



# Reconstructing tephra fall deposits via ensemble-based data assimilation techniques

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Abstract. In recent years, there has been a growing interest in ensemble approaches for modelling volcanic plumes. The development of such techniques enables the exploration of novel methods for incorporating real observations into tephra dispersal models. However, traditional data assimilation algorithms, including ensemble Kalman filter methods, can yield suboptimal state estimates for positive-definite variables such as volcanic aerosols and tephra deposits. This study proposes two new ensemble-based data assimilation techniques for semi-positive-definite variables with highly skewed uncertainty distributions, including aerosol concentrations and tephra deposit mass loading. The proposed methods are applied to reconstruct the tephra fallout deposit resulting from the 2015 Calbuco eruption using an ensemble of 256 runs performed with the FALL3D dispersal model. Two datasets of deposit thickness measurements are considered: an assimilation dataset including 161 observations, and a validation dataset for an independent assessment of the methods. Results show that the assimilation leads to a significant improvement over the first-guess results, obtained from the simple ensemble forecast. The spatial distribution of the tephra fallout deposit thickness and the ashfall volume according to the analyses are in good agreement with estimations based on field measurements and isopach maps reported in previous studies. Both assimilation methods show a similar performance in terms of evaluation metrics and spatial distribution of the deposit. Finally, the potential application of the methodologies for the improvement of ash-cloud forecasts produced for operational models is also discussed.

## 5 1 Introduction

Multiple hazards are associated with volcanic eruptions including lava flows, pyroclastic density currents, lahars, volcanic plumes, and tephra fallout. Specifically, the dispersal of volcanic plumes poses a serious threat to flight safety (e.g. Clarkson et al., 2016) and the subsequent fallout of tephra can cause structural damage to buildings and infrastructures due to excessive loading, disrupt communication networks, airports, power plants, and water and energy distribution networks (Wilson et al., 2014). Additionally, fresh fallout deposits may be resuspended by aeolian processes affecting the air quality and prolonging the impacts of an eruption many years afterwards (Folch et al., 2014; Mingari et al., 2020). The characterisation and quantification of past eruptive events is also of paramount importance for volcano hazard and risk assessment studies, which infer the likelihood of future eruption scenarios based on the past volcano behaviour.

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Past explosive volcanic eruptions are often characterised and classified by means of tephra deposits (Bonadonna et al., 2015) and provide critical information to infer Eruption Source Parameters (ESP) relevant to hazards, such as eruption column height, mass eruption rate, or total erupted volume (Martí et al., 2016; Constantinescu et al., 2022), especially for poorly monitored volcanoes. Traditionally, volcanologists rely on simple field-based models to obtain certain ESP (e.g. erupted volume) assuming an exponential decay with distance for some deposit parameters such as deposit thickness (Pyle, 1989; Bonadonna and Costa, 2013). However, it is well recognised that this simplistic approach is inappropriate for tephra fall deposits with complex distribution patterns (e.g. Bonadonna et al., 1998; Martí et al., 2016). In fact, many deposits exhibit abrupt thickness variations over short distances, display well-developed secondary maxima, grain size bimodalidity (Durant et al., 2009), stratified deposit with alternating layer characteristics, and include other complexities that make the reconstruction of tephra fallout deposits challenging (Scasso et al., 1994). In contrast, physics-based approaches built upon Volcanic Ash Transport and Dispersal (VATD) models include multiple physical parameterisations and are a much more powerful tool for representing the distribution of tephra deposits. However, the accuracy of deterministic models is highly sensitive to uncertain model input parameters (e.g. eruption column height or physical properties of particles) and the underlying meteorological fields. Alternatively, probabilistic modelling approaches provide a framework to incorporate uncertainties associated with model input data. Specifically, ensemble-based modelling strategies allow one to characterise and quantify model uncertainties and have been proven to enhance VATD model skills (Bonadonna et al., 2012; Madankan et al., 2014; Stefanescu et al., 2014). For example, several VATD models have been used to conduct ensemble simulations, including ASH3D (Denlinger et al., 2012), COSMO-ART (Vogel et al., 2014), HYSPLIT (Dare et al., 2016; Zidikheri et al., 2018), NAME (Dacre and Harvey, 2018; Beckett et al., 2020), and FALL3D (Sandri et al., 2016; Folch et al., 2022; Martinez et al., 2022). Furthermore, different inversion modelling techniques based on ensemble approaches have shown to produce improved volcanic ash forecasts consistent with observations by constraining ash emission estimates and model parameters (Pelley et al., 2015; Zidikheri et al., 2017; Harvey et al., 2020).

The incorporation of ensemble capabilities in VATD models set the foundations for developing and implementing ensemble-based data assimilation and inversion techniques (see Folch and Mingari, 2022, for a recent detailed review). Two main approaches have been explored in the literature to assimilate volcanic aerosol observations from satellites: ensemble Kalman filters (Fu et al., 2016, 2017; Osores et al., 2020; Pardini et al., 2020; Mingari et al., 2022) and ensemble particle filter methods (Zidikheri and Lucas, 2021a, b). Specifically, ensemble Kalman filter (EnKF) methods, used for sequential data assimilation, are based on the Kalman filter (Kalman, 1960) and approximate the probability distributions by an ensemble of system states and assume that the prior model errors and the observation noise are Gaussian. However, lower-bounded variables such as water-vapour mixing ratio (Kliewer et al., 2016), rainfall (Husak et al., 2007) and aerosol concentrations (O'Neill et al., 2000) frequently have skewed and near-zero distributions and are not well described by Gaussian distributions. As a result, traditional EnKF methods in VATD models often yield suboptimal state estimates (Folch and Mingari, 2022).

This study explores two new ensemble-based data assimilation techniques for positive-definite variables and its implementation in VATD models, the Gaussian with Nonnegative Constraints (GNC) method and the Gamma, Inverse-Gamma, and Gaussian Ensemble Kalman Filter (GIGG-EnKF), a sequential method proposed by Bishop (2016) for highly skewed nonnegative distributions. Posselt and Bishop (2018) applied this approach for the nonlinear data assimilation of precipitation rate





observations and compared the results with the analysis produced by a classical EnKF algorithm. It was concluded that the analysis ensemble of precipitation rates produced by the GIGG-EnKF bears a closer resemblance with the Bayesian posterior when the distribution is skewed.

This paper aims to reconstruct the tephra fall deposit of the 2015 Calbuco eruption from a scattered set of observations. The rich existing dataset available for this eruption, consisting of deposit samples collected up to 500 km downwind from the volcano, provides an excellent test case to evaluate the proposed methodology. The Gaussian with nonnegative constraints (GNC) method and the Gamma Inverse-Gamma (GIG) method, based on the GIG equation set proposed by Bishop (2016), are used to assimilate deposit thickness data. Both methods are used here to reconstruct a complete map of the tephra-fall deposit from a dataset of uncertain observations and an ensemble of model realisations based on numerical simulations performed with the FALL3D dispersal model. In addition, a technique for emission source inversion based on the GNC method is also presented and discussed. As an initial step, this manuscript is focused on the assimilation of tephra deposits, which is crucial for long-term tephra hazard assessment, leaving to future studies the assimilation of volcanic clouds and the potential use of these two methods in operational ash forecast contexts.

The manuscript is organised as follows. The ensemble-based data assimilation methods are introduced in Sect. 2. A brief description of the 2015 Calbuco eruption is outlined in Sect. 3 where details about the observational datasets are given. Subsequently, Sect. 3 describes the numerical experiments and shows the results obtained by both methods. In Sect. 4, the GNC method is used to invert the Calbuco source term. Section 5 dwells on potential implications of the proposed methodology and possible future applications and limitations are further discussed. Finally, conclusions are drawn in Sect. 6.

#### 2 Methodology

Data assimilation (DA) techniques have been widely used to study and forecast geophysical systems and have been applied in a variety of research and operational settings (Carrassi et al., 2018). Commonly, DA methods aim at obtaining an optimal state of a dynamical system by combining model forecasts with observations. Specifically, ensemble-based methods comprise two consecutive steps. In the forecast step, a prior ensemble is estimated by running an ensemble of m forward models until the observation time. In the analysis step, observational information is included to obtain a posterior estimate.

Assume that the state of the physical system is represented by a model state vector  $\boldsymbol{x} \in \mathbb{R}^n$ , where n is the system dimension, and that the observations are given by a vector  $\boldsymbol{y}_o \in \mathbb{R}^p$ , where p is the number of observations (see Table 1 for the list of symbols). Suppose that exists a linear observation operator  $\mathbf{H} \in \mathbb{R}^{p \times n}$  which translates the model state  $\boldsymbol{x}$  into the observation space:

$$y = \mathbf{H}x \tag{1}$$

where y represent a p-dimensional vector. In our particular case, the model state vector x is constructed from the two-dimensional tephra deposit load (in  $kg m^{-2}$ ) and the components of x represent the mass load at the n grid points of the computational domain. On the other hand, the observations vector  $y_o$  represents a list of scattered deposit thickness observations (in cm). The deposit mass load and the deposit thickness are related by a proportionality factor, i.e. the bulk density of





the tephra deposit, assumed, for the sake of simplicity for our application, to be constant and equal to  $800 \text{ kg m}^{-3}$ . In addition, the translation of x into the observations space requires horizontal bi-linear interpolations for each measurement site.

In a probabilistic framework, the Probability Density Function (PDF) of the state y conditioned to the observation  $y_o$  is relevant to the assimilation techniques, and can be computed via the Bayes's theorem (Jazwinski, 1970):

$$P(\mathbf{y}|\mathbf{y}_o) = \frac{P(\mathbf{y}_o|\mathbf{y})P(\mathbf{y})}{P(\mathbf{y}_o)}$$
(2)

where  $P(y|y_o)$  is the a posterior PDF, P(y) and  $P(y_o)$  are the prior PDFs, and  $P(y_o|y)$  (also known as the observation likelihood or observational PDF) is the probability of observing  $y_o$  given a true y.

In consequence, the determination of the posterior PDF requires the specification of both the prior and the observational PDFs. This paper proposes two ensemble-based assimilation strategies which rely on (2) and differ on the assumptions made about these PDFs. The GNC method (Sect. 2.1) uses an all-at-once assimilation approach looking for the model state that maximises the vectorial form of (2) and observations are assimilated all at once. In contrast, the GIG method (Sect. 2.2) uses a serial assimilation approach in which the univariate version of (2) is explicitly written for each single observation and the full dataset of *p* observations is assimilated in a sequential way.

## 105 2.1 The GNC method

The Gaussian with Nonnegative Constraints (GNC) method assumes a multi-dimensional Gaussian probability distribution for y, defined in (1), given as:

$$P(\boldsymbol{y}) \propto \exp\left\{-\frac{1}{2}(\boldsymbol{y} - \overline{\boldsymbol{y}})^{\mathsf{T}} \mathbf{P}^{-1} (\boldsymbol{y} - \overline{\boldsymbol{y}})\right\}$$
(3)

where  $\mathbf{P} \in \mathbb{R}^{p \times p}$  is the error covariance matrix and  $\overline{y} \in \mathbb{R}^p$  is the average model state vector in the observation space. Similarly, for the sake of simplicity, measurements are assumed to be normally distributed with observation error covariance matrix  $\mathbf{R} \in \mathbb{R}^{p \times p}$ , i.e.

$$P(\boldsymbol{y}_o|\boldsymbol{y}) \propto \exp\left\{-\frac{1}{2}(\boldsymbol{y}_o - \boldsymbol{y})^{\mathsf{T}} \mathbf{R}^{-1}(\boldsymbol{y}_o - \boldsymbol{y})\right\}$$
(4)

The most likely state is the one that maximises the posterior PDF, Eq. (2), or equivalently the one that minimises the GNC cost function J:

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$$J(y) \propto (y - \overline{y})^{\mathsf{T}} \mathbf{P}^{-1} (y - \overline{y}) + (y_o - y)^{\mathsf{T}} \mathbf{R}^{-1} (y_o - y)$$
 (5)

Note that the expression above is actually very similar to the cost function used in classical variational methods (e.g. 3DVAR, Carrassi et al., 2018) with the difference that  $\overline{y}$  plays the role of the model background state in the VAR methods, and the first term in (5) is computed in the observations space rather than in the model space as usual. This is justified because expressing the functional J in the observations space is advantageous in those cases where observations are localised and/or nearly zero, i.e. circumscribed to portion of the computational domain (this is what typically occurs for volcanic clouds and fallout





Table 1. List of symbols using the following convention: matrices in upper case bold, vectors in lower case bold, scalars in italics.

Symbol	Description				
	General definitions				
m	ensemble size				
n	dimension of model state vector				
p	number of observations				
$oldsymbol{x} \in \mathbb{R}^n$	Model state vector				
$oldsymbol{y_o} \in \mathbb{R}^p$	Observations vector				
$oldsymbol{y} \in \mathbb{R}^p$	Model state vector in the observation space				
$\mathbf{H} \in \mathbb{R}^{p \times n}$	Observation operator				
	GNC method				
$\mathbf{P} \in \mathbb{R}^{p \times p}$	Model covariance matrix				
$\mathbf{R} \in \mathbb{R}^{p \times p}$	Observation error covariance matrix				
$oldsymbol{w} \in \mathbb{R}^m$	Vector of weight factors				
$\overline{oldsymbol{y}} \in \mathbb{R}^p$	Average model state vector (obs. space)				
$\mathbf{Y} \in \mathbb{R}^{p \times m}$	Ensemble model state matrix (obs. space)				
$\mathbf{Y}' \in \mathbb{R}^{p \times m}$	Ensemble perturbations matrix (obs. space)				
	GIG method				
$y_j$	$j$ -th component of $oldsymbol{y}$				
$y^o_j$	$j$ -th component of $oldsymbol{y}_o$				
$\overline{y}_{j}^{f}$	Mean of prior distribution of $y_j$				
$\overline{y}_{j}^{a}$	Mean of analysis distribution of $y_j$ (14a)				
$P_j^r$	Type 1 relative error variance of prior				
	$P_j^r := \operatorname{var}(y_j^f) / (\overline{y}_j^f)^2$				
$\Pi_j^r$	Type 1 relative error variance of analysis (14b)				
	$\Pi_j^r := \operatorname{var}(y_j^a) / (\overline{y}_j^a)^2$				
$R_j^r$	Type 1 relative error variance of observation <sup>†</sup>				
	$R_j^r := \operatorname{var}(y_j^o) / (y_j^t)^2$				
$\tilde{P}_i^r$	Type 2 relative error variance of prior				
J					
J	$(\tilde{P}_j^r)^{-1} = (P_j^r)^{-1} + 1$				
$ ilde{R}_{j}^{r}$	$(P_j^r)^{-1} = (P_j^r)^{-1} + 1$ Type 2 relative error variance of observation				

 $<sup>^{\</sup>dagger}$  where  $\boldsymbol{y}_{j}^{t}$  is the true value of the j-th observation.





deposits). Moreover, the functional in (5) yields to a much reduced system when compared to classical VAR methods because the observation space normally has a much lower dimension (p << n).

In the context of an ensemble modelling formulation, the forecast step provides an ensemble of m independent prior states vectors  $\mathbf{x}_i$  representing m model realisations at the analysis time. The GNC method looks for the best linear estimate of the system state in the subspace spanned by the ensemble of vectors  $\mathbf{x}_i$ :

$$\boldsymbol{x} = w_1 \boldsymbol{x}_1 + \dots + w_m \boldsymbol{x}_m \tag{6}$$

where  $w_i \ge 0$  (i=1...m) is a set of nonnegative weight factors for each ensemble member. The important point here is that the nonnegative constraints on  $w_i$  relax the Gaussian hypotheses and avoid the occurrence of non-physical solutions. The linearity of the observation operator  $\mathbf{H}$  allows the analysis to be expressed in the observation space as:

$$\mathbf{130} \quad \mathbf{y} = w_1 \mathbf{y}_1 + \dots + w_m \mathbf{y}_m = \mathbf{Y} \mathbf{w} \tag{7}$$

where  $y_i = \mathbf{H}x_i$  and the matrix  $\mathbf{Y} \in \mathbb{R}^{p \times m}$  is defined by  $[y_1, \dots, y_m]$ . The ensemble mean is used to approximate the average model state vector  $\overline{y} \in \mathbb{R}^p$ , i.e.:

$$\overline{\boldsymbol{y}} \approx \frac{1}{m} \sum_{i=1}^{m} \boldsymbol{y}_{i} \tag{8}$$

whereas the ensemble-based error covariance matrix is used to approximate P according to:

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$$\mathbf{P} \approx \frac{1}{m-1} \mathbf{Y}' \mathbf{Y}'^{\mathsf{T}}$$
 (9)

where the matrix of ensemble perturbations  $\mathbf{Y}' \in \mathbb{R}^{p \times m}$  is given by  $[y_1 - \overline{y}, \dots, y_m - \overline{y}]$ . Replacing (7), (8) and (9) in (5), the GNC cost function J can be expressed as the equivalent quadratic form:

$$J(\boldsymbol{w}) = \frac{1}{2} \boldsymbol{w}^{\mathsf{T}} \mathbf{Q} \boldsymbol{w} + \boldsymbol{b}^{\mathsf{T}} \boldsymbol{w} + \dots$$
 (10)

with

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$$\mathbf{Q} = \mathbf{Y}^{\mathsf{T}}(\mathbf{P}^{-1} + \mathbf{R}^{-1})\mathbf{Y}$$
 (11a)

$$\boldsymbol{b} = -\mathbf{Y}^{\mathsf{T}}(\mathbf{P}^{-1}\overline{\boldsymbol{y}} + \mathbf{R}^{-1}\boldsymbol{y}_{o}) \tag{11b}$$

In order to find the optimal vector of weight factors w, the optimisation problem  $\min_{w\geq 0} J(w)$  must be solved. Then, the analysis vector state  $x^a$  can be computed by replacing the optimal w in (6). The minimisation of the quadratic form (10) subject to the constraints  $w_i \geq 0 \ \forall i$  is a nonnegative quadratic programming problem and there is no analytical solution for the global minimum due to the nonnegativity constraint. However, it can be solved using the iterative approach proposed by Sha et al. (2007):

$$w_i \leftarrow w_i \left[ \frac{-b_i + \sqrt{b_i^2 + a_i c_i}}{a_i} \right] \tag{12}$$



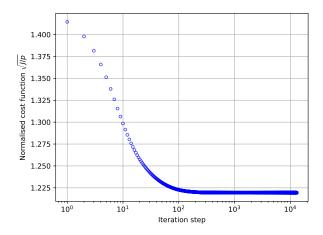


Figure 1. Iterative approach to minimise the GNC cost function J subject to the non-negativity constraints. Under the multiplicative updates in (12) the cost function decreases monotonically. In this particular example, the convergence required more than  $10^4$  iterative steps.

as long as the matrix  $\mathbf{Q}$  is symmetric and semipositive definite, as can be easily verified from (11a). Under the multiplicative updates (12), the cost function decreases monotonically to the value of its global minimum as shown by Sha et al. (2007). The vectors  $\mathbf{a} = \mathbf{A}_+ \mathbf{w}$  and  $\mathbf{c} = \mathbf{A}_- \mathbf{w}$  must be updated in each iterative step, where  $\mathbf{A}_+ = |\mathbf{Q}| + \mathbf{Q}$  and  $\mathbf{A}_- = |\mathbf{Q}| - \mathbf{Q}$ . For illustrative purposes, Fig. 1 shows the decrease of the normalised cost function, defined as  $\sqrt{J/p}$ , under the multiplicative updates in (12) for the case study presented later in Sect. 3. More than  $10^4$  iterative steps were required to get low enough residuals to satisfy the convergence criteria.

## 2.2 The GIG method

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Bishop (2016) introduced a variation of the Ensemble Kalman Filter (EnKF) that solves the univariate Bayes' theorem for non-negative variables with skewed (asymmetrical) probability distributions. The so-called GIGG-EnKF (with GIGG standing for Gamma, Inverse-Gamma and Gaussian) allows non-negative variables typically involving near-zero values (i.e. with right-skewed probability distributions), such as aerosol, water vapour, cloud, and precipitation concentrations, to be directly assimilated, thus avoiding the use of Gaussian anamorphosis nonlinear transformations (e.g. Amezcua and Leeuwen, 2014). The GIGG-EnKF algorithm is based on the generalised two-stage multivariate ensemble filter described by Anderson (2003). The first stage involves the univariate GIGG-EnKF in which an ensemble-based estimate of the posterior distribution of the observed variable is generated from a single observation and a prior ensemble of state estimates. In the second stage, the univariate method is extended by propagating this information to the complete model state vector using a linear regression approximation.

According to the strategy proposed by Bishop (2016), in the first step, the Bayes' theorem is solved in the univariate form of (2) assuming a distribution pair for the prior probability and the likelihood PDF of the error-prone observations given truth, respectively. A single observation is assimilated using an appropriate equation set depending on three different cases: GIG,



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IGG and G. This paper will only consider the GIG equation set, aimed at situations in which the prior can be described by a Gamma distribution and the observation likelihood can be represented by an Inverse-Gamma distribution. In addition, Bishop (2016) introduced the IGG equation set (Inverse-Gamma prior and Gamma observation likelihood) and the G equation set (Gaussian distributions).

In the GIG case, the posterior probability is given by a gamma PDF:

$$P(y_j|y_j^o) \propto y_j^{(\Pi_j^r)^{-1}-1} \exp\left\{-\frac{y_j}{\Pi_j^r \ \overline{y}_j^a}\right\}$$

$$\tag{13}$$

where  $y_j$  and  $y_j^o$  are the j-th components of the vectors  $\boldsymbol{y}$  and  $\boldsymbol{y}_o$  respectively. The posterior univariate gamma PDF is characterised by two parameters, namely, the analysis mean  $\overline{y}_j^a$  and the type 1 relative error variance of the analysis  $\Pi_j^r := \text{var}(y_j^a)/(\overline{y}_j^a)^2$ , that in the GIG method are given by:

$$\frac{1}{\overline{y}_j^a} = \frac{1}{\overline{y}_j^f} + \frac{\tilde{P}_j^r}{\tilde{R}_j^r + \tilde{P}_j^r} \left\{ \frac{1}{y_j^o} - (\tilde{R}_j^r + 1) \frac{1}{\overline{y}_j^f} \right\}$$
(14a)

$$(\Pi_i^r)^{-1} = (\tilde{R}_i^r)^{-1} + (\tilde{P}_i^r)^{-1} \tag{14b}$$

where  $\tilde{P}_{j}^{r}$  and  $\tilde{R}_{j}^{r}$  are the type 2 relative error variance of the prior and observations respectively (see Table 1 for details). Bishop (2016) provides a stochastic equation to generate individual members of the analysis ensemble (i.e.  $y_{ji}^{a}$ ) with the low-order moments of the posterior distribution being consistent with (14). See Appendix A and Eq. (A1) for further details. In addition, the ensemble generated in this way is ensured to converge to the true posterior PDF for large ensembles.

In the second step, the univariate case is extended according to the second-stage linear regression step proposed by Anderson (2003) in order to find the analysis ensemble for the complete model state vector. The update of the k-th model state vector variable of the i-th ensemble member due to the j-th observation is computed according to:

$$x_{ki}^{a} = x_{ki}^{f} + \delta_{kj} \frac{\operatorname{cov}(x_{k}^{f}, y_{j}^{f})}{\operatorname{var}(y_{j}^{f})} (y_{ji}^{a} - y_{ji}^{f})$$
(15)

where  $\delta_{kj}$  is a localisation weight satisfying  $0 \le \delta_{kj} \le 1$ . In the original formulation, Bishop (2016) assumed  $\delta_{kj} = 1$ . However, we found that (15) can introduce artificial results when a poor correlation exist between the  $x_k$  and  $y_j$  variables, including non-physical values (e.g. negative mass load). By introducing the localisation weight  $\delta_{kj}$  in this work, the effect of poorly correlated observations is intended to be limited. Instead of the traditional distance-dependent localisation, we use here the correlation-cutoff method to localise the analysis (Yoshida and Kalnay, 2018). This method localises the observation impacts based on their background error correlations and only cross covariances between variables that have strong background error correlation are considered in (15). In this work, a quadratic function has been chosen for defining the cutoff function (see Chang and Kalnay, 2022, for a similar approach). Specifically, the localisation weight assigned for the j-th observation at the k-th analysis grid cell was defined as:

$$\delta_{kj} = \begin{cases} 1 & \text{if } |\rho_{kj}| \ge \rho_o \\ (\rho_{kj}/\rho_o)^2 & \text{if } |\rho_{kj}| < \rho_o \end{cases}$$
 (16)



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where  $\rho_{kj} = \operatorname{corr}(x_k^f, y_j^f)$  is the prior ensemble correlation and  $\rho_o$  is a tunable parameter controlling the intensity of the localisation. Localisation is enabled when  $0 < \rho_o < 1$  and  $\rho_o = 0$  implies no localisation at all, as in the original formulation.

The inverse-gamma PDFs assign non-zero probability densities only for positive observations and, as a result, zero observations cannot be properly assimilated using the GIG equation set (e.g. see 14a). This problem is addressed here by redefining zero observation data according to:

$$y_j^o \leftarrow r * \epsilon_{min} \tag{17}$$

where  $r \in (0,1]$  is a random number and  $\epsilon_{min}$  is the minimum error expected for the deposit thickness observations (assumed to be  $\epsilon_{min} = 1$  mm in this work where we deal only with visible tephra deposits).

The GIG method is a sequential procedure: a single observation is assimilated in order to update the prior ensemble forecast using the GIG equation set; subsequently, this procedure is repeated until all observations have been sequentially assimilated. In contrast, the GNC method described in Sec. 2.1 represents an all-at-once assimilation technique. To summarise, a pseudocode of the sequential procedure used to implement the GIG method is detailed in the Algorithm 1.

**Algorithm 1** Pseudocode of the GIG method based on the Bishop (2016) algorithm for the case in which the prior is a gamma distribution and the observation likelihood is an inverse-gamma distribution.

**Require:** List of observations  $\{y_i^o\}$  with their relative errors

**Ensure:** Analysis ensemble  $x_{ki}^a$ 

```
1: get x_{ki}^f

⊳ Generate the prior ensemble

2: procedure GIG METHOD
3:
         randomly shuffle observation list
         for all y_i^o do
4:
                                                                                                                                    if y_i^o = 0 then
5:
                 r \leftarrow \text{random number} \in (0, 1]
6:
7:
                 y_j^o \leftarrow r * \epsilon_{min}
                                                                                                                                 ▶ Redefine zero observations
             end if
8:
             \overline{y}_i^a \leftarrow \text{Eq. } (14a)
9:
10:
             y_{ii}^a \leftarrow \text{Eq.} (A1)
             x_{ki}^a \leftarrow \text{Eq.} (15)
11:
                                                                                                                            12:
             x_{ki}^f \leftarrow x_{ki}^a
                                                                                                                                  ▶ Update the prior ensemble
13:
         end for
14: end procedure
```





Table 2. Calbuco deposit datasets considered in this study.

Reference	Data type	Purpose	
Van Eaton et al. (2016)	Thickness at	Data assimilation	
	163 locations		
Reckziegel (unpublished)	Thickness at	Data validation	
	45 locations		
Romero et al. (2016)	Isopachs map	Data validation	

# 3 Reconstruction of the 2015 Calbuco deposit

In this section, the procedures described in Sect. 2 are applied to the 2015 Calbuco eruption in order to obtain the analysed deposit thickness. With this in mind, the field measurements reported by Van Eaton et al. (2016) will be considered for assimilation purposes.

## 3.1 Fallout deposit and datasets

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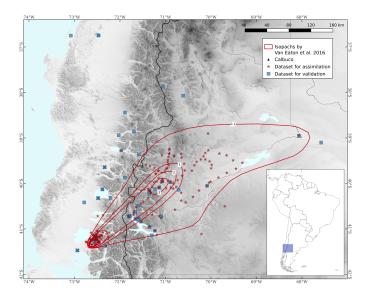
The 2015 eruption of the Calbuco stratovolcano (41.33°S, 72.61°W) in Southern Chile involved two major eruptive pulses on 22–23 April along with a third minor pulse on 30 April (Romero et al., 2016). During the most energetic phase on 23 April, stratospheric eruption columns higher than 15 km above the vent level (~17 km above sea level) were developed. Regions over Southern Chile and the Argentinian Patagonia were severely affected by tephra fall. According to different estimations based on field studies, deposit volume ranges between 0.27 and 0.58 km³ (Romero et al., 2016; Van Eaton et al., 2016).

The availability of independent and comprehensive datasets of field observations makes the Calbuco tephra deposit an excellent case study. Van Eaton et al. (2016) reported the thickness and stratigraphy of the fall deposits at 163 sampling sites within a 500-km radius from the volcano summit. This dataset is considered in this work for assimilation purposes. Figure 2 shows the location of sampling sites and the isopachs of fall deposit thickness (in centimetres) reported by Van Eaton et al. (2016). It is interesting to note the presence of a secondary thickness maximum ~200 km downwind from the vent, located around two major cities of the Argentinian Patagonia: Junín de los Andes (39.95°S, 71.07°W) and San Martín de los Andes (40.16°S, 71.35°W), likely due to ash aggregation processes (e.g. Costa et al., 2010).

A complementary dataset composed of 45 independent measurements of deposit thickness (F. Reckziegel, pers. comm.) is also considered for validation purposes using the evaluation metrics described in the following section. Figure 2 shows the location of the corresponding sampling sites (square symbols). Finally, a hand-drawn isopachs map built from a third independent dataset (Romero et al., 2016) is also used to evaluate the tephra deposit distribution considering the isopachs for 0.1, 0.5, 1.0 and 2.0 mm. The three datasets are summarised in Table 2. The assimilation methods require a dataset of measurements along with the corresponding absolute or relative errors. Specifically, the GNC method requires the absolute







**Figure 2.** Isopachs of fall deposit thickness in centimetres and location of the 161 sampling sites (two ambiguous measurements were discarded) of the assimilation dataset reported by Van Eaton et al. (2016) (red stars). The map also shows the locations of an independent dataset composed by 45 measurements (blue squares) used for validation purposes.

error  $\epsilon_j$  associated with the j-th measurement  $y_j^o$  (the observation error covariance matrix  ${\bf R}$  is assumed diagonal with elements  $\epsilon_j^2$ ). On the other hand, the GIG method requires the relative error  $\epsilon_j^r = \epsilon_j/y_j^t$ , where  $y_j^t$  is the true value of the j-th observation. Observation error standard deviations are assumed to be dependent on the measured value. The procedure for determining observation errors is as follows: the full dataset of observations is classified in 7 groups using a spectral clustering algorithm (Pedregosa et al., 2011), as shown in Fig. 3. Three feature were used for clustering: latitude, longitude and deposit thickness of every measurement. The error for the j-th measurement  $y_j^o$  is approximated by standard deviation associated with the corresponding cluster; the true value  $y_j^t$ , required to estimate relative errors, is approximated by the cluster mean value.

#### 3.2 Validation metrics

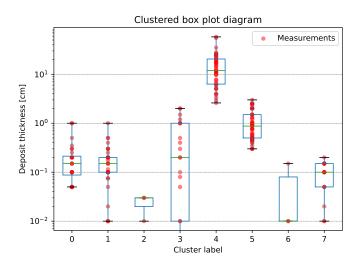
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As validation metrics, we consider the Mean Bias Error (MBE) and the Root-Mean-Square Error (RMSE), defined as usual:

$$MBE = \frac{1}{p} \sum_{j=1}^{p} y_j^o - y_j$$
 (18a)

$$RMSE = \sqrt{\frac{1}{p} \sum_{j=1}^{p} (y_j^o - y_j)^2}$$
 (18b)

The MBE (in cm) quantifies the tendency to overestimate or underestimate observations for the overall dataset whereas the RMSE (in cm) measures the average magnitude of the errors. These two metrics are suitable when a uniform distribution of the errors is expected. However, the datasets in this work contain measurements spanning four orders of magnitude and the



**Figure 3.** Observational data are classified in 7 groups using a spectral clustering algorithm taking into account three features: latitude, longitude, and deposit thickness of each measurement. The observation absolute error is assumed to be the standard deviation for the corresponding cluster and the relative error is approximated by the ratio of standard deviation to the cluster mean.

assumption of a constant absolute error seems to be inappropriate in this case because only proximal data (i.e. the largest measurements of deposit thickness) contributes significantly to MBE and RMSE. The weighted MBE and RMSE represent alternative evaluation metrics that are more meaningful in our case. These dimensionless metrics are defined according to:

$$\text{wMBE} = \frac{1}{p} \sum_{j=1}^{p} \frac{y_j^o - y_j}{\epsilon_j}$$
 (19a)

250 wRMSE = 
$$\sqrt{\frac{1}{p} \sum_{j=1}^{p} \frac{(y_j^o - y_j)^2}{\epsilon_j^2}}$$
 (19b)

where the deviations from observation  $(y_j^o - y_j)$  are weighted by the observation uncertainties  $(\epsilon_j)$ . The impact of the assimilation can be better characterised by means of these metrics. Ideally, the wMBE should be close to zero and wRMSE should approach one.

## 3.3 Ensemble modelling

Numerical simulations were carried out using the latest version release of FALL3D (v8.2), an open-source offline Eulerian model for atmospheric passive transport and deposition of aerosols and particles, including tephra species. FALL3D solves the so-called Advection-Diffusion-Sedimentation (ADS) equation (Folch et al., 2020; Prata et al., 2021). The new FALL3D version has been designed to support increasingly larger scientific workloads and prepare the code for the transition to extreme-scale computing systems. Specifically, the code version v8.x has been released with several improvements over previous versions, including improvements in the model physics, numerical algorithmic methods, and computational efficiency. In addition, from



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Table 3. FALL3D model configuration parameters for the 2015 Calbuco runs.

Parameter	Value			
Ensemble size	256			
Resolution	$0.05^{\circ} \times 0.05^{\circ}$			
Number of grid points	$180 \times 160 \times 45$			
Species	40 tephra bins			
Grain Size Distribution	bi-Gaussian			
Run time	100 h			
Emission source	Suzuki source (Pfeiffer et al., 2005)			
Mass emission rate	Estimated from column height			
	(Degruyter and Bonadonna, 2012)			

version v8.1 onwards, the FALL3D model enables ensemble simulations to be performed very efficiently by means of a single parallel task (Folch et al., 2022). Ensemble modelling allows one to characterise and quantify model uncertainties due to poorly constrained input parameters and errors in the model physics parameterisations or the underlying model-driving meteorological data. In addition, the ability to generate ensemble runs makes it possible to improve forecasts by incorporating observations using different ensemble-based data assimilation techniques.

The configuration of the FALL3D model used in this work is summarised in Table 3. A three-dimensional computational domain with a horizontal resolution of  $\sim$ 4 km  $(0.05^{\circ})$  and  $180 \times 160 \times 45$  grid points was defined. The Total Grain-Size Distribution (TGSD) of tephra injected into the atmosphere consists of the sum of two log-normal distributions (i.e. bi-Gaussian in  $\Phi$ -units) including 40 tephra bins. The modes and standard deviations of the bimodal distribution were computed using the parameterisation proposed by Costa et al. (2016), which estimates them from the eruption intensity and magma viscosity. The mode of the coarser population was located at -1.2 $\Phi$  with a standard deviation of 1.71 $\Phi$ , while the mode of the finer population was 3.49 $\Phi$  with a standard deviation of 1.46 $\Phi$ . The weight of each subpopulation was set to  $p_c = 0.15$  and  $p_f = 0.85$  for the coarse and fine population, respectively. The vertical mass distribution of the source term depends on the eruptive column top height (H) according to the following parameterisation (Pfeiffer et al., 2005):

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$$\frac{dm}{dz} \propto \left\{ \left( 1 - \frac{z}{H} \right) \exp \left[ A_s \left( \frac{z}{H} - 1 \right) \right] \right\}^{\lambda_s}$$
 (20)

where  $A_s$  and  $\lambda_s$  are the so-called Suzuki parameters (Suzuki, 1983; Pfeiffer et al., 2005). Finally, the mass emission rate was computed from the eruptive column top height using the expression proposed by Degruyter and Bonadonna (2012).

An ensemble size of 256 members is considered in this paper. The prior ensemble was generated by perturbing the Eruption Source Parameters (ESP) and the horizontal wind components around a reference value using either uniform or truncated normal distributions with the range of parameter uncertainties. The Latin Hypercube Sampling (LHS, McKay et al., 1979) was used to efficiently sample the parameter space. Table 4 lists the perturbed model parameters along with the corresponding reference values and sampling uncertainty ranges. The central member of the ensemble was defined considering two eruptive





**Table 4.** Ensemble configuration. The perturbed model parameters are: eruption column height (H), eruption phase start time  $(T_i)$ , phase duration  $(\Delta T)$ , parameters  $A_s$  and  $\lambda_s$  of the Suzuki vertical mass distribution, the fine mode of the bi-Gaussian TGSD, the density of aggregates, and the wind components.

Parameter	Reference value	Distribution	Sampling range
Н	15 & 15 km avl <sup>†</sup>	Uniform	±25%
$T_i$	$21~\&~28~h^\ddagger$	Uniform	15 min
$\Delta T$	1.5 & 6 h	Uniform	15 min
$A_s$	6	Uniform	3
$\lambda_s$	3	Uniform	2
Fine mode	$3.49~\Phi$	Uniform	$2 \Phi$
Density agg.	$450~\rm kgm^{-3}$	Uniform	$100~\rm kgm^{-3}$
U wind	ERA5§	Gaussian	$\pm 20\%$
V wind	ERA5§	Gaussian	$\pm 20\%$

<sup>&</sup>lt;sup>†</sup> Two eruptive phases are considered. Heights are given in km above the vent level;

phases with column top heights at H=15 km avl (above vent level) for each phase. The ensemble was generated by perturbing H for each phase independently and assuming a sampling range of  $\pm 25\%$ . The source start time for the central member was defined at 21:00 UTC on 22 April 2015 (1st phase) and at 04:00 UTC on 23 April 2015 (2nd phase) assuming a duration of 1.6 h and 6 h for each phase. Source start times and phase duration were also perturbed.

## 3.4 Prior ensemble distribution

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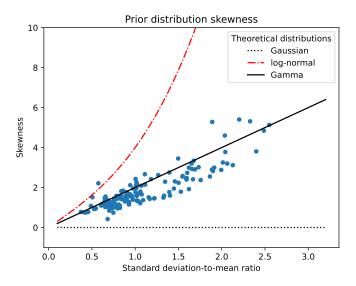
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Before showing how the assimilation methods perform, it is worth to consider the prior (i.e. forecast) ensemble distribution and to check weather the PDF P(y) fulfils the assumption of the GIG method (i.e. a gamma prior distribution). To this purpose, the skewness  $\tilde{\mu}_3$  (i.e. the third standardised moment,  $\mu_3/\sigma^3$ ) of the prior distribution was computed from the random samples  $y_{ji}^f$ , i.e. the forecasted deposit thickness according to the i-th ensemble member ( $i=1\dots m=256$ ) at the sampling site of the j-th observation ( $j=1\dots p=161$ ). Figure 4 shows the results for each observation point by plotting  $\tilde{\mu}_3$  as a function of the standard deviation-to-mean ratio, i.e.  $\sqrt{P_j^r}$ . For illustrative purposes, the skewness of three different theoretical distributions (Gaussian, log-normal and gamma) is shown. As expected, the symmetric Gaussian distribution, characterised by  $\tilde{\mu}_3=0$ , does not reproduce the positively skewed prior distribution. The log-normal family of probability distributions represent an example of distributions for positive-definite variables with a lower bound. However, as shown in Fig. 4, the log-normal distribution cannot properly represent the prior distribution because the theoretical skewness is extremely large in this case. In contrast, the skewness computed from the prior distributions (blue dots) is well approximated by the relationship  $\tilde{\mu}_3=2\sqrt{P_j^r}$ , which is the theoretical expression corresponding to the gamma distribution (solid black line).

<sup>&</sup>lt;sup>‡</sup> In hours since 22 April 2015 at 00:00 UTC;

<sup>§</sup> ECMWF atmospheric reanalysis in model levels.





**Figure 4.** Skewness as a function of the standard deviation-to-mean ratio for the prior distribution at the sampling locations (blue circles). Results for some theoretical distributions (Gaussian, log-normal and gamma) are also shown for comparison.

In order to dig further into the similarities between the gamma and prior distributions, Fig. 5 explicitly shows histograms of sampled prior distributions along with the corresponding theoretical gamma distribution for some observation sites. The theoretical gamma distributions were constructed using the sampled first and second moments. As observed, a good agreement is found between both distributions in almost all cases. Note that when the type 1 relative error variance is greater than 1  $(P_j^r > 1)$  the gamma probability density decreases monotonically (Fig. 5a–b) and the mode becomes zero. In contrast, when  $P_j^r < 1$  or, equivalently, when  $\overline{y}_j^2 > \text{var}(y_j)$ , the mode becomes positive (Fig. 5c–p). The results obtained is this section justify the suitability of the GIG method to deal with the assimilation volcanic deposit data.

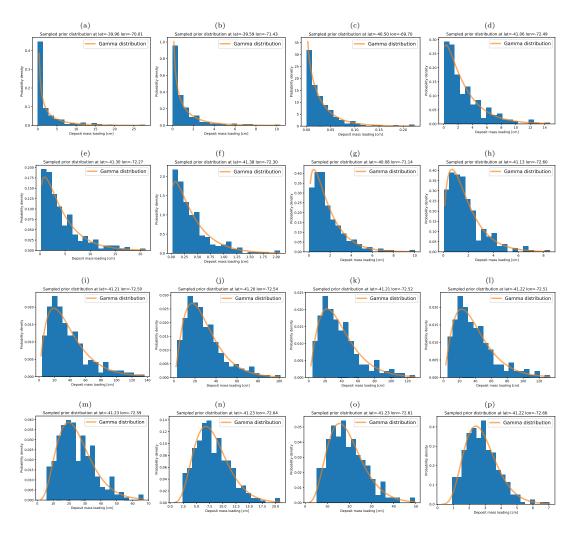
## 3.5 Analyses

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The GNC method gives a set of weight factors for each ensemble member  $(w_i, i = 1, ..., m)$  and the best estimate of the system state is obtained by replacing the optimal weight factors in (6). On the other hand, the GIG method produces an ensemble of analysis states. In this section, the analysis ensemble mean corresponding to the GIG method is used for comparison purposes. Figure 6 compares the analysis results at each sampling site with observations (from the same assimilated dataset) according to the GNC (Fig. 6a) and the GIG (Fig. 6b) methods. Given that the GIG method also provides the mean and standard deviation of the analysis, this information is used in Fig. 6b to represent the analysis error through error bars. Unfortunately, an estimate of the analysis error is not provided by the GNC method and, consequently, error bars cannot be shown in Fig. 6a. Note that most of the data lies in the 1:10 ratio band, with only 4 (2) data points laying outside this band for the GNC (GIG) method despite the challenging fact that measurements span a range of four orders of magnitude between  $10^{-2}$  and  $10^2$  cm, approximately.



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**Figure 5.** Histograms of sampled prior distributions along with the corresponding theoretical gamma distribution at some selected observation sites.

Table 5 reports the different evaluation metrics in order to quantify deviations from observations in the prior ensemble mean (forecast) and the analyses for both GNC and GIG methods. The bias computed using the assimilation dataset for the prior mean was negative (MBE=-4.31 cm), meaning that measurements were overestimated systematically by the first guess. In contrast, the analysis underestimates observations according to GNC method (MBE=1.32 cm) and slightly overestimates according to GIG methos (MBE=-0.30 cm). In both cases, the magnitude of the forecast bias was clearly reduced. The prior RMSE (RMSE=8.53 cm) was also reduced by the GNC (RMSE=4.23 cm) and the GIG (RMSE=2.02 cm) methods. The impact of the assimilation is better characterised by the relative or weighted metrics. The assimilation methods lead to a strong reduction of both wMBE and wRMSE, as reported in Table 5. For example, the weighted RMSE was reduced from wRMSE=3.96 to wRMSE=1.15 (GNC method) and wRMSE=1.31 (GIG method). Note that the analysis are close to the ideal value wRMSE=1.





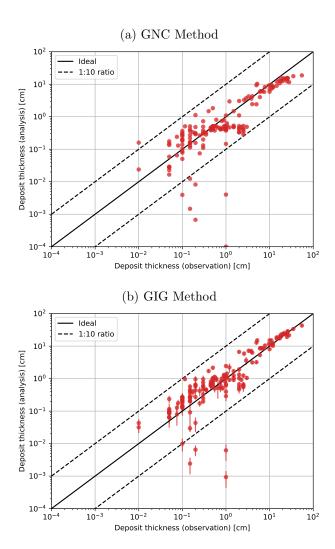


Figure 6. Comparison between analysis and observations from the assimilation dataset (161 points, see Table 2).

For the validation dataset, the metrics for the prior mean yield values similar to those of the assimilation dataset, i.e. wMBE=3.77 and wRMSE=6.47. In terms of the weighted metrics, the GNC method showed the best performance (wMBE=-0.04 and wRMSE=1.81). Regarding the GIG method, the metrics were computed for different values of the parameter  $\rho_o$ , required by (16), i.e.  $\rho_o = 0.0$  (no localisation),  $\rho_o = 0.01$ ,  $\rho_o = 0.05$ ,  $\rho_o = 0.1$  and  $\rho_o = 0.5$ . Since the GIG method is stochastic, five realisations were performed for each  $\rho_o$ . Table 5 reports the metrics corresponding to the best realisation (i.e. the one with the least wRMSE) and the average over the realisations. Taking into account both weighted bias and RMSE, the best performance was obtained for  $\rho_o = 0.0$ , i.e. wMBE=-0.55 and wRMSE=1.83, suggesting that the localisation procedure did not have a significant impact. This is also confirmed by looking at the average wRMSE: the best result is obtained when localisation is disabled (wRMSE=2.32 for  $\rho_o = 0$ ).





**Table 5.** Metrics computed for each method considering the assimilation and validations datasets. The parameter  $\rho_o$  used for the GIG method is specified in parentheses (see Eq. 16).

	MBE	RMSE	wMBE	wRMSE
	(cm)	(cm)		
	Assimilation dataset			
Prior mean	-4.31	8.53	-2.15	3.96
GNC Method	1.32	4.23	0.30	1.15
GIG Method <sup>†</sup>	-0.30	2.02	-0.18	1.31
	Validation dataset			
Prior mean	-2.11	5.39	-3.77	6.47
GNC Method	1.43	6.21	-0.04	1.81
GIG Method $(0.00)^{\dagger}$	0.11	2.51	-0.55	1.83
GIG Method $(0.01)^{\dagger}$	0.22	2.97	-0.64	2.05
GIG Method $(0.05)^{\dagger}$	0.23	3.28	-0.60	1.96
GIG Method $(0.10)^{\dagger}$	0.00	1.90	-0.65	1.94
GIG Method $(0.50)^{\dagger}$	-0.07	2.73	-0.68	2.34
GIG Method $(0.00)^{\S}$	0.17	2.85	-0.72	2.32
GIG Method $(0.01)^{\S}$	0.21	2.98	-0.68	2.36
GIG Method $(0.05)^{\S}$	0.25	3.23	-0.73	2.52
GIG Method (0.10)§	0.11	2.51	-0.78	2.44
GIG Method $(0.50)^{\S}$	0.05	2.83	-0.77	2.55

<sup>†</sup> Best realisation

Figure 7 compares the analyses with measurements from the validation dataset. While most of the points lie within the 1:10 ratio band, an almost zero thickness was found according to the analysis at four sampling sites with positive measurements. These points were excluded from the log-log plot in Fig. 7 in order to facilitate the visualisation, but they were included in the calculation of the evaluation metrics reported in Table 5. The sampling sites for the excluded data are found west of the Andes as indicated by Fig. 2 (cross symbols), suggesting that the tephra deposit distribution was not properly modelled over this area with strong gradients of mass loading.

Finally, the tephra fall deposit reconstruction according to GNC/GIG methods is shown in Fig. 8. For comparative purposes, the Romero et al. (2016) deposit contours (hand-drawn isopachs for 0.1, 0.5, 1.0 and 2.0 mm) are over imposed in Fig. 8, providing an alternative estimate of the fallout deposit from an independent dataset of thickness measurements. Note that a remarkable feature of the GNC/GIG reconstructed deposits is the presence of a distal secondary thickness maximum  $\sim$ 250 km downwind from vent, indicating that this volcanic eruption involved a complex plume dynamics. Both reconstructions success-

<sup>§</sup> Average over five realisations



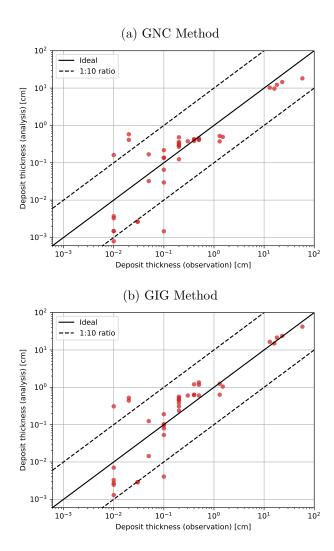


Figure 7. Comparison between analysis and observations from the validation dataset (41 points, see Table 2 and text for explanation).

fully captured the location of this maximum due to ash aggregation processes which were described using the parameterisation proposed (Cornell et al., 1983; Folch et al., 2010) assuming a aggregate particle class having a diameter of 200  $\mu m$  and a density sampled in a range centred around 450  $kg m^{-3}$  (see Table 4).

## 4 Source term inversion: Application to the 2015 Calbuco eruption

A major advantage of the GNC method is that it allows estimating the Eruption Source Parameters (ESP) in a straightforward way, with inverse modelling coming at no extra computational cost. This is because FALL3D solves an almost linear problem with weak non-linearity effects (e.g. due to gravity current, wet deposition, or aggregation) and, consequently, a re-scaling of the emission source term  $s_i$  associated with the *i*-th ensemble member according to  $s_i \rightarrow w_i s_i$  leads to a deposit mass loading



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re-scaled correspondingly, being  $w_i$  the weight factor provided by the GNC method. In this case, the best estimation of the total source term is given by:

$$s^a = \sum_i w_i s_i \tag{21}$$

where  $\{s_i\}$  are the emission source terms of the prior ensemble members (in kg m<sup>-3</sup> s<sup>-1</sup>). Figure 9 shows the emission rate profiles resulting from the source inversion, expressed in terms of the linear source emission strength (in kg m<sup>-1</sup> s<sup>-1</sup>), i.e.  $s^a \times dA$ , where dA is the area of the cell grid.

As most of the 256 weight factors converge to zero,  $w_i \rightarrow 0$ , the profiles in Fig. 9 reflect only those ensemble members that effectively contribute to the analysed deposit mass loading. According to the GNC inverse modelling results, each eruptive phase is characterised by different vertical mass distribution and emission rates. The first phase results on higher cloud top heights, reaching almost 20 km (a.s.l.) whereas the column heights during the second phase remain at around 16 km (a.s.l.). Note also that the prior ensemble was defined assuming the same emission source and sampling parameters for both eruptive phases. For instance, the ensemble configuration was defined assuming the same reference value for column heights, i.e. 15 km above vent level for both phases (Table 4) or 17 km (a.s.l.), approximately. Therefore, the resulting asymmetry between both pulses observed in Fig. 9 is a result of the GNC inversion procedure. In other words, the GNC method can discard inappropriate ensemble members and pick out those that are consistent with observations. The vertical mass distribution for the first phase results on a larger spread and most of the total erupted mass stems from the second phase. Assuming a bulk density of  $\rho_b = 800 \text{ kg m}^{-3}$ , it is possible to estimate the total erupted volume. The time-series for mass eruption rate and total erupted volume are also depicted in Fig. 9. In particular, the final total erupted volume was around 0.33 km³. According to the inversion, the erupted volume corresponding to the first and second phases were 0.11 km³ (32.9%) and 0.22 km³ (67.1%), respectively. These results are in good agreement with the estimations reported by Romero et al. (2016), which give a total bulk tephra fall deposit volume of  $0.27 \pm 0.007 \text{ km}^3$  with 62% of the total volume corresponding to the second phase.

## 375 5 Discussion

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Traditional ensemble-based DA methods such as the ensemble Kalman filter (EnKF) are based on the Gaussian hypothesis. However, it is well known that analyses produced by these methods are suboptimal when either the model state variables or the observation errors are not Gaussian distributed. Volcanic aerosol concentrations and tephra deposit mass loading are two remarkable examples of non-Gaussian-distributed variables with highly-skewed distributions. This explains why the application of EnKF-like methods in VATD models often leads to nonphysical results and oscillations (e.g. the occurrence of negative concentrations). The ensemble-based GNC and GIG methods introduced in this paper are promising alternatives for dealing with assimilation of volcanic data. These methods differ in the assumptions made about the prior distribution and the likelihood of the observation conditioned on the true state. Both methods have been applied to the assimilation of volcanic deposit data and have shown to have a similar performance in terms of the evaluation metrics.



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The GNC method assumes a multi-dimensional Gaussian distribution and solves an optimization problem with nonnegative constraints to ensure plausible physical solutions. The GNC method, constrained here to assimilate deposit observations, can be easily extended to other observables as long as the observation operator is linear. For example, VATD models could use it to assimilate column mass observations of volcanic aerosols, but the assimilation of other satellite-retrieved variables (e.g. Aerosol Optical Depth) would require of an alternative approach. Since the linear estimator (6) includes the prior ensemble mean  $(w_i = 1/m)$  and any specific ensemble member  $(\mathbf{w} = [0, \dots, 1, \dots, 0]^\mathsf{T})$  among the possible solutions, the solution obtained through the minimisation process (12) converges to an analysis state which, by construction, improves the prior ensemble mean and any individual ensemble member. Furthermore, the GNC method ensures that the prior wRMSE is reduced by the analysis state. This can be checked from (5) by noting that, if the iterative solving procedure is started from the uniform vector with components  $w_i = 1/m$ , the normalised cost function (defined as  $\sqrt{J/p}$ ) coincides with the weighted RMSE of the prior ensemble mean before the first iteration (as long as the matrix R is diagonal). In fact, the metrics reported in Table 5 illustrate this property of the solution, e.g. the wRMSE is reduced from wRMSE=8.53 (prior ensemble mean) to wRMSE=1.32 (analysis). This is not the case for individual ensemble members; the weighted RMSE associated with a specific member may be less than the analysis relative error. This is a desirable property of the solution for statistically non-significant members. In fact, note that the first term in (5) penalises deviations from the ensemble mean (i.e. statistically non-significant members), while the second term penalises deviations from the observations. As a result, the solution provided by the GNC method satisfies two properties: (i) the analysis is statistically significant and, (ii) deviations from the observations are small.

The GIG method is a sequential assimilation procedure proposed by Bishop (2016), in which single observations are sequentially assimilated. The GIG method is based on the GIG equation set for the special case where the prior distribution can be described by a gamma PDF and an inverse-gamma observation likelihood. This is a stochastic method providing an ensemble of analyses and does not require a linear observation operator. These reasons make the GIG method a better candidate for implementation in VATD models as it would allow performing multiple assimilation cycles by restarting a corrected ensemble forecast. The GIG method enables near-zero semi-positive-definite variables with highly skewed uncertainty distributions to be assimilated, and avoids the occurrence of negative mass loading at the observation site. However, a linear regression approximation is used to propagate the observation information to the extended model-observation state vector in the assimilation step. Consequently, negative values can be introduced by (15) when the analysis in computed at the grid points of the computational domain. In order to solve this potential issue, we explored a localisation technique which limits the observation impacts based on their background error correlations using (16). Unfortunately, this localisation approach did not show a clear improvement according to the evaluation metrics. Probably, the linear regression approximation went well in this work due to the linearity of the observation operator. However, for more general problems involving non-linear observation operators, the regression (15) may need to be reformulated and the localisation approach proposed here could be more relevant.



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#### 6 Conclusions

This paper has proposed two ensemble-based data assimilation methods for semi-positive-definite variables. The methods were applied to reconstruct the tephra fallout deposit of the 2015 Calbuco eruption in Chile by assimilating measurements of deposit thickness. An assessment based on an independent observational dataset yielded similar results for both methods in terms of the evaluation metrics. We conclude that both methods are a promising alternative for assimilation of volcanic fallout data. In addition, the techniques presented in this work can be extended to other volcanic observations, including satellite-retrieved data of different aerosol types, and used in operational forecast contexts.

The GNC method provides an ensemble of weight factors and can be used also for source term inversion in a straightforward way. Unlike the majority of source term inversion methods (e.g. Folch and Mingari, 2022), which focus on determining specific ESP associated with oversimplified parameterisations of the source term, this approach reconstructs the overall space-time distribution of the source and it is not constrained by any specific parameterisation of the emission source term.

The GIG method is a second-order method and provides an ensemble of analyses without the linear observation operator assumption. Consequently, it represents an attractive alternative for assimilation of volcanic aerosol observations from satellite retrievals. To this purpose, the analysis ensemble from the GIG method could be used to perform multiple assimilation cycles by restarting an ensemble forecast. This approach has the potential to improve the accuracy of operational forecasts of volcanic clouds.

## Appendix A: Generation of the analysis ensemble in the GIG method

In order to generate an ensemble  $y_{ji}^a$  with the correct posterior density from the forecast ensemble for the case of an univariate gamma prior and an inverse-gamma observation-likelihood, Bishop (2016) proposed the following stochastic equation:

$$\frac{y_{ji}^{a} - \overline{y}_{j}^{a}}{\overline{y}_{j}^{a}} = \frac{y_{ji}^{f} - \overline{y}_{j}^{f}}{\sqrt{\left(\overline{y}_{j}^{f}\right)^{2} + \operatorname{var}(y_{ji}^{f})}} + \tilde{P}_{j}^{r} \left(\tilde{P}_{j}^{r} + \tilde{R}_{j}^{r}\right)^{-1} \times \left\{ \frac{z_{ji} - \overline{z}_{j}}{\sqrt{\left(\overline{z}_{j}\right)^{2} - 2\operatorname{var}(z_{ji})}} - \frac{y_{ji}^{f} - \overline{y}_{j}^{f}}{\sqrt{\left(\overline{y}_{j}^{f}\right)^{2} + \operatorname{var}(y_{ji}^{f})}} \right\}$$
(A1)

where  $\overline{y}_j^a$  can be computed using Eq. (14a). Equation (A1) ensures that the analysis ensemble  $y_{ji}^a$  is consistent with the type 1 relative error variance of the posterior given by Eq. (14b) provided  $z_{ji}$  is randomly sampled from a gamma PDF with type 1 relative error variance  $R_j^z$  and mean  $\overline{z}_j$  given by:

$$\left(R_j^z\right)^{-1} = \left(\tilde{R}_j^r\right)^{-1} + 2\tag{A2a}$$

$$\overline{z}_j = \frac{\left(\tilde{R}_j^r\right)^{-1} + 2}{\left(\tilde{R}_j^r\right)^{-1}} y_j^o \tag{A2b}$$





Code and data availability. FALL3D-8.2.0 is available under the version 3 of the GNU General Public License (GPL) at https://doi.org/10. 5281/zenodo.6343786 (Folch et al., 2020; Prata et al., 2021). Observational datasets, code used for the GNC/GIG methods and input model parameter file along with the pre- and post-processing scripts have been archived on Zenodo at https://doi.org/10.5281/zenodo.7259531.

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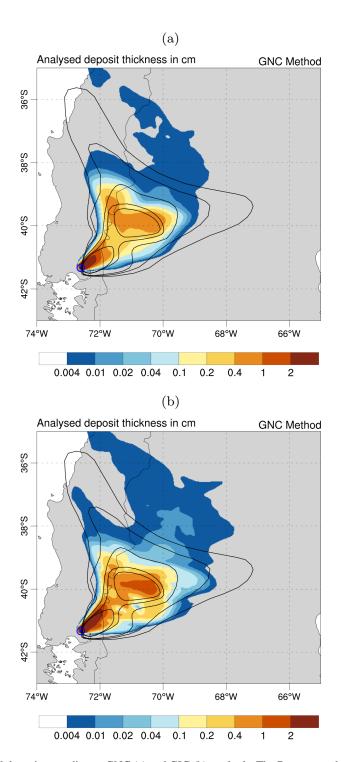


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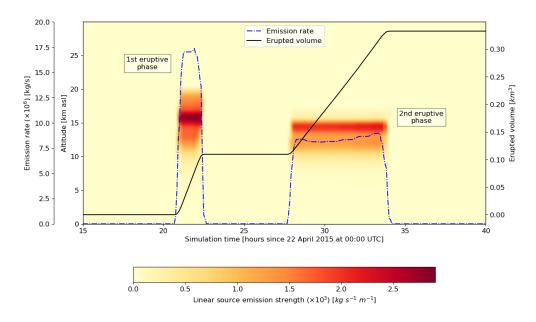




**Figure 8.** Reconstructed tephra fall deposit according to GNC (a) and GIG (b) methods. The Romero et al. (2016) deposit contours (isopachs for 0.1, 0.5, 1.0 and 2.0 mm) are also over imposed for comparative purposes.







**Figure 9.** Profiles of emission rate and time-series of Eruption Source Parameters (ESP) for the 2015 Calbuco eruption according to the GNC inverse modelling approach.