

1 **Improving Madden–Julian Oscillation Simulation in Atmospheric General**  
2 **Circulation Models by Coupling with Snow–Ice–Thermocline One-dimensional**  
3 **Ocean Model**

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13

**14 Abstract**

15 A one-column turbulent kinetic energy-type ocean mixed-layer model snow–ice–  
16 thermocline (SIT) when coupled with three atmospheric general circulation models  
17 (AGCMs) yield superior Madden–Julian Oscillation (MJO) simulation. SIT is designed  
18 to have fine layers similar to those observed near the ocean surface; therefore, it can  
19 realistically simulate the diurnal warm layer and cool skin. This refined discretization of  
20 the near ocean surface in SIT provides accurate sea surface temperature (SST) simulation;  
21 thus, facilitating realistic air–sea interaction. Coupling SIT with European Centre  
22 Hamburg Model, Version 5, Community Atmosphere Model, Version 5, High-Resolution  
23 Atmospheric Model significantly improved MJO simulation in three coupled AGCMs  
24 compared with the AGCM driven with prescribed SST. This study suggests two major  
25 improvements to the coupling process. First, during the preconditioning phase of MJO  
26 over the Maritime Continent (MC), the over underestimated surface latent heat bias in  
27 AGCMs can be corrected. Second, during the phase of strongest convection over MC, the  
28 change in intraseasonal circulation in the meridional circulation enhancing near-surface  
29 moisture convergence is the dominant factor in the coupled simulations relative to the  
30 uncoupled experiments. The study results show that a fine vertical resolution near the  
31 surface, which better captures temperature variations in the upper few meters of the ocean,  
32 considerably improves different models with different configurations and physical  
33 parameterization schemes; this could be an essential factor for accurate MJO simulation.

34 **Keywords:** Madden–Julian Oscillation, coupling, warm layer

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36 **Short summary (plain text)**

37       We show that coupling a high-resolution one-column ocean model to three  
38 atmospheric general circulation models dramatically improves Madden–Julian  
39 Oscillation (MJO) simulations. It suggests two major improvements to the coupling  
40 process in the preconditioning and strongest convection phases over the Maritime  
41 Continent. Our results demonstrate a simple but effective way to significantly improve  
42 MJO simulation and potentially seasonal to subseasonal prediction.

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44

## 45 **1 Introduction**

46       The Madden–Julian Oscillation (MJO) is the dominant pattern of atmospheric  
47 intraseasonal variability in the tropics (Madden and Julian, 1972; Zhang, 2005; Jiang et  
48 al., 2020). It has been reported that the MJO convection is often observed over sea surface  
49 temperature (SST) of greater than 28°C in the Indo-Pacific warm pool (Salby and Hendon,  
50 1994). MJO is an eastward-propagating ocean–atmosphere and convection-circulation  
51 coupled phenomenon that lasts for 20–100 days. On these timescales, low-level moisture  
52 convergence, warm SST, and shallow upper-ocean mixed-layer depth precede the  
53 eastward propagation of organized deep convection by approximately ten days; opposite  
54 conditions followed by approximately ten days (Krishnamurti et al., 1988; Hendon and  
55 Salby, 1994; Woolnough et al., 2000). Heat flux exchange between the atmosphere and  
56 ocean modulates the intraseasonal oscillation (Shinoda and Hendon, 1998). Studies have  
57 emphasized the importance of moisture and heat flux feedback in MJO (Sobel et al., 2008,  
58 2010; DeMott et al., 2015). Besides, oceanic wave dynamics are suggested to be  
59 associated with MJO, for example, zonal wind stress anomalies driven by the MJO force  
60 eastward-propagating oceanic equatorial Kelvin wave (Hendon et al., 1998; Webber et  
61 al., 2010), and the signals could extend as deep as 1500 m in the ocean (Matthews et al.,  
62 2007). Furthermore, the westward-propagating oceanic equatorial Rossby wave can  
63 initiate the next MJO in the Indian Ocean (Webber et al., 2010; Webber et al., 2012).  
64 Evidence of oceanic intraseasonal signals coupling with atmospheric signals was  
65 observed in terms of the sea level, surface heat flux, salinity, and temperature during field  
66 experiments and in situ monitoring (Oliver and Thompson, 2011; Drushka et al., 2012;  
67 Wang et al., 2013; Chi et al., 2014; Matthews et al., 2014; DeMott et al., 2015; Fu et al.,  
68 2015).

69       Recent modeling studies have demonstrated that most coupled models could  
70 improve MJO simulations but that the ocean driven by the atmosphere contributes  
71 indirectly by improving the mean state, heat flux, fresh water, and momentum. DeMott  
72 et al. (2016) estimated that direct SST-driven ocean feedback contributes to the MJO  
73 propagation up to 10% by a change in column moisture. A comparison of the direct and  
74 indirect effects of SST indicated that direct effects, such as SST-driven surface fluxes,  
75 tend to offset wind-driven fluxes (DeMott et al., 2015; DeMott et al., 2016; DeMott et al.,  
76 2019). The factor of indirect ocean feedback on the atmospheric physical process includes  
77 strong MJO convection can amplify the radiative feedback to MJO convections  
78 associated with large cloud systems (Del Genio and Chen, 2015). The SST gradients can  
79 drive the MJO low-level convergence (Hsu and Li, 2012; Li and Carbone, 2012) and  
80 destabilize lower tropospheric to further enhance low-level convergence to the east of  
81 MJO convergence (Wang and Xie, 1998; Marshall et al., 2008; Benedict and Randall,  
82 2011; Fu et al., 2015). Many observational and model studies have reported that coupled  
83 feedback enhances the MJO with strong horizontal moisture advection, driven by sharp  
84 mean near-equatorial meridional moisture gradients (DeMott et al., 2015; Jiang et al.,  
85 2018; DeMott et al., 2019; Jiang et al., 2020). These findings suggest that high-frequency  
86 SST perturbations could improve moisture convergence efficiency and enhance MJO  
87 propagation through relatively smooth background moisture distribution.

88       Tseng et al. (2015) identified the key role of the upper-ocean warm layer in  
89 improving the MJO eastward propagation simulation using the European Centre  
90 Hamburg Model, Version 5 (ECHAM5), coupled with the one-column ocean model  
91 named snow–ice–thermocline (SIT). Many observational (Drushka et al., 2012; Chi et al.,  
92 2014) and modeling (Klingaman and Woolnough, 2013; DeMott et al., 2019; Klingaman

93 and Demott, 2020) studies have supported this hypothesis. However, coupling the SIT to  
94 only one atmospheric general circulation model (AGCM) may be insufficient to prove  
95 the effectiveness of the coupling. In this study, we coupled the SIT to three AGCMs:  
96 ECHAM5, Community Atmosphere Model, Version 5 (CAM5), and High-Resolution  
97 Atmospheric Model (HiRAM). We also discussed the coupling mechanism that leads to  
98 simulation improvement.

99 The remainder of the paper is organized as follows. In section 2, we describe the  
100 models, experimental designs, and observational data. Section 3 and 4 present the results  
101 and discussion, respectively.

## 102 **2 Data, model experiments and methodology**

### 103 **2.1 Observation and atmospheric/oceanic data**

104 Observational data used in this study include precipitation from Global Precipitation  
105 Climatology Project V1.3 (GPCP, 1° resolution, 1997–2010; Adler et al., 2003), outgoing  
106 longwave radiation (OLR, 1° resolution, 1997–2010; Liebmann, 1996) and daily SST  
107 (Optimum Interpolated SST, 0.25° resolution, 1989–2010; Banzon et al., 2014) from the  
108 National Oceanic Atmosphere Administration. The in situ ocean temperature profiles  
109 from 1989 to 2010 were obtained from the Tropical Ocean Global Atmosphere program  
110 (McPhaden et al., 2010).

111 Atmospheric variables were obtained from the European Centre for Medium-range  
112 Weather Forecast Reanalysis-interim (Dee et al., 2011) from 1989 to 2010. The variables  
113 include zonal wind, meridional wind, temperature, specific humidity, sea level pressure,  
114 geopotential high, latent heat, sensible heat, shortwave and longwave radiation. Oceanic

115 temperature data from 1989 to 2010 were obtained from the NCEP Global Ocean Data  
116 Assimilation System (GODAS) (Behringer and Xue, 2004) provided by the  
117 NOAA/OAR/ESRL PSL, Boulder, Colorado, USA  
118 (<https://psl.noaa.gov/data/gridded/data.godas.html>).

## 119 **2.2 Model experiments**

120 In this study, we coupled the one-column ocean model SIT (Tu and Tsuang, 2005;  
121 Tsuang et al., 2009) to three AGCMs. SIT simulates variations in the SST and upper-  
122 ocean temperature, including the diurnally-varying cool skin and warm layer in the upper  
123 few meters of the ocean and the turbulent kinetic energy (TKE; Gaspar et al., 1990) in  
124 the water column (Tsuang et al., 2001; Tu and Tsuang, 2005; Tu, 2006; Tsuang et al.,  
125 2009; Tu and Tsuang, 2014; Tseng et al., 2015; Lan et al., 2021). Cool skin is a very thin  
126 layer that has a direct contact with the atmosphere and warm layer is the warmer sea water  
127 immediately below the cool skin in the top few meters of the ocean. They fluctuate  
128 diurnally in response to atmospheric forcing. SIT with high vertical resolution  
129 realistically simulates the warm-layer (within top 10 m) and cool-skin (the top layer with  
130 0.001 m thickness), and improve the simulation of upper ocean temperature (Tu and  
131 Tsuang, 2005; Tsuang et al., 2009). The model has been verified at a tropical ocean site  
132 (Tu and Tsuang, 2005), in the South China Sea (Lan et al., 2010), and Caspian Sea  
133 (Tsuang et al., 2001). The melt and formation of snow and ice above a water column have  
134 been introduced (Tsuang et al., 2001). The three AGCMs used in this study are as follows.  
135 ECHAM5, the fifth-generation AGCM developed at the Max Planck Institute for  
136 Meteorology (Roeckner, 2003; Roeckner et al., 2006) is a spectral model that employs  
137 the Nordeng (Nordeng, 1994) cumulus convective scheme. We used a horizontal  
138 resolution of T63 (approximately 2°) with 31 vertical layers and a model top at 10 hPa

139 (approximately 30 km). The second one is NCAR Community Atmospheric Model  
140 version 5 (Hurrell et al., 2013) from the National Center for Atmospheric Research. We  
141 used a horizontal resolution of approximately  $1.875^\circ$  latitude  $\times$   $2.5^\circ$  longitude and 30  
142 vertical layers with the Zhang–McFarlane deterministic convection scheme (Zhang and  
143 McFarlane, 1995) and the University of Washington Shallow Convection (Park and  
144 Bretherton, 2009). HiRAM was developed based on Geophysical Fluid Dynamical  
145 Laboratory global atmosphere and land model AM2 (Team et al., 2004; Zhao et al., 2009)  
146 with few modifications (Chen et al., 2019). We used a horizontal resolution of  
147  $0.5^\circ$  latitude  $\times$   $0.5^\circ$  longitude with 32 vertical levels. For boundary layer and free  
148 atmospheric turbulence, the model adopted the 2.5 order parameterization of Mellor and  
149 Yamada (1982). Surface fluxes are computed based on the Monin–Obukhov similarity  
150 theory, given the atmospheric model’s lowest level of wind, temperature, and moisture.

151       There are 42 vertical layers in SIT, with 12 layers in the upper 10 m: the surface,  
152 0.05 mm, 1 m, 2 m, 3 m, 4 m, 5 m, 6 m, 7 m, 8 m, 9 m, and 10 m below ocean surface.  
153 The fine resolution was designed to realistically simulate the upper-ocean warm layer,  
154 including a layer at 0.05 mm, reproducing the cool skin of the ocean surface. It is worth  
155 noting that coupling a high-vertical-resolution TKE ocean model with an AGCM is  
156 unconventional. To account for neglected horizontal processes, the model ocean was  
157 weakly nudged (with a 30-day time scale) to the observed GODAS monthly mean ocean  
158 temperature below a depth of 10 m. Nudging was not applied in the upper 10 m. The  
159 timestep of SIT and AGCMs exchange ocean surface fluxes varying associated with the  
160 model resolution, which is 720, 1800, and 900 seconds in ECHAM-SIT, CAM5-SIT, and  
161 HiRAM-SIT, respectively. AGCMs were coupled with the SIT in the tropical region



162 between 30°S and 30°N and forced by prescribed monthly mean OISST outside this  
163 tropical belt.

164 The experiments comprised three sets of coupled AGCM simulations (ECHAM5-  
165 SIT, CAM5-SIT, and HiRAM-SIT) and standalone AGCM simulations forced by  
166 observed monthly mean OISST (ECHAM5, CAM5, and HiRAM) from 1985 to 2005.  
167 The experiments were designed to evaluate the effect of atmosphere–ocean coupling on  
168 MJO simulations. Table 1 presents the model and experiment details.

### 169 **2.3 Methodology**

170 The analysis focused on the boreal cool season (November–April) when the  
171 eastward propagation tendency of the MJO is the most prominent. We used the CLIVAR  
172 MJO Working Group diagnostics package (CLIVAR, 2009) and a 20–100-day filter to  
173 analyze intraseasonal variability. The MJO phase composites were computed using the  
174 real-time multivariate MJO index (Wheeler and Hendon, 2004), defined as the leading  
175 pair of principal components of intraseasonal OLR, and 850 and 200 hPa zonal winds in  
176 the tropics.

177 The vertically integrated MSE budget was diagnosed based on the following  
178 equation:

$$179 \left\langle \frac{\partial h}{\partial t} \right\rangle = - \left\langle u \frac{\partial h}{\partial x} \right\rangle - \left\langle v \frac{\partial h}{\partial y} \right\rangle - \left\langle \omega \frac{\partial h}{\partial p} \right\rangle + \langle LW \rangle + \langle SW \rangle + \langle LH \rangle + \langle SH \rangle \quad (1)$$

180 where  $h$  is the MSE ( $h = cpT + gz + Lq$ );  $u$  and  $v$  are the zonal and meridional velocities,  
181 respectively;  $\omega$  is the vertical pressure velocity;  $LW$  and  $SW$  are the longwave and  
182 shortwave radiation fluxes, respectively;  $LH$  and  $SH$  are the latent and sensible surface  
183 heat fluxes, respectively. The mass–weighted vertical integration from the surface to 200  
184 hPa is denoted as  $\langle \cdot \rangle$ , and intraseasonal anomalies are represented as  $\langle \cdot \rangle'$ , which were  
185 isolated using a 20–100-day bandpass Lanczos filter (Duchon, 1979).

186

187 **3 Results**188 **3.1 MJO simulations: ECHAM5-SIT, CAM5-SIT, and HiRAM-SIT**189 **3.1.1 General structure**

190 We compared simulated MJO characteristics using three coupled and uncoupled  
191 AGCMs. Figure 1 shows the wavenumber–frequency spectra of simulated 850-hPa zonal  
192 wind (shading) and precipitation (contours). All three uncoupled AGCMs (hereafter  
193 referred to as AGCMs) simulated intraseasonal signals with lower frequency than the  
194 observed and overestimated the westward propagation with periods greater than 80 days  
195 (Figs. 1e–g). The ECHAM5 and HiRAM simulated signals of wavenumbers 1–3 instead  
196 of the observed wavenumber 1 in 850-hPa zonal wind. These results show that all three  
197 AGCMs simulated stationary fluctuations with low frequency that were not consistent  
198 with the observation. By contrast, coupled AGCMs realistically reproduce the observed  
199 spectral characteristics and strength of the eastward propagation at wavenumbers 1 to 2  
200 in 850-hPa zonal wind (Figs. 1b–d). Although HiRAM simulated eastward propagation  
201 in a wider frequency spectrum than the observed, the coupled model clearly displays  
202 improvements in the MJO simulation compared with the stationary intraseasonal  
203 fluctuation in the uncoupled simulation. Hovmöller diagrams presented in Fig. 2 illustrate  
204 the temporal evolution of 850-hPa zonal wind and precipitation in the tropics in  
205 observation and simulations. All three models simulated either stationary (CAM5 and  
206 HiRAM) or weak eastward-propagating (ECHAM5) signals in AGCMs, but more  
207 realistically simulated the eastward propagation of the MJO in the coupled models.

208 However, the propagation in the ECHAM5-SIT is still slightly slower than the observed.  
209 The improvement obtained in coupled models suggests that active ocean–atmosphere  
210 interaction is crucial for successful MJO simulation.

### 211 **3.1.2 Atmospheric and oceanic profiles**

212 The composite MJO life cycle featuring intraseasonal OLR and 10-m surface wind  
213 anomalies for boreal winter in eight phases following Wheeler and Hendon (2004) is  
214 displayed in Fig. 3. All three coupled models simulated realistic MJO with enhanced  
215 circulations and propagation tendency compared with the uncoupled ones. The MJO in  
216 phase 4, when deep convection is the strongest over the Maritime Continent (MC),  
217 demonstrates the large-scale zonally overturning circulation coupling with the convection  
218 (Fig. 4). The positive heating region in the coupled experiment is significantly enlarged,  
219 deepened, and westward-tilted with increasing height compared with those in the  
220 uncoupled experiment. Correspondingly, the convective-circulation envelope of the MJO  
221 is thicker and longitudinally wider in coupled experiments. The strong convection is  
222 associated with much enhanced low-level moisture convergence (green contours).  
223 Furthermore, the area of positive rainfall anomaly in the coupled experiment becomes  
224 larger, and the sea level pressure anomaly is meridionally more confined, exhibiting the  
225 characteristics of intensified Kelvin wave-like perturbations to the east of the deep  
226 convection. This enhancement of low-level moisture convergence is consistent with the  
227 frictional wave-conditional instability of the second kind mechanism (Wang and Rui,  
228 1990; Kang et al., 2013). The enhancement of the Kelvin wave can be observed in the  
229 symmetric wavenumber–frequency spectra (Fig. 5). The spectra between 0 and  $0.35 \text{ day}^{-1}$   
230 are presented to highlight the MJO and equatorial Kelvin waves. The coherence at

231 wavenumbers of 2–4 for the 10–20-day period is simulated stronger in three coupled than  
232 uncoupled models.

233 In addition to the atmospheric structure, the SST (Fig. S1) and vertical profile of  
234 ocean temperature examined are presented in Fig. 6. The observed SST variation in MJO  
235 variability is well reproduced in all three coupled models (Fig. S1). The warm SST leads  
236 the main MJO convection by approximately 5–10 days, followed by the cold SST  
237 approximately 5–10 days later (Flatau et al., 1997; DeMott et al., 2015; Tseng et al., 2015).  
238 Moreover, the observed amplitude fluctuation (approximately 0.5° to 1°C) is realistically  
239 simulated. The observed ocean temperature profiles, characterized by the warm layer,  
240 along the equator from the Indian Ocean to the western Pacific are well simulated in the  
241 three coupled models (Fig. 6). Meanwhile, simulated temperature anomalies are larger in  
242 ECHAM5-SIT than in CAM5-SIT and HiRAM-SIT. Figure. S2 shows the fluctuations of  
243 observed SST and simulated SST in three sets of coupled and uncoupled model. There is  
244 no fluctuation as expected in uncoupled simulations, whereas the simulated SST  
245 fluctuates with phases similar to the observed at different locations. The amplitudes in  
246 ECHAM5-SIT and CAM5-SIT are similar to the observed, whereas those in HiRAM-SIT  
247 seems to be smaller in the western Pacific. The differences between models are likely due  
248 to the different atmospheric model configurations, because they were coupled to the same  
249 1-D ocean model. Since the atmosphere is the main driver to extract heat form the ocean,  
250 different responses of atmospheric models likely have different effects on SST. The cause  
251 of quantitative differences between models needs further detailed analysis to pinpoint.  
252 The consistent results in all three coupled models support the conclusion of Tseng et al.  
253 (2015) that resolving fine vertical resolution in the upper ocean improves the simulation  
254 of the warm layer and MJO propagation and variability. Our results further demonstrate

255 that the effect of atmosphere–ocean coupling on the MJO could be independent of  
256 AGCMs with different configurations and atmospheric physical parameterizations, and  
257 that coupling seems to be a more fundamental approach.

### 258 **3.1.3 Performance comparison**

259 Model performance is summarized in Fig. 7. The scatter plot shows the power ratio  
260 of east–west propagating waves (X-axis) versus the pattern correlation between the  
261 simulated and observed precipitation anomaly in Hovmöller diagrams (Fig. 2; Y-axis).  
262 The east:west ratio was calculated by dividing eastward-propagating power by westward-  
263 propagating power of 850–hPa zonal wind summed over wavenumbers of 1–2 and a  
264 period of 30–80 days. Compared with the observation, coupled simulations (marked by  
265 circles) exhibit better simulation than uncoupled simulations (marked by asterisks). A  
266 comparison of combined explained variance using RMM1 and RMM2 (Fig. 7b) based on  
267 Wheeler and Hendon (2004) shows marked increases after coupling. The comparison  
268 demonstrates that coupling is essential for realistic MJO simulations.

### 269 **3.2 Mechanism discussion**

270 We applied the MSE budget to diagnose the moisture budget associated with the  
271 MJO. Figure 8 shows a Hovmöller diagram of MSE tendency averaged by 10°S–EQ  
272 overlaying precipitation anomalies. MSE tendency derived from reanalysis fluctuates in  
273 quadrature with precipitation anomaly with positive (negative) MSE tendency, leading  
274 (lagging) major convection by approximately one to two phases (DeMott et al., 2015;  
275 DeMott et al., 2016; DeMott et al., 2019). Coupled models simulate stronger eastward  
276 propagation in the MSE tendency and precipitation anomalies and realistic phase lag

277 between the two. Stronger MSE tendencies in coupled simulations are observed in  
278 ECHAM5 and HiRAM but are less clear in CAM5. Figures 8d, g, and j show the  
279 differences between coupled and uncoupled simulations. One notable feature is the  
280 positive (negative) MSE tendency preceding positive (negative) precipitation anomaly  
281 and preconditions an environment for eastward propagation of active (inactive)  
282 convection and associated circulation. Next, we diagnosed the relative contribution of  
283 each term in Equation 1 to the MSE tendency with the focus on the MC, where the largest  
284 positive MSE tendency and precipitation anomaly were found.

### 285 **3.2.1 Preconditioning phase**

286 Following the peak MSE tendency over the MC (120°E–150°E) during phase 2 (Figs.  
287 8d, g, and j), values of each term contributing to the column-integrated MSE tendency in  
288 Equation 1 preceding the deep convection over the MC area (10°S–EQ, 120°E–150°E)  
289 are shown in Fig. 9. Vertical advection is the dominant term with the major compensation  
290 from longwave radiation during phase 2 when convection is still in the eastern Indian  
291 Ocean, as identified by Wang et al. (2017). Moreover, the LH term is consistent within  
292 all three models to contribute less negative MSE tendency in coupled models than  
293 ACGMs. The results show that the contribution comes from the LH term in this early  
294 phase stage. The LH effect was overlooked in Tseng et al. (2015) because of the weak  
295 MJO variability in coupled simulations. However, this negative LH bias becomes one of  
296 the key factors in enhancing the leading MSE tendency during the MJO preconditioning  
297 phases. This suggests that the surface latent flux bias in AGCMs can be corrected by  
298 involving the coupling process in the preconditioning phase. Generally, coupling  
299 improves the budget simulation. The positive contribution of vertical advection and

300 negative contribution of LH in MSE tendency is closer to realistic in the coupled  
301 simulations during the initial phase of the MJO.

### 302 **3.2.2 Phase of strongest convection over MC**

303 We compared the spatial distribution of MSE and precipitation in phase 4 when  
304 convection was the strongest in the MC (Fig. 10). In the observation, the main convection  
305 occurs in the MC from 90°E to 150°E. A positive MSE tendency with a maximum value  
306 near 10°N and 10°S is identified in the east of the MJO convection centered near the  
307 equator. Meanwhile, a negative integrated MSE tendency is found in the west of the MJO  
308 convection, and the meridionally confined structure near the equator exhibits the  
309 characteristics of the equatorial Kelvin wave embedded in the MJO. Clearly, coupled  
310 models outperform uncoupled models in reproducing these signals. To quantify the  
311 contribution of coupling to the improvement, we follow Jiang et al. (2018) to project all  
312 MSE terms to the observations (Fig. 11). The dominant contribution of horizontal  
313 advection to the MSE tendency in observation (Fig. 11a) is well simulated in the coupled  
314 simulations but not in uncoupled simulations by ECHAM5 and CAM5 (Figs. 11b and c).  
315 Although a similar dominant effect was observed in both simulation types in HiRAM, it  
316 is enhanced in the coupled simulation (Fig. 11d). The horizontal advection term is further  
317 decomposed into zonal and meridional components (Figs. 11e–h); both components have  
318 a positive contribution, but the meridional component has a larger amplitude.  
319 Furthermore, the uncoupled ECHAM5 and CAM5 simulate unrealistic features: positive  
320 contribution from zonal advection but negative contribution from meridional advection.  
321 In contrast, coupled models well simulate the dominance of meridional advection. In  
322 HiRAM, the uncoupled model simulates almost equally positive contributions from both





### 343 3.3 Discussion: mean state and intraseasonal variance

344 We examined the simulated mean state, which has been suggested a key factor  
345 affecting MJO simulations (Inness et al., 2003; Watterson and Syktus, 2007; Kim et al.,  
346 2009; Kim et al., 2011; Kim et al., 2014; Jiang et al., 2018; Jiang et al., 2020). The three  
347 models exhibited different tropical SST responses to coupling (Fig. S3e). Over the warm  
348 pool area, CAM-SIT and HiRAM-SIT underestimate the SST, whereas ECHAM5-SIT  
349 overestimates the SST. Note that warm SST bias in the eastern tropical Pacific was  
350 simulated in the three models due to the lack of oceanic circulation in the SIT. The  
351 simulated zonal wind in the three models (Figs. S3b–d) demonstrated different responses  
352 to coupling. Figures S2c show the 850-hPa zonal wind differences between coupled and  
353 uncoupled models (shading) and the total field in uncoupled models (contours). Figures  
354 S3f–h show the 10°S–EQ averaged 850-hPa zonal in the coupled and uncoupled models.  
355 In ECHAM5-SIT, the westerly wind is slightly enhanced in the eastern Indian Ocean but  
356 decreases in the western Indian Ocean and western Pacific. In CAM5-SIT, westerly wind  
357 reduces in the Indian Ocean but enhances over the western Pacific. The HiRAM-SIT has  
358 similar changes as in ECHAM5-SIT, e.g., decreases over the MC area but increases in  
359 the western Indian Ocean and Pacific. Generally, the three models disagree on the zonal  
360 wind mean state changes in response to coupling.

361 The mean moisture changes are substantially enhanced over the tropical areas in  
362 ECHAM5 after coupling (Figs. S4b and e). However, in CAM5 and HiRAM, no clear  
363 change was observed to the south of the equator, but strong drying was observed to in the  
364 north (Figs. S4c, d, f, and g). The only common feature among the three models e is the  
365 enhanced meridional gradient of mean moisture, which is consistent with many previous  
366 studies (Kim et al., 2014; Jiang et al., 2018; Ahn et al., 2020). Our budget analysis

367 demonstrated that the meridional transport by the intraseasonal meridional circulation is  
368 the dominant term. It also showed that the meridional gradient of mean moisture is the  
369 secondary effect in enhancing MJO simulations by coupling. After coupling, the mean  
370 precipitation changes are more consistent among the three models (Fig. S5). One of the  
371 major changes is the southward shift of the major precipitation zone, resulting in  
372 precipitation increases over the regions south of the equator, except in the MC. Similarly,  
373 the precipitation intraseasonal variance (20–100 days filtered) was markedly enhanced in  
374 these regions (Fig. S6). The ECHAM5-SIT exhibits a relatively minor increase over the  
375 western MC. In contrast, the HiRAM-SIT exhibits the strongest enhancement,  
376 particularly in the Indian Ocean. Generally, all three coupled models enhance the  
377 intraseasonal signals over the tropics with discrepancies in detail. Meanwhile, the model  
378 mean state does not substantially improve after coupling. Thus, in this study, the mean  
379 state is not the main contribution to enhancing the MJO simulation after coupling. Instead,  
380 coupling leading to rigorous atmosphere–ocean interaction in intraseasonal time scale is  
381 likely the reason for improving MJO simulation.

382

#### 383 **4 Discussion**

384 This study used a one-column TKE-type ocean mixed-layer model SIT coupled  
385 with AGCMs to improve MJO simulation. SIT that is designed to have fine layers near  
386 the surface well simulates warm layer, cool skin, and their diurnal fluctuations. This  
387 refined discretization under the ocean surface in SIT provides improved SST simulation;  
388 thus, improving realistic air–sea interaction. Coupling SIT with ECHAM5, CAM5, and  
389 HiRAM significantly improves the MJO simulation in the three AGCMs compared with  
390 that in the prescribed SST-driven AGCMs. The vertical cross-section indicates that the

391 strengthened low-level convergence during the preconditioning phase is better simulated  
392 in the coupled experiment. Furthermore, the phase variation and amplitude of the SST  
393 and ocean temperature under the surface can be realistically simulated. Our results reveal  
394 that the MJO can be realistically simulated in terms of strength, period, and propagation  
395 speed by increasing the vertical resolution of the one-column ocean model to better  
396 resolve the upper-ocean warm layer.

397         The MSE budget analysis revealed that the coupling effects during the  
398 preconditioning and mature phases exhibit different contributions. During the  
399 preconditioning phase, the positive contribution of vertical advection and negative  
400 contribution of LH in MSE tendency are closer to realistic values in coupled simulations  
401 during the initial phase of the MJO. Additionally, the meridional component of the  
402 horizontal advection term is the dominant term during the mature phase of the strongest  
403 convection in the MC to enhance the simulation after coupling. Improved meridional  
404 circulation is essential in the coupled simulations that outperformed uncoupled  
405 experiments. The results confirm that the dominance of dynamic influence over  
406 thermodynamic influence in response to the atmosphere–ocean coupling is the key  
407 process in improving MJO simulations.

408         In summary, this study suggests two major enhancements of the coupling process.  
409 First, the underestimated surface LH bias in AGCMs can be corrected during the  
410 preconditioning phase of the MJO over MC. Second, during the strongest convection  
411 phase over MC, the change in intraseasonal circulation in the meridional circulation is the  
412 dominant factor in coupled simulations relative to uncoupled experiments. Although  
413 many studies have indicated the key role played by the mean state, the mean state in our

414 simulations provides only a secondary contribution to enhancing MJO simulation, with  
415 coupling being the main contributor. For example, zonal wind and precipitation changed  
416 inconsistently among the three models after coupling. Instead, the meridional gradient of  
417 the mean moisture and intraseasonal precipitation variance has a better relationship after  
418 coupling. Therefore, coupling leading to rigorous atmosphere–ocean interaction in the  
419 intraseasonal time scale, but no change in mean states, is likely the reason for MJO  
420 simulation improvement. This study supports previous findings (Tseng et al., 2015) that  
421 enhancing atmosphere–ocean coupling by considering an extremely high vertical  
422 resolution in the first few meters of the ocean model improves MJO simulations. It also  
423 supports that this improvement is independent of AGCMs with different configurations  
424 and physical parameterization schemes. Resolving the atmosphere–ocean coupling may  
425 be more beneficial than modifying the atmospheric physical parameterization schemes in  
426 GCM. In brief, this study suggested the effectiveness of air–sea coupling for improving  
427 MJO simulation in a climate model and demonstrated the critical effect of being able to  
428 simulate warm layer. Additionally, the findings presented here enhance our understanding  
429 of the physical processes that shape the characteristics of the MJO.

430

431 **Code and data availability.** The model code of CAM5–SIT, ECHAM5-SIT and  
432 HiRAM-SIT is available at <https://doi.org/10.5281/zenodo.5701538>,  
433 <https://doi.org/10.5281/zenodo.5510795> and <https://doi.org/10.5281/zenodo.5701579>.  
434 Observational data used in this study include precipitation from Global Precipitation  
435 Climatology Project V1.3 (GPCP, 1° resolution), OLR (1° resolution), and daily SST  
436 (Optimum Interpolated SST, 0.25° resolution) from the National Oceanic and  
437 Atmosphere Administration, and variables were obtained from the European Centre for  
438 Medium-range Weather Forecast Reanalysis-interim. All experiments were conducted at  
439 the National Center for High-performance Computing. All model codes and data  
440 availability presented here can be obtained by contacting the first author, Dr. Wan-Ling  
441 Tseng ([wtseng@ntu.edu.tw](mailto:wtseng@ntu.edu.tw)).

442

443 **Author contributions.** HHH and WLT have responsibility for conceptualization,  
444 including analyzing the data and writing the manuscript. YYL, WLL, PHK, BJT, CYT,  
445 and HCL developed the model and provided the simulations.

446

447 **Competing interests.** The authors declare that they have no conflict of interest.

448

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456

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- 663

	<b>ECHAM5-SIT</b>	<b>CAM5-SIT</b>	<b>HiRAM-SIT</b>
<b>AGCM</b>	ECHAM5	CAM5	HiRAM
<b>Horizontal resolution</b>	T63(~2°)	1.9°×2.5°	1°×1°
<b>SST</b>	OISST	OISST	OISST
<b>BC SIC</b>	OISST	OISST	OISST
<b>OT/OS</b>	GODAS	GODAS	GODAS
<b>Atmosphere vertical resolution</b>	L31	L30	L32
<b>Ocean vertical resolution</b>	42	42	42
<b>Coupled region</b>	30°S–30°N	30°S–30°N	30°S–30°N
<b>Time</b>	1985–2005 (21 years)		

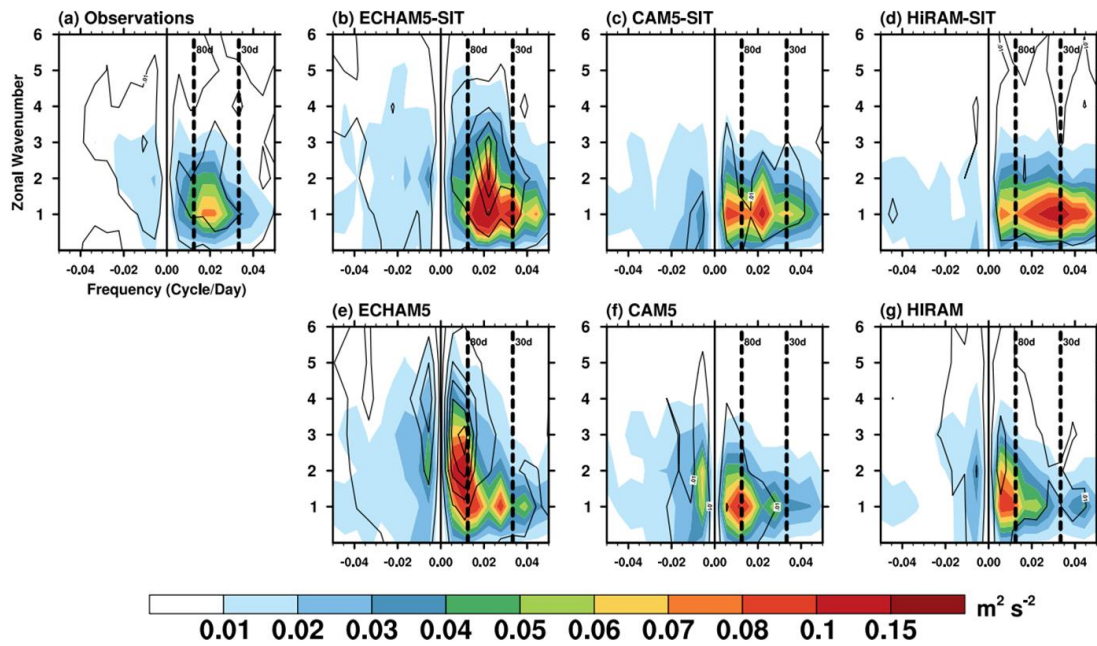
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665 **Table 1.** Detailed information of models and experiments

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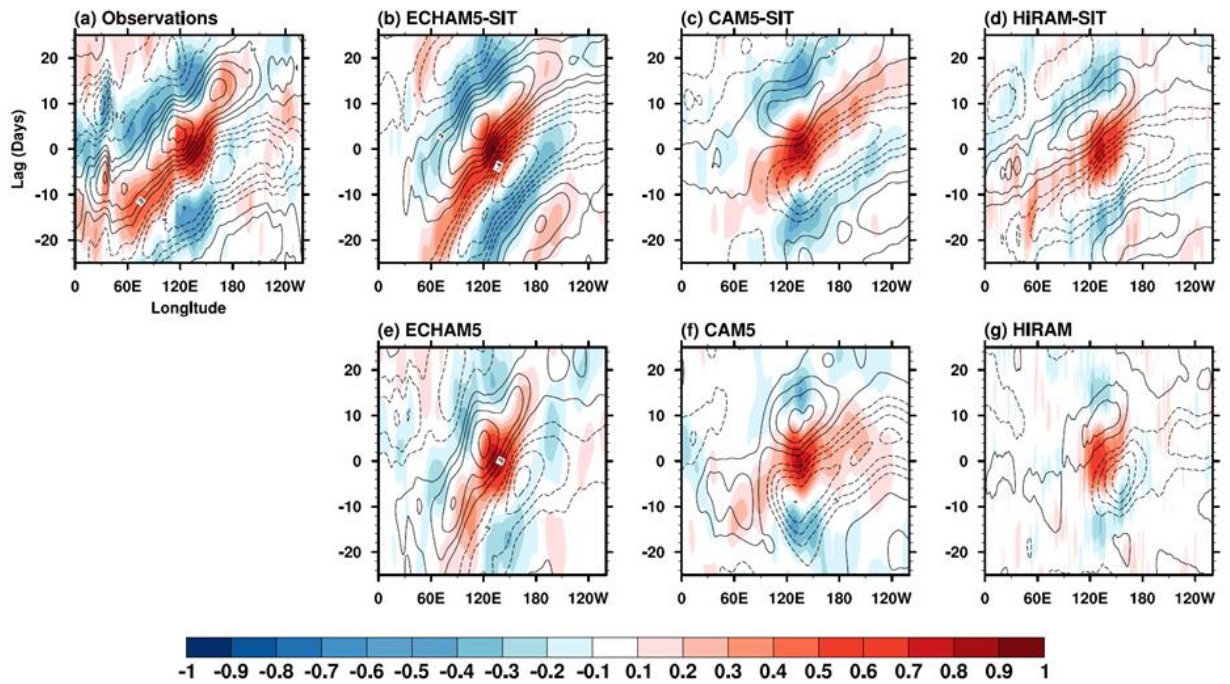


668

669 **Figure 1.** Wavenumber–frequency spectra for equatorial 850-hPa zonal wind (shading;  
 670  $\text{m}^2 \text{s}^{-2}$ ) and precipitation (contours;  $\text{mm}^2 \text{day}^{-2}$ ) over  $10^\circ\text{S}$ – $10^\circ\text{N}$  from (a) observations  
 671 and simulations using the (b–d) coupled and (e–g) uncoupled AGCM.

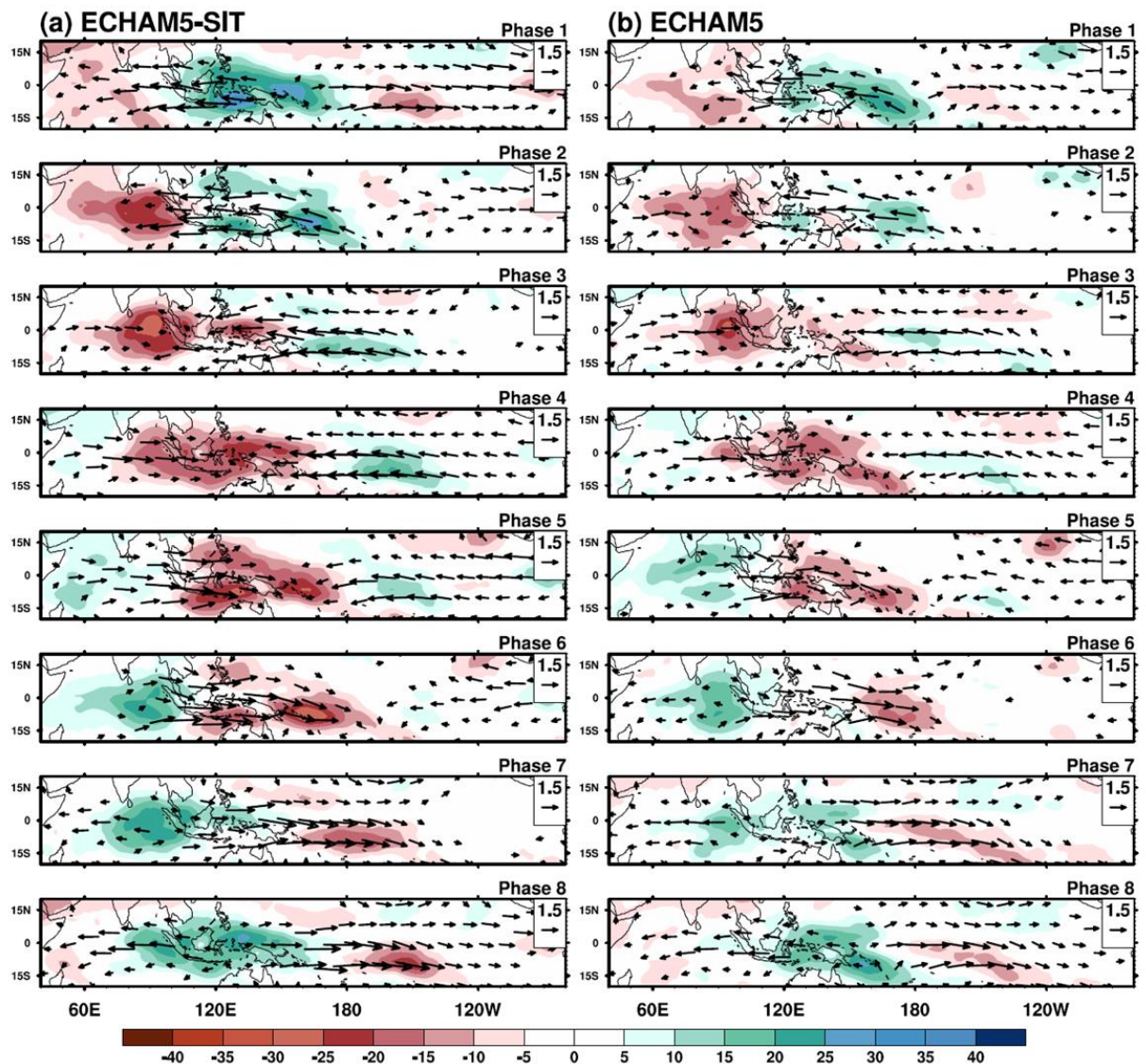
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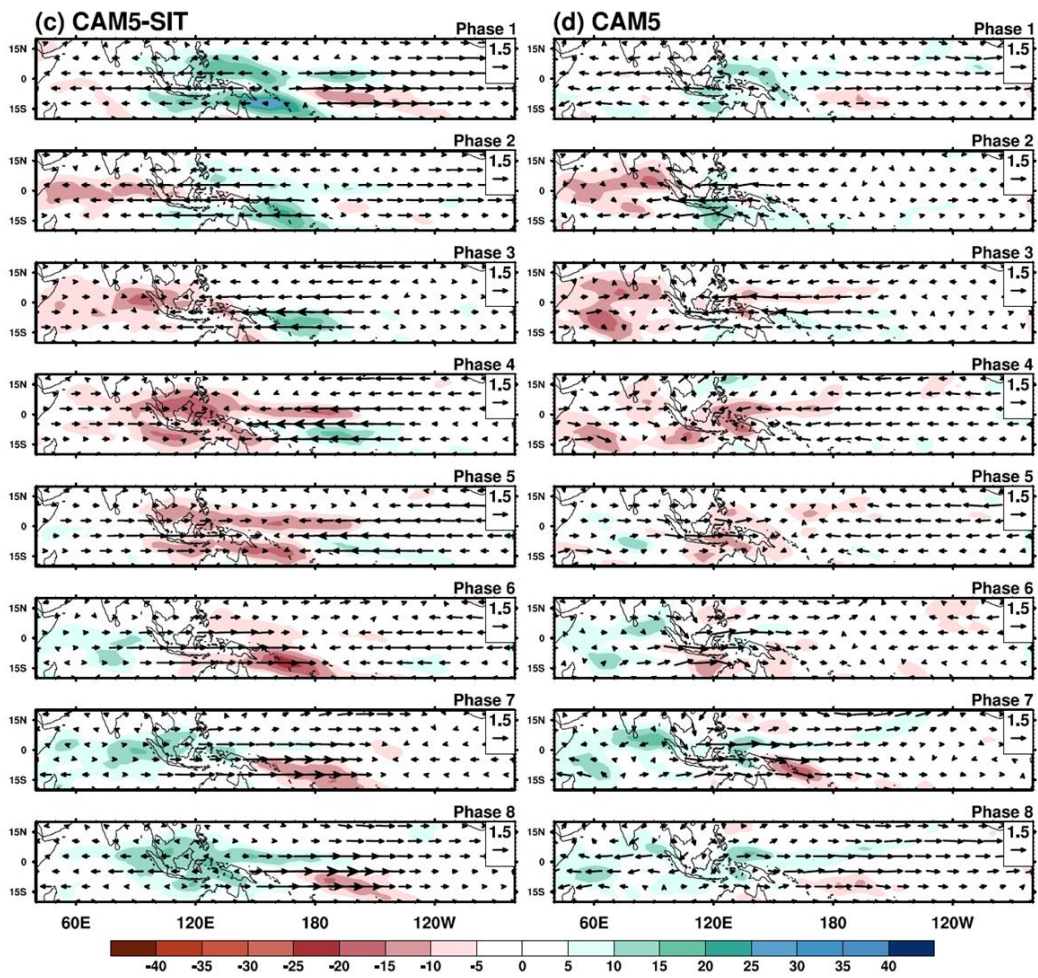
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675 **Figure 2.**  $10^{\circ}\text{S}$ – $10^{\circ}\text{N}$  averaged lag–longitude diagrams of intraseasonal precipitation  
 676 (shading) and 10-m zonal wind (contour) correlated against precipitation at region ( $10^{\circ}\text{S}$ –  
 677  $5^{\circ}\text{N}$ ,  $120^{\circ}\text{E}$ – $150^{\circ}\text{E}$ ) from (a) observations and simulations using the (b–d) coupled and  
 678 (e–g) uncoupled AGCM. The contour interval is 0.1.



679  
 680 **Figure 3.** Composite November–April 20–100-day OLR ( $W m^{-1}$ ; color) and 10-m surface  
 681 wind anomalies ( $m s^{-1}$ ; vectors) as a function of the MJO phase in (a) ECHAM5-SIT and  
 682 (b) ECHAM5. Vectors  $<0.6 m s^{-1}$  are not shown. The reference vector in units of  $m s^{-1}$  is  
 683 shown at the bottom right of each panel. The number of days used to generate the composite for each  
 684 phase is shown to the right of each panel.

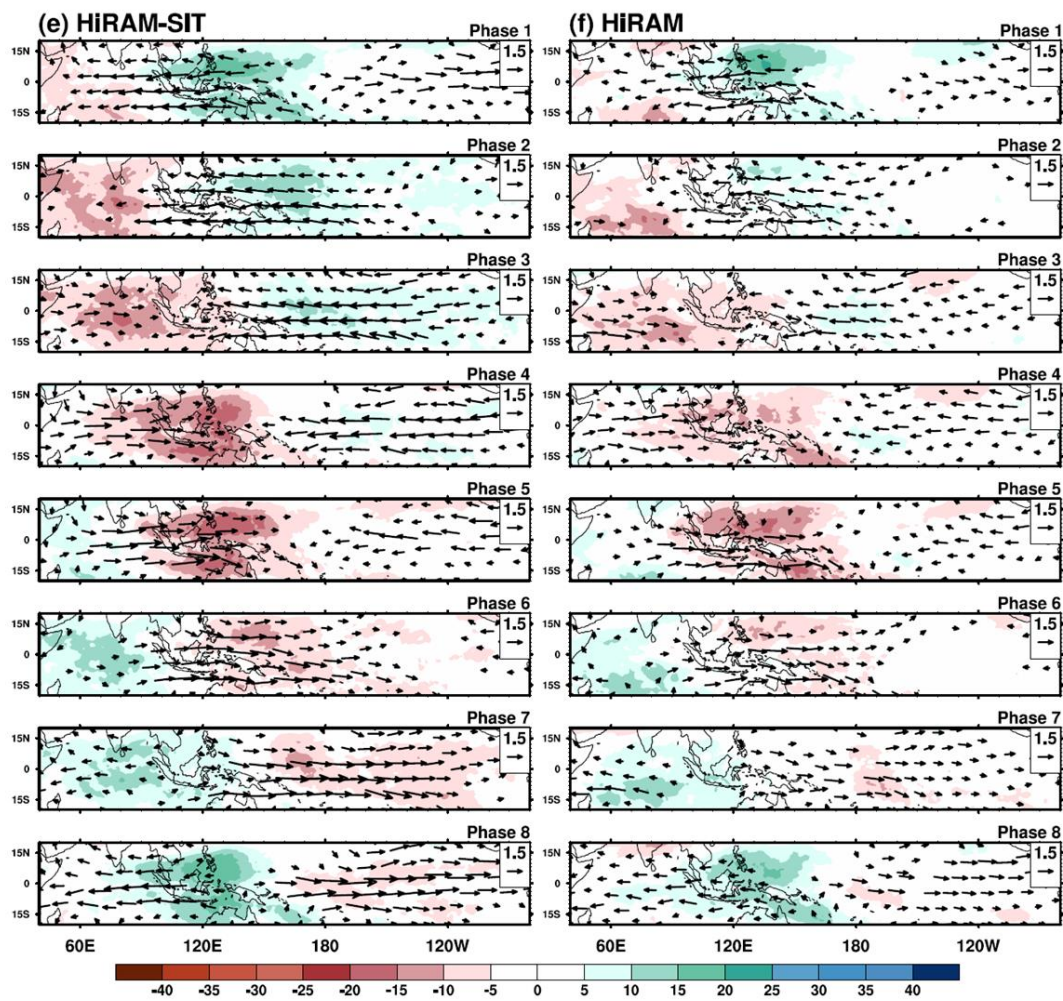
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687 **Figure 3.** continued, (c) CAM5-SIT and (d) CAM5.

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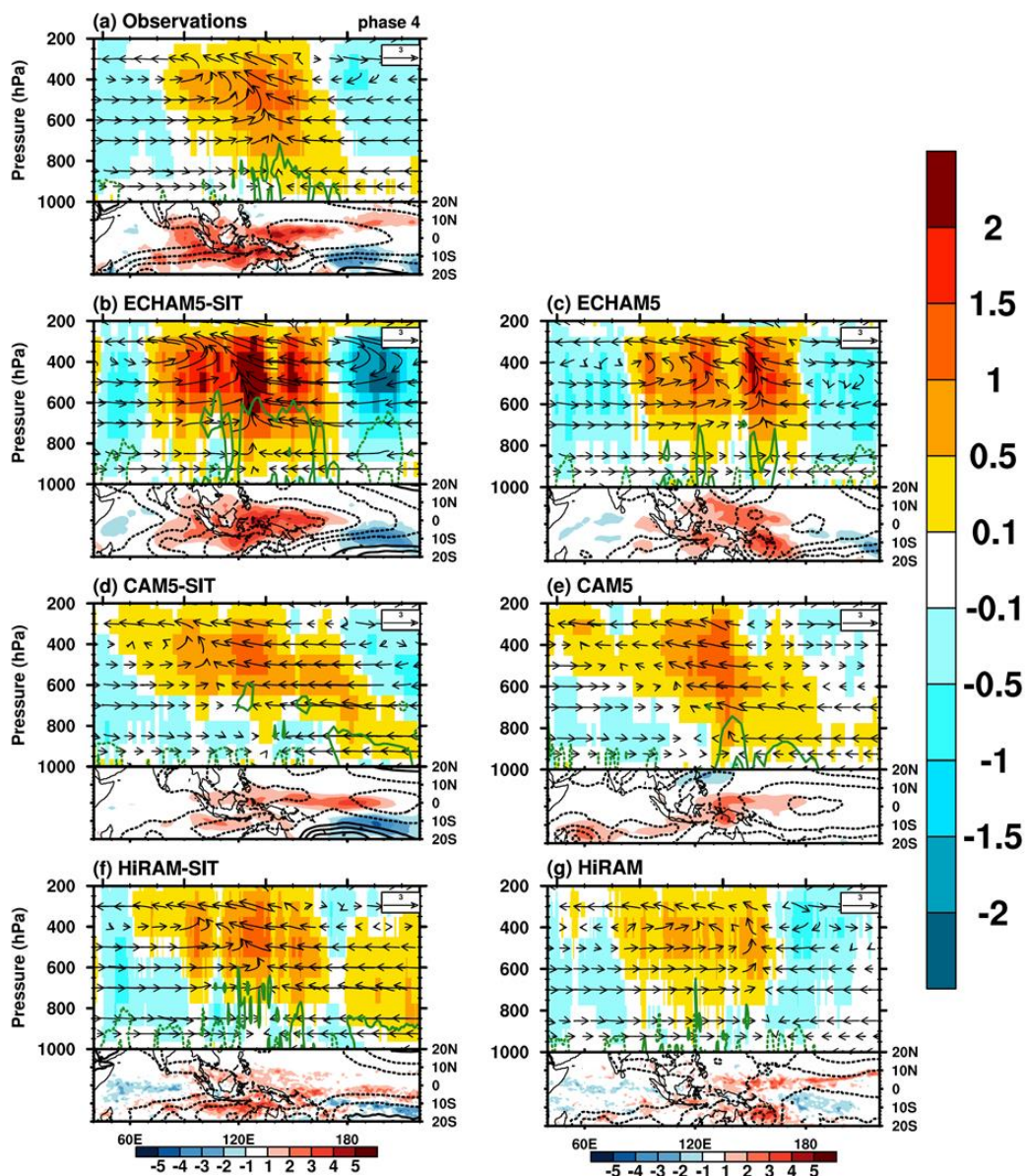


689

690 **Figure 3.** continued, (e) HiRAM-SIT and (f) HiRAM.

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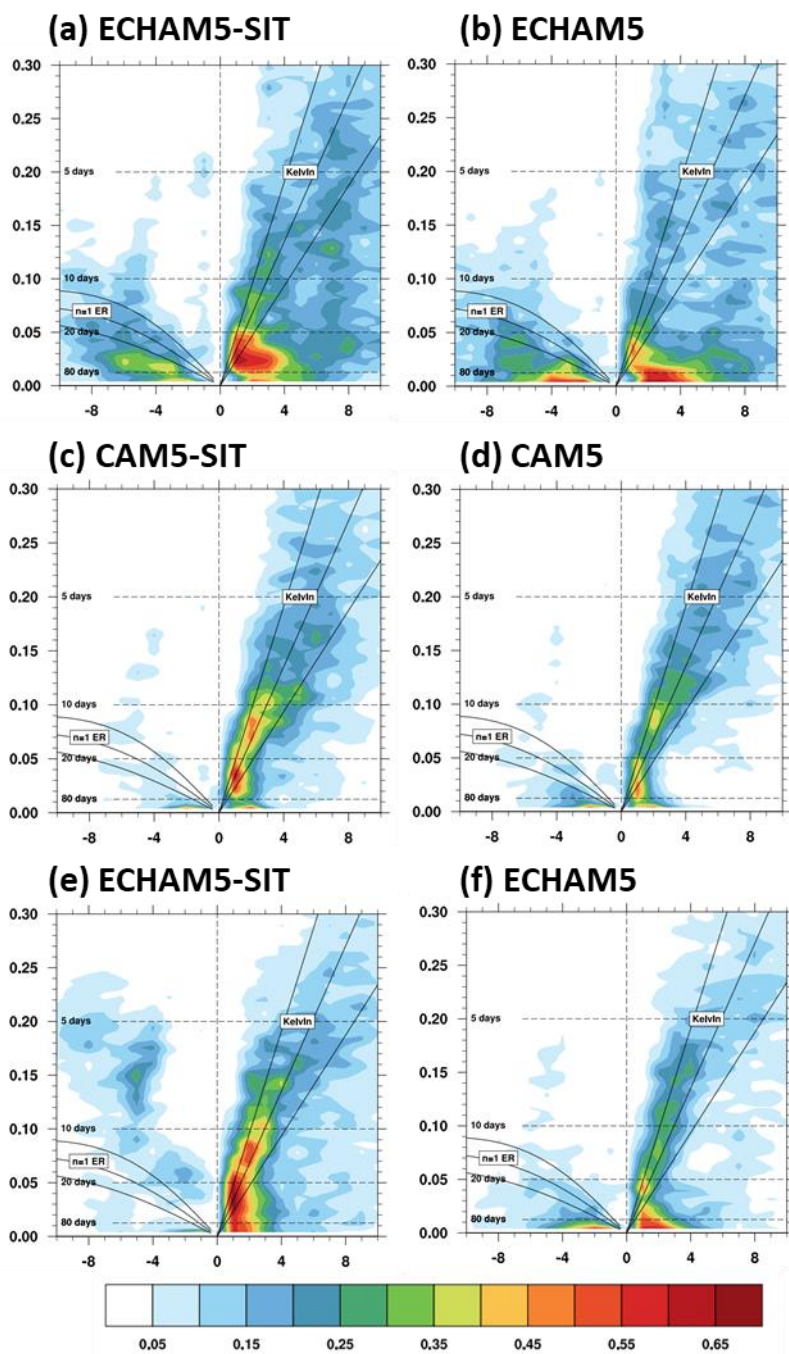
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693

694 **Figure 4.** Structure of simulated MJO in phase 4. The longitude–height cross-sections  
 695 (averaged over 10°S–EQ) of the MJO scaled wind circulation (vector,  $u$ :  $\text{m s}^{-1}$ ,  $\omega$ :  
 696  $10^{-2} \text{ Pa s}^{-1}$ ), Q1 (shading, unit:  $\text{K day}^{-1}$ ), and the horizontal moisture convergence (green  
 697 contour, unit:  $10^{-6} \text{ g kg}^{-1} \text{ s}^{-1}$ ) from (a) observations and simulations using the (b–d)  
 698 coupled and (e–g) uncoupled AGCMs. The contour interval of the moisture convergence  
 699 is  $8 \times 10^{-6} \text{ g kg}^{-1} \text{ s}^{-1}$ ; solid line is positive. Precipitation (shading, unit:  $\text{mm day}^{-1}$ ) and

- 700 sea level pressure (contour, unit: hPa). Contour interval of sea level pressure is 30 hPa;
- 701 dashed line indicates negative.
- 702

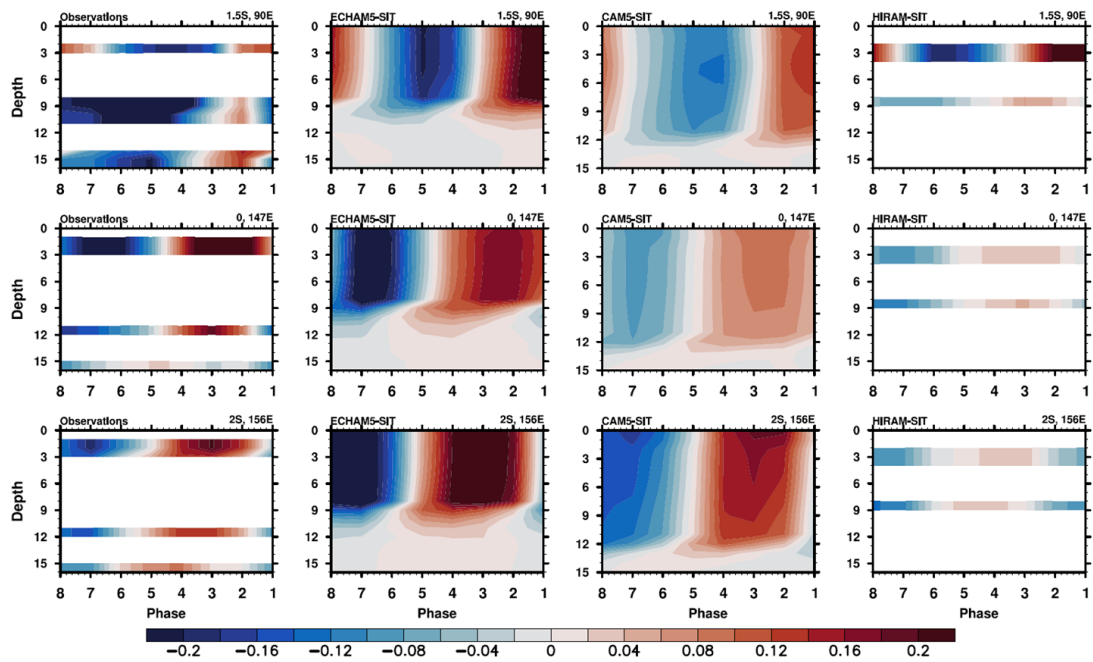


703

704 **Figure 5.** Symmetric wavenumber–frequency spectra of 10°N–10°S-averaged 850-hPa705 zonal wind using the (a, c, e) coupled and (b, d, f) uncoupled AGCMs. Units:  $\text{m}^2 \text{s}^{-2}$ .

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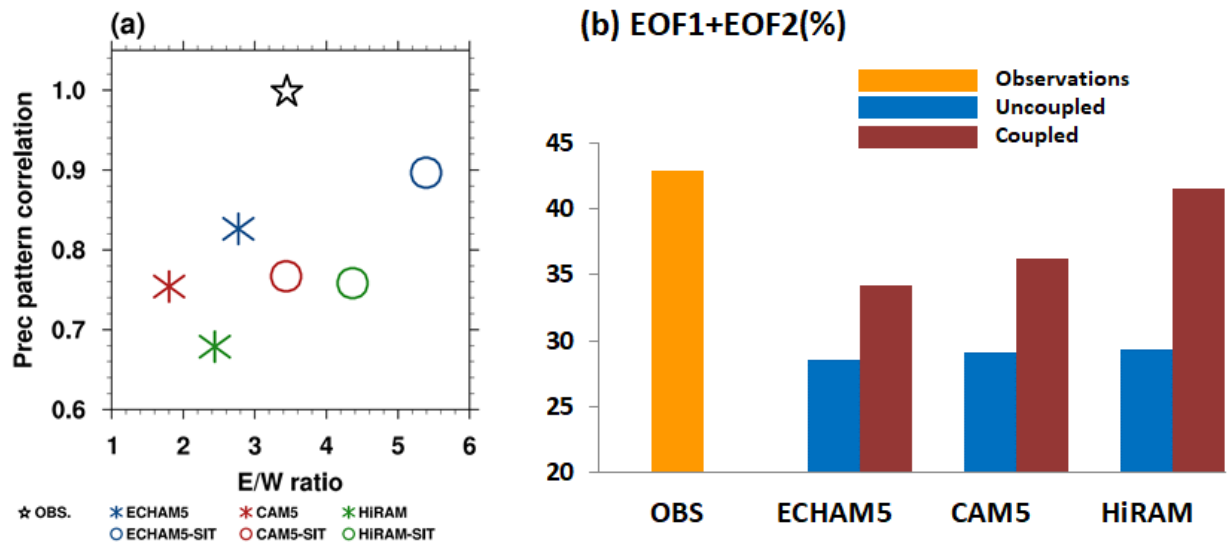
707

708 **Figure 6.** Vertical ocean temperature ( $^{\circ}\text{C}$ ) profiles with respect to MJO phases for  
 709 intraseasonal anomalies (i.e., with 20–100-day filtering) in observations and simulations  
 710 using coupled models. Observations are in suit with data from TAO. Because of storage  
 711 limitations, only 3 and 10 m water temperatures are presented in the HiRAM-SIT  
 712 simulation.

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717 **Figure 7.** Scatter plots of various MJO indices based on observation and experiments

718 (Table 1). (a) X-axis is the power ratio of east–west propagating waves. The east–west

719 ratio was calculated by dividing the sum eastward-propagating power by the westward-

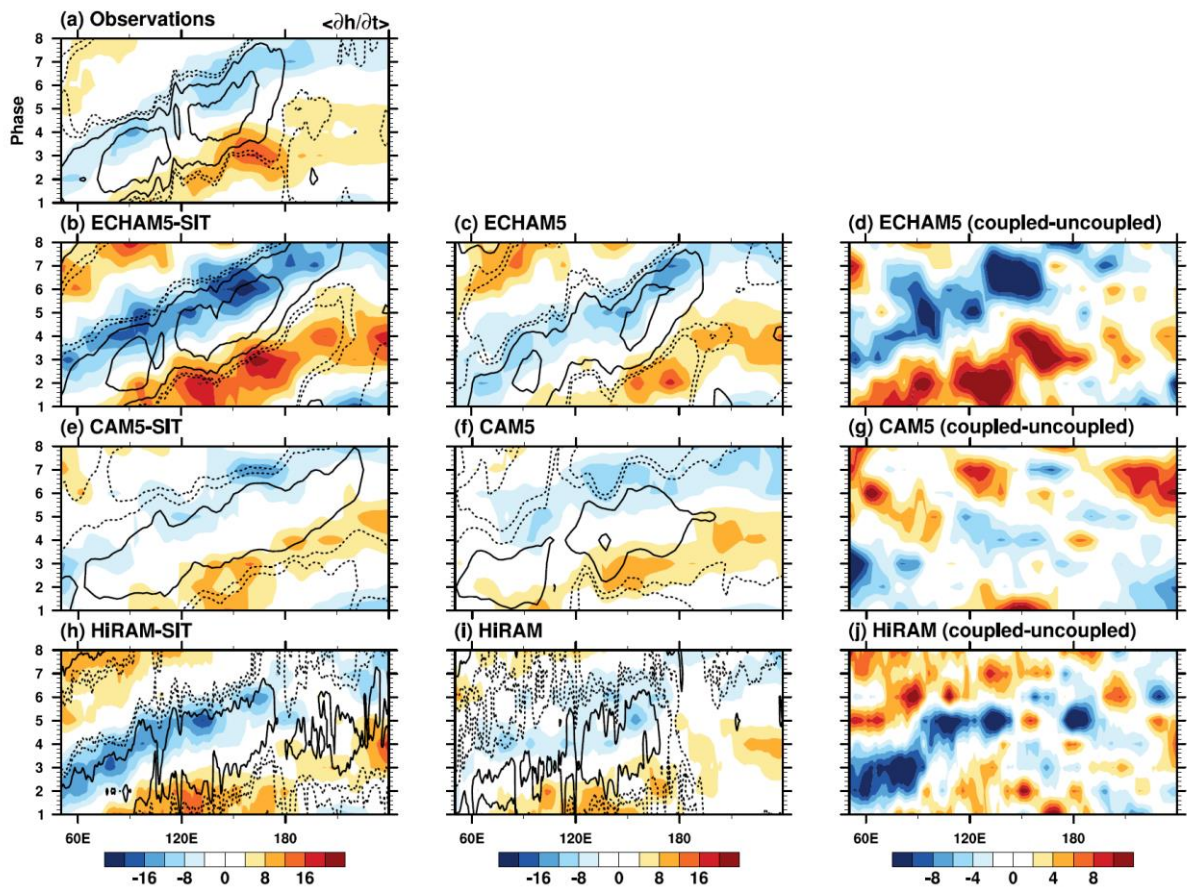
720 propagating counterpart within wavenumbers 1–3 (1–2 for zonal wind), period 30–80

721 days. Y-axis is the pattern correlation of precipitation eastward propagation, as shown in

722 Fig. 2. (b) Sum of RMM1 and RMM2 variances based on Wheeler and Hendon (2004).

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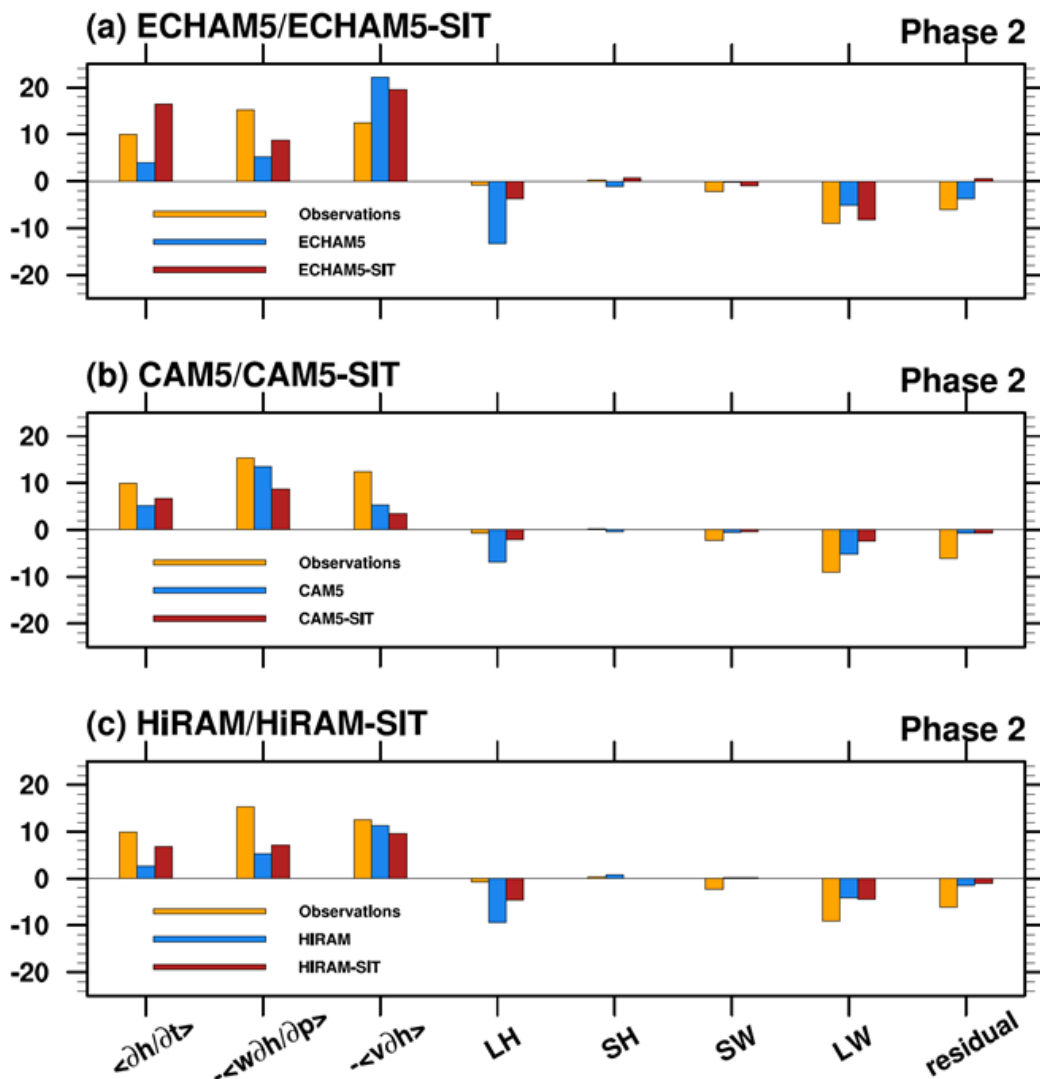
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726 **Figure 8.** 10°S–EQ averaged Hovmöller diagrams of MSE (shading;  $\text{J kg}^{-1}$ ) and  
 727 precipitation (contour;  $\text{mm day}^{-1}$ ) composite followed the RMM index from (a)  
 728 observations and simulations using the (b, e, j) coupled and (c, f, k) uncoupled AGCMs  
 729 and (d, i, l) their difference. The contour interval is precipitation anomalies.

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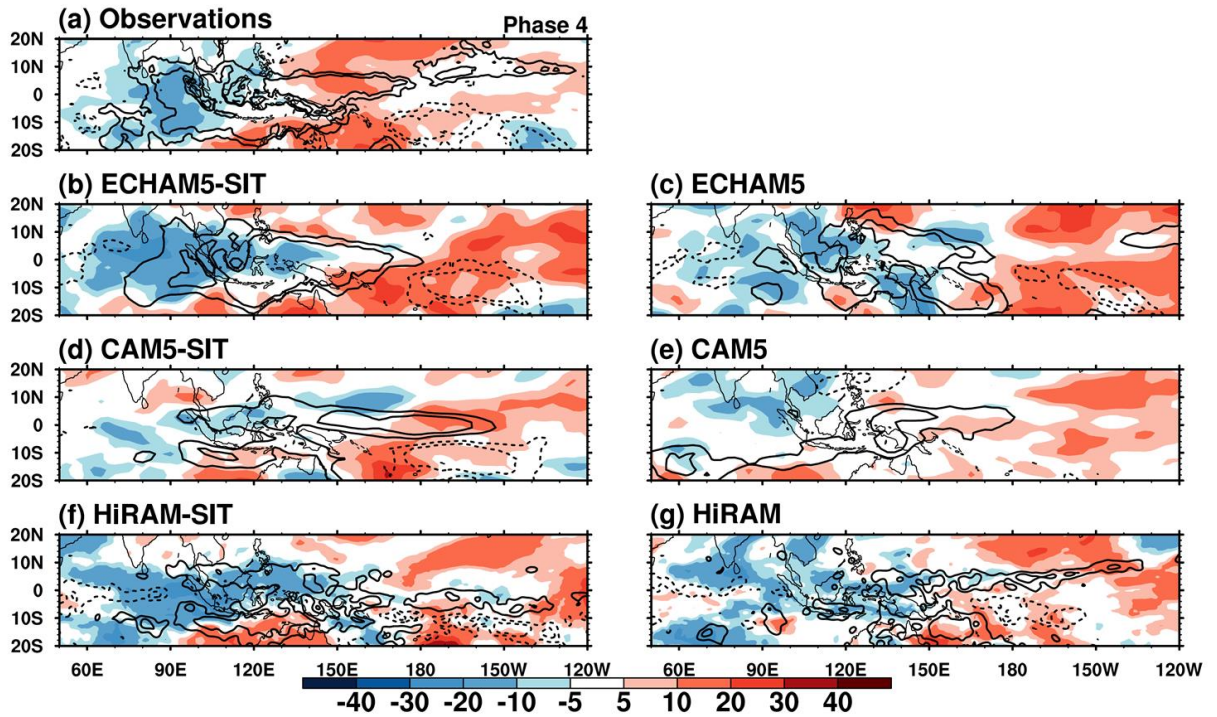
731

732 **Figure 9.** Model-simulated column-integrated MSE budget terms ( $\text{J kg}^{-1} \text{s}^{-1}$ ) during  
 733 phase 2 of the MJO. Black, red, and blue represents the data from the observations,  
 734 Nordeng scheme simulation, and Tiedtke scheme simulation, respectively. The averaged  
 735 domain is  $10^\circ\text{S}$ –EQ and  $120^\circ$ – $150^\circ\text{E}$ .

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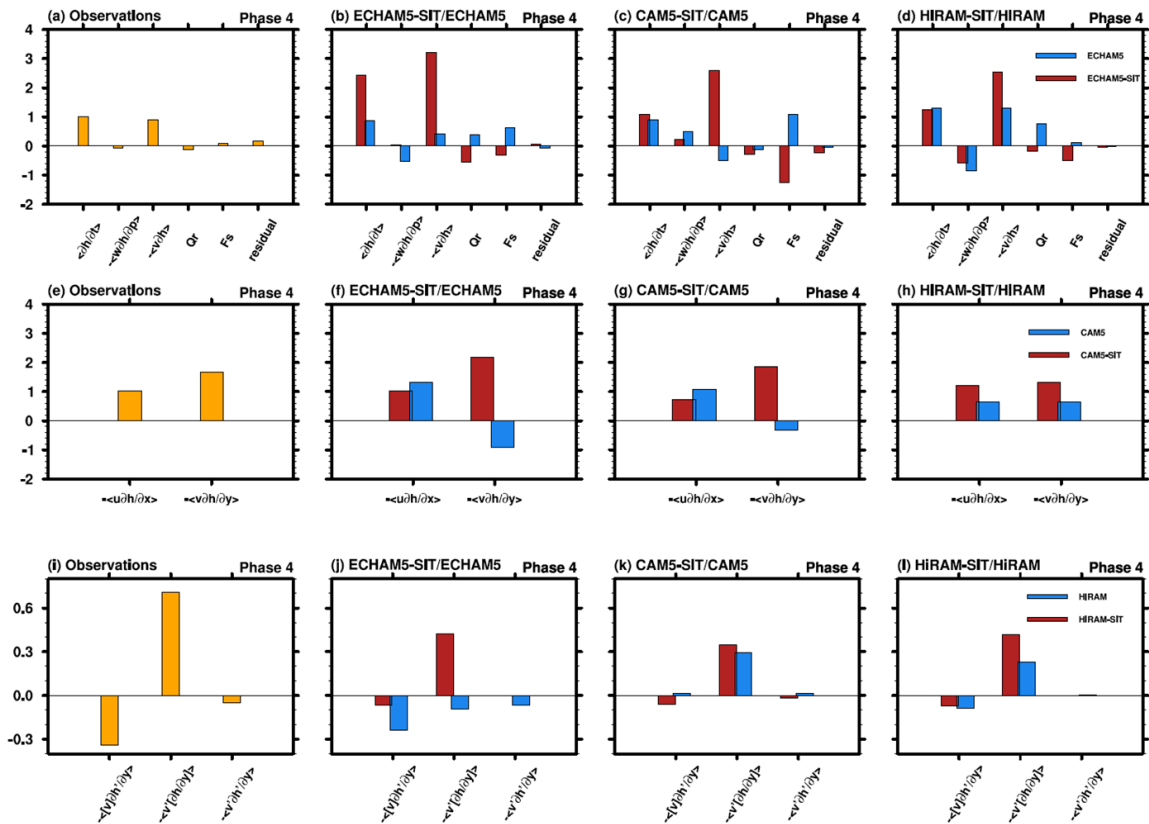
739

740 **Figure 10.** Phase 4 of the column-integrated MSE tendency (shading;  $\text{J kg}^{-1} \text{s}^{-1}$ ) and  
 741 precipitation (contours;  $\text{mm day}^{-1}$ ) based on (a) observation, (b) ECHAM5-SIT, (c)  
 742 ECHAM5, (d) CAM5-SIT, (e) CAM5, (g) HiRAM-SIT, and (f) HiRAM. The nine-point  
 743 local smoothing is applied in the intraseasonal precipitation variance of HiRAM here  
 744 (contours only).

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749 **Figure 11.** (a–d) Relative role of each MSE component of phase 4 through the projection

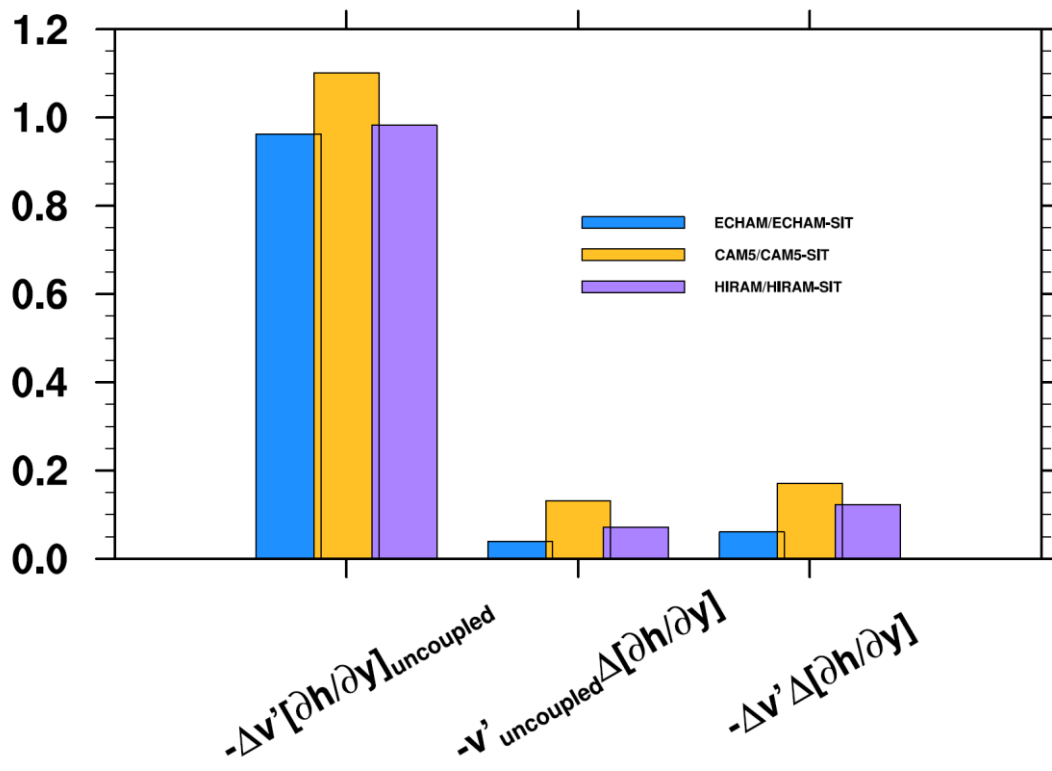
750 of the spatial pattern of each MSE budget over the MC (domain) onto the total MSE

751 tendency pattern (Fig. 8a). (e–h) Decomposite of the total horizontal MSE advection

752 based on zonal and meridional components. (i–l) Decomposite of the meridional

753 horizontal MSE advection based on the MJO circulation and the mean state of moisture.

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756 **Figure 12.** Bar chart of relative contribution of intraseasonal convergence and

757 background moisture between the coupled and uncoupled changes in MJO phase 4.

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