011; Limber et al., 2014).

83 Existing models of planform erosion of rocky beaches include 1) a mesoscale (1 to 100 84 years) alongshore-coupled cross-shore profile model, SCAPE (Walkden and Hall, 2005), in 85 which waves erode the substrate when the substrate is not armored by sediment and sediment is transported by waves using linear wave theory; 2) a numerical model of sea-cliff retreat that 86 87 focuses on the mechanical abrasion of a notch at the cliff toe and subsequent failure of the cliff 88 and sediment comminution in the surf zone (Kline et al., 2014); and 3) a numerical model of 89 headlands and pocket beaches that takes into account wave energy convergence/divergence and 90 the processes of sediment production and redistribution by waves (Limber at al., 2014).

91 Previous work on marsh-shoreline erosion considers the heterogeneity of substrate 92 erodibility using a percolation theory model (Leonardi & Fagherazzi, 2015). In this system, low 93 wave energy conditions lead to patchy failure of large marsh portions, resulting in a strong 94 dependence on the spatial distribution of substrate resistance. In contrast, high-wave-energy 95 conditions cause the shoreline to erode uniformly, such that the spatial heterogeneity in marsh 96 erodibility does not influence the erosion rate (Leonardi & Fagherazzi, 2015). This ignores 97 variations in fetch, which can be important for rocky coastal systems.

These previous process-based models are all computationally expensive and require
 specific knowledge of sediment and wave characteristics to accurately apply at local scales. To
 model systems for which minimal field data are available, or to explore the general behavior of
 planform erosion in rocky coasts under a broad range of conditions, a reduced-complexity model
 (Ranasinghe, 2020) is necessary.

103 3 Model

104 We developed the Numerical model of coastal Erosion by Waves and Transgressive 105 Scarps, V1.0 (NEWTS1.0) (Palermo et al., 2023) to study the planform-shoreline erosion of detachment-limited coasts by waves, uniform erosion, or a combination of these processes. This 106 reduced-complexity model can be used to explore long-term (1000s-10,000s of years) trends in 107 landscape evolution that result from these processes. Uniform erosion includes dissolution or 108 109 mass backwasting and is modeled with a spatially uniform rate of shoreline retreat, which generally smooths the coastline but generates cuspate points where promontories are eroded. 110 111 Wave erosion occurs in proportion to the wave energy that the coastline is exposed to and to the 112 angle of incidence of the incoming waves, such that the erosion rate depends on the wave energy 113 in the cross-shore direction per unit of length along the coast (Komar, 1997; Ashton & Murray, 2009; Huppert et al., 2020). Coastlines that have larger exposure to the lake (larger fetch) 114 115 experience higher wave energy and therefore faster wave erosion. We model this energy-116 dependent erosion by computing the fetch of every incident wave angle that may impact a given 117 point on the shoreline and weighting this fetch by the cosine of the angle between the incident 118 wave crests and the shoreline. Mathematically, this is equivalent to the dot product of the 119 direction of wave travel and the direction normal to the shoreline. 120

121 3.1 General description and model setup

122 3.1.1 Model domain and structure





123 3.1.1.1 Model domain

124 The domain of the model (Fig. 1) is a grid discretized into N_x cells in the x direction and N_y cells in the y direction, with cell spacings Δx and Δy , such that $x_i = i\Delta x$ and $y_j = j\Delta y$. The 125 126 value of each grid cell, $z_{i,i}$, corresponds to the landscape elevation. The boundaries of the grid 127 are periodic. Each cell in the domain is defined as either liquid or land based on its elevation 128 relative to sea level. Cells below sea level are fixed and do not erode. Shoreline cells, defined as 129 land cells directly adjacent to liquid, may be eroded by coastal processes through uniform 130 erosion and wave erosion. Sea level is an input to the system that the user can vary throughout a 131 model run. 132

133



134

Figure 1: Example model domain with a sea level of a) 0 m, b) 40 m, and c) 80 m. This domainis used in Figs. 4 and 5.

137 3.1.1.2 Identification of lake and shoreline cells

138 Boundaries in the grid are identified using pixel connection definitions of either 4-connected 139 or 8-connected cells. Liquid cells that are 8-connected to each other comprise the same lake. 140 Islands are defined as groups of land cells that are surrounded by liquid cells. Lakes can also 141 occur inside islands and islands inside these lakes, so we define a lake hierarchy to identify and 142 model each lake individually. The first level in this hierarchy is the land that is connected to the 143 border of the domain. First order lakes are lakes that are immediately surrounded by this land 144 that extends to the border of the domain. A first order island is immediately surrounded by a first 145 order lake. A second order lake is surrounded by a first order island, and so on. This continues 146 such that Nth-order islands are surrounded by Nth-order lakes, and Nth-order lakes are 147 surrounded by N-minus-one-order islands. This hierarchy allows us to identify and isolate unique 148 lakes, which will be important when we consider wave-driven erosion.

- 149
- 150 3.1.1.3 Cellular grid erosion

151 Each cell starts with an initial strength, S_{init}, (see Sections 3.1.3 to 3.3) which is depleted 152 according to a rate law associated with each coastal process until reaching 0 (see Sections 3.2 153 and 3.3), at which point the cell erodes. Coastal erosion occurs on shoreline cells, defined as land 154 cells adjacent to liquid cells, and decreases the elevation of those cells by a specified depth of 155 erosion, d_e , which is user specified. For cells eroded by coastal processes, z(t) = z(t-1) - z(t-1)156 d_e , where t is model time. For uniform erosion, d_e is conceptualized as the scarp dissolution 157 depth. For wave erosion, d_e is conceptualized as a wave base. Shoreline cells become lake cells 158 once eroded. To avoid numerical artifacts associated with the time discretization, the timestep 159 must be set such that the amount of erosion per iteration is a small fraction of the total cell size.





- 160 In practice, we set the time step to erode less than $1/10^{\text{th}}$ of a cell at a given time given the cell 161 spacing and rate law. The model run terminates if a lake cell becomes adjacent to a boundary cell
- 162 because the wave erosion model requires a closed coastline.
- 163
- 164 3.1.1.4 Order of operations

165 During each timestep, erosion occurs according to three steps, if enabled: 1) Sea-level

166 Change, 2) Wave Erosion, and 3) Uniform Erosion (Fig. 2). Here we describe the general model

167 components and simulation procedure. The governing equations for Uniform Erosion and Wave

- 168 erosion are outlined in more detail in sections 3.2 and 3.3, respectively.
- 169



170

- Figure 2: Model structure showing the time loop in which the model 1) updates sea-level change,then calculates shoreline erosion due to 2) waves and 3) uniform erosion processes.
- 173

The first operation of the model is sea-level change. The sea level changes as an input rate or
according to an input sea-level curve. The new sea level is used to define the lake(s) and
shoreline(s) (Section 3.1.1.2 and 3.1.2).

Next, wave erosion of the shoreline(s) occurs as a function of the fetch—the open-water 177 178 distance wind and waves travel before reaching a point on the coast-and the angle between the wave crests of the incident waves, $\boldsymbol{\varphi}$, and the azimuth of the shoreline, $\boldsymbol{\theta}$ (Section 3.3). In this 179 180 module, the shoreline is first identified and traced such that shoreline cells are ordered in a 181 counterclockwise direction. The shoreline is then used to calculate the shoreline angle, incident 182 wave angle, and associated fetch at each cell along the shoreline (Section 3.3.1). The elevation of 183 eroded shoreline cells is lowered, their labels are changed to liquid cells as appropriate, and the 184 shoreline is updated (See Section 3.4, Fig 5). This approach considers sediment removal as





- instantaneous. Future variations of the model could consider the erosion also as a function of theheight of the material being eroded or the excavation rate of weathered rubble.
- 187 Finally, uniform erosion of the updated shoreline occurs (Section 3.2). Here, the shoreline
- 188 erodes as a function of the alongshore length of the shoreline as measured along cell boundaries
- 189 (Section 3.1.2 and 3.2). And again, the elevation of eroded shoreline cells is lowered, the labels
- 190 of eroded cells are changed to liquid cells, and the shoreline is updated.
- 191
- 192 3.1.2 Defining the shoreline
- 193 There are two options for defining shoreline cells: land cells that are either 4-
- 194 connected to the lake (Fig. 3a) or 8-connected to the lake (Fig. 3b). We choose Δx and Δy to be
- small enough to represent the relevant features of the shoreline.



196

- 197 Figure 3: Shoreline cells and associated strength loss weighting for a shoreline that is a) 4-
- 198 connected to liquid cells or b) 8-connected to liquid cells. Arrows point in the direction of
- erosion into each shoreline cell from neighboring lake cells. Increasing darkness of shorelinecells indicate increasing strength loss weighting.

The shoreline cells need to be ordered so that the lake can be represented as a polygon for the fetch computation. To order the shoreline cells in closed loops, we start at the first indexed shoreline cell of the longest shoreline and move counterclockwise to find the next shoreline cell. Once a sequence of the first 3 cells is repeated, the loop is closed and the shoreline is deemed complete. Any remaining shoreline cells that do not lie on this loop represent the shoreline of a separate first-order lake, or of an island or higher order lake contained within the lake. Next,





207 ordering the shorelines of the islands contained within the current lake begins on the first 208 remaining shoreline cell. We repeat this process until all land cells bordering liquid are included

209 in a closed shoreline. When there are multiple first-order lakes in a landscape domain, the

210 shorelines for each lake and its enclosed islands are ordered one at a time.

211 3.1.3 Cell strength and coastal erosion processes

212 All cells start with an initial strength, S_{init}, which represents how difficult it is to erode 213 the land (Equation 1). We model the domain as having uniform strength, but this could easily be 214 extended to a scenario with heterogeneous strength. The strength of a cell is initialized as a 215 reference strength, S_0 , multiplied by the ratio between the cell area, $A = \Delta x \Delta y$, and a reference 216 cell area, $A_0 = \Delta x_0 \Delta y_0$, with reference spacing Δx_0 and Δy_0 (Equation 1). The reference 217 strength and area nondimensionalize strength and maintain proportions that mitigate 218

discretization bias. The magnitude of these values can be chosen by the user.

$$S_{init} = S_0 \frac{A}{A_c} \tag{1}$$

220 Strength is lost from each shoreline cell at a rate that depends on the exposed perimeter of 221 the cell and an erosion rate law specific to either uniform erosion or wave erosion processes. 222 Change in strength is grid-independent for grids sufficiently fine to satisfy model stability 223 because the strength is initialized with a reference cell area in proportion to the parameterized 224 cell area. To mitigate discretization bias, Δx , Δy , and Δt must be sufficiently small that Δt is less 225 than the time to completely erode a cell (See Sections 3.2 and 3.3), and that Δx and Δy properly 226 represent the shoreline morphology. In practice, we choose Δx to be equal to Δy .

227 As time progresses, each shoreline cell loses strength until failure, $S_{i,j} = 0$, at which 228 point the cell has eroded. It is possible for the strength loss in one time step to exceed the 229 remaining strength of the cell. When this occurs, the excess time spent eroding the cell is passed 230 along to all new shoreline neighbors of the eroded cell, representing the time of erosion that 231 neighboring cell will incur after the erosion of the original shoreline. If a new shoreline cell is 232 inheriting excess time from multiple neighbors, the mean excess time is used to compute the 233 strength loss. In our simulations, taking the mean of the excess time resulted in the least grid 234 bias.

235 Modeled erosion could be underestimated or redistributed improperly if the strength loss 236 for an eroding cell is consistently large relative to the initial strength of the domain. The 237 shoreline would then not update with the newly exposed cells, rather constantly passing strength 238 loss to its neighbors, and inaccurately characterizing the morphology. We implement a sub-239 timestep routine to capture the effect of the changing shoreline within a single timestep when the 240 strength loss of any shoreline cell in the domain exceeds a certain threshold of the initial 241 strength, α , which ranges between 0 and 1. In the modified time-step routine, the damage is 242 computed and the shoreline updated in sub-timesteps, which segments the time-step and allows 243 erosion to occur in smaller increments.

244 3.2 Uniform erosion model

245 The rate of shoreline retreat by uniform erosion is set by an erodibility coefficient, 246 $k_{uniform}$ (Eq. 2). Strength loss due to uniform erosion occurs as a function of the amount of shoreline in contact with the lake for a given cell, represented as the number of 4-connected 247





sides, s_c , and 8-connected corners, c, in contact with lake cells (Eq. 3; Fig. 3). Because the

249 diagonal of the cell is longer than the side by a factor of $\sqrt{2}$, it would take $\sqrt{2}$ times longer for a

shoreline to retreat across a cell diagonal than in the perpendicular direction. To correct for this

in our model, the strength loss computed from an exposed corner is $\sqrt{2}/2$ as much as the strength lost from an exposed side.

$$\frac{dx}{dt} = k_{uniform},$$
(2)

254
$$\frac{\Delta S_{i,j}}{S_0} = -k_{uniform} \left(S_c + \frac{\sqrt{2}c}{2} \right) \frac{\Delta x}{\Delta x_0} \Delta t, \tag{3}$$

255 3.3 Wave erosion model

256 Wave erosion occurs at a rate determined by a wave erodibility coefficient, k_{wave} [m·yr⁻ 257 ¹], and the wave energy flux in the cross-shore direction, *E* (Eq. 4). The wave energy flux 258 depends on the wave height, *H*, and the angle between the wave crests of the incident waves, φ , 259 and the azimuth of the shoreline, θ (Eq. 5). Wave height scales with fetch, *F*, such that $H \propto \sqrt{F}$ 260 (Hasslemann, 1973; Smith and Waseda, 2008). Therefore, we use fetch to approximate the wave 261 energy density for a wave from a given direction on a coastline (Eq. 6).

$$\frac{dx}{dt} = k_{wave}E \quad , \tag{4}$$

263
$$E = \frac{1}{16} \rho g H^2 \cos \left(\varphi - \theta\right), \tag{5}$$

$$E \propto \rho g F \cos \left(\varphi - \theta \right), \tag{6}$$

265 The strength loss of a cell due to waves can be described as

266
$$\frac{\Delta S_{i,j}}{S_0} = -k_{wave} \left(S_c + \frac{\sqrt{2}c}{2}\right) \int_{\varphi=0}^{2\pi} F(\varphi) \cos\left(\varphi - \theta\right) d\varphi \frac{\Delta x}{\Delta x_0} \Delta t.$$
(7)

267 If the strength loss in a time step exceeds a parameter-set threshold, a sub-timestep 268 routine is implemented. Because the fetch calculation is the costliest step of the model, in this 269 sub-timestep routine, we estimate the fetch weighting by interpolating the fetch of the nearest 270 neighbor shoreline cells. This avoids additional costly fetch computations during the sub-271 timestep updates and allows us to approximate erosion driven by waves in a way that limits error 272 without slowing down the model simulation.

273 3.3.1 Modeling wave energy density

The rate of strength loss of each shoreline cell is proportional to the wave energy density. We model the wave energy density to be proportional to the fetch and the cosine of the angle between the incident wave crest and the shoreline (Fig. 4). To compute this quantity, we measure the fetch in all directions around the shoreline, in increments of $d\varphi$, for each shoreline cell. For each direction, we extend a ray from the cell center in the direction 90° – φ and step along the ray in increments of a distance δ until reaching the opposite shore. When the ray extends past the opposite shoreline, we take one step back and define this point as the intersection. The distance





between this intersection and the originating shoreline cell center is the fetch in the direction

from which a wave would propagate (Fig 4b). To calculate the amount of strength loss each cell incurs, we compute the area of a polygon defined by the ray-shoreline intersections for that cell

(Fig. 4a). We call this area the "fetch area." The length of the ray in each direction is then

(Fig. 4a). We call this area the field area. The length of the ray in each direction is then weighted by the cosine of the angle between the shoreline and the incident wave crest, $\varphi - \theta$

286 (Fig. 4a). The area of the polygon defined by these cosine-weighted fetch lengths is computed

287 and called the "wave area." The wave area for each point on the shoreline approximates the

288 integral in Eq. 7.



289

Figure 4: a) Fetch area (black) and wave area (white) computed for a point (red circle) on a
typical model shoreline (blue). The area shown in b) is outlined in red. b) Zoomed-in view of
fetch line-of-sight rays (black) and angle-weighted line-of-sight rays (white) computed for the

same point. In this example, $d\varphi = 2^{\circ}$ and the ray step size, $\delta = 0.05$ m.

294 3.4 Model output

295 The model can be initialized with any user defined topographic model. In the simulations 296 presented here, we initialize the grid with a synthetic topography consisting of a pseudo-fractal 297 surface with variance of 10,000 superimposed on an elliptical depression with a depth of 25% of 298 the domain relief and eroded by river incision to 95% of the initial terrain relief using a 299 landscape evolution model (Perron et al., 2008, 2009, 2012). We then flood the domain by 300 raising sea level by 40 m. The model of shoreline retreat by uniform and wave erosion is then applied to the domain. Here, we show examples of an initial landscape eroded by either wave 301 302 erosion or uniform erosion, to illustrate separately the effects of the two erosional mechanisms in 303 the model (Fig. 5). However, all model components may be run in combination. We do not provide examples of combined uniform and wave erosion models here. 304





305 The initial shoreline exhibits a dendritic shape due to flooding of the incised river valleys 306 (Fig. 5). Through time, the uniform erosion model drives shoreline retreat at the same rate 307 everywhere around the perimeter of the lake, resulting in widening valleys and increasing the 308 pointedness of promontories or headlands (Fig. 5). The overall shape of the lake is maintained, 309 but becomes smoother and tends toward circular. In the case of wave erosion, the river valleys 310 erode slowly while the exposed parts of the coast erode more rapidly (Fig. 5). The embayed river 311 valleys largely maintain their shapes, whereas the central, high-fetch portion of the lake grows 312 larger and smoother.



313

314 Figure 5: Shaded relief maps of example model simulations of uniform erosion and wave erosion 315 through time, starting from the same initial condition. Blue color indicates liquid cells, with darker blues indicating deeper depths. Gold color indicates land cells, with lighter shades 316 indicating higher elevations. Black lines trace shorelines. Erodibility coefficients are $k_{wave} =$ 317 $k_{uniform} = 0.00001 \text{ m} \cdot \text{yr}^{-1}$. Uniform erosion (top) results in greater overall smoothness that is 318 punctuated by pointy headlands, whereas wave erosion (bottom) results in blunted headlands, 319 320 smooth open sections of coast, and preservation of sharp features in sheltered areas. Landscape 321 time-steps shown correspond to similar amounts of erosion between wave and uniform 322 examples.

323 Because the long-term retreat of bedrock coastlines is generally too slow to be 324 measurable with historical aerial and satellite images, the data needed to fully validate this model 325 are not presently available. Nonetheless, a visual comparison can be drawn between coastal 326 features found on Earth and the coastline shapes generated by each end-member erosional 327 mechanism in the model, which is the main goal of our modeling approach. Instances of 328 dissolution and backwasting include karst lakes found in the Plitvice Lakes, Croatia (Fig. 6d) and 329 in Florida, USA (Fig. 6a) as well as scarp retreat due to weathering and backwasting, such as 330 Caineville Mesa, Utah, USA (Fig. 6b). These shorelines exhibit the same overall smoothness,





- punctuated by sharp headlands, as is seen in the shorelines formed by uniform erosion in our $\frac{222}{100}$ and $\frac{1}{100}$
- 332 model (Fig. 5).

A bedrock lake that has been eroded recently by waves is exemplified by Lake Rotoehu,
 New Zealand (Fig. 6c). In these examples, we observe blunted headlands and smooth, rounded
 stretches in open sections of coast, and crenulated shorelines in more protected areas of coast –

similar to the shorelines formed by wave erosion in our model (Fig. 5).

- 337
- 338 339



340

353

341 Figure 6: a) Karst lakes in Florida, USA (Map Data: © Google Earth, Landsat/Copernicus). Lake

342 Butler and the surrounding region. b) Caineville Mesa, Utah, USA (Map Data: © Google Earth,

343 Landsat/Copernicus). c) Lake Rotoehu, New Zealand (Map Data: © Google Earth, CNS/Airbus).

- d) Plitvice Lakes, Croatia (Map Data: © Google Earth, DigitalGlobe).
- 345 4 Model tests

346 4.1 Comparison with analytical solution and sensitivity to shoreline connectedness

For the simple case of an initially circular shoreline, we compute the shoreline evolution analytically and compare this known solution with our numerical model results. For the uniform erosion case, the rate at which the radius of a circle increases, \dot{r} , is equal to the constant of erosion, in this case $k_{uniform}$.

$$\dot{r}(t) = k_{uniform} \tag{8}$$

352 Therefore, the radius, r, at time, t, and initial radius, r_0 , for uniform erosion is:

$$r(t) = r_0 + k_{uniform}t \tag{9}$$

For wave erosion, the rate of increase of the radius, \dot{r} , depends on the constant of erosion, k_{wave} , and the integral of the fetch, *F*, at each angle between the incoming wave crest and the shoreline, $(\varphi - \theta)$ in all directions around the circle:

357
$$F(\varphi) = r\sqrt{2(1 + \cos(2(\varphi - \theta)))}$$
 (10)





(12)

358

$$\dot{r}(t) = \frac{k_{wave}}{2} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (F(\varphi)\cos(\varphi - \theta))^2 d\varphi$$
(11)

359 Computing this integral simplifies to:

361 Therefore, the radius, r, at time, t, for wave erosion is:

362
$$r(t) = \frac{r_0}{1 - r_0 k_{wave} \frac{3\pi}{4}t}$$
 (13)

 $\dot{r(t)} = k_{wave} \frac{3\pi}{4} r(t)^2$

We use the analytical solution for the radius through time for each case to calculate the
shoreline position and area of the circular lake as it is eroded by either uniform or wave erosion.
To compute the relative error of the numerical model, a test circular lake is eroded for 17,400
years, resulting in approximately 20% and 25% increase in lake area for wave and uniform
erosion, respectively, and compare this to the analytical solution.

368 Because the model operates on a rectangular grid, some amount of distortion of a circle is 369 expected. While this distortion cannot be avoided entirely by increasing the grid resolution, 370 increasing it can reduce the error in the shoreline shape by allowing the shoreline to retreat in 371 finer increments. A fine grid, however, comes at increased computational cost. The spatial 372 resolution, Δx and Δy , should be chosen to be small enough to represent the features of the 373 shoreline, but large enough to keep computational costs reasonable.

We perform these simulations for uniform and wave erosion with both 8-connected and 4-connected versions of the model (Fig. 3). The 8-connected model performs significantly better than the 4-connected model, as shown by the relative error in lake area. The 8 connected case maintains relative error less than 2% throughout the simulation whereas the error in the 4connected model increases roughly linearly with time, ending at approximately 7% (Fig. 7a). The distortion is worse in the 4-connected case for both uniform erosion and wave erosion, and systematically worse in the diagonal directions (Fig. 7b,c). This analysis suggests that grid bias is

a more important source of error in the model than spatial discretization.







382

Figure 7: a) The error in lake area through time of an initially circular lake relative to the
analytical solution for 8-connected (solid) and 4-connected (dotted) models of uniform erosion
(black) and wave erosion (blue). The initial condition (dashed), analytical solution (red), and
modeled 8-connected and 4-connected shorelines at time=17400 are shown for b) uniform
erosion and c) wave erosion, with zoomed in results shown for d) uniform erosion and e) wave
erosion.

389 4.2 Resolution sensitivity

390 4.2.1 Grid resolution

391 Although the grid resolution affects the size of the features that can be resolved in the 392 landscape, it does not substantially affect the amount of coastal erosion. As discussed above, the 393 strength loss in this model is insensitive to grid resolution, Δx , and time step, Δt , assuming that 394 Δx is fine enough to resolve the features of interest and that Δt is small enough to limit erosion 395 to less than the maximum cell strength in a single time step. The total amount of strength in the domain is independent of Δx because the number of cells is proportional to Δx^{-2} and the 396 strength of each cell is proportional to Δx^2 . The damage in each time step is independent of Δx 397 because the number of cells on the shoreline is proportional to Δx^{-1} and the damage per cell is 398 399 proportional to Δx .

400 4.2.2 Threshold strength parameter





401 The threshold strength parameter, α , was introduced to prevent excess strength reduction 402 from being neglected when a cell has less strength than is depleted in a timestep. A smaller

403 threshold strength parameter results in a more frequent application of the sub-timestep routine

404 and smaller sub-timesteps. With a less stringent threshold strength parameter (>0.05), the

405 shoreline may erode more than the analytical solution in a time step, leading to a positive slope

406 in the relative error in strength against the threshold strength parameter (Fig. 8).



407

408 Figure 8: Error in total strength reduction as a function of the threshold strength parameter,

409 expressed as a percentage of the error for the smallest value of the threshold strength parameter,

410 for a typical model lake (the initial condition in Fig. 5) eroded over one time step by uniform

411 erosion (black) and wave erosion (blue).

412 4.3 Fetch ray angular and distance increments

413 We test the sensitivity of the fetch-area calculation to the angle between rays, $d\varphi$, and the 414 ray step size, δ . This test allows us to analyze the error in fetch of a typical model due to these 415 parameters. The error measurements provide a basis for selecting an angle between rays and a 416 ray step size that optimize the trade between computational time and model accuracy.

417 We compute the error in fetch area over a range of ray angles and step sizes. With a fixed 418 ray step size of $0.05\Delta x$ (the nominal step sized used in our simulations), we compute the fetch 419 error for each shoreline cell over a range of 0.012° to 10°, corresponding to 30,000 and 36 rays, 420 respectively. With a fixed ray angle of 2° (the nominal ray angle used in our simulations), we 421 compute the relative fetch error over a range of ray step sizes between $0.01\Delta x$ to Δx . The fetch-422 area error of each cell is computed relative to the fetch area of the finest resolution in each parameter: 2° between rays and a ray step size of $0.05\Delta x$ (Fig. 9). The error, as well as the 423 424 standard deviation in errors, in each scenario converges to zero, indicating that as the angle 425 between rays and the ray step size become small the fetch area converges to a constant value.







426 427

427 Figure 9: Relative error in fetch area for a range of step sizes with ray angle of 2° (black) and for 428 a range of ray angles with step size of $0.05\Delta x$ (red).

429 5 Discussion and Conclusions

In this paper, we present NEWTS1.0, a cellular model of coastline erosion in detachmentlimited environments by uniform erosion and by wave erosion. For uniform erosion, the
coastline erodes at a constant rate everywhere along the shoreline. For wave-driven erosion, the
coastline erodes as a function of the fetch and the angle between the incident waves and the
shoreline.

435 While our uniform erosion rate law is similar to that of Howard (1995), our modeling 436 approach is different. Because there are multiple mechanisms that may erode a coast in our 437 model, memory of the strength loss of the substrate is necessary. Rather than rays extending at a 438 constant rate from the interior points representing retreat as is done in Howard's 1995 model, the 439 strength of shoreline (or scarp edge) points is reduced by an amount proportional to the number 440 and direction of neighboring lake cells. Our wave erosion model contains a dependence on wave 441 power like in other models (Walkden and Hall, 2005; Limber et al., 2014), but simplifies the 442 influence of sediment and other factors to a constant. This simplification is useful for locations 443 without readily available grain size or sediment cover data, and to investigate the long-term 444 influence of these processes. Our model is also unusual among coastal erosion models in that it 445 evaluates multiple closed coastlines (or lakes) in a landscape domain rather than a single reach of 446 open coastline, and that it focuses on the planform morphology of eroding rocky closed-basin 447 shorelines. A limitation of this model is that sediment redistribution is not included in the erosion 448 rate laws and there is no sedimentation along the coast. Sediment abrasion and cover could be 449 incorporated in future versions of our model through a spatially heterogeneous and time-450 dependent erodibility coefficient, k; however, this would likely require parameterization from 451 field data.

As a reduced-complexity model, NEWTS1.0 can be applied to investigate coastal
systems in remote environments where field work is difficult or impossible. This includes
locations such as the arctic or Saturn's moon Titan, home to the only other active coastlines in





455 456 457 458 459 460	our solar system. The simplicity of our model allows for efficient, long-term simulations of coupled landscape evolution and coastal erosion in detachment-limited systems. Among coastal systems on Earth, investigations of fetch dependence and the resulting morphology given a combination of erosional mechanisms would be particularly relevant to the carbonate geomorphology community, as dissolution and wave activity are both often acting simultaneously along these coasts.
461	Acknowledgments
462	We thank David Mohrig, Di Jin, Heidi Nepf, Jorge Lorenzo-Trueba, Santiago Benavides,
463	and Paul Corlies for helpful discussions. Any use of trade, firm, or product names is for
464	descriptive purposes only and does not imply endorsement by the U.S. Government.
465	
466	Funding:
467	National Science Foundation Graduate Research Fellowship grant 1745302 (RVP)
468	NASA Cassini Data Analysis Program grants 80NSSC18K1057 and 80NSSC20K0484
469	(RVP, JTP, ADA, JMS, SPDB, AGH).
470	United States Geological Survey, Coastal and Marine Hazards Research Program (RVP)
471	Heising-Simons Foundation (SPDB)
472	
473	Author contributions:
474	Conceptualization: RVP, JTP, ADA, JMS, SPDB, AGH
475	Methodology: RVP, JTP, ADA, JMS
476	Investigation: RVP, JTP, ADA
477	Visualization: RVP
478	Supervision: JTP, ADA, AGH
479	Writing—original draft: RVP
480	Writing—review & editing: RVP, JTP, ADA, JMS, SPDB, AGH
481	
482	Competing interests: Authors declare that they have no competing interests.
483	
484	Code/Data availability: NEWTS1.0 model (Palermo et al., 2023) code is available at
485	https://doi.org/10.5066/P9Q6GDGP.
486	
487	
488	6 References
489	Adams, P.N., 2004, Assessing coastal wave energy and the geomorphic evolution of rocky coasts
490	[Ph.D. thesis]: Santa Cruz, California, University of California–Santa Cruz, 175 p.
491	Adams, P.N., Anderson, R.S., and Revenaugh, J., 2002. Microseismic measurement of wave
492	energy delivery to a rocky coast: Geology, v. 30, p. 895–898, doi:10.1130/0091-
493	7613(2002)030 <0895:MMOWED>2.0.CO;2.
494	Adams, P.N., Storlazzi, C.D., Anderson, R.S., 2005. Nearshore wave-induced cyclical flexing of
495	sea cliffs. Journal of Geophysical Research Earth Surface 110, 1–19,

- 496 https://doi.org/10.1029/2004JF000217.
- 497 Ashton, A. D., Murray, A. B., Littlewood, R., Lewis, D. A., & Hong, P., 2009. Fetch-limited
 498 self-organization of elongate water bodies. Geology, 37(2), 187-190.





499	Bramante, J. F., Perron, J. T., Ashton, A. D., and Donnelly, J. P., 2020. Experimental
500	quantification of bedrock abrasion under oscillatory flow, Geology, 48, 541–545,
501	https://doi.org/10.1130/G47089.1.
502	Emery, K. O., and Kuhn, G. G., 1980. Erosion of rock coasts at La Jolla, California. Marine
503	Geology, 37, 197–208.
504	Günther, F., Overduin, P.P., Sandakov, A.V., Grosse, G., Grigoriev, M.N., 2013. Short and long-
505	term thermo-erosion of ice-rich permafrost coasts in the Laptev Sea region.
506	Biogeosciences 10, 4297–4318.
507	Hasselmann, K., Barnett, T.P., Bouws, E., Carlson, H., Cartwright, D.E., Enke, K., Ewing, J.A.,
508	Gienapp, A., Hasselmann, D.E., Kruseman, P. and Meerburg, A., 1973. Measurements of
509	wind-wave growth and swell decay during the Joint North Sea Wave Project
510	(JONSWAP). Ergaenzungsheft zur Deutschen Hydrographischen Zeitschrift, Reihe A.
511	Howard A. D. (1995) Simulation modeling and statistical classification of escarpment planforms.
512	Geomorphology 12.3, 187–214, 61–78.
513	Huppert, K. L., Perron, J. T., & Ashton, A. D., 2020. The influence of wave power on bedrock
514	sea-cliff erosion in the Hawaiian Islands. Geology, 48(5), 499-503.
515	https://doi.org/10.1130/G47113.1
516	Kline, S.W., Adams, P.N., Limber, P.W., 2014. The unsteady nature of sea cliff retreat due to
517	mechanical abrasion, failure and comminution feedbacks. Geomorphology 219, 53-67,
518	https://doi.org/10.1016/j.geomorph.2014.03.037.
519	Komar P. D., 1998. Prentice-Hall, Englewood Cliffs, New Jersey, 429 pp.
520	Lamont-Smith T, Waseda T., 2008. Wind Wave Growth at Short Fetch. Journal of Physical
521	Oceanography, 38(7), 1597-1606. doi:10.1175/2007JPO3712.1
522	Limber, P.W., Murray, A.B., Adams, P.N., Goldstein, E.B., 2014. Unraveling the dynamics that
523	scale cross-shore headland relief on rocky coastlines: 1. Model development. Journal of
524	Geophysical Research Earth Surface 119, 854–873,
525	https://doi.org/10.1002/2013jf002950.
526	Limber, P.W., Murray, A.B., 2011. Beach and sea-cliff dynamics as a driver of long-term rocky
527	coastline evolution and stability. Geology 39, 1147–1150,
528	https://doi.org/10.1130/g32315.1.
529	Palermo, R. V., Piliouras, A., Swanson, T. E., Ashton, A. D., & Mohrig, D., 2021. The effects of
530	storms and a transient sandy veneer on the interannual planform evolution of a low-relief
531 532	coastal cliff and shore platform at Sargent Beach, Texas, USA. Earth Surface Dynamics,
533	9(5), 1111-1123. Palermo, R.V., Perron, J.T., Soderblom, J.M., Birch, S.P.D., Hayes, A.G., Ashton, A.D., 2023,
534	Numerical model of coastal Erosion by Waves and Transgressive Scarps (NEWTS)
535	Version 1.0: U.S. Geological Survey software release,
536	https://doi.org/10.5066/P9Q6GDGP.
537	Perron, J. T., Dietrich, W. E., & Kirchner, J. W., 2008. Controls on the spacing of first-order
538	valleys. Journal of Geophysical Research: Earth Surface, 113(4), 1–21.
539	https://doi.org/10.1029/2007JF000977
540	Perron, J. T., J. W. Kirchner, and W. E. Dietrich, 2009. Formation of evenly spaced ridges and
541	valleys, Nature, 460, 502–505, doi:10.1038/ nature08174.
542	Perron, J. T., P. W. Richardson, K. L. Ferrier, and M. Lapôtre, 2012. The root of branching river
543	networks, Nature, 492, 100–103, doi:10.1038/ nature11672.





- Ranasinghe, R., 2020. On the need for a new generation of coastal change models for the 21st century. Scientific reports, 10(1), p.2010.
- Robinson, L. A., 1977. Marine erosive processes at the cliff foot, Marine Geology, 23, 257–271, https://doi.org/10.1016/0025- 853227(77)90022-6.
- Sunamura, T. (2018). A fundamental equation for describing the rate of bedrock erosion by
 sediment-laden fluid flows in fluvial, coastal, and aeolian environments. Earth Surface
 Processes and Landforms, 43(15), 3022-3041.
- Sunamura, T., 1976. Feedback relationship in wave erosion of laboratory rocky coast, The
 Journal of Geology, 84, 427–437, 115 https://doi.org/10.1086/628209.
- 553 Sunamura, T., 1992. Geomorphology of Rocky Coasts. Wiley, Chichester, UK.
- 554 Trenhaile, A. S., 1987. The Geomorphology of Rock Coasts, Oxford University Press, Oxford.
- 555 Trenhaile, A.S., 2001. Modeling the effect of weathering on the evolution and morphology of
- shore platforms. Journal of Coastal Research, 17, 398–406.
- Trenhaile, A. S., 2002. Rock coasts, with particular emphasis on shore platforms,
 Geomorphology, 48, 7–22, https://doi.org/10.1016/S0169-555X(02)00173-3.
- Trenhaile, A.S., 2011. Cliffs and Rock Coasts. In: Treatise on Estuarine and Coastal Science
 Vol. 3, eds. Flemming, B.W. and Hansom, J.D., Elsevier, p. 171-192
- Trenhaile AS., 2015. Coastal notches: Their morphology, formation, and function. Earth-Science
 Reviews. 150, 285-304. doi:10.1016/j.earscirev.2015.08.003
- Trenhaile, A.S., 2016. Rocky coasts—Their role as depositional environments. Earth-Science
 Reviews. 159, 1–13.
- Walkden, M. J. A. and Hall, J. W., 2005. A predictive Mesoscale model of the erosion and profile development of soft rock shores, Coastal Engineering, 52, 535–563, 20
 https://doi.org/10.1016/j.coastaleng.2005.02.005.
- Young, A. P., R. E. Flick, W. C. O'Reilly, D. B. Chadwick, W. C. Crampton, and J. J. Helly,
 2014. Estimating cliff retreat in southern California considering sea level rise using a
- 570 sand balance approach, Marine Geology, 348,15–26,
- 571 https://doi.org/10.1016/j.margeo.2013.11.007.