

1 **Accelerated Estimation of Sea Spray-Mediated Heat Flux**
2 **Using Gaussian Quadrature: Case Studies with a Coupled**
3 **CFSv2.0-WW3 System**

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8

9 **Abstract.** Sea spray-mediated heat flux plays an important role in air-sea heat transfer. Heat flux
10 integrated over droplet size spectrum can well simulate total heat flux induced by sea spray droplets.
11 Previously, a fast algorithm of spray-flux assuming single-radius droplets (A15) was widely used since
12 the full-size spectrum integral is computationally expensive. Based on the Gaussian Quadrature (GQ)
13 method, a new fast algorithm (SPRAY-GQ) of sea spray-mediated heat flux is derived. The performance
14 of SPRAY-GQ is evaluated by comparing heat fluxes with those estimated from the widely-used A15.
15 The new algorithm shows a better agreement with the original spectrum integral. To further evaluate the
16 numerical errors of A15 and SPRAY-GQ, the two algorithms are implemented into a coupled CFSv2.0-
17 WW3 system, and a series of 56-day simulations in summer and winter are conducted and compared.
18 The comparisons with satellite measurements and reanalysis data show that the SPRAY-GQ algorithm
19 could lead to more reasonable simulation than the A15 algorithm by modifying air-sea heat flux. For
20 experiments based on SPRAY-GQ, the sea surface temperature at mid-high latitudes of both hemispheres,
21 particularly in summer, is significantly improved compared with the experiments based on A15. The
22 simulation of 10-m wind speed and significant wave height at mid-low latitudes of the Northern
23 Hemisphere after the first two weeks is improved as well. These improvements are due to the reduced
24 numerical errors. The computational time of SPRAY-GQ is about the same as that of A15. Therefore,
25 the newly-developed SPRAY-GQ algorithm has a potential to be used for calculation of spray-mediated
26 heat flux in coupled models.

27

28 **1 Introduction**

29 Sea spray droplets, ejected from oceans, include film drops, jet drops and spume drops (Veron, 2015).
30 The first two types of droplets are generated from bubble bursting caused by ocean surface wave breaking,
31 with radius ranging from 0.5 μm to 50 μm (Resch and Afeti, 1991; Thorpe, 1992; Melville, 1996; Spiel,
32 1997; Andreas, 1998; Lhuissier and Villermaux, 2012). Spume drops are generated by strong winds ($>$
33 7-11 m/s) which directly tear the wave crests, with larger radius ranging from tens to hundreds of
34 micrometers (Koga, 1981; Andreas et al., 1995; Andreas, 1998). Sea spray droplets play an important
35 role in weather and climate processes (Fox-Kemper et al., 2022). On one hand, sea spray droplets
36 contribute to local marine aerosols and subsequently modify the local radiation balance (Fairall et al.,
37 1983; Burk, 1984; Fairall and Larsen, 1984). On the other hand, sea spray droplets affect the fluxes of
38 heat, momentum, salt, and freshwater between atmosphere and ocean (Andreas, 1992; Andreas et al.,
39 2008; Andreas, 2010; Andreas et al., 2015; Ling and Kao, 1976; Fairall et al., 1994; Andreas and
40 Decosmo, 2002).

41 The sea spray-mediated heat transfer mainly occurs within the droplet evaporation layer (DEL) near
42 the sea surface (Andreas and Decosmo, 1999, 2002; Fairall et al., 1994). Sea spray droplets with the same
43 temperature as ocean surface can lead to sensible heat flux in DEL, while water evaporated from these
44 droplets can further release latent heat to the atmosphere (Andreas, 1992; Borisenkov, 1974; Bortkovskii,
45 1973; Wu, 1974; Monahan and Van Patten, 1988; Ling and Kao, 1976). Part of the sea spray-mediated
46 sensible heat is absorbed by droplet evaporation, which further increases the air-sea temperature
47 difference, and thus increases the sea spray-mediated sensible heat flux (Fairall et al., 1994; Andreas and
48 Decosmo, 2002). Since strong winds produce more sea spray droplets with larger radius, sea spray-
49 mediated heat fluxes increase with wind speed (Fairall et al., 1994), and contribute to more than 10% of
50 the total surface heat flux after reaching the threshold speed ($>$ 11 m/s for sensible heat flux and $>$ 13 m/s
51 for latent heat flux)(Andreas et al., 2008). In addition, when a droplet is released into the air, it is
52 accelerated due to surface winds (Edson and Andreas, 1997; Fairall et al., 1994; Van Eijk et al., 2011;
53 Wu et al., 2017). If the droplet could fall back into the ocean, additional momentum would be injected
54 into the ocean from the atmosphere (Andreas, 1992, 2004).

55 The usual bulk parameterizations in numerical models for surface fluxes only include the interfacial

56 (turbulent) fluxes (e.g., Fairall et al., 1996), while neglecting the significant contributions of sea spray
57 droplets in DEL (Andreas et al., 2008; Fairall et al., 1994; Smith, 1997; Emanuel, 1995). Andreas and
58 Emanuel (2001) implemented sea spray-mediated heat flux and momentum flux parameterizations into
59 a simple tropical cyclone model, and found that the sea spray-mediated heat flux can significantly
60 enhance tropical cyclone intensity. The similar enhancement of tropical cyclone intensity was also
61 noticed in recent regional coupling systems by including sea spray-mediated heat flux (Xu et al., 2021a;
62 Liu et al., 2012; Garg et al., 2018; Zhao et al., 2017). In the First Institute of Oceanography Earth System
63 Model, Bao et al. (2020) first incorporated the sea spray-mediated heat flux in global climate simulation.
64 Following Bao et al. (2020), Song et al. (2022) found that the sea spray-mediated heat flux can lead to
65 cooling at the air-sea interface and westerlies strengthening in the Southern Ocean, and thus improves
66 estimates of sea surface temperature (SST).

67 Since the parameterization of sea spray-mediated heat flux derived from observations requires full-
68 size spectral integral and thus is computationally expensive for large-scale models (Table 1, details in
69 Section 4.2; Andreas, 1989, 1990, 1992; Andreas et al., 2015), a simplified algorithm based on a single
70 radius of sea spray droplets (Andreas et al., 2015; Andreas et al., 2008) is widely used in atmosphere-
71 ocean coupling systems (Xu et al., 2021a; Liu et al., 2012; Garg et al., 2018; Zhao et al., 2017; Song et
72 al., 2022; Bao et al., 2020), and apt to produce numerical errors. To reduce these numerical errors induced
73 by the single radius of sea spray droplets, we develop a new fast algorithm of sea spray-mediated heat
74 flux based on the Gaussian Quadrature (GQ) method, a fast and accurate way to calculate spectral integral.
75 The GQ method has been successfully used for the estimation of domain-averaged radiative flux profiles
76 (Li and Barker, 2018). The performance of the GQ-based fast algorithm of the sea spray-mediated heat
77 flux is evaluated and compared with the simplified algorithm for single radius of Andreas et al. (2015),
78 referred to as A15 hereafter. The results are first compared with the original parameterization using full-
79 size spectral integral (A92, hereafter). Then the parameterizations with different algorithms are
80 implemented in a global coupled atmosphere-ocean-wave system (Shi et al., 2022), and the results are
81 compared with global satellite measurements and reanalysis data.

82 The rest of the paper is structured as follows: observation and reanalysis data for comparisons are
83 introduced in Section 2; the derivation of the GQ-based fast algorithm and the global coupling system

84 are described in Section 3; the performance of the new fast algorithm is evaluated in Section 4. Finally,
85 a summary and discussion are given in Section 5.

86 **2 Data**

87 The fifth generation European Centre for Medium-Range Weather Forecasts (ECMWF) Reanalysis
88 (ERA5; Hersbach et al., 2020) 10-m wind speed (WSP10), 2-m air temperature (T02), 2-m dewpoint
89 temperature, surface pressure and significant wave height (SWH) with a spatial resolution of 0.5° are
90 used. Additionally, WSP10, T02 and 2-m specific humidity (SPH) data from the Objectively Analyzed
91 air-sea Fluxes (OAFflux) products (Yu et al., 2008) are also applied for comparison, with $1^\circ \times 1^\circ$ resolution.
92 The daily average satellite Optimum Interpolation SST (OISST) data are obtained from the National
93 Oceanic and Atmospheric Administration (NOAA) with a spatial resolution of 0.25° (Reynolds et al.,
94 2007). The global monthly mean salinity observations from European Space Agency (ESA;
95 https://climate.esa.int/sites/default/files/SSS_cci-D1.1-URD-v1r4_signed-accepted.pdf) are applied.
96 Besides, we also use the monthly global ocean RSS Satellite Data Products for WSP10
97 (https://data.remss.com/wind/monthly_1deg/) and the Reprocessed L4 Satellite Measurements for SWH
98 (<https://doi.org/10.48670/moi-00177>), to validate the simulation results and ERA5 data.

99 **3 Methods**

100 **3.1 Development of a Fast Algorithm Based on GQ**

101 The effects of sea spray droplets on sensible and latent heat fluxes ($H_{S,SP}$, $H_{L,SP}$) contribute to the total
102 turbulent sensible and latent heat fluxes ($H_{S,T}$, $H_{L,T}$) at the air-sea interface. That is,

$$H_{S,T} = H_S + H_{S,SP}, \quad (1)$$

$$H_{L,T} = H_L + H_{L,SP}. \quad (2)$$

103 where H_S and H_L are the sensible and latent heat fluxes at the air-sea interface due to the air-sea
104 differences of temperature and humidity. Based on observations of total turbulent heat fluxes and the
105 COARE algorithm (Andreas et al., 2015; Fairall et al., 1996), A92 integrates the sea spray-mediated
106 sensible and latent heat flux spectrums over initial droplet radius ($Q_S(r_0)$ and $Q_L(r_0)$) to estimate $H_{S,SP}$

107 and $H_{L,SP}$ (details in Appendix A; Andreas, 1989, 1990, 1992; Andreas and Decosmo, 2002). The
108 distributions of $Q_S(r_0)$ and $Q_L(r_0)$ spectrums as functions of initial droplet radius r_0 under various
109 atmosphere and ocean states are shown in Fig. 1, indicating that Q_S and Q_L spectrums are more
110 sensitive to the change of WSP10, and less sensitive to other variables, including T02, 2-m relative
111 humidity, SST, surface air pressure and sea surface salinity.

112 Since the calculation of $H_{S,SP}$ and $H_{L,SP}$ in A92 is computationally expensive due to full-size
113 spectral integral (Eqn. A5-A6 of Appendix A), it is difficult to apply A92 directly in coupled modeling
114 systems. A15 (Andreas et al., 2015) developed a fast algorithm by using a single representative droplet
115 radius (details in Appendix B), which was widely adopted in recent regional and global coupling systems
116 (Xu et al., 2021a; Liu et al., 2012; Garg et al., 2018; Zhao et al., 2017; Song et al., 2022; Bao et al., 2020).
117 In this study, we apply a 3-node GQ method (details in Appendix C) to develop a new fast algorithm to
118 approximate the full-size spectral integral of A92. Notably, GQ can converge exponentially to the actual
119 integral only for a smooth function, which is a prerequisite for GQ (McClarren, 2018). Since as functions
120 of r_0 , $Q_S(r_0)$ and $Q_L(r_0)$ are not smooth (Fig. 1), a data sorting from largest to smallest is required.
121 After sorting, local $Q_S(r_0)$ and $Q_L(r_0)$ become $Q_{S_sort}(m)$ and $Q_{L_sort}(m)$, and then GQ can be used
122 to estimate the integral of $Q_{S_sort}(m)$ and $Q_{L_sort}(m)$. Note that the independent variable m is not
123 equivalent to the original r_0 , but only indicates the position. In this way, according to Appendix C,
124 $m_1=443$, $m_2=251$, $m_3=58$ are three GQ nodes of $Q_{S_sort}(m)$ and $Q_{L_sort}(m)$, and we can get the
125 corresponding r_0 for local $Q_S(Q_L)$, denoted as $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$. However, the
126 sorting leads to high complexity of GQ comparable to A92, and the values of $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and
127 $r_{S3}(r_{L3})$ vary under various atmosphere and ocean states in the globe. Therefore, it is necessary to find
128 the general approximate values of $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$ via global statistical analyses, to
129 avoid the sorting in application.

130 To derive the general approximate values of $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$, we calculate the
131 distribution of the sea spray-mediated heat flux spectral following A92, based on the global daily WSP10,
132 T02, 2-m dewpoint temperature, surface pressure and SWH of ERA5 and OISST from August 1, 2018
133 to August 31, 2018. Since the sea spray-mediated heat flux is not sensitive to salinity (Fig. 1e&f) and
134 only monthly observational data is available, the ESA monthly salinity is applied. From the global

135 spectrums, we sort Q_S and Q_L from largest to smallest to obtain local r_{S1} , r_{S2} and r_{S3} (r_{L1} , r_{L2} and
 136 r_{L3}) for every grid point, whose global distribution of occurrence frequency in percentage is shown in
 137 Fig. 2. It is noted that except for r_{L3} , all other five nodes have frequency roughly concentrated at a
 138 constant (peak frequency >65% in Fig. 2a, b, d-f; Eqn. 3&4), while for r_{L3} , there is a 92.53%
 139 concentration between 55 and 90 μm (Fig. 2c). Then we found that r_{L3} (55-90 μm) is related to WSP10
 140 (Fig. S1 in supplementary), thereby we set the approximate values as

$$r_{S1} = 459.056, r_{S2} = 294.185, r_{S3} = 166.771, \quad (3)$$

$$r_{L1} = 443.914, r_{L2} = 251.0498, \quad (4)$$

$$r_{L3} = \begin{cases} 60.310WSP10^{0.1161}, & WSP10 \geq 2 \text{ m/s} \\ 58.086, & WSP10 < 2 \text{ m/s} \end{cases}, \quad (5)$$

141 where the unit of the radius is micrometer. Afterwards, we directly use Eqn. 3-5 to approximate the full-
 142 size spectral integral of A92 without sorting as

$$\int_a^b Q_S(r_0)dr_0 \approx \frac{b-a}{2} \sum_{i=1}^3 \omega_i Q_S(r_{Si}), \quad (6)$$

$$\int_a^b Q_L(r_0)dr_0 \approx \frac{b-a}{2} \sum_{i=1}^3 \omega_i Q_L(r_{Li}). \quad (7)$$

143 Here a and b are the lower and upper limits of r_0 , which are set to $2\mu m$ and $500\mu m$ based on Andreas
 144 (1990), and ω_i is the corresponding weight ($\omega_1=\omega_3=0.556$, $\omega_2=0.889$), obtained from McClarren
 145 (2018). The new fast algorithm for approximations of $H_{S,SP}$ and $H_{L,SP}$ is referred to as SPRAY-GQ
 146 hereafter.

147 3.2 CFSv2.0-WW3 Coupling System

148 A coupled system based on Climate Forecast System model version 2.0 (CFSv2.0) and
 149 WAVEWATCH III (WW3) is employed to evaluate and compare the effects of sea spray-mediated heat
 150 flux parameterized by A15 and SPRAY-GQ. The CFSv2.0-WW3 has three components, the Global
 151 Forecast System (GFS; <http://www.emc.ncep.noaa.gov/GFS/doc.php>) as the atmosphere component of
 152 CFSv2.0, the Modular Ocean Model version 4 (MOM4; Griffies et al., 2004) as the ocean component of
 153 CFSv2.0, and the WW3 (WAVEWATCH III Development Group, 2016) as the ocean surface wave
 154 component. The variables between CFSv2.0 and WW3 are interpolated and passed using the Chinese

155 Community Coupler version 2.0 (C-Coupler2; Liu et al., 2018).

156 The CFSv2.0 is mainly applied for intraseasonal and seasonal prediction (e.g., Saha et al., 2014). The
157 atmosphere component GFS uses a spectral triangular truncation of 382 waves (T382) in the horizontal,
158 equivalent to a grid resolution of nearly 35 km, and 64 sigma-pressure hybrid layers in the vertical. The
159 MOM4 is integrated on a nominal 0.5° horizontal grid with enhanced horizontal resolution to 0.25° in
160 the tropics, and there are 40 levels in the vertical. The CFSv2.0 initial fields at 00:00 UTC of the first
161 day for experiments were generated by the real time operational Climate Data Assimilation System
162 (Kalnay et al., 1996), downloaded from the CFSv2.0 official website
163 (<http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod>). The latitude range of WW3 is 78°S – 78°N
164 with a spatial resolution of $1/3^\circ$. The initial wave fields were generated from 10-day simulations starting
165 from rest in a stand-alone WW3 model, forced by ERA5 10-m winds and ice concentration. The open
166 boundary conditions of WW3 were also obtained by the global simulation of the stand-alone WW3 model.

167 In the coupling system, the WW3 obtains 10-m wind and ocean surface current from CFSv2.0, and
168 then provides wave parameters to CFSv2.0. Several wave-mediated processes, including upper ocean
169 mixing modified by Stokes drift-related processes, air-sea fluxes modified by surface current and Stokes
170 drift, and momentum roughness length, are considered. Details of this system are referred to Shi et al.
171 (2022).

172 A series of numerical experiments is conducted to evaluate the effects of the two fast algorithms (A15
173 and SPRAY-GQ) of sea spray-mediated heat flux on ocean, atmosphere and waves in two 56-day periods,
174 from January 3 to February 28, 2017 and from August 3 to September 28, 2018 for boreal winter and
175 boreal summer, respectively. For each period, two sensitivity experiments are carried out. The first is the
176 SPRAY-A15 experiment, in which A15 is used with two-way fully coupling. The second is the SPRAY-
177 GQ experiment, in which SPRAY-GQ fast algorithm is used instead of A15. In addition, we also carry
178 out another 7-day experiment using A92 (SPRAY-A92) to test the runtime.

179 4 Results

180 4.1 Comparison with A92

181 Based on the daily global WSP10, T02, 2-m dewpoint temperature, surface pressure and SWH of
182 ERA5, the daily global OISST, and the ESA monthly global salinity, $H_{S,SP}$ and $H_{L,SP}$ from A15,
183 SPRAY-GQ and A92 are calculated (Fig. 3). The computational time for SPRAY-GQ is about the same
184 as that for A15, and about 36 times less than the time for A92. Compared with A92 (the black dotted
185 line), A15 (red) overestimates $H_{S,SP}$ for low $H_{S,SP}$ (<50 W/m²) and underestimates $H_{S,SP}$ for high
186 $H_{S,SP}$ (>50 W/m²) with a root mean square error (RMSE= $\sqrt{\sum_{i=1}^n(\hat{y}_i - y_i)^2/n}$, \hat{y}_i is A15 value, y_i is
187 A92 value, and n is the total number of grid points) of 3.40 W/m² (Fig. 3a), while A15 shows consistent
188 overestimations with a RMSE of 2.98 W/m² for $H_{L,SP}$ (Fig. 3b). Overall, the RMSE of A15 is about
189 2.69 W/m² for sea-spray mediated total heat flux ($TH_{SP} = H_{S,SP} + H_{L,SP}$; Fig. 3c). Andreas et al. (2015)
190 derived A15 from A92 using single-radius droplets as bellwethers and wind functions, and extrapolated
191 the wind functions at high wind speeds >25 m/s. Since the wind speeds in the study are less than 25 m/s
192 (Fig. S1), the large difference between A15 and A92 is mainly due to the use of single-radius droplets.
193 Compared with A15, SPRAY-GQ (blue) has less deviation from A92 for both $H_{S,SP}$ and $H_{L,SP}$ (Fig.
194 3a&b). The corresponding RMSEs of SPRAY-GQ for $H_{S,SP}$, $H_{L,SP}$ and TH_{SP} are 0.83 W/m², 0.92
195 W/m² and 0.62 W/m², all significantly lower ($P<0.05$ in Student's t-test) than those of A15.

196 To test robustness of the results, we also use WSP10, T02 and SPH of OAFlux dataset to estimate
197 $H_{S,SP}$ and $H_{L,SP}$. As shown in Fig. 4, SPRAY-GQ has significantly ($P<0.05$ in Student's t-test) lower
198 deviations and RMSEs than A15, consistent with Fig. 3. Note that the values of $H_{S,SP}$ and $H_{L,SP}$ in
199 Fig.4 are larger than those in Fig. 3. It is because OAFlux only provides neutral wind speeds, calculated
200 from wind stress and the corresponding roughness by assuming air is neutrally stratified. The neutral
201 winds from OAFlux are larger than winds in ERA5 as indicated by previous studies (Lindemann et al.,
202 2021; Seethala et al., 2021).

203 In addition, since it is common to derive SWH from empirical equations (e.g., Andreas et al., 2008;
204 Andreas et al., 2015; Andreas and Decosmo, 2002; Andreas, 1992), we also use SWH generated by
205 empirical equations of WSP10 (Andreas, 1992) instead of ERA5 SWH to estimate $H_{S,SP}$ and $H_{L,SP}$

206 (Fig. 5). Again, the RMSEs decrease significantly ($P < 0.05$ in Student's t-test) in SPRAY-GQ compared
207 to A15, though the RMSEs become higher for all estimates due to the enhanced biases of SWH. The
208 difference between SPRAY-GQ and A92 is always smaller than that between A15 and A92. Next, we
209 will evaluate and compare the two fast algorithms in an atmosphere-ocean-wave coupled system
210 (CFSv2.0-WW3).

211 **4.2 Comparison in the CFSv2.0-WW3 Coupling System**

212 To compare the computational time of different parameterizations in the large-scale modeling system,
213 the runtime of the fully coupled experiments for 7-day forecast is given in Table 1 as an example. It is
214 shown that the runtime is about the same for SPRAY-GQ and SPRAY-A15. Both experiments run about
215 17 times faster than SPRAY-A92.

216 To illustrate the numerical errors of the two fast algorithms discussed in the context of the coupled
217 system, comparisons are made for simulated SSTs, WSP10s as well as SWHs against OISST and ERA5
218 reanalysis. The results in the first three days are excluded in the comparison, since the wave influences
219 are weak at the beginning of the simulations. Overall, the WSP10s of simulations are generally in the
220 range of 0-25 m/s globally. At mid-high latitudes, the WSP10s generally exceed 10 m/s (Fig. S2&S3 of
221 the supplementary), at which the effects of sea spray can become significant (Andreas et al., 2015;
222 Andreas et al., 2008).

223 **4.2.1 Sea Surface Temperature (SST)**

224 In the austral summer, compared with OISST, large SST biases (>1 °C or <-1 °C) of SPRAY-A15
225 occur in the Southern Hemisphere (SH; Fig. S4a in supplementary), especially in the Southern Ocean. It
226 is always a challenge to reduce the large SST biases in the Southern Ocean for climate models (e.g.,
227 Alessandro et al., 2019; Wang et al., 2014; Li et al., 2013; Bodas-Salcedo et al., 2012; Ceppi et al., 2012).
228 In Fig. 6a, SSTs north (south) of 50°S in experiment SPRAY-A15 are mainly underestimated
229 (overestimated). The domain-averaged RMSE (0-360°E, 40-75°S) in experiment SPRAY-A15 increases
230 in the first month and then levels off (red solid line in Fig. 6c), while the domain-averaged RMSE in
231 experiment SPRAY-GQ levels off about a week earlier (black solid line in Fig. 6c). The mean RMSE in

232 SPRAY-GQ is significantly lower than that in SPRAY-A15 ($P < 0.05$ in Student's t-test). The increased
 233 (decreased) SSTs north (south) of 50°S in SPRAY-GQ compared to those in SPRAY-A15 (Fig. 6b)
 234 reduce the RMSE of SST in SPRAY-GQ. We also calculate the mean absolute error,
 235 $\text{MAE} = \sum_{i=1}^n |\hat{y}_i - y_i|/n$, where \hat{y}_i is simulated value and y_i is OISST data, and n is the total number
 236 of grid points. The MAEs are consistent with RMSEs (dotted line in Fig. 6c). Furthermore, the mean
 237 errors, $\text{ME} = \sum_{i=1}^n (\hat{y}_i - y_i)/n$ (Fig. S5a in the supplementary), are smaller in SPRAY-GQ than SPRAY-
 238 A15.

239 To understand the effects of sea spray droplets on SST, we calculate the total heat flux ($\text{TH} = H_{S,T} + H_{L,T}$)
 240 differences between SPRAY-GQ and SPRAY-A15 (Fig. 7a). The TH differences are significantly
 241 correlated with SST differences (Fig. S4b in the supplementary), with the spatial correlation coefficient
 242 of -0.41 ($P < 0.05$ in Student's t-test). We further decompose direct and indirect effects of sea spray
 243 droplets on heat fluxes following Song et al. (2022). The direct effect ($H_{S,SP}$ and $H_{L,SP}$) is induced
 244 directly by sea spray droplets, calculated from A15 (Eqn. B1-B4 of Appendix B) and SPRAY-GQ
 245 (Section 3.1). The indirect effect (H_S and H_L) is the heat flux variation induced by changes of
 246 atmosphere and ocean variables (including wind, pressure, humidity and temperature) caused by direct
 247 effect, estimated by subtracting $H_{S,SP}$ and $H_{L,SP}$ from the output heat fluxes ($H_{S,T}$ and $H_{L,T}$) of
 248 experiment SPRAY-A15 and SPRAY-GQ.

249 In the Southern Ocean, although direct differences of $H_{S,SP}$ and $H_{L,SP}$ are relatively small (< 10
 250 W/m^2 , Fig. 7b, e, & h), the resulting changes of temperature and humidity lead to relatively large
 251 differences in indirect effects of H_S and H_L (Fig. 7c, f, & i). Enhanced (reduced) TH_{SP} from ocean to
 252 atmosphere in the summer leads to increased (decreased) air-sea temperature difference and thus
 253 enhances (weakens) H_S . Meanwhile the warmer (cooler) air also causes more (less) evaporation and thus
 254 more (less) H_L . Finally, the enhanced (reduced) TH cools (warms) SST.

255 In the boreal summer, large SST biases ($> 1^\circ\text{C}$ or $< -1^\circ\text{C}$) of SPRAY-A15 mainly occur at mid-high
 256 latitudes of the Northern Hemisphere (NH; Fig. S6a in supplementary). Significant underestimations
 257 occur in the western and northern part of the North Pacific and at mid latitudes of the North Atlantic,
 258 while large positive SST biases mainly occur in the eastern part of the North Pacific and at high latitudes
 259 of the North Atlantic (Fig. 8a). In experiment SPRAY-GQ, SSTs are warmer (cooler) in the previously

260 underestimated (overestimated) regions (Fig. 8b). Therefore, the domain-averaged RMSE and MAE (0-
261 360°E, 20-75°N) in SPRAY-GQ are significantly lower ($P < 0.01$ in Student's t-test) than in SPRAY-A15
262 after the first three weeks (Fig. 8c). Compared to SPRAY-A15, the overall underestimation is reduced in
263 SPRAY-GQ (Fig. S5b). The spatial correlation coefficient between TH differences and SST differences
264 (Fig. 9a&Fig. S6b) is -0.32 ($P < 0.05$ in Student's t-test). Consistent with the austral summer, the SST
265 changes are related to the changes of heat flux (Fig. 9). The indirect effects of latent heat flux (Fig. 9f)
266 play a major role in TH differences, which are modified by the direct effects (Fig. 9b, e, &h). In addition,
267 the changes of surface wind also contribute to the changes of SST. The reduced winds weaken the upper
268 ocean mixing, the water becomes more stratified, and then the SST tends to be warmer, and vice versa
269 (Fig. S7&S8).

270 **4.2.2 10-m Wind Speed (WSP10) and Significant Wave Height (SWH)**

271 Compared with experiment SPRAY-A15, significant differences of WSP10 in SPRAY-GQ occur at
272 mid-low latitudes of the NH (0-360°E, 0-60°N) in both winter and summer (Fig.S7b&S8b). As we know,
273 satellite scatterometer and altimeter data are usually used to validate WSP10 and SWH for short term
274 weather forecast (e.g., Accadia et al., 2007; Djurdjevic and Rajkovic, 2008; Myslenkov et al., 2021).
275 However, due to the spatial and temporal coverage of satellite data, we can only obtain the monthly
276 averaged satellite data for the globe. Therefore, we compare the monthly averaged WSP10 and SWH
277 from simulations with the corresponding satellite data (Fig. S9-S12). The comparison results (Fig.
278 S9a&c-S12a&c) are consistent with those compared with ERA5 (Fig. S9b&d-S12b&d). From Fig. S9e-
279 S12e, the differences of WSP10s between ERA5 and the satellite data are always less than 1 m/s and the
280 differences of SWHs are always less than 0.3 m. Since ERA5 provides daily data for comparison, we
281 will use ERA5 for validation in the following.

282 The ME of WSP10 (SPRAY-A15 minus ERA5) is 0.28 m/s and 0.47 m/s in winter and summer (red
283 in Fig. S5c&d), respectively, mainly due to the overestimations over the Pacific and the Atlantic Ocean
284 (red in Fig.10a&11a). Whereas in SPRAY-GQ, the ME (SPRAY-GQ minus ERA5) is 0.15 m/s and 0.33
285 m/s in winter and summer respectively (black in Fig. S5c&d). The domain-averaged RMSEs and MAEs
286 of WSP10s increase with time in the first two weeks and then gradually level off (Fig. 10c&11c). The

287 differences of WSP10 RMSEs and MAEs between SPRAY-GQ (black) and SPRAY-A15 (red) are very
288 small in the first two weeks. Afterwards the mean values of RMSE and MAE in SPRAY-GQ are lower
289 than those in SPRAY-A15 significantly at 95% confidence level in both boreal winter (Fig. 10c) and
290 boreal summer (Fig. 11c).

291 The simulated SWHs changes are closely related to the changes of WSP10s (Shi et al., 2022).
292 Therefore, the differences of SWHs (Fig.12&13) are consistent with those of WSP10s (Fig.10&11), with
293 overestimated (underestimated) WSP10s corresponding to overestimated (underestimated) SWHs
294 compared with ERA5. The SWHs in SPRAY-GQ are significantly different from those in SPRAY-A15
295 (Fig. 12b&13b). In winter (summer), the SWH RMSE averages for SPRAY-A15 and SPRAY-GQ are
296 1.31 m (0.98 m) and 1.23 m (0.87 m), and after the first two weeks the RMSE and MAE in SPRAY-GQ
297 are lower than those in SPRAY-A15 significantly at 95% confidence level in both winter (Fig. 12c) and
298 summer (Fig. 13c).

299 The direct and indirect effects of sea spray droplets on heat fluxes can influence estimates of WSP10
300 and then SWH. The changes of WSP10s are related to the direct effects ($H_{S,SP}$ and $H_{L,SP}$; Fig. 7b, e, &h;
301 Fig. 9b, e, &h). The spatial correlation coefficients between WSP10 differences (Fig. S7b&S8b) and
302 TH_{SP} differences (Fig. 7b&9b) are 0.51 and 0.69 ($P<0.01$ in Student's t-test) in winter and summer,
303 respectively, because TH_{SP} differences can influence the sea level pressure (SLP) distribution (Fig.
304 S15&S16), and subsequently surface winds. For example, compared with SPRAY-A15, the decreased
305 TH_{SP} of SPRAY-GQ in the Northwest Pacific in summer (Fig. 9b) leads to higher SLP and smaller
306 pressure gradient (Fig. S16), and thus decreased WSP10 (Fig. 11b); while the increased TH_{SP} in the
307 Gulf of Alaska (Fig. 9b) leads to lower SLP and larger pressure gradient (Fig. S16), and thus enhanced
308 WSP10 (Fig. 11b). The accelerated (decelerated) WSP10s further result in increased (decreased)
309 interfacial heat transport (H_S , H_L), as well as increased (decreased) SWHs.

310 **5 Conclusions and Discussion**

311 Based on a GQ method, we develop a new fast algorithm based on Andreas's (1989, 1990, 1992) full-
312 size microphysical parameterization (A92) for sea spray-mediated heat fluxes. Using global satellite
313 measurements and reanalysis data, we found that the difference between SPRAY-GQ and A92 is

314 significantly smaller than that between A15 and A92 (Andreas et al., 2015). To evaluate the numerical
315 error of SPRAY-GQ/A15 fast algorithm, we implement them in the two-way coupled CFSv2.0-WW3
316 system. A series of 56-day simulations from January 3 to February 28, 2017 and from August 3 to
317 September 28, 2018 are conducted. The results are compared against satellite measurements and ERA5
318 reanalysis. The comparison shows that the sea spray-mediated heat flux in SPRAY-GQ can reasonably
319 modulate total heat flux compared with SPRAY-A15, and significantly reduce the SST biases in the
320 Southern Ocean (mid-high latitudes of the NH) for the austral (boreal) summer, as well as WSP10 and
321 SWH after the first two weeks at mid-low latitudes of the NH for both boreal winter and summer. Overall,
322 our fast algorithm based on GQ is applicable to sea spray-mediated heat flux parameterization in coupled
323 models.

324 To investigate the effects of spray-mediated heat flux on simulations, two 56-day experiments without
325 sea spray effect (CTRL) in boreal winter and summer are conducted respectively, and the differences of
326 simulated SST, WSP10, SWH, T02 and SPH between SPRAY-GQ and CTRL are compared in Fig. S17-
327 S21 in the supplementary. The introduction of sea spray cannot significantly reduce the global overall
328 errors of simulations, but it leads to regional improvements (blue in Fig. S17e&f-S21e&f). For example,
329 compared with CTRL in Jan-Feb, 2017, SST MAE of SPRAY-GQ in the southeast of Australia decreases
330 (Fig. S17e), because of warmer SST (Fig. S17c) related to reduced wind (Fig. S18c). The reduced wind
331 here also leads to lower SWH (Fig. S19c) and thus reduced SWH overestimation (Fig. S19e). Meanwhile,
332 SPRAY-GQ reduces MAE of T02 and SPH (Fig. S20e&S21e) by increasing temperature and moisture
333 (Fig. S20c&S21c). The reduced errors are related to the relatively large WSP10s over the areas (Fig.
334 S2&S3), since the effects of sea spray become important at wind speeds larger than 10 m/s.

335 In addition to the variables aforementioned, the changes of simulated cloud fraction were also
336 compared. However, the effects of sea spray-mediated heat flux on cloud fraction are non-significant for
337 the 2-month simulation, so the results are not shown. Besides, the lack of other processes related to sea
338 spray may be one of the reasons why the global overall error cannot be reduced effectively. For example,
339 for simulated WSP10 and SWH in SPRAY-GQ, the significant overestimations in the SH still exist
340 especially in Aug-Sep, 2018 (Fig. S18&S19 in supplementary). As Andreas (2004) indicated, sea spray
341 droplets also influence the surface momentum flux by injecting more momentum into the ocean from the

342 atmosphere, which might further decrease the surface wind speed. We will consider this process in the
343 future study.

344 Sea spray-mediated heat fluxes are related to the sea spray generation function (SSGF). Based on a
345 number of laboratory and field observations, varieties of SSGF were derived (e.g., Koga, 1981; Monahan
346 et al., 1982; Troitskaya et al., 2018; Andreas, 1992, 1998, 2002; Fairall et al., 1994; Veron, 2015),
347 whereas their differences can reach six orders of magnitude (Andreas, 1998). There is currently no
348 consensus on the most suitable choice. In this study, we use SSGF of Fairall et al. (1994), recommended
349 by Andreas (2002), to get a mean bias of 3.70 W/m² and 0.095 W/m² for latent and sensible heat flux
350 respectively (Andreas et al., 2015), consistent with recent observations of Xu et al. (2021b). However,
351 the improved SST and other variables cannot be reliably assigned to the usage of the GQ method, due to
352 the uncertainties of the coupled model itself and SSGF.

353 When wind speed is larger than 10 m/s, spray-mediated heat flux can become as important as the
354 interfacial heat flux (Andreas and Decosmo, 1999, 2002). Particularly, even in the absence of air-sea
355 temperature difference, the spray-mediated sensible heat flux is still present (Andreas et al., 2008). As
356 indicated by previous studies (e.g., Garg et al., 2018; Song et al. 2022), it is necessary to superimpose
357 the spray-mediated heat flux on the bulk formula to complete the physics of turbulent heat transfer for
358 coupled simulation. Since the full microphysical parameterization (A92) is computationally expensive,
359 an efficient algorithm that captures the main features of A92 can be beneficial to large-scale climate
360 systems or operational storm models. The GQ method proposed in the study can efficiently calculate the
361 spray-mediated heat flux, and agree better with A92 than A15. Therefore, the GQ based spray-mediated
362 heat flux is promising to be widely applied in large-scale climate systems and operational storm models.

363

364 **Appendix A**

365 **Microphysical Parameterization of A92**

366 Based on the cloud microphysical parameterization of Pruppacher and Klett (1978), Andreas (1989,
367 1990, 1992) proposed a parameterization of sea spray-related heat fluxes for droplets with different radius,

368 from formation at sea surface to equilibrium with environment, that is,

$$Q_S = \rho_w C_{ps} (T_w - T_{eq}) \left[1 - \exp\left(-\frac{\tau_f}{\tau_T}\right) \right] \left(\frac{4\pi r_0^3}{3} \frac{dF}{dr_0} \right), \quad (A1)$$

$$Q_L = \begin{cases} \rho_w L_v \left\{ 1 - \left[\frac{r(\tau_f)}{r_0} \right]^3 \right\} \left(\frac{4\pi r_0^3}{3} \frac{dF}{dr_0} \right), & \tau_f \leq \tau_r, \\ \rho_w L_v \left\{ 1 - \left(\frac{r_{eq}}{r_0} \right)^3 \right\} \left(\frac{4\pi r_0^3}{3} \frac{dF}{dr_0} \right), & \tau_f > \tau_r. \end{cases} \quad (A2)$$

369 Here Q_S , Q_L are sensible heat flux and latent heat flux resulted from sea spray droplets with initial
 370 radius r_0 , ρ_w is the sea water density, C_{ps} is the specific heat, L_v is the latent heat of vaporization of
 371 water, T_w is the water temperature, T_{eq} is the temperature of droplet when it reaches thermal
 372 equilibrium with ambient condition, r_{eq} is the radius of droplet when it reaches moisture equilibrium
 373 with ambient condition, τ_f is the residence time for droplets in the atmospheric, $r(\tau_f)$ is the
 374 corresponding radius, τ_T is the characteristic e-folding time of droplet temperature, and τ_r is the
 375 characteristic e-folding time of droplet radius. The detailed calculation of these microphysical quantities
 376 can be found in Andreas (1989, 1990, 1992). dF/dr_0 is the sea spray generation function, which
 377 represents the produced number of droplets with initial radius r_0 (Andreas, 1992). For this term, the
 378 function of Fairall et al. (1994) was recommended by Andreas (2002). According to the review in
 379 Andreas (2002), the dF/dr_0 of Fairall et al. (1994) is related on that of Andreas (1992) as

$$\frac{dF}{dr_0} = 38 \times 3.84 \times 10^{-6} U_{10}^{3.41} r_0^{-0.024} \left. \frac{dF_{A92}}{dr_{80}} \right|_{U_{10}=11 \text{ m/s}}, \quad (A3)$$

$$\left. \frac{dF_{A92}}{dr_{80}} \right|_{U_{10}=11 \text{ m/s}} = \begin{cases} e^{(4.405-2.646(\log r_{80})-3.156(\log r_{80})^2+8.902(\log r_{80})^3-4.482(\log r_{80})^4)}, & r_{80} \leq 15 \mu\text{m}; \\ 1.02 \times 10^4 r_{80}^{-1}, & 15 \leq r_{80} \leq 37.5 \mu\text{m}; \\ 6.95 \times 10^6 r_{80}^{-2.8}, & 37.5 \leq r_{80} \leq 100 \mu\text{m}; \\ 1.75 \times 10^{17} r_{80}^{-8}, & r_{80} \geq 100 \mu\text{m} \end{cases} \quad (A4)$$

380 Here U_{10} is the 10-m wind, $r_{80} = 0.518 r_0^{0.976}$.

381 The total sea spray fluxes are obtained by integrating Q_S and Q_L corresponding to all r_0 . Based on
 382 Andreas (1990), the lower and upper limits of r_0 is $2 \mu\text{m}$ and $500 \mu\text{m}$, that is,

$$\overline{Q_S} = \int_2^{500} Q_S(r_0) dr, \quad (A5)$$

$$\overline{Q_L} = \int_2^{500} Q_L(r_0) dr. \quad (A6)$$

383 Note that $\overline{Q_S}$ and $\overline{Q_L}$ are nominal sea spray fluxes but not the actual $H_{S,SP}$ and $H_{L,SP}$ (Andreas and
 384 Decosmo, 1999, 2002), because there are interactions between these two terms and the microphysical

385 functions also lead to uncertainties (Fairall et al., 1994). Therefore, $\overline{Q_S}$ and $\overline{Q_L}$ are tuned by non-
 386 negative constants α , β and γ (Andreas and Decosmo, 2002; Andreas et al., 2008; Andreas et al., 2015;
 387 Andreas, 2003) as

$$H_{S,SP} = \beta \overline{Q_S} - (\alpha - \gamma) \overline{Q_L}, \quad (\text{A7})$$

$$H_{L,SP} = \alpha \overline{Q_L}. \quad (\text{A8})$$

388 In Eqn. (A8), the α term indicates the sea spray-mediated latent heat flux from the top of DEL to
 389 atmosphere. Because the evaporation of droplets absorbs heat, which is provided by sea spray-mediated
 390 sensible heat (Fairall et al., 1994), the negative α term appears in Eqn. (A7). The evaporation also cools
 391 DEL and thus increases the air-sea temperature difference, therefore it contributes to a positive γ term
 392 in Eqn. (A7). Different values of α , β and γ were given in Andreas and Decosmo (2002), Andreas
 393 (2003), Andreas et al. (2008) and Andreas et al. (2015), to minimize the bias between estimations and
 394 observations of turbulent heat fluxes measured by eddy correlation. And Andreas et al. (2015) validated
 395 the most observation data, which are 4000 sets, to derive $\alpha = 2.46, \beta = 15.15, \gamma = 1.77$.

396 Appendix B

397 Fast Algorithm of A15

398 Andreas (2003) and Andreas et al. (2008, 2015) developed a fast algorithm to approximate $H_{S,SP}$,
 399 $H_{L,SP}$ by a characteristic radius, that is,

$$H_{S,SP} = \beta \overline{Q_S} - (\alpha - \gamma) \overline{Q_L} \approx \rho_w C_{ps} (T_W - T_{eq,100}) V_S(u_*), \quad (\text{B1})$$

$$H_{L,SP} = \alpha \overline{Q_L} \approx \rho_w L_v \left\{ 1 - \left[\frac{r(\tau_{f,50})}{50 \mu\text{m}} \right]^3 \right\} V_L(u_*). \quad (\text{B2})$$

400 Here $T_{eq,100}$ is T_{eq} of droplets with $r_0=100 \mu\text{m}$, $\tau_{f,50}$ is τ_f of droplets with $r_0=50 \mu\text{m}$, V_S and
 401 V_L are functions of the bulk friction velocity u_* . As indicated by Andreas et al. (2008, 2015), the
 402 characteristic radiuses of 100 μm and 50 μm for sensible and latent heat fluxes are chosen,
 403 respectively, because Q_S and Q_L show a large peak in the vicinity of these values (Fig. 1). V_S and V_L
 404 are calculated in Andreas et al. (2015) as

$$V_S = \begin{cases} 3.92 \times 10^{-8}, & 0 \leq u_* \leq 0.1480 \text{ m/s} \\ 5.02 \times 10^{-6} u_*^{2.54}, & u_* \geq 0.1480 \text{ m/s} \end{cases} \quad (\text{B3})$$

$$V_L = \begin{cases} 1.76 \times 10^{-9}, & 0 \leq u_* \leq 0.1358 \text{ m/s} \\ 2.08 \times 10^{-7} u_*^{2.39}, & u_* \geq 0.1358 \text{ m/s} \end{cases} \quad (\text{B4})$$

405 Appendix C

406 Gaussian Quadrature (GQ)

407 GQ is a method to approximate the definite integral of a function $f(x)$ via the function values at a
 408 small number of specified nodes (Gauss, 1815; Jacobi, 1826). In this study we use the form of n-node
 409 Gauss–Legendre quadrature on $[-1, 1]$ as

$$\int_{-1}^1 f(x) dx \approx \sum_{i=1}^n \omega_i f(x_i). \quad (\text{C1})$$

410 Here x_i is the specified node, and ω_i is the corresponding weight. For $n=3$, $x_1=-0.775$, $x_2=0$,
 411 $x_3=0.775$, $\omega_1=\omega_3=0.556$, $\omega_2=0.889$.

412 For a function $g(\xi)$ on $[a, b]$, Eqn. (C1) can be transformed to

$$\begin{aligned} \int_a^b g(\xi) d\xi &= \int_{-1}^1 g\left(\frac{b-a}{2}x + \frac{a+b}{2}\right) \frac{d\xi}{dx} dx \\ &\approx \frac{b-a}{2} \sum_{i=1}^n \omega_i g\left(\frac{b-a}{2}x_i + \frac{a+b}{2}\right). \end{aligned} \quad (\text{C2})$$

413 Code and data availability

414 The code of sea spray can be found under <https://doi.org/10.5281/zenodo.7100345> or
 415 <https://zenodo.org/record/7100345#.Y66vRtVByHt> (Shi and Xu, 2022). The code for CFSv2.0-WW3
 416 system can be found under <https://doi.org/10.5281/zenodo.5811002> (Shi et al., 2021) including the
 417 coupling, preprocessing, run control and postprocessing scripts. The initial fields for CFSv2.0 are
 418 generated by the real time operational Climate Data Assimilation System, downloaded from the CFSv2.0
 419 official website (<http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod>). The daily average satellite
 420 Optimum Interpolation SST (OISST) data are obtained from NOAA (<https://www.ncdc.noaa.gov/oisst>).
 421 The fifth generation European Centre for Medium-Range Weather Forecasts (ECMWF) Reanalysis
 422 (ERA5) are available at the Copernicus Climate Change Service (C3S) Climate Data Store
 423 (<https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels>). The daily Objectively
 424 Analyzed air-sea Fluxes (OAFflux) products are available at <https://oafux.whoi.edu/heat-flux>. The global
 425 monthly mean salinity observations of European Space Agency (ESA) are from <https://climate.esa.int>.

426 The monthly global ocean RSS Satellite Data Products for 10-m wind speed are from
427 https://data.remss.com/wind/monthly_1deg/, and the Reprocessed L4 Satellite Measurements for
428 significant wave height are from <https://doi.org/10.48670/moi-00177>.

429 **Author contribution**

430 FX and RS designed the experiments and RS carried them out. RS developed the code of coupling
431 parametrizations and produced the figures. RS prepared the manuscript with contributions from all co-
432 authors. FX contributed to review and editing.

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438 **Competing Interests**

439 The contact author has declared that neither they nor their co-authors have any competing interests.

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Table 1. The runtime of CFSv2.0-WW3 global experiments for 7-day forecast with different parameterizations.

7-day Forecast	Runtime (h)
SPRAY-A92	126.94
SPRAY-A15	7.60
SPRAY-GQ	7.67

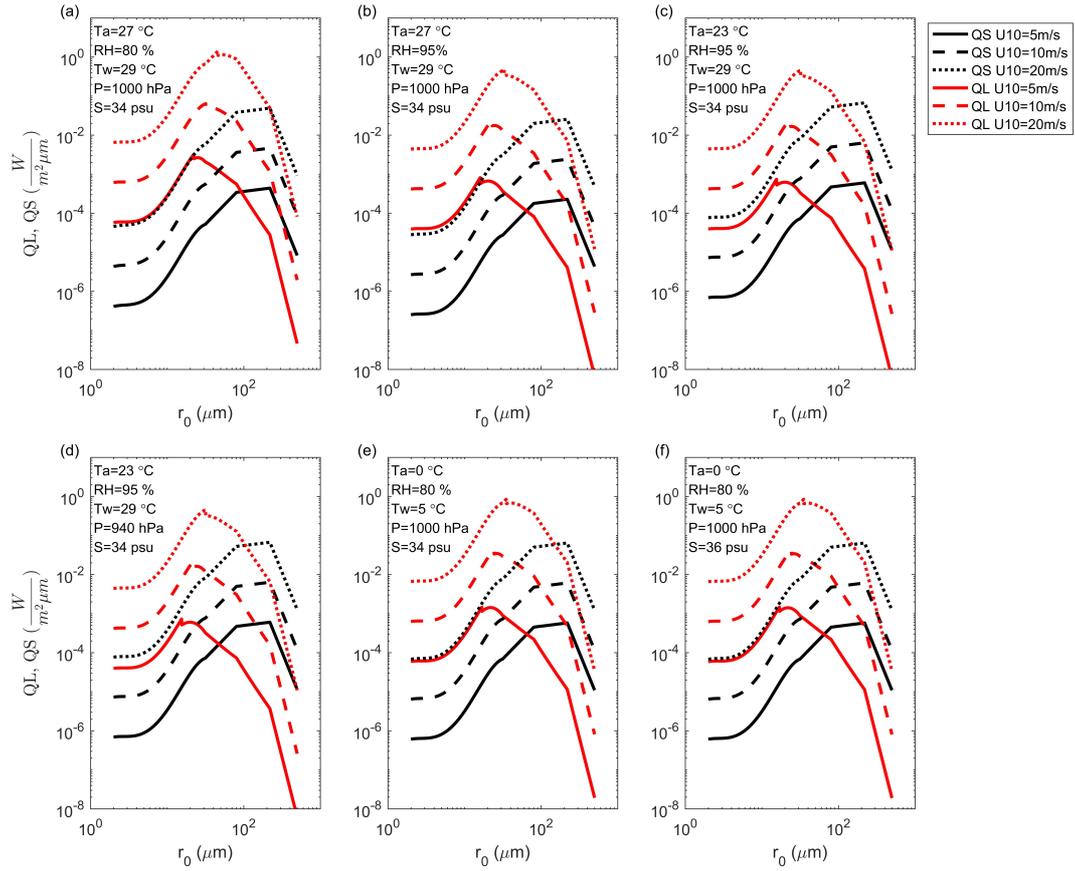


Figure 1. The radius-specific sea spray-mediated sensible (Q_S ; black) and latent (Q_L ; red) heat fluxes as functions of initial radius r_0 : U_{10} , T_a , RH , T_w , P and S are 10-m wind speed, 2-m air temperature, 2-m relative humidity, sea surface temperature, surface air pressure and surface salinity, respectively.

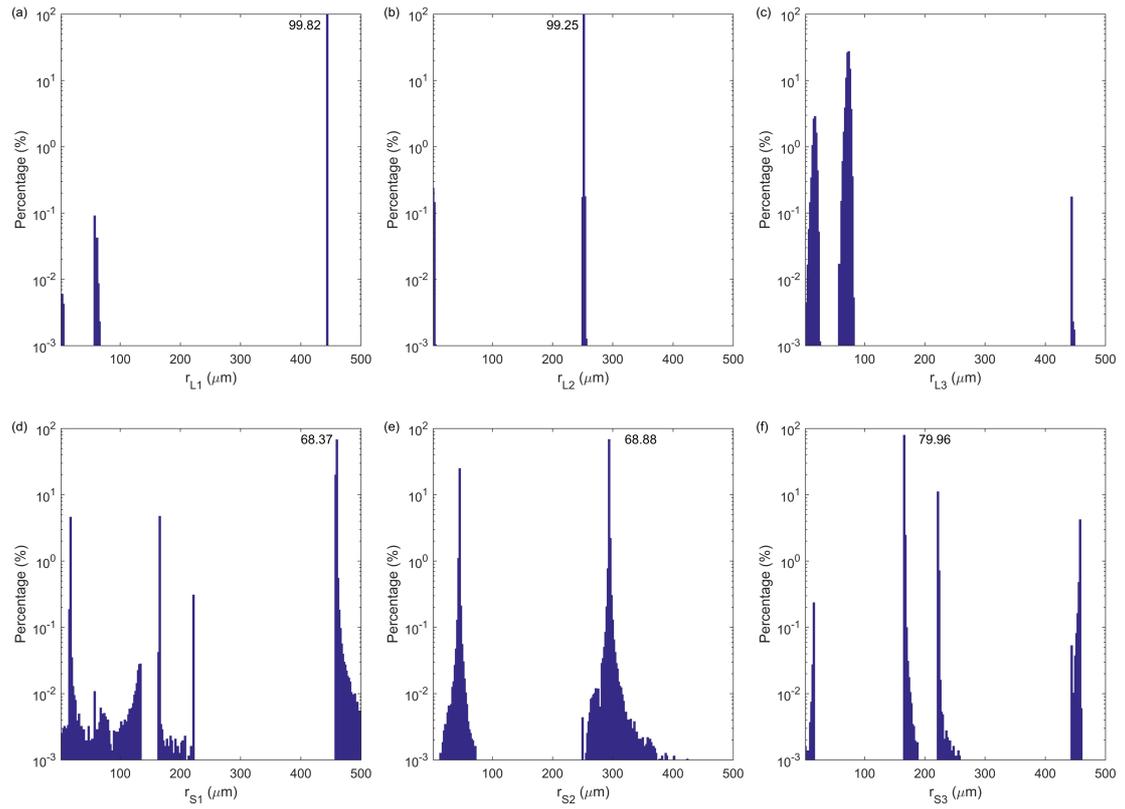


Figure 2. The distribution of occurrence frequency in percentage for GQ radius nodes: (a) the first node of latent heat flux; (b) the second node of latent heat flux; (c) the third node of latent heat flux; (d) the first node of sensible heat flux; (e) the second node of sensible heat flux; (f) the third node of sensible heat flux. The peak frequencies are marked.

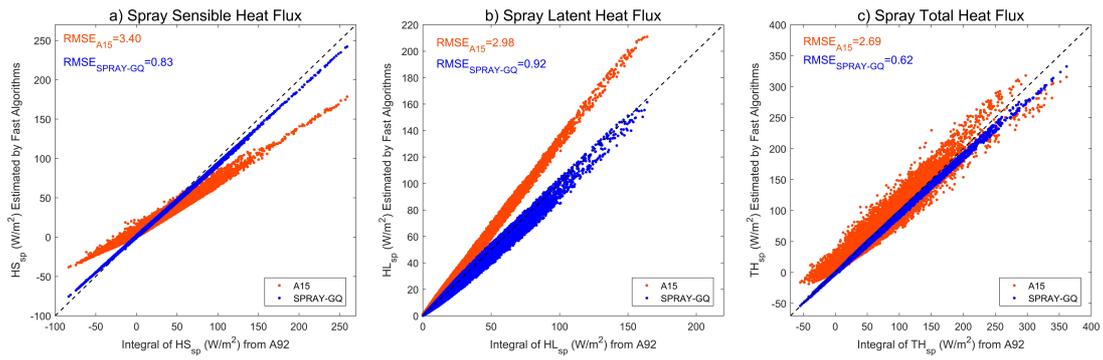


Figure 3. Scatter plots of $H_{S,SP}$ (a), $H_{L,SP}$ (b) and total heat flux $TH_{SP} = H_{S,SP} + H_{L,SP}$ (c) estimated by fast algorithms (y-axis) vs those estimated by spectral integral in microphysical parameterization (x-axis): The dotted black line is $y=x$. The corresponding RMSEs are marked in the upper left corner.

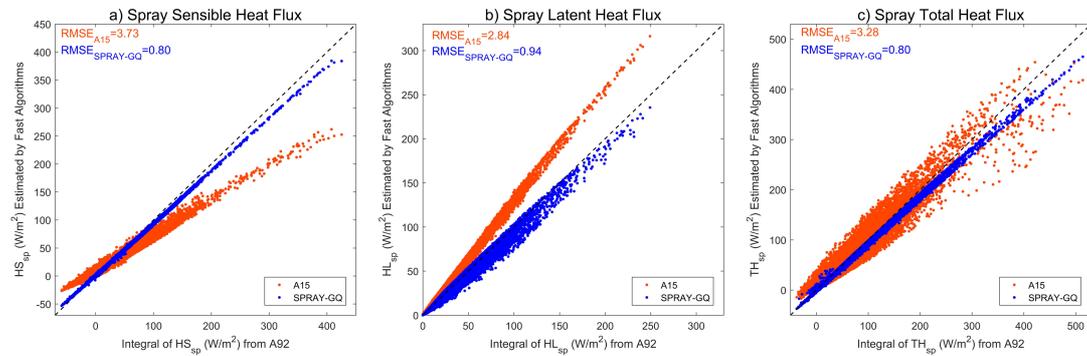


Figure 4. The same as Figure 3, but WSP10, 2-m air temperature and 2-m specific humidity of OAF flux are used.

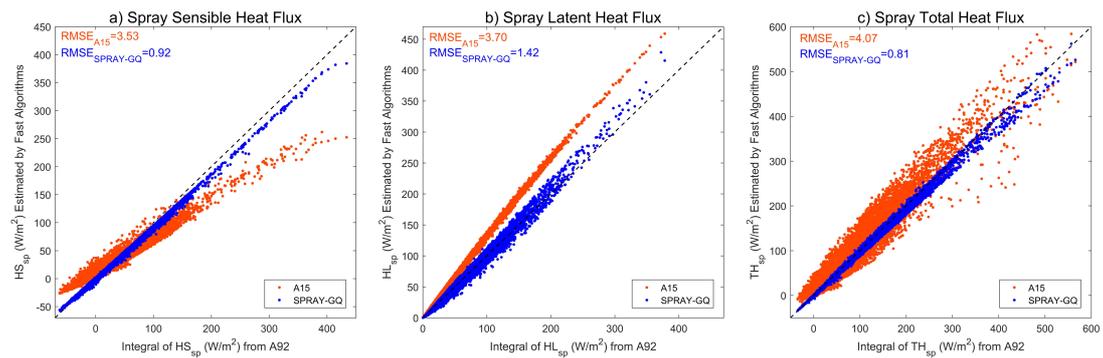


Figure 5. The same as Figure 4, but SWH is derived by WSP10 instead of ERA5 SWH.

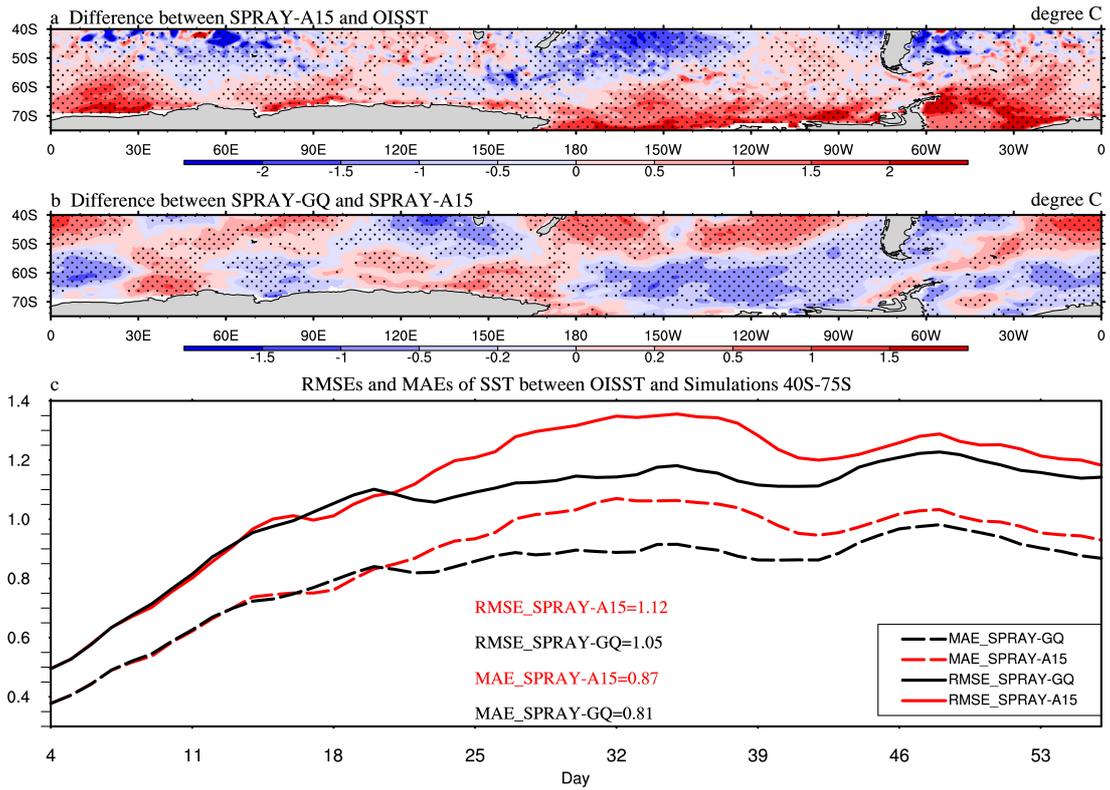


Figure 6. The 53-day average SST ($^{\circ}\text{C}$) differences between SPRAY-A15 and OISST (a; SPRAY-A15 minus OISST), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360 $^{\circ}\text{E}$, 40-75 $^{\circ}\text{S}$) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.

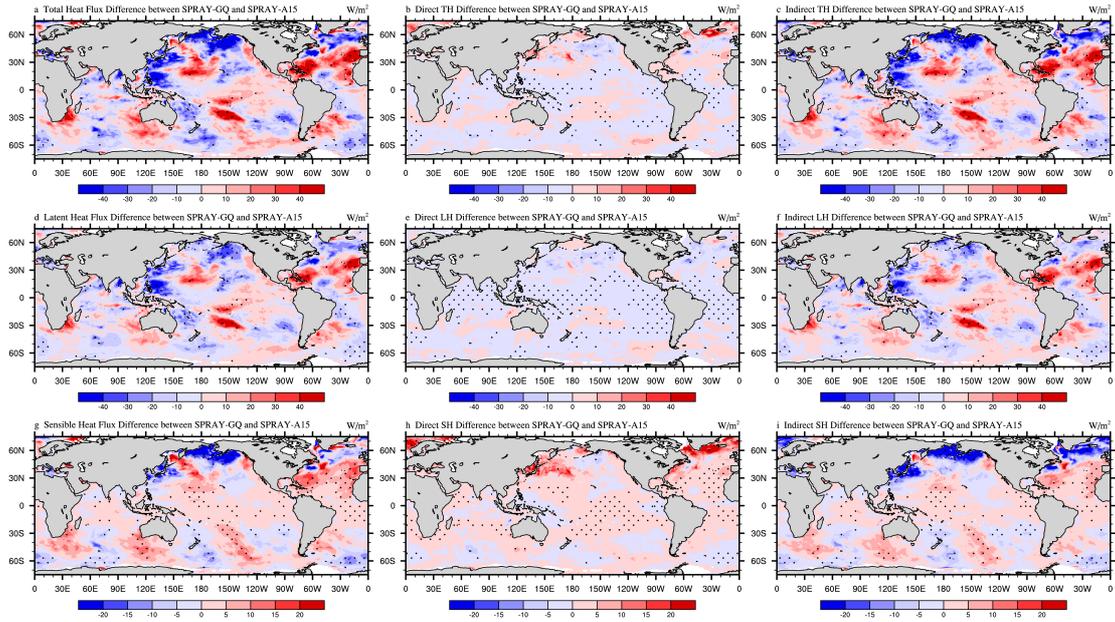


Figure 7. The 53-day average differences of total heat flux (a-c), latent heat flux (d-f), and sensible heat flux (g-i) between SPRAY-GQ and SPRAY-A15 (SPRAY-GQ minus SPRAY-A15) in Jan-Feb, 2017. The direct differences indicate sea spray-mediated heat flux differences (b, e, h), and the indirect differences indicate interfacial (bulk) heat flux differences resulted by sea spray (c, f, i). The dotted areas are statistically significant at 95% confidence level. A positive value of flux indicates an upward direction.

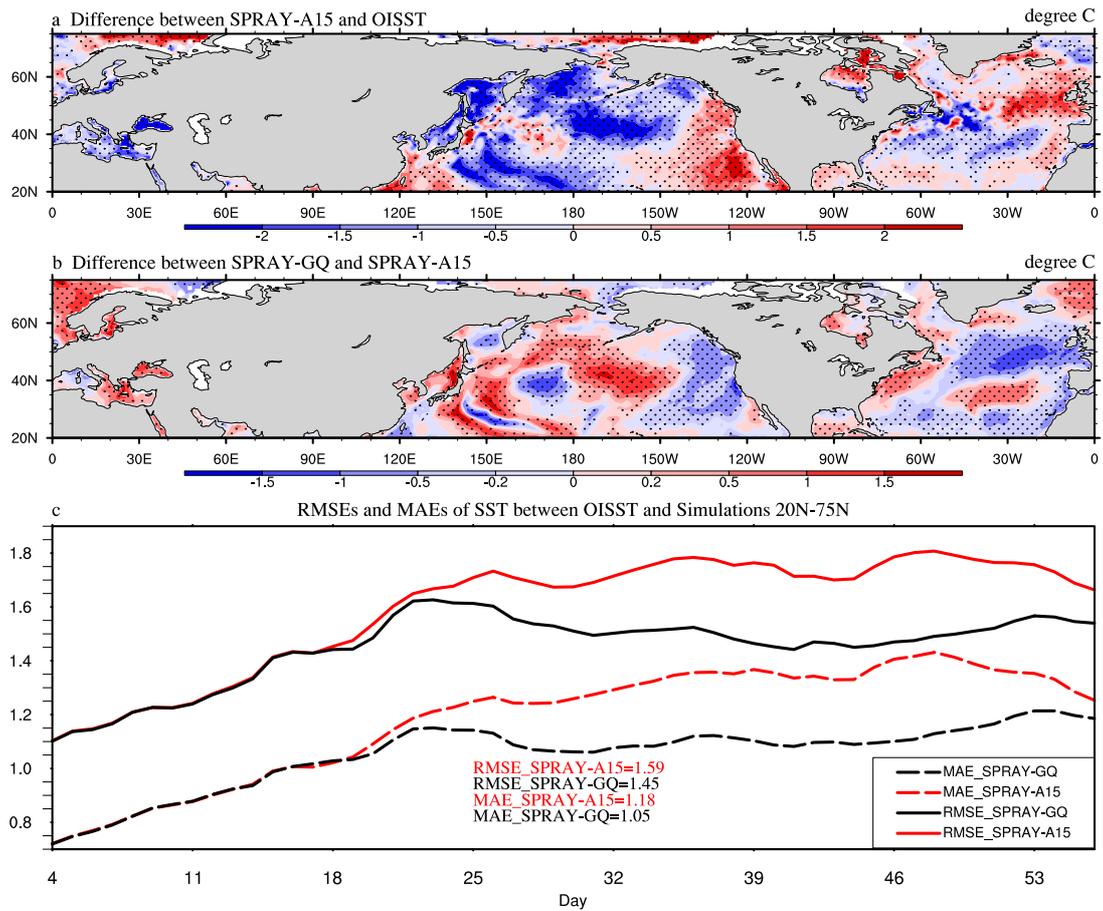


Figure 8. The same as Figure 6, but for Aug-Sep, 2018 in 0-360°E, 20-75°N.

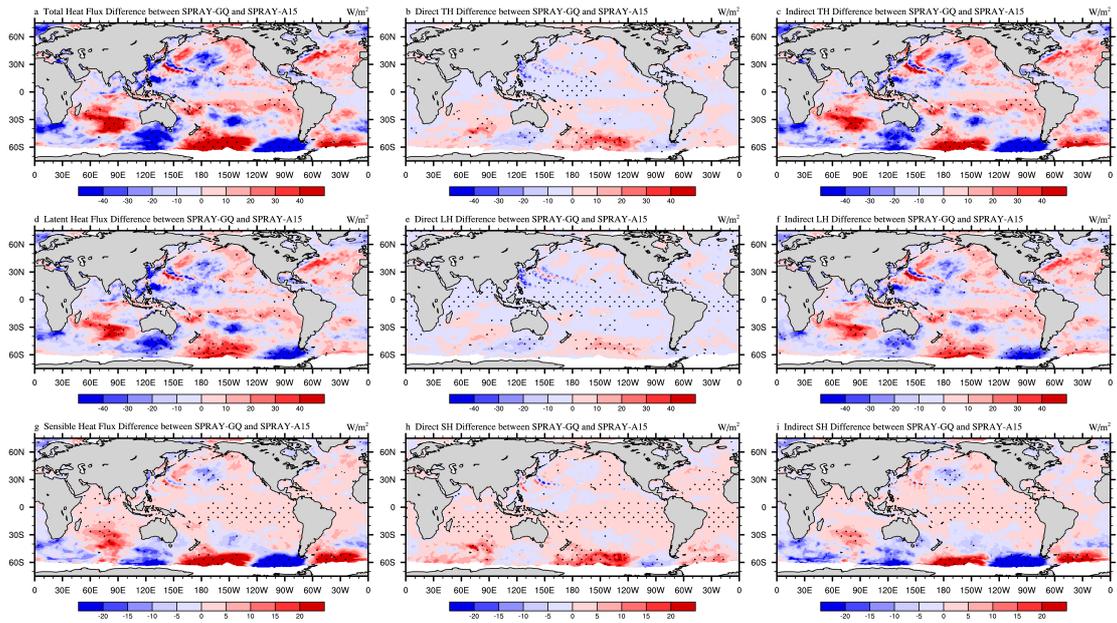


Figure 9. The same as Figure 7, but for Aug-Sep, 2018.

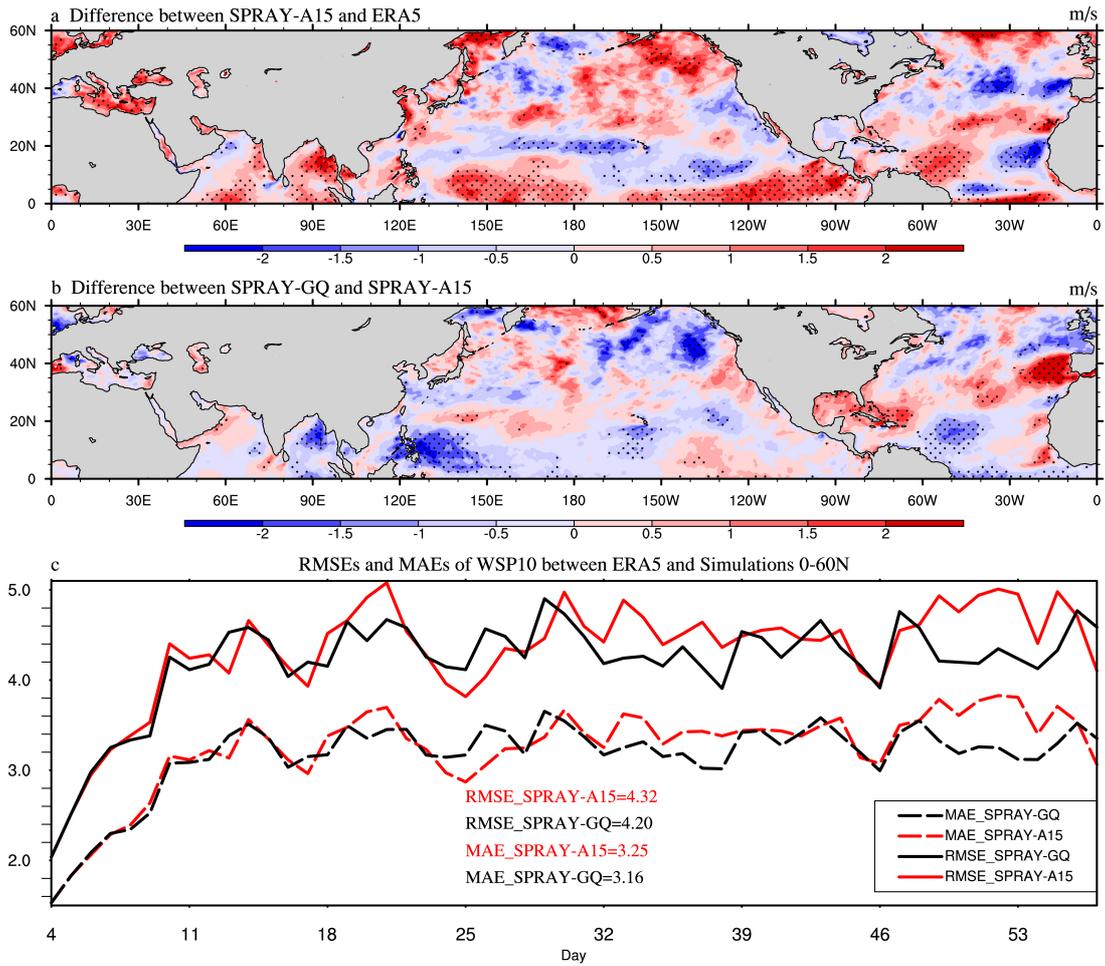


Figure 10. The 53-day average WSP10 (m/s) differences between SPRAY-A15 and ERA5 (a; SPRAY-A15 minus ERA5), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360°E, 0-60°N) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.

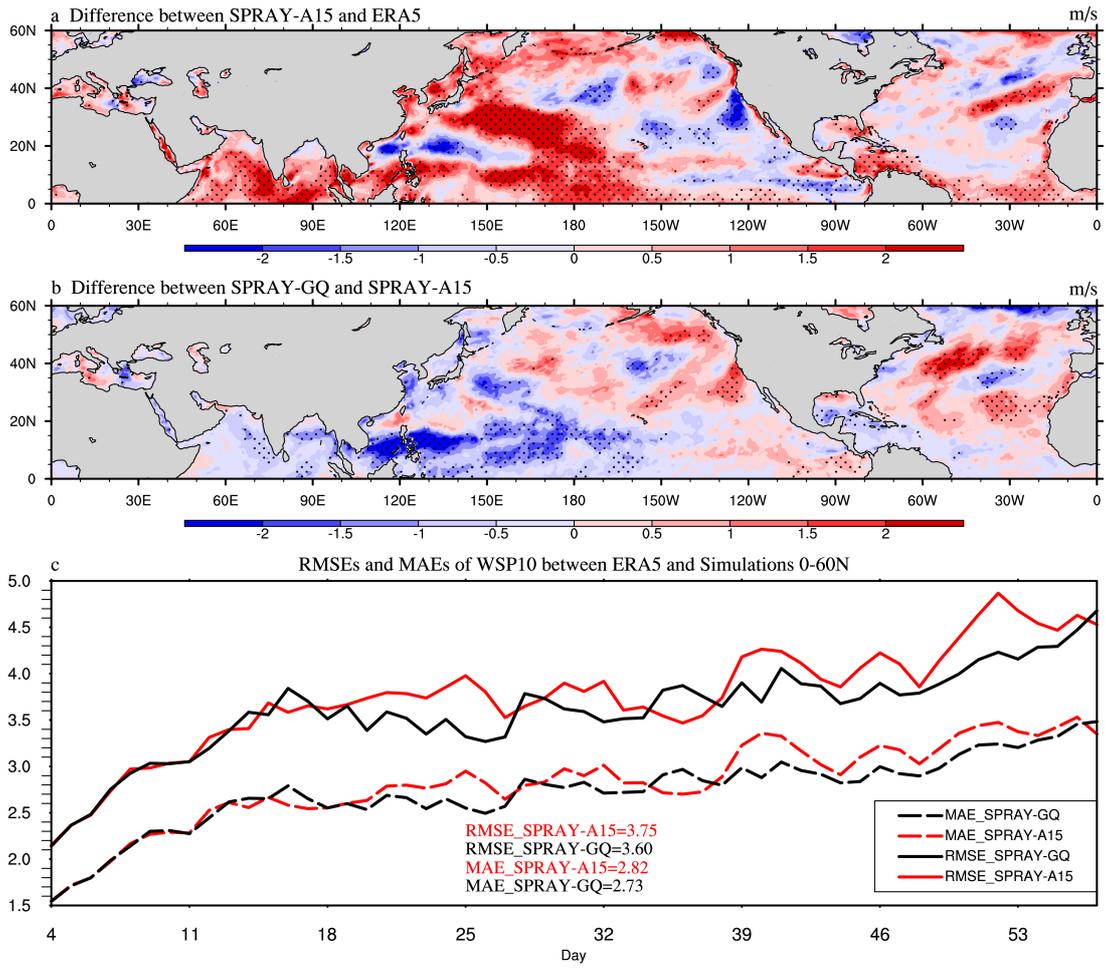


Figure 11. The same as Figure 10, but for Aug-Sep, 2018.

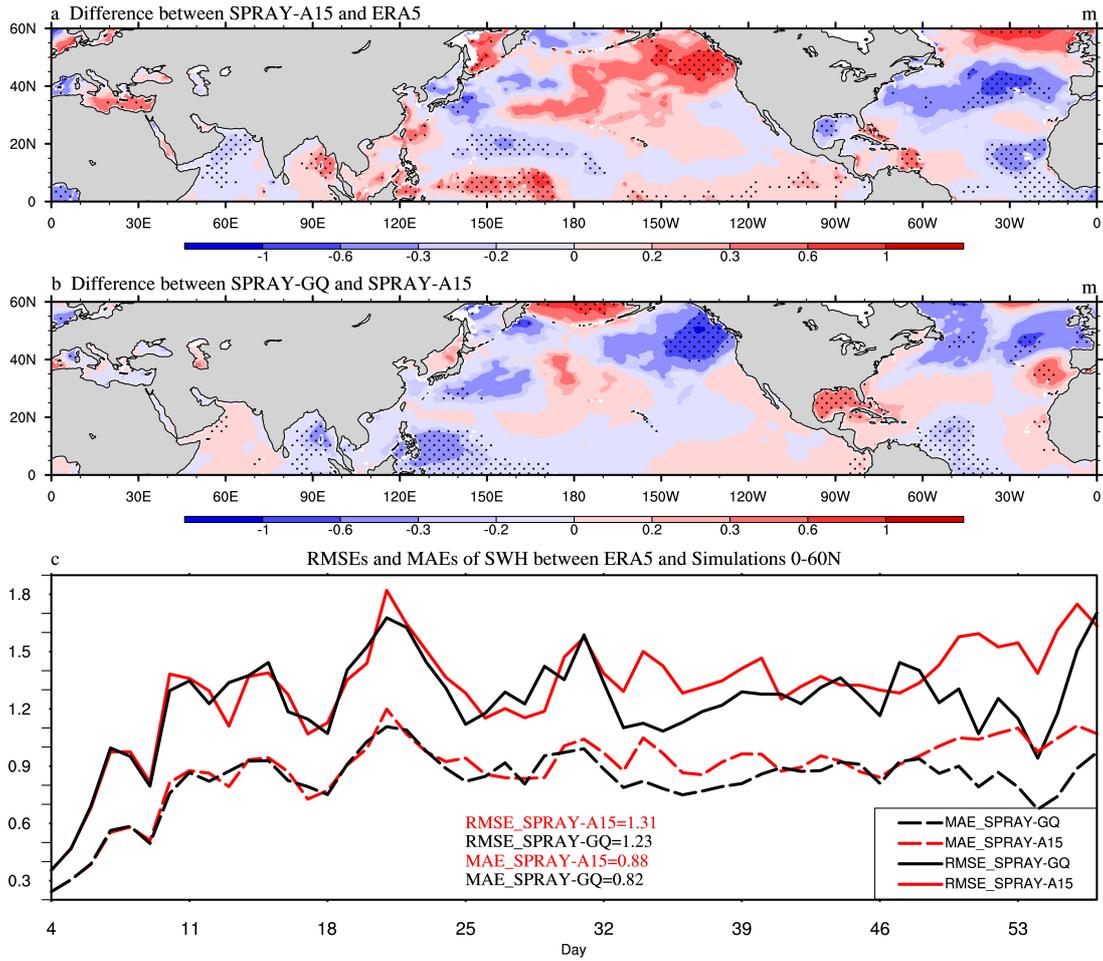


Figure 12. The 53-day average SWH (m) differences between SPRAY-A15 and ERA5 (a; SPRAY-A15 minus ERA5), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360°E, 0-60°N) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.

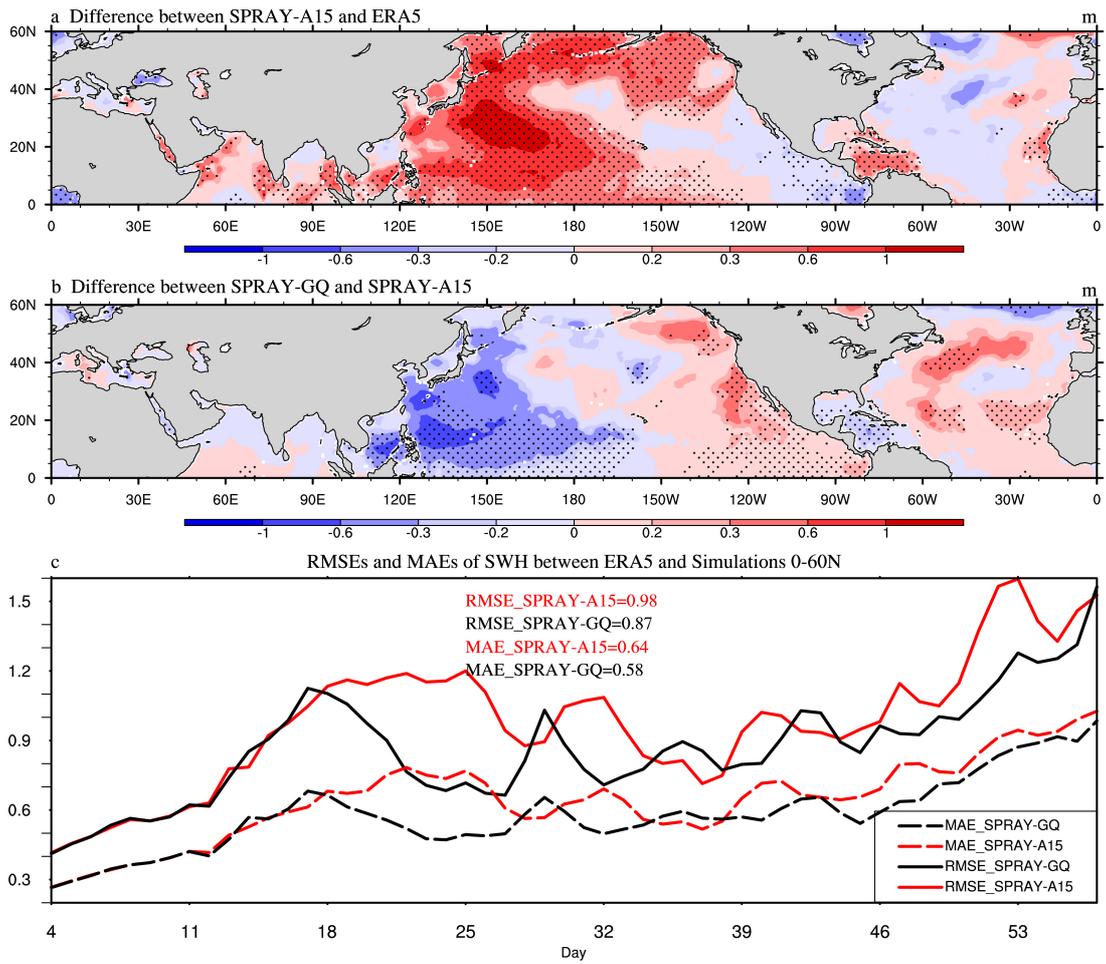


Figure 13. The same as Figure 12, but for Aug-Sep, 2018.