1 Accelerated Estimation An Improved Parameterization

of Sea Spray-Mediated Heat Flux Using Gaussian Quadrature: Case Studies with a Coupled CFSv2.0-WW3

4 System

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10 Abstract. Sea spray-mediated heat flux plays an important role in air-sea heat transfer. Heat flux 11 integrated over droplet size spectrum can well simulate total heat flux induced by sea spray droplets. 12 Previously, a fast algorithm of spray-flux scheme assuming single-radius droplets (A15) was widely used 13 since the full-size spectrum integral is <u>computationally</u> emputational expensive. Based on the Gaussian 14 Quadrature (GQ) method, a new fast algorithm scheme (SPRAY-GQ) of sea spray-mediated heat flux is 15 derived. The performance of SPRAY-GQ is evaluated by comparing heat fluxes with those estimated 16 from the widely-used A15. The new algorithm scheme shows a better agreement with the original 17 spectrum integral. To further evaluate the numerical errors performance of A15 and SPRAY-GQ, the 18 two algorithms schemes are implemented into a coupled CFSv2.0-WW3 system, and a series of 56-day 19 simulations in summer and winter are conducted and compared. The comparisons with satellite 20 measurements and reanalysis data show that the SPRAY-GQ algorithmscheme could simulate lead to 21 air-sea heat flux-more reasonable simulationy than the A15 algorithm scheme by modifying air-sea heat 22 flux. For experiments based on SPRAY-GQ, the sea surface temperature at mid-high latitudes of both 23 hemispheres, particularly in summer, is significantly improved compared with the experiments based on 24 A15. The simulation of 10-m wind speed and significant wave height at mid-low latitudes of the Northern 25 Hemisphere after the first two weeks is improved as well. The computational time of SPRAY-GQ is 26 about the same as that of A15. Thereby, the newly-developed SPRAY-GQ algorithm scheme-has a 27 potential to be used for improving air-seacalculation of ealeulate spray-mediated heat flux in coupled 28 models.

29

30 1 Introduction

31 Sea spray droplets, ejected from oceans, include film drops, jet drops and spume drops (Veron, 2015). 32 The first two types of droplets are generated from bubble bursting caused by ocean surface wave breaking, 33 with radius ranging from 0.5 µm to 50 µm (Resch and Afeti, 1991; Thorpe, 1992; Melville, 1996; Spiel, 34 1997; Andreas, 1998; Lhuissier and Villermaux, 2012). Spume drops are generated by strong winds (> 35 7-11 m/s) which directly tear the wave crests, with larger radius ranging from tens to hundredshundreds 36 of micrometers um (Koga, 1981; Andreas et al., 1995; Andreas, 1998). Sea spray droplets play an 37 important role in weather and climate processes (Fox-Kemper et al., 2022). On one hand, sea spray 38 droplets contribute to local marine aerosols and subsequently modify the local radiation balance (Fairall 39 et al., 1983; Burk, 1984; Fairall and Larsen, 1984). On the other hand, sea spray droplets affect the fluxes 40 of heat, momentum, salt, and freshwater between atmosphere and ocean (Andreas, 1992; Andreas et al., 41 2008; Andreas, 2010; Andreas et al., 2015; Ling and Kao, 1976; Fairall et al., 1994; Andreas and 42 Decosmo, 2002).

43 The sea spray-mediated heat transfer mainly occurs within the droplet evaporation layer (DEL) near 44 the sea surface (Andreas and Decosmo, 1999, 2002; Fairall et al., 1994). Sea spray droplets with the same 45 temperature as ocean surface can lead to sensible heat flux in DEL, while water evaporated from these 46 droplets can further release latent heat to the atmosphere (Andreas, 1992; Borisenkov, 1974; Bortkovskii, 47 1973; Wu, 1974; Monahan and Van Patten, 1988; Ling and Kao, 1976). Part of the sea spray-mediated 48 sensible heat is absorbed by droplet evaporation, which further increases the air-sea temperature 49 difference, and thus increases the sea spray-mediated sensible heat flux (Fairall et al., 1994; Andreas and 50 Decosmo, 2002). Since strong winds produce more sea spray droplets with larger radius, sea spray-51 mediated heat fluxes increase with wind speed (Fairall et al., 1994), and contribute more than 10% of the 52 total surface heat flux after reaching the threshold speed (> 11 m/s for sensible heat flux and > 13 m/s 53 for latent heat flux)(Andreas et al., 2008). In addition, when a droplet is released into the air, it is 54 accelerated due to surface winds (Edson and Andreas, 1997; Fairall et al., 1994; Van Eijk et al., 2011; 55 Wu et al., 2017). If the droplet could fall back into the ocean, additional momentum would be injected 56 into the ocean from the atmosphere (Andreas, 1992, 2004).

57 The usual bulk parameterizations in numerical models for surface fluxes only include the interfacial

58 (turbulent) fluxes (e.g., Fairall et al., 1996), while neglecting the significant contributions of sea spray 59 droplets in DEL (Andreas et al., 2008; Fairall et al., 1994; Smith, 1997; Emanuel, 1995). Andreas and 60 Emanuel (2001) implemented sea spray-mediated heat flux and momentum flux parameterizations into 61 a simple tropical cyclone model_, and found that the sea spray-induced heat flux significantly enhances 62 the tropical cyclone intensity, offsetting the negative effect of enhanced surface drag by strong wind and 63 waves. They, and found that the sea spray-mediated induced heat flux can significantly enhance tropical 64 cyclone intensity. It is well known that The strong winds and high waves induced by tropical cyclones 65 can enhance sea surface roughness and thus surface drag coefficients, which tend to reduce tropical 66 cyclone intensity (Emanuel, 1995). In additionFurthermore, the accelerated sea spray droplets by 67 surface winds also lead to more dissipation of tropical cyclone kinetic energy (Andreas, 1992, 2004). 68 Andreas and Emanuel (2001) found that the sea spray-induced heat flux significantly enhances the 69 tropical cyclone intensity, offsetting Tthese negative effects could be offset by the sea spray-70 mediatedinduced heat flux. The similar enhancement of tropical cyclone intensity was also shown 71 noticed in recent regional coupling systems by including sea spray-mediated heat flux (Xu et al., 2021b; 72 Liu et al., 2012; Garg et al., 2018; Zhao et al., 2017). In the First Institute of Oceanography Earth System 73 Model, Bao et al. (2020) first incorporated the sea spray-mediated heat flux in global climate simulation. 74 Following Bao et al. (2020), Song et al. (2022) found that the sea spray-mediated heat flux can lead to 75 cooling at the air-sea interface and strengthening westerlies in the Southern Ocean, and thus improves 76 estimates of sea surface temperature (SST).

77 Since the parameterization of sea spray-mediated heat flux derived from observations requires full-78 size spectral integral and thus is too computationallyer expensive intensive for large-scale 79 modelsdemands huge amount of computational time (Table 1, details in Section 4.2; Andreas, 1989, 80 1990, 1992; Andreas et al., 2015), a simplified algorithm parameterization-based on a single radius of 81 sea spray droplets (Andreas et al., 2015; Andreas et al., 2008) is widely used in atmosphere-ocean 82 coupling systems (Xu et al., 2021b; Liu et al., 2012; Garg et al., 2018; Zhao et al., 2017; Song et al., 83 2022; Bao et al., 2020), and apt to produce numerical errorsignificant biases. To reduce these numerical 84 errors biases induced by the single radius of sea spray droplets, we develop a new fast algorithm 85 parameterization of sea spray-mediated heat flux based on the Gaussian Quadrature (GQ) method, a fast 86 and accurate way to calculate spectral integral. The GQ method has been successfully used for the 87 estimation of domain-averaged radiative flux profiles (Li and Barker, 2018). The performance of the 88 GQ-based fast algorithm parameterization of the sea spray-mediated heat flux is evaluated and compared 89 with the simplified algorithm parameterization for single radius of Andreas et al. (2015), referred to as 90 A15 hereafter. The results are first compared with the original parameterization using full-size spectral 91 integral (A92, hereafter). Then the parameterizations with different algorithms, are implemented in a 92 global coupled atmosphere-ocean-wave system (Shi et al., 2022), and the results are compared with 93 global satellite measurements and reanalysis data.

The rest of the paper is structured as follows: observation and reanalysis data for comparisons are introduced in Section 2; the derivation of the GQ-based <u>fast algorithm parameterization</u> and the global coupling system are described in Section 3; the performance of the new <u>fast algorithm parameterization</u> is evaluated in Section 4. Finally, a summary and discussion are given in Section 5.

98 2 Data

99 The fifth generation European Centre for Medium-Range Weather Forecasts (ECMWF) Reanalysis 100 (ERA5; Hersbach et al., 2020) data assimilated huge amounts of historical data and thus provided reliable 101 hourly estimates. ERA5-10-m wind speed (WSP10), 2-m air temperature (T02), 2-m dewpoint 102 temperature, surface pressure and significant wave height (SWH) with a spatial resolution of 0.5° are 103 used. Additionally, WSP10, T022-m air temperature and 2-m specific humidity (SPH) data from the 104 Objectively Analyzed air-sea Fluxes (OAFlux) products (Yu et al., 2008) are also applied for comparison, 105 with 1°×1° resolution. The daily average satellite Optimum Interpolation SST (OISST) data are obtained 106 from the National Oceanic and Atmospheric Administration (NOAA) with a spatial resolution of 0.25° 107 (Reynolds et al., 2007). The global monthly mean salinity observations from European Space Agency 108 (ESA; https://climate.esa.int/sites/default/files/SSS cci-D1.1-URD-v1r4 signed-accepted.pdf—) are 109 applied. Besides, we also use the monthly global ocean RSS Satellite Data Products for WSP10 (https://data.remss.com/wind/monthly_1deg/) for WSP10 and the Reprocessed L4 Satellite 110 111 Measurements for SWH (https://doi.org/10.48670/moi-00177), to validate theour simulation results and 112 ERA5 data.

113 **3 Methods**

114 **3.1 Development of a Fast Algorithm Based on GQ**

115 The effects of sea spray droplets on sensible and latent heat fluxes $(H_{S,SP}, H_{L,SP})$ contribute to the total 116 <u>turbulent</u> sensible and latent heat fluxes $(H_{S,T}, H_{L,T})$ at the air-sea interface. That is,

$$H_{S,T} = H_S + H_{S,SP},\tag{1}$$

$$H_{L,T} = H_L + H_{L,SP}.$$
 (2)

where H_S and H_L are the sensible and latent heat fluxes at the air-sea interface due to the air-sea 117 118 differences of temperature and humidity. Based on observations of total turbulent heat fluxes and the 119 COARE algorithm (Andreas et al., 2015; Fairall et al., 1996)Based on eddy correlation observations, 120 A92 (Andreas, 1989, 1990, 1992; Andreas et al., 2015) integrates the sea spray-mediated sensible and 121 latent heat flux spectrums over initial droplet radius ($Q_S(r_0)$ and $Q_L(r_0)$) to estimate $H_{S,SP}$ and $H_{L,SP}$ 122 (details in Appendix A)(details in Appendix A; Andreas, 1989, 1990, 1992; Andreas and Decosmo, 2002). 123 The distributions of $Q_S(r_0)$ and $Q_L(r_0)$ spectrums as functions of initial droplet radius r_0 under 124 various atmosphere and ocean state are shown in Fig. 1, indicating that Q_S and Q_L spectrums are more 125 sensitive to the change of WSP1010 m wind speed, and less sensitive to other variables, including T022-126 m air temperature, 2-m relative humidity, sea surface temperatureSST, surface air pressure and sea 127 surface salinity.

The calculation of $H_{S,SP}$ and $H_{L,SP}$ in A92 requires huge amount of computational time is 128 129 computationally expensive due to full-size spectral integral (Eqn. A5-A6 of Appendix A), therefore it is 130 difficult to apply A92 directly in coupled modeling systems. A15 (Andreas et al., 2015) developed a fast 131 algorithm by using a single representative droplet radius (details in Appendix B), which was widely 132 adopted in recent reginalregional and global coupling systems (Xu et al., 2021b; Liu et al., 2012; Garg 133 et al., 2018; Zhao et al., 2017; Song et al., 2022; Bao et al., 2020). In this study, we apply a 3-node GQ 134 method (details in Appendix C) to develop a new fast algorithm to approximate the full-size spectral 135 integral of A92. Notably, GQ can converge exponentially to the actual integral only for a smooth function, 136 which is athe prerequisite for GQ (Mcclarren, 2018). Since as functions of r_0 , $Q_S(r_0)$ and $Q_L(r_0)Q_S$ 137 and Q_{E} are not smooth (Fig. 1), a data sorting from largest to smallest is required. After sorting, local

138 $Q_{S}(r_{0})$ and $Q_{L}(r_{0})Q_{S}$ and Q_{L} become $Q_{S_{sort}}(m)$ and $Q_{L_{sort}}(m)Q_{SS}$ and Q_{LS} , and then GQ can 139 be used to estimate the integral of $Q_{S_sort}(m)$ and $Q_{L_sort}(m)Q_{ss}$ and Q_{Ls} . Note that the independent 140 variable <u>m</u> is not equivalent to the original r_0 , but only indicates the position. In this way, according to 141 <u>Appendix C, m_1 =443, m_2 =251, m_3 =58 are three GQ nodes of $Q_{S \text{ sort}}(m)$ and $Q_{L \text{ sort}}(m)$, and we</u> 142 can get the corresponding r_0 for local $Q_S(Q_L)$, denoted as $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$. 143 However, the sorting leads to high complexity of GQ comparable to A92₂- and the values of r_{S1} (r_{L1}). 144 $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$ vary under various atmosphere and ocean environments in the globe. Therefore, 145 it is necessary to find the general approximate values law of $r_{S1}(r_{L1}), r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$ viaby 146 <u>global statistical analyses, GQ nodes for Q_{SS} and Q_{LS} to avoid the sorting in application.</u>

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148 To derive the general approximate values law of $r_{S1}(r_{L1})$, $r_{S2}(r_{L2})$ and $r_{S3}(r_{L3})$ GQ nodes, we 149 calculate the distribution of the sea spray-mediated heat flux spectral following A92, based on the global 150 daily WSP10, T022-m air temperature, 2-m dewpoint temperature, surface pressure and SWH of ERA5 151 and OISST from August 1, 2018 to August 31, 2018. Since the sea spray-mediated heat flux is not 152 sensitive to salinity (Fig. 1e&f) and only monthly observational data is available, the ESA monthly 153 salinity is applied. The ESA monthly salinity is also applied since the sea spray mediated heat flux is the 154 least sensitive to salinity (Fig. 1e&f) and only monthly salinity observation data is available. From the 155 global spectrums, we sort Q_s and Q_L from largest to smallest to obtain local. The GQ nodes 156 corresponding to r_0 of the sensible (latent) heat flux after sorting are denoted as r_{S1} , r_{S2} and r_{S3} (r_{L1} , 157 r_{L2} and r_{L3} for every each grid point, whose global distribution of occurrence frequency in percentage 158 is shown in Fig. 2. It is noted that except for that r_{L3} is related to WSP10 (Fig. 2c), all other five nodes have frequency roughly concentrated at a constant (peak frequency >65% in Fig. 2a, b, d-f; Eqn. 3&4), 159 160 while for r_{L3} , there is a 92.53% concentration between 55 and 90 μm (Fig. 2c). And tThen we found 161 that r_{L3} (55-90 µm) is related to WSP10 (Fig. S1 in supplementary), , that is thereby we set the 162 approximate values as

 $r_{s1} = 459.056, r_{s2} = 294.185, r_{s3} = 166.771,$ (3)

$$r_{L1} = 443.914, r_{L2} = 251.0498, \tag{4}$$

$$r_{L3} = \begin{cases} 60.310 \text{WSP10}^{0.1161}, \text{ WSP10} \ge 2 \text{ } m/s \\ 58.086, \text{ WSP10} < 2 \text{ } m/s \end{cases},$$
(5)

163 where the unit of the radius is micrometer. <u>Afterwards, we don't sort anymore, aAnd directly use Eqn.</u>

164 <u>3-5 then the 3-node GQ</u> to approximate the full-size spectral integral of A92 without sorting as-are

$$\int_{a}^{b} Q_{S}(r_{0}) dr_{0} \approx \frac{b-a}{2} \sum_{i=1}^{3} \omega_{i} Q_{S}(r_{Si}), \tag{6}$$

$$\int_{a}^{b} Q_{L}(r_{0}) dr_{0} \approx \frac{b-a}{2} \sum_{i=1}^{3} \omega_{i} Q_{L}(r_{Li}).$$
⁽⁷⁾

Here a and b are the lower and upper limits of r_0 , which are set to $2\mu m$ and $500\mu m$ based on Andreas (1990), and ω_i is the corresponding weight ($\omega_1 = \omega_3 = 0.556$, $\omega_2 = 0.889$), obtained from Mcclarren (2018). Thus, we can directly use Eqn. (3-7) to estimate the GQ-based $H_{S,SP}$ and $H_{L,SP}$ approximations, avoiding sorting. The new fast algorithm to for approximations estimate the GQ-based of $H_{S,SP}$ and $H_{L,SP}$ approximations is referred to as SPRAY-GQ hereafter.

170 **3.2 CFSv2.0-WW3 Coupling System**

171 A coupled system based on Climate Forecast System model version 2.0 (CFSv2.0) and 172 WAVEWATCH III (WW3) is employed to evaluate and compare the effects of sea spray-mediated heat 173 flux parameterized by A15 and SPRAY-GQ. The CFSv2.0-WW3 has three components, the Global 174 Forecast System (GFS; http://www.emc.ncep.noaa.gov/GFS/doc.php) as the atmosphere component of 175 CFSv2.0, the Modular Ocean Model version 4 (MOM4; Griffies et al., 2004) as the ocean component of 176 CFSv2.0, and the WW3 (WAVEWATCH III Development Group, 2016) as the ocean surface wave 177 component. The variables between CFSv2.0 and WW3 are interpolated and passed using the Chinese 178 Community Coupler version 2.0 (C-Coupler2; Liu et al., 2018).

The CFSv2.0 is mainly applied for intraseasonal and seasonal prediction (e.g., Saha et al., 2014). The atmosphere component GFS uses a spectral triangular truncation of 382 waves (T382) in the horizontal, equivalent to a grid resolution of nearly 35 km, and 64 sigma-pressure hybrid layers in the vertical. The MOM4 is integrated on a nominal 0.5° horizontal grid with enhanced horizontal resolution to 0.25° in the tropics, and there are 40 levels in the vertical. <u>The CFSv2.0 initial fields at 00:00 UTC of the first</u> day for experiments were generated by the real time operational Climate Data Assimilation System (Kalnay et al., 1996), downloaded from the CFSv2.0 official website
(http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod). The latitude range of WW3 is 78°S–78°N
with a spatial resolution of 1/3°. The initial wave fields arewere generated from 10-day simulation
starting from rest in a stand-alone WW3 model, forced by ERA5 10-m winds and ice concentration. The
open boundary conditions of WW3 arewere also obtained by the global simulation of the stand-alone
WW3 model.

In the coupling system, the WW3 obtains 10-m wind and ocean surface current from CFSv2.0, and then provides wave parameters to CFSv2.0. Several wave-mediated processes, including upper ocean mixing modified by Stokes drift-related processes, air-sea fluxes modified by surface current and Stokes drift, and momentum roughness length, are considered. Details of this system are referred to Shi et al. (2022).

196 A series of numerical experiments is conducted to evaluate the effects of the two fast algorithms 197 parameterizations (A15 and SPRAY-GQ) of sea spray-mediated heat flux on ocean, atmosphere and 198 waves in two 56-day periods, from January 3 to February 28, 2017 and from August 3 to September 28, 199 2018 for boreal winter and boreal summer, respectively. For each period, two sensitivity experiments are 200 carried out. The first is the SPRAY-A15 experiment, in which A15 is used with two-way fully coupling. 201 The second is the SPRAY-GQ experiment, in which SPRAY-GQ fast algorithm parameterization-is used 202 instead of A15. In addition, we also carry out another 7-day experiment using A92 (SPRAY-A92) to test 203 the runtime.

204 4 Results

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206 **4.1 Comparison with A92**

Based on the daily global WSP10, <u>T022-m air temperature</u>, 2-m dewpoint temperature, surface pressure and SWH of ERA5, the daily global OISST, and the ESA monthly global salinity, $H_{S,SP}$ and $H_{L,SP}$ from A15, SPRAY-GQ and A92 are calculated (Fig. 3). The computational time for SPRAY-GQ is about the same as that for A15, and about 36 times less than the time for A92. Compared with A92

211 (the black dotted line), A15 (red) overestimates $H_{S,SP}$ for low $H_{S,SP}$ (<50 W/m²) and underestimates 212 $H_{S,SP}$ for high $H_{S,SP}$ (>50 W/m²) with a root mean square error (RMSE= $\sqrt{\sum_{i=1}^{n} (\hat{y}_i - y_i)^2 / n}, \hat{y}_i$ is A15 213 value, y_i is A92 value, and *n* is the total number of grid points) of 3.40 W/m² (Fig. 3a), while A15 214 shows consistent overestimations with a RMSE of 2.98 W/m² for H_{LSP} (Fig. 3b). Overall, the RMSE 215 of A15 is about 2.69 W/m² for sea-spray mediated total heat flux ($TH_{SP} = H_{S,SP} + H_{L,SP}$; Fig. 3c). And reas 216 et al. (2015) derived A15 from A92 using single-radius droplets as bellwethers and wind functions, and 217 extrapolated the wind functions at high wind speeds >25 m/s. Since here the wind speeds in the study are 218 less than 25 m/s (Fig. S1), the large difference between A15 and A92 is mainly due to the use of single-219 radius droplets. Compared with A15, SPRAY-GQ (blue) has less deviation from A92 for both $H_{S,SP}$ and 220 $H_{L,SP}$ (Fig. 3a&b). The corresponding RMSEs of SPRAY-GQ for $H_{S,SP}$, $H_{L,SP}$ and TH_{SP} are 0.83 W/m², 0.92 W/m² and 0.62 W/m², all significantly lower (P<0.05 in Student's t-test) than those of A15. 221 222 To test robustness of the results, we also use WSP10, T022-m air temperature and 2-m SPHspecific 223 humidity of OAFlux dataset to estimate $H_{S,SP}$ and $H_{L,SP}$. As shown in Fig. 4, SPRAY-GQ has 224 significantly (P<0.05 in Student's t-test) lower deviations and RMSEs than A15, consistent with Fig. 3. 225 Note that the values of $H_{S,SP}$ and $H_{L,SP}$ in Fig.4 are larger than those in Fig. 3, since the equivalent 226 neutral wind speed from OAFlux is generally overestimated compared to the observed wind speed 227 (Seethala et al., 2021; Praveen Kumar et al., 2012). Because It is because OAFILux only provides neutral 228 wind speeds, calculated from wind stress and the corresponding roughness by assuming air is neutrally 229 stratified., Twhile previous studies indicated the neutral winds from OAFlux are larger than winds in 230 ERA5 as indicated by previous studies (Lindemann et al., 2021; Seethala et al., 2021). 231 In addition, since it is common to derive SWH from empirical equations (e.g., Andreas et al., 2008; 232 Andreas et al., 2015; Andreas and Decosmo, 2002; Andreas, 1992), we also use SWH generated by empirical equations of WSP10 (Andreas, 1992) instead of ERA5 SWH to estimate $H_{S,SP}$ and $H_{L,SP}$ 233 234 (Fig. 5). Again, the RMSEs decrease significantly (P<0.05 in Student's t-test) in SPRAY-GQ compared 235 to A15, though the RMSEs become higher for all estimates due to the enhanced biases of SWH. 236 Thereby, it is clear that the performance of SPRAY GQ is always better than A15_tThe difference 237 between SPRAY-GQ and A92 is always smaller than that between A15 and A92. Next, we will evaluate

and compare the two fast algorithms in an atmosphere-ocean-wave coupled system (CFSv2.0-WW3).

239 4.2 Comparison in the CFSv2.0-WW3 Coupling System

240 To compare the computational time of different parameterizations in the large-scale modeling system, 241 the runtime of the fully coupled experiments for 7-day forecast is given in Table 1 as an example. It is 242 shown that the runtime is about the same for SPRAY-GQ and SPRAY-A15. Both experiments run about 243 17 times faster than SPRAY-A92. 244 To illustrate the numerical errors of the two fast algorithms discussed in the context of the coupled 245 systemIn this section, comparisons are made for simulated SSTs, WSP10s as well as SWHs against 246 OISST and ERA5 reanalysis (Figs. 6-11), to present the numerical errors of the two fast algorithms 247 discussed in the context of the coupled system. The results in the first three days are excluded in the 248 comparison, since the wave influences are weak at the beginning of the simulations. Overall, the WSP10s 249 of simulations are generally in the range of 0-25 m/s globally. As shown in Fig. S2&S3 of 250 supplementary, Aat mid-high latitudes, the WSP10s generally can-exceed 10 m/s (Fig. S2&S3 of the 251 supplementary), the threshold at which the effects of sea spray can become significant (Andreas et al., 252 2015; Andreas et al., 2008). Overall, WSP10s of simulations are in the range of 0.25 m/s. Besides, we 253 calculate the runtime of the fully coupled experiments with different parameterizations for 7 day forecast 254 (Table 1). The runtime computational time is about the same for experiments SPRAY-GQ and SPRAY-255 A15, while the runtime of SPRAY GQ experiment is about 17 times less than the runtime of SPRAY-256 A92 experiment.

257 **4.2.1 Sea Surface Temperature (SST)**

258 In the austral summer, compared with OISST, large SST biases (>1 \degree C or <-1 \degree C) of SPRAY-A15 259 occur in the Southern Hemisphere (SH; Fig. S44a in supplementary), especially in the Southern Ocean. 260 It is always a challenge for reducing the large SST biases in the Southern Ocean for climate models (e.g., 261 Alessandro et al., 2019; Wang et al., 2014; Li et al., 2013; Bodas-Salcedo et al., 2012; Ceppi et al., 2012). 262 In Fig. 6a, SSTs north (south) of 50°S in experiment SPRAY-A15 are mainly underestimated 263 (overestimated). The domain-averaged RMSE (0-360°E, 40-75°S) increases in the first month and then 264 levels off (red solid line in Fig. 6c). While the domain-averaged RMSE in experiment SPRAY-GQ levels 265 off about a week earlier (black solid line in Fig. 6c). The time series of mean RMSE in SPRAY-GQ is 266 significantly lower than that in SPRAY-A15 (P<0.05 in Student's t-test). The increased (decreased) SSTs 267 north (south) of 50°S in SPRAY-GQ compared to those in SPRAY-A15 (Fig. 6b) reduce the RMSE of 268 <u>SST in SPRAY-GQ</u>. And the results of We also calculate the mean absolute error, $(MAE = \sum_{i=1}^{n} |\hat{y}_i - y_i|)$ 269 y_i / n , where \hat{y}_i is simulated value and y_i is OISST data, and n is the total number of grid points. The 270 MAEs) are consistent with RMSEs (dotted line in Fig. 6c). Furthermore, tThe eorresponding positive 271 mean errors, $(ME = \sum_{i=1}^{n} (\hat{y}_i - y_i)/n)$ in (-Fig. S5a in the supplementary), indicates the overall 272 overestimation, which areis reduced smaller in SPRAY-GQ than SPRAY-A15. The decreased SST 273 RMSE in SPRAY-GQ is resulted from the increased (decreased) SSTs north (south) of 50°S (Fig. 6b). 274 To understand the effects of sea spray droplets on SST, we calculate the total heat flux $(TH=H_{S,T}+H_{L,T})$ 275 differences between SPRAY-GQ and SPRAY-A15 (Fig. 12g7a). The TH differences are significantly 276 correlated with SST differences (Fig. S1b-S4b in the supplementary), with the spatial correlation 277 coefficient of -0.41 (P<0.05 in Student's t-test). We further decompose direct and indirect effects of sea 278 spray droplets on heat fluxes following Song et al. (2022). The direct effect ($H_{S,SP}$ and $H_{L,SP}$) is induced 279 directly by sea spray droplets, calculated from A15 (Eqn. B1-B4 of Appendix B) and SPRAY-GQ 280 (Section 3.1). The indirect effect (H_S and H_L) is the heat flux variation induced by changes of 281 atmosphere and ocean variables (including wind, pressure, humidity and temperature) caused by direct 282 effect, estimated by subtracting $H_{S,SP}$ and $H_{L,SP}$ from the output heat fluxes ($H_{S,T}$ and $H_{L,T}$) of 283 experiment SPRAY-A15 and SPRAY-GQ.

In the Southern Ocean, although direct differences of $H_{S,SP}$ and $H_{L,SP}$ are relatively small (<10 W/m², Fig. <u>12b7b</u>, e, &h), the resulting changes of temperature and humidity lead to relatively large differences in indirect effects of H_S and H_L (Fig. <u>12e7c</u>, f, &i). Enhanced (reduced) $TH_{S,SP}$ from ocean to atmosphere in the summer leads to increased (decreased) air-sea temperature difference and thus enhances (weakens) H_S . Meanwhile the warmer (cooler) air also causes more (less) evaporation and thus more (less) H_L . Finally, the enhanced (reduced) TH cools (warms) SST.

In the boreal summer, large SST biases (>1 °C or <-1 °C) of SPRAY-A15 mainly occur at mid-high
latitudes of the Northern Hemisphere (NH; Fig. <u>S2a-S6a</u> in supplementary). Significant underestimations
occur in the western and northern part of the North Pacific and at mid latitudes of the North Atlantic,
while large positive SST biases mainly occur in the eastern part of the North Pacific and at high latitudes

294 of the North Atlantic (Fig. 7a8a). In experiment SPRAY-GQ, SSTs are warmer (cooler) in the previously 295 underestimated (overestimated) regions (Fig. 7-b8b). Therefore, the domain-averaged RMSE and MAE 296 (0-360°E, 20-75°N) in SPRAY-GQ is are significantly lower (P<0.01 in Student's t-test) than in SPRAY-297 A15 after the first three weeks (Fig. 7e8c). Compared to And SPRAY-A15, the overall underestimation 298 is reduced in SPRAY-GQ-than SPRAY-A15 (Fig. S5b). The spatial correlation coefficient between TH 299 differences and SST differences (Fig. 13g9a&Fig. S2bS6b) is -0.32 (P<0.05 in Student's t-test). 300 Consistent with the austral summer, the SST changes are related to the changes of heat flux (Fig. 139). 301 The indirect effects of latent heat flux (Fig. 13e9f) play a major role in TH differences, which are 302 modified by the direct effects (Fig. 13b9b, e, &h). In addition, the changes of surface wind also contribute 303 to the changes of SST. The enhanced (reduced) winds lead to stronger (weaker) ocean mixing, and thus 304 eooler (warmer) SST (Fig. S3&S4). The reduced winds weaken the upper ocean mixing, so the water 305 becomes more stratified, and then the SST tends to be warmer, and vice versa (Fig. S7&S8).

306 4.2.2 10-m Wind Speed (WSP10) and Significant Wave Height (SWH)

307 Compared with experiment SPRAY-A15, significant improvements differences of WSP10 in SPRAY-308 GQ occur at mid-low latitudes of the NH (0-360°E, 0-60°N) in both winter and summer 309 (Fig. S7b& S8b9). As we know, satellite scatterometer and altimeter data are usually used to validate 310 WSP10 and SWH for short term weather forecast (e.g., Accadia et al., 2007; Djurdjevic and Rajkovic, 311 2008; Myslenkov et al., 2021). However, dDue to the spatial and temporal limitations coverage of satellite 312 data, we can only obtain the monthly averaged satellite data for the globe. So we compare the monthly 313 differences of averaged WSP10 and SWH from simulations over the periods between simulations 314 and with the corresponding satellite data (Fig. S9-S12). The average WSP10 and SWH differences 315 compared with satellite datacomparison results (Fig. S9a&c-S12a&c) are consistent with those compared 316 with ERA5 (Fig. S9b&d-S12b&d). From Fig. S9e-S12e, Besides, the differencesdomain averaged 317 RMSEs of WSP10s between ERA5 and the satellite data are always less than 1 m/s and are 0.48 m/s, 318 0.53 m/s, 0.15 m and 0.12 m (Fig. S9e S12e the differences of). SWHs are always less than 0.3 m. Since 319 ERA5 provides daily data for comparison and the differences between ERA5 and satellite data are small, 320 we will use ERA5 to validate simulations for validation in the following.

321 The domain averaged biasME of WSP10 (SPRAY-A15 minus ERA5) is 0.37-28 m/s and 0.24 47 m/s 322 in winter and summer (red in Fig. S5c&d), respectively, mainly due to the overestimations over the 323 Pacific and the Atlantic Ocean (red in Fig. 108a&119a). Whereas in SPRAY-GQ, the domain-averaged 324 biasME (SPRAY-GQ minus ERA5) is 0.26-15 m/s and 0.03-33 m/s in winter and summer respectively 325 (black in Fig. S5c&d). The domain-averaged RMSEs and MAEs of WSP10s increase with time in the first two weeks and then gradually level off (Fig. 108c&119c). The differences of WSP10 RMSEs and 326 327 MAEs between SPRAY-GQ (black) and SPRAY-A15 (red) are very small in the first two weeks. 328 Afterwards the mean values of time series of RMSE and MAE in SPRAY-GQ is are lower than that those 329 in SPRAY-A15 significantly at 959% confidence level in both boreal winter (Fig. 8e10c) and boreal 330 summer (Fig. 119c).

331 The simulated SWHs changes are closely related to the changes of WSP10₅ (Shi et al., 2022). 332 Therefore, the differences of SWHs (Fig. 120 ± 120 are consistent with those of WSP10s (Fig. 108 ± 119), 333 with overestimated (underestimated) WSP10s corresponding to overestimated (underestimated) SWHs 334 compared with ERA5. The SWHs in SPRAY-GQ showare significantly differencents with those in 335 SPRAY-A15The SWHs in SPRAY-GQ improve compared with those in SPRAY-A15 (Fig. 336 10b12b&131b), especially in summer. In winter (summer), the SWH RMSE averages for SPRAY-A15 337 and SPRAY-GQ are 1.31 m (0.98 m) and 1.23 m (0.87 m), and after the first two weeks the__time series 338 of RMSE and MAE in SPRAY-GQ is are lower than that those in SPRAY-A15 significantly at 959% 339 confidence level in both winter (Fig. 10e12c) and summer (Fig. 11e13c).

340 The direct and indirect effects of sea spray droplets on heat fluxes can influence estimates of WSP10 341 and then SWH. The changes of WSP10s are related to the direct effects ($H_{S,SP}$ and $H_{L,SP}$; Fig. 12b7b, e, 342 &h; Fig. 13b9b, e, &h). The spatial correlation coefficients between WSP10 differences (Fig. 343 S3bS7b&S4bS8b) and TH_{SP} differences (Fig. 12h7b&13h9b) are 0.51 and 0.69 (P<0.01 in Student's t-344 test) in winter and summer, respectively. Because TH_{SP} differences can influence the sea level pressure 345 (SLP) distribution (Fig. S15&S16), and thus influence subsequently surface winds. For example, 346 compared with SPRAY-A15, the decreased TH_{SP} of SPRAY-GQ in the Northwest Pacific in summer 347 (Fig. 9b) leads to higher SLP and smaller pressure gradient (Fig. S16), and thus decreased WSP10 (Fig. 348 11b); while the increased TH_{SP} in the Gulf of Alaska (Fig. 9b) leads to lower SLP and larger pressure

349 gradient (Fig. S16), and thus enhanced WSP10 (Fig. 11b). the directly increased (decreased) heat fluxes 350 enhance (reduce) turbulence, promote (hinder) the downward transmission of momentum from the upper 351 layer of atmosphere, and then accelerate (decelerate) the surface wind speed (Wallace et al., 1989). While 352 tThe accelerated (decelerated) WSP10s further result in increased (decreased) interfacial heat transport 353 (H_S , H_L), as well as increased (decreased) SWHs.

354 5 Conclusions and Discussions

355 Based on a GQ method, we develop a new fast algorithm based on Andreas's (1989, 1990, 1992) full-356 size microphysical parameterization (A92) for sea spray-mediated heat fluxes. Using global satellite 357 measurements and reanalysis data, we found that the difference between SPRAY-GQ and A92 is 358 significantly smaller than that between A15 and A92SPRAY-GQ parameterization is validated to 359 approximate A92 more accurately than the A15 fast algorithm (Andreas et al., 2015). To evaluate the 360 numerical error of SPRAY-GQ/A15 fast algorithmparameterization, we implement them in the two-way 361 coupled CFSv2.0-WW3 system. A series of 56-day simulations from January 3 to February 28, 2017 and 362 from August 3 to September 28, 2018 are conducted. The results are compared against OISST satellite 363 measurements and ERA5 reanalysis. The comparison shows that the sea spray-mediated heat flux in 364 SPRAY-GQ can reasonably modulate total heat flux compared with SPRAY-A15, and significantly 365 improve reduce the SST biases in the Southern Ocean (mid-high latitudes of the NH) for the austral 366 (boreal) summer, as well as WSP10 and SWH after the first two weeks at mid-low latitudes of the NH 367 for both boreal winter and summer. Overall, our fast algorithm based on GQ is applicable to sea spray-368 mediated heat flux parameterization in coupled models.

- 369 To investigate the effects of spray-mediated heat flux on simulations, two 56-day experiments without
- 370 sea spray effect (CTRL) in boreal winter summer and summer respectively winter are conducted added,
- 371 and the differences of simulated SST, WSP10, SWH, T02 and SPH between SPRAY-GQ and CTRL are
- 372 compared in Fig. S17-S21 in the supplementary. The introduction of sea spray cannot significantly
- 373 reduce the global overall errors of simulations, but it leads to regional improvements (blue in Fig.
- 374 S17e&f-S21e&f). For example, compared with CTRL in Jan-Feb, 2017, SST MAE of SPRAY-GQ in
- 375 the southeast of Australia decreases (Fig. S17e), because of warmer SST (Fig. S17c) related to reduced

376 wind (Fig. S18c). The reduced wind here also leads to lower SWH (Fig. S19c) and thus reduced SWH
377 overestimation (Fig. S19e). Meanwhile, T02 and SPH in CTRL are underestimated in this area (Fig.
378 S20a&S21a), while SPRAY-GQ reducesdeereases MAE of T02 and SPH (Fig. S20e&S21e) by
379 increasing temperature and moisture (Fig. S20c&S21c). Besides, tThe reduced errors are related to the
380 relatively large WSP10s over the corresponding areas (Fig. S2&S3), since the effects of sea spray become
381 significant important at wind speeds larger than 10 m/s.

382 In addition to the variables aforementioned, the changes of simulated cloud fraction were also 383 compared. However, the effects of sea spray-mediated heat flux on cloud fraction are non-significant for 384 the 2-month simulation, so the results are not shown. Besides, the lack of other processes related to sea 385 spray may be one of reasons why the global overall error cannot be reduced effectively. For example, for 386 simulated WSP10 and SWH in SPRAY-GQ, the SPRAY-GQ parameterization used in the study mainly 387 improves the biases at mid-low latitudes of the NH, while the significant overestimations in the SH are 388 only slightly improved still exist especially in Aug-Sep, 2018 (Fig. S3S18&-S196 in supplementary). As 389 Andreas (2004) indicated, sea spray droplets also influence the surface momentum flux by injecting more 390 momentum into the ocean from the atmosphere, which might further decrease the surface wind speed. 391 We will consider this process in the future study.

392 Sea spray-mediated heat fluxes are sensitive torelated to the sea spray generation function 393 $(SSGF)dF/dr_n$. Based on a number of laboratory and field observations, varieties of $SSGFdF/dr_n$ 394 were derived (e.g., Koga, 1981; Monahan et al., 1982; Troitskaya et al., 2018; Andreas, 1992, 1998, 2002; 395 Fairall et al., 1994; Veron, 2015), whereas their differences can reach six orders of magnitude (Andreas, 396 1998). There is currently no consensus on the most suitable choice. In this study, we use $SSGF \frac{dF}{dr_{\pi}}$ 397 of Fairall et al. (1994), recommended by Andreas (2002), to get a mean bias of 3.70 W/m² and 0.095 398 W/m² for latent and sensible heat flux respectively (Andreas et al., 2015). It is also, consistent with recent 399 observations of Xu et al. (2021a). (Andreas et al., 2015)Since the new scheme based on GQ is 400 independent of sea spray generation function, the new scheme can also be applied to sea spray mediated heat fluxes estimation with different dF/dro.Besides, considering the uncertainty of SSGF, the sea 401 402 spray mediated heat fluxes in A92 have been tuned by non-negative constants based on observations and 403 the COARE algorithm to reduce the uncertainties (Andreas and Decosmo, 2002; Andreas et al., 2008;

404 Andreas et al., 2015; Andreas, 2003). In this study, we use the constants (Eqn. A7 A8 in Appendix A) 405 for the SSGF (Fairall et al., 1994) to get a mean bias of 3.70 and 0.095 W/m² for latent and sensible heat 406 flux respectively in A92 compared to observations (Andreas et al., 2015). Therefore, a few W/m² 407 improvements of numerical errors in this study are relevant. Even though, the improved SST and other 408 variables cannot be reliably assigned to the usage of the GQ method, due to the uncertainties of the 409 coupled model itself and SSGF. 410 When wind speed is larger than 10 m/s, spray-mediated heat flux can become as important as the 411 interfacial heat flux (Andreas and Decosmo, 1999, 2002). Particularly, even in the absence of air-sea 412 temperature difference, the spray-mediated sensible heat flux is still present (Andreas et al., 2008). As 413 indicated by previous studies (e.g., Garg et al., 2018; Song et al. 2022), it is necessary to superimpose 414 the spray-mediated heat flux on the bulk formula to complete the physics of turbulent heat transfer for 415 coupled simulation. Since the full microphysical parameterization (A92) is computationally expensive, 416 an efficient algorithm that captures the main features of A92 can be beneficial to large-scale climate 417 systems or operational storm models. The GQ method proposed in the study can efficiently calculate the 418 spray-mediated heat flux, and agree better with A92 than A15. Thereby, Thus, the GQ based spray-419 mediated heat flux method has a great potential tois promising to be widely applied in large-scale climate 420 systems and operational storm models.

421

422 Appendix A

423 Microphysical Parameterization of A92

424 Based on the cloud microphysical parameterization of Pruppacher and Klett (1978), Andreas (1989,

425 1990, 1992) proposed a parameterization of sea spray-related heat fluxes for droplets with different radius,

426 from formation at sea surface to equilibrium with environment, that is,

$$Q_S = \rho_w C_{ps} \left(T_w - T_{eq} \right) \left[1 - \exp\left(-\frac{\tau_f}{\tau_T} \right) \right] \left(\frac{4\pi r_0^3}{3} \frac{dF}{dr_0} \right),\tag{A1}$$

$$Q_{L} = \begin{cases} \rho_{w}L_{v}\left\{1 - \left[\frac{r(\tau_{f})}{r_{0}}\right]^{3}\right\} \left(\frac{4\pi r_{0}^{3}}{3}\frac{dF}{dr_{0}}\right), \tau_{f} \leq \tau_{r}, \\ \rho_{w}L_{v}\left\{1 - \left(\frac{r_{eq}}{r_{0}}\right)^{3}\right\} \left(\frac{4\pi r_{0}^{3}}{3}\frac{dF}{dr_{0}}\right), \tau_{f} > \tau_{r}. \end{cases}$$
(A2)

Here Q_S , Q_L are sensible heat flux and latent heat flux resulted by sea spray droplets with initial radius 427 428 r_0 , ρ_w is the sea water density, C_{ps} is the specific heat, L_v is the latent heat of vaporization of water, 429 T_w is the water temperature, T_{eq} is the temperature of droplet when it reaches thermal equilibrium with 430 ambient condition, r_{eq} is the radius of droplet when it reaches moisture equilibrium with ambient 431 condition, τ_f is the residence time for droplets in the atmospheric, $r(\tau_f)$ is the corresponding radius, 432 τ_T is the characteristic e-folding time of droplet temperature, and τ_r is the characteristic e-folding time 433 of droplet radius. The detailed calculation of these microphysical quantities can be found in Andreas 434 (1989, 1990, 1992). And dF/dr_0 is the sea spray generation function, which represents the number 435 produced of droplets with initial radius r_0 (Andreas, 1992). For this term, the function of Fairall et al. 436 (1994) was recommended by Andreas (2002). According to the review in Andreas (2002), the dF/dr_0 437 of Fairall et al. (1994) is related on that of Andreas (1992) as

$$\frac{dF}{dr_{0}} = 38 \times 3.84 \times 10^{-6} U_{10}^{3.41} r_{0}^{-0.024} \frac{dF_{A92}}{dr_{80}} \Big|_{U_{10}=11 \text{ m/s}}, \quad (A3)$$

$$\frac{dF_{A92}}{dr_{80}} \Big|_{U_{10}=11 \text{ m/s}} =$$

$$\left(e^{(4.405-2.646(logr_{80})-3.156(logr_{80})^{2}+8.902(logr_{80})^{3}-4.482(logr_{80})^{4})}, r_{80} \le 15\mu m; \\ 1.02 \times 10^{4} r_{80}^{-1}, 15 \le r_{80} \le 37.5\mu m; \\ 6.95 \times 10^{6} r_{80}^{-2.8}, 37.5 \le r_{80} \le 100\mu m; \\ 1.75 \times 10^{17} r_{80}^{-8}, r_{80} \ge 100\mu m \right)$$

438 Here U_{10} is the 10-m wind, $r_{80} = 0.518 r_0^{0.976}$.

439 The total sea spray fluxes are obtained by integrating Q_s and Q_L corresponding to all r_0 . Based on 440 Andreas (1990), the lower and upper limits of r_0 is $2\mu m$ and $500\mu m$, that is,

$$\overline{Q_S} = \int_2^{500} Q_S(r_0) dr, \tag{A5}$$

$$\overline{Q_L} = \int_2^{500} Q_L(r_0) dr.$$
(A6)

While $\overline{Q_S}$ and $\overline{Q_L}$ are nominal sea spray fluxes but not the actual $H_{S,SP}$ and $H_{L,SP}$ (Andreas and Decosmo, 1999, 2002), because there are interactions between these two terms and the microphysical functions also lead to uncertainties (Fairall et al., 1994). Therefore, $\overline{Q_S}$ and $\overline{Q_L}$ are tuned by nonnegative constants α , β and γ (Andreas and Decosmo, 2002; Andreas et al., 2008; Andreas et al., 2015; Andreas, 2003) as

$$H_{S,SP} = \beta \overline{Q_S} - (\alpha - \gamma) \overline{Q_L}, \tag{A7}$$

$$H_{L,SP} = \alpha \overline{Q_L}.$$
 (A8)

446 In Eqn. (A8), the α term indicates the sea spray-mediated latent heat flux from the top of DEL to 447 atmosphere. Because the evaporation of droplets absorbs heat, which is provided by sea spray-mediated 448 sensible heat (Fairall et al., 1994), the negative α term appears in Eqn. (A7). Whereas the evaporation 449 also cools DEL and thus increases the air-sea temperature difference, therefore it contributes to a positive 450 γ term in Eqn. (A7). Different values of α , β and γ were given in Andreas and Decosmo (2002), 451 Andreas (2003), Andreas et al. (2008) and Andreas et al. (2015), to minimize the bias between 452 estimations and observations of turbulent heat fluxes measured by eddy correlation. And Andreas et al. 453 (2015) validated the most observation data, which are 4000 sets, to derive $\alpha = 2.46, \beta = 15.15, \gamma =$ 454 1.77.

455 Appendix B

456 Fast Algorithm of A15

457 Andreas (2003) and Andreas et al. (2008, 2015) developed a fast algorithm to approximate $H_{S,SP}$, 458 $H_{L,SP}$ by a characteristic radius, that is,

$$H_{S,SP} = \beta \overline{Q_S} - (\alpha - \gamma) \overline{Q_L} \approx \rho_w C_{ps} (T_W - T_{eq,100}) V_s(u_*),$$
(B1)
$$H_{L,SP} = \alpha \overline{Q_L} \approx \rho_w L_v \left\{ 1 - \left[\frac{r(\tau_{f,50})}{50 \mu m} \right]^3 \right\} V_L(u_*).$$
(B2)

Here $T_{eq,100}$ is T_{eq} of droplets with $r_0=100 \ \mu\text{m}$, $\tau_{f,50}$ is τ_f of droplets with $r_0=50 \ \mu\text{m}$, V_s and V_L are functions of the bulk friction velocity u_* . As indicated by Andreas et al. (2008, 2015), the characteristic radiuses of 100 μm and 50 μm for sensible and latent heat fluxes are chosen, respectively, because Q_s and Q_L show a large peak in the vicinity of these values (Fig. 1). V_s and V_L are calculated in Andreas et al. (2015) as

$$V_{S} = \begin{cases} 3.92 \times 10^{-8}, & 0 \le u_{*} \le 0.1480 \ m/s\\ 5.02 \times 10^{-6} u_{*}^{2.54}, & u_{*} \ge 0.1480 \ m/s, \end{cases}$$
(B3)

$$V_L = \begin{cases} 1.76 \times 10^{-9}, & 0 \le u_* \le 0.1358 \ m/s \\ 2.08 \times 10^{-7} u_*^{2.39}, & u_* \ge 0.1358 \ m/s \end{cases}$$
(B4)

464 Appendix C

465 Gaussian Quadrature (GQ)

GQ is a method to approximate the definite integral of a function f(x) via the function values at a small number of specified nodes (Gauss, 1815; Jacobi, 1826). In this study we use the form of n-node Gauss-Legendre quadrature on [-1, 1] as

$$\int_{-1}^{1} f(x) dx \approx \sum_{i=1}^{n} \omega_i f(x_i).$$
(C1)

469 Here x_i is the specified node, and ω_i is the corresponding weight. For n=3, x_1 =-0.775, x_2 =0, 470 x_3 =0.775, ω_1 = ω_3 =0.556, ω_2 =0.889.

471 While for a function
$$g(\xi)$$
 on [a, b], Eqn. (C1) can be transformed to

$$\int_{a}^{b} g(\xi) d\xi = \int_{-1}^{1} g\left(\frac{b-a}{2}x + \frac{a+b}{2}\right) \frac{d\xi}{dx} dx$$

$$\approx \frac{b-a}{2} \sum_{i=1}^{n} \omega_{i} g\left(\frac{b-a}{2}x_{i} + \frac{a+b}{2}\right).$$
(C2)

472 Code and data availability

473 The code of sea spray can be found under https://doi.org/10.5281/zenodo.7100345 or 474 https://zenodo.org/record/7100345#.Y66vRtVByHt (Shi and Xu, 2022). The code for CFSv2.0-WW3 475 system can be found under https://doi.org/10.5281/zenodo.5811002 (Shi et al., 2021) including the 476 coupling, preprocessing, run control and postprocessing scripts. The initial fields for CFSv2.0 are 477 generated by the real time operational Climate Data Assimilation System, downloaded from the CFSv2.0 478 official website (http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod). The daily average satellite 479 Optimum Interpolation SST (OISST) data are obtained from NOAA (https://www.ncdc.noaa.gov/oisst). 480 The fifth generation European Centre for Medium-Range Weather Forecasts (ECMWF) Reanalysis 481 (ERA5) are available at the Copernicus Climate Change Service (C3S) Climate Date Store 482 (https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels). The daily Objectively 483 Analyzed air-sea Fluxes (OAFlux) products are available at https://oaflux.whoi.edu/heat-flux. The global 484 monthly mean salinity observations of European Space Agency (ESA) are from https://climate.esa.int. 485 The monthly global ocean RSS Satellite Data Products for 10-m wind speed are from 486 https://data.remss.com/wind/monthly ldeg/, and the Reprocessed L4 Satellite Measurements for

487 significant wave height are from https://doi.org/10.48670/moi-00177.

488 Author contribution

489 FX and RS designed the experiments and RS carried them out. RS developed the code of coupling 490 parametrizations and produced the figures. RS prepared the manuscript with contributions from all co-491 authors. FX contributed to review and editing.

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497 Competing Interests

498 The contact author has declared that neither they nor their co-authors have any competing interests.

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7-day Forecast	Runtime (h)
SPRAY-A92	126.94
SPRAY-A15	7.60
SPRAY-GQ	7.67

Table 1. The runtime of CFSv2.0-WW3 global experiments for 7-day forecast with different parameterizations.



Figure 1. The radius-specific sea spray-mediated sensible (Q_S ; black) and latent (Q_L ; red) heat fluxes as functions of initial radius r_0 : U₁₀, Ta, RH, Tw, P and S are 10-m wind speed, 2-m air temperature, 2-m relative humidity, sea surface temperature, surface air pressure and surface salinity, respectively.



Figure 2. The distribution of occurrence frequency in percentage for GQ radius nodes: (a) the first node of latent heat flux; (b) the second node of latent heat flux; (c) the third node of latent heat flux; (d) the first node of sensible heat flux; (e) the second node of sensible heat flux; (f) the third node of sensible heat flux. The peak frequencies are marked.



Figure 3. Scatter plots of $H_{S,SP}$ (a), $H_{L,SP}$ (b) and total heat flux $TH_{SP} = H_{S,SP} + H_{L,SP}$ (c) estimated by fast algorithms (y-axis) vs those estimated by spectral integral in microphysical parameterization (x-axis): The dotted black line is y=x. The corresponding RMSEs are marked in the upper left corner.



Figure 4. The same as Figure 3, but WSP10, 2-m air temperature and 2-m specific humidity of OAFlux are used.



Figure 5. The same as Figure 4, but SWH is derived by WSP10 instead of ERA5 SWH.



Figure 6. The 53-day average SST (°C) differences between SPRAY-A15 and OISST (a; SPRAY-A15 minus OISST), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360°E, 40-75°S) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.



Figure 7. The 53-day average differences of total heat flux (a-c), latent heat flux (d-f), and sensible heat flux (g-i) between SPRAY-GQ and SPRAY-A15 (SPRAY-GQ minus SPRAY-A15) in Jan-Feb, 2017. The direct differences indicate sea spray-mediated heat flux differences (b, e, h), and the indirect differences indicate interfacial (bulk) heat flux differences resulted by sea spray (c, f, i). The dotted areas are statistically significant at 95% confidence level. A positive value of flux indicates an upward direction.



Figure 8. The same as Figure 6, but for Aug-Sep, 2018 in 0-360°E, 20-75°N.



Figure 9. The same as Figure 7, but for Aug-Sep, 2018.



Figure 10. The 53-day average WSP10 (m/s) differences between SPRAY-A15 and ERA5 (a; SPRAY-A15 minus ERA5), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360°E, 0-60°N) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.



Figure 11. The same as Figure 10, but for Aug-Sep, 2018.



Figure 12. The 53-day average SWH (m) differences between SPRAY-A15 and ERA5 (a; SPRAY-A15 minus ERA5), the differences between SPRAY-GQ and SPRAY-A15 (b; SPRAY-GQ minus SPRAY-A15), and the time series of domain-averaged RMSE and MAE (c; 0-360°E, 0-60°N) in Jan-Feb, 2017. The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.



Figure 13. The same as Figure 12, but for Aug-Sep, 2018.