- 1 Embedding a One-column Ocean Model (SIT 1.06) in the
- 2 Community Atmosphere Model 5.3 (CAM5.3; CAM5-
- 3 SIT v1.0) to Improve Madden–Julian Oscillation
- **4 Simulation in Boreal Winter**

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Abstract

12	The effect of the air-sea interaction on the Madden-Julian Oscillation (MJO)
13	was investigated using the one-column ocean model Snow-Ice-Thermocline (SIT
14	1.06) embedded in the Community Atmosphere Model 5.3 (CAM5.3; hereafter
15	CAM5–SIT v1.0). The SIT model with 41 vertical layers was developed to simulate
16	sea surface temperature (SST) and upper-ocean temperature variations with a high
17	vertical resolution that resolves the cool skin and diurnal warm layer and the upper
18	oceanic mixed layer. A series of 30-year sensitivity experiments were conducted in
19	which various model configurations (e.g., coupled versus uncoupled, vertical
20	resolution and depth of the SIT model, coupling domains, and absence of the diurnal
21	cycle) were considered to evaluate the effect of air-sea coupling on MJO simulation.
22	Most of the CAM5–SIT experiments exhibited exhibit higher fidelity than the CAM5-
23	alone experiment in characterizing the basic features of the MJO such as
24	spatiotemporal variability and the eastward propagation in boreal winter. The overall
25	MJO simulation performance of CAM5–SIT benefited benefits from (1) better
26	resolving the fine vertical structure of upper-ocean temperature and therefore the air-
27	sea interaction that resulted result in more realistic intraseasonal variability in both
28	SST and atmospheric circulation and (2) the adequate thickness and vertical
29	resolution of the oceanica vertically-gridded ocean mixed layer. The sensitivity
30	experiments demonstrated demonstrate the necessity of coupling the tropical eastern
31	Pacific in addition to the tropical Indian Ocean and the tropical western Pacific.
32	Coupling is more essential in the south than north of the equator in the tropical
33	western Pacific. Enhanced MJO could be obtained without considering the diurnal
34	cycle in coupling.

1. Introduction

36	The Madden-Julian Oscillation (MJO) is a tropical large-scale convection
37	circulation system that propagates eastward across the warm pool region from the
38	tropical Indian Ocean (IO) to the western Pacific (WP) on an intraseasonal time scale
39	(Madden and Julian, 1972). The MJO is not just an atmospheric phenomenon. The
40	findings of the multination joint from a multi-nation field campaign called the
41	<u>Dynamics of MJO/</u> Cooperative Indian Ocean Experiment on Intraseasonal Variability
42	in the Year 2011/Dynamics of the MJO ((DYNAMO/CINDY2011; de Szoeke et al.,
43	2017; Johnson and Ciesielski, 2017; Pujiana et al., 2018; Yoneyama et al., 2013;
44	Zhang and Yoneyama, 2017) revealed vigorous air-sea coupling during the evolution
45	of the MJO (Chang et al., 2019; DeMott et al., 2015; Jiang et al., 2015, 2020; Kim et
46	al., 2010; Li et al., 2016; Li et al., 2020; Newman et al., 2009; Pei et al., 2018; Tseng
47	et al., 2014). During the suppression of convection, the MJO propagates eastward
48	with light winds, which is accompanied by enhanced downwelling shortwave
49	radiation absorption, weaker upward latent and sensible fluxes, less cloudiness and
50	precipitation, and weaker vertical mixing in the upper ocean, thus causing an increase
51	in the upper-ocean temperature. In the following active phase when deep convection
52	occurs, downwelling shortwave radiation is reduced and stronger westerly winds
53	enhance evaporation and latent/sensible heat flux (LHF/SHF) loss from the ocean
54	surface, thus causing a decrease in the upper-ocean temperature (DeMott et al., 2015;
55	Madden and Julian, 1972, 1994; Zhang, 2005).
56	In addition to the ocean surface, the structure of the upper ocean also evolves.
57	Alappattu et al. (2017) reported that during an MJO event, surface flux perturbations
58	cause changes in the ocean thermohaline structure, thus affecting the mixed-layer
59	temperature. The following change in sea surface temperature (SST) can further affect
60	atmospheric circulation of the MJO. Variations in SST mediate heatLHF and SHF

61 exchange across the air-sea interface. Although SST responds to atmospheric forcing, 62 itsthe modulation of surface heat fluxesLHF and SHF provides feedback to the 63 atmosphere (DeMott et al., 2015; Jiang et al., 2020). Li et al. (2008, 2020) proposed 64 that the phase relationship between SST and convection implies a delayed air-sea 65 interaction mechanism whereby a preceding active-phase MJO may trigger an 66 inactive-phase MJO through the delayed effect of the induced SST anomaly over the 67 IO. The reduction in SST caused by a preceding active-phase MJO may, in turn, 68 vieldyields delayed ocean feedback that initiates a suppressed-phase MJO, and vice 69 versa. The nonnegligible by-no-means negligible effect of intraseasonal SST 70 variations caused by surface heat fluxes suggests that the ocean state can affect the 71 MJO (DeMott et al., 2015, 2019; Hong et al., 2017; Li et al., 2020). 72 Since its discovery almost five decades ago, the MJO remains a phenomenon 73 that poses a challenge to the capacity of state-of-the-art atmospheric general 74 circulation models (AGCMs) and climate models such as those participating in the 75 Coupled Model Intercomparison Project phase 5 and 6 to generate successful 76 simulations (Ahn et al., 2017, 2020; Bui and Maloney 2018; Jiang et al., 2020; Hung 77 et al., 2013; Kim et al., 2011). 78 Recent studies have reported that air–sea coupling improves the representation of 79 the MJO in numerical simulation (Bernie et al., 2008; Crueger et al., 2013; DeMott et 80 al., 2015; Li et al., 2016; Li et al., 2020; Tseng et al., 2014; Woolnough et al., 2007). 81 Tseng et al. (2014) indicated that effectively resolving the upper-ocean warm layer to 82 capture temperature variations in the upper few meters of the ocean could improve 83 MJO simulation. DeMott et al. (2015) suggested that the tropical atmosphere—ocean 84 interaction may sustain or amplify the pattern of the enhanced and suppressed 85 atmospheric convection of the eastward propagation. DeMott et al. (2019) demonstrated that the improved MJO eastward propagation in four coupled models 86

resulted from enhanced low-level convective moistening for a rainfall rate of >5 mm 88 day⁻¹ due to air–sea coupling. In addition, numerical experiments have been 89 performed to investigate the effect of the diurnal cycle on the MJO (Hagos et al., 90 2016; Oh et al., 2013), with the results suggesting that the strength and propagation of 91 the MJO through the Maritime Continent (MC) were enhanced when the diurnal cycle 92 was ignored. 93 Although previous studies have demonstrated the importance of considering the 94 air—sea interaction in a numerical model to improve MJO simulation, additional 95 details regarding model configuration (e.g., vertical resolution, and depth of the ocean 96 mixed layer, coupling domain, and absence of the diurnal cycle in air-sea coupling) 97 have not been systematically explored. Tseng et al. (2014) coupled the one-column 98 ocean model Snow-Ice-Thermocline (SIT; Tu and Tsuang, 2005) to the fifth 99 generation of the ECHAM AGCM (ECHAM5-SIT) and indicated that a vertical 100 resolution of 1 m was essential to yield an improved simulation of the MJO with a 101 realistic strength and eastward propagation speed. 102 In this study, we coupled the SIT model to the Community Atmosphere Model 103 version 5.3 (CAM5.3; Neale et al., 2012)—the atmosphere component of the 104 Community Earth System Model version 1.2.2 (CESM1.2.2; Hurrell et al., 2013)—to 105 explore how the air sea interaction in AGCMs can improve improvement of MJO 106 simulation by coupling SIT model to another AGCM is reproducible in modeling 107 science. The CAM5.3, which has been widely used for the long-term simulation of the 108 climate system, could not efficiently simulate the eastward propagation of the MJO; 109 instead, the model simulated a tendency for the MJO to move westward in the IO 110 (Boyle et al., 2015, Jiang et al, 2015). By contrast, the updated CESM2 with the new CAM6 could realistically simulate the MJO (Ahn et al., 2020; Danabasoglu et 111 112 al., 2020). Thus, the well-explored CAM5, which does not produce a realistic MJO,

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113	appears to be a favorable choice for exploring how coupling a simple one-dimensional
114	(1-D) ocean model, such as the SIT model, can improve MJO simulation, as well as
115	the effects of model configuration. on the degree of the improvement. Such a study
116	can also enhance our understanding regarding the <u>effect of</u> air–sea coupling's <u>effect</u>
117	on the MJO.
118	This study examined how air sea coupling can improve MJO simulation,
119	especially that of the eastward propagation that has been poorly simulated in many
120	climate models. The MJO that exhibits a more substantial eastward propagation in
121	boreal winter than in other seasons was the targeted feature in this study. WeTo
122	examine the sensitivity of MJO simulations to different configurations of air-sea
123	coupling, we conducted a series of 30-year numerical experiments by considering
124	various model configurations (e.g., coupled versus uncoupled, vertical resolution and
125	depth of the SIT model, coupling domains, and absence of the diurnal cycle) to
126	investigate the effect of air-sea coupling. This paper is organized as follows. Section
127	2 describes the data, methodology for validation, the model used for simulation, and
128	model setup. Section 3 presents the the design of coupled model numerical
129	experiments. Section 43 describes the effect of various modelair-sea coupling
130	configurations on the MJO simulation determined through detailed MJO diagnostics.
131	A discussion Discussion and conclusions are provided in Section 54.
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133	2. <u>2.</u> Data, methodology, and model description, and experimental designs
134	2.1 Observational data and analysis methods
135	2.1 Data and methodology
136	The data analyzed in this study include precipitation from the Global
137	Precipitation Climatology Project, (GPCP), outgoing longwave radiation (OLR) and
138	daily SST (Optimum Interpolation SST; OISST) from the National Oceanic and

139	Atmosphere Administration (NOAA), and parameters from the ERA-Interim (ERA-I)
140	reanalysis (Adler et al., 2003; Dee et al., 2011; Lee et al., 2011; Reynolds and Smith,
141	1995; Schreck et al., 2018). The initial SST data for the SIT model were obtained
142	from the Hadley Centre Sea Ice and Sea Surface Temperature dataset (Rayner et al.,
143	2003; HadISST1) and the ocean subsurface data (40-layer climatological ocean
144	temperature, salinity, and currents) for nudging were retrieved from the National
145	Centers for Environmental Prediction (NCEP) Global Ocean Data Assimilation
146	System (GODAS; Behringer and Xue, 2004). Ocean bathymetry was derived from the
147	NOAA ETOPO1 data (Amante and Eakins, 2009) and interpolated into 1.9° × 2.5°
148	horizontal resolution.
149	We used the CLIVAR MJO Working Group diagnostics package (CLIVAR,
150	2009) and a 20-100-day filter (Kaylor, 1977; Wang et al., 2014) to determine
151	intraseasonal variability. MJO phases were defined following the index (namely,
152	RMM1 and RMM2) proposed by Wheeler and Hendon (2004), which considers the
153	first two principal components of the combined near-equatorial OLR and zonal winds
154	at 850 and 200 hPa. The band-passed filtered data were used for calculating the index
155	and defining phases.
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157	2.2 Model description
158	2.2.1 CAM5.3
159	The CAM5.3 used in this study has a horizontal resolution of 1.9 $^{\circ}$ latitude \times
160	2.5° longitude and 30 vertical levels with the model top at 0.1 hPa. The MJO could
161	not be realistically simulated in the CAM5.3. Boyle et al. (2015) demonstrated that
162	although making the deep convection dependent on SST improved the simulation of
163	the MJO variance, it exerted a significant negative effect on the mean-state climate of
164	low-level cloud and absorbed shortwave radiation. By comparing the simulation 24

results of an uncoupled and coupled CAM5.3, Li et al. (2016) suggested that air-sea coupling and the convection scheme most significantly affected the MJO simulation in the climate model.

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2.2.2 1-D high-resolution TKE ocean model

169 170 The 1-D high-resolution turbulence kinetic energy (TKE) ocean model SIT was 171 used to simulate the diurnal fluctuation of SST and surface energy fluxes.— (Lan et 172 al., 2010; Tseng et al., 2014; Tu and Tsuang, 2005). The model was well verified 173 against surface and subsurface observations in in situ measurements on board the R/V 174 Oceanographic Research Vessel 1 and 3 over the South China Sea (Lan et al., 2010) 175 and on R/V Vickers over the tropical WP (Tu and Tsuang, 2005). Variations in sea 176 water temperature (T), current ($-\vec{u}$), and salinity (S) were determined (Gaspar et al., 177 1990) using the following equations. $\frac{\partial T}{\partial t} = (k_h + v_h) \frac{\partial^2 T}{\partial z^2} + \frac{R_{SR}}{\partial v_{tot} G_{tr}} \frac{\partial F}{\partial z}$ 178 $\frac{\partial \vec{u}}{\partial t} = -f \hat{k} \times \vec{u} + (k_m + v_m) \frac{\partial^2 \vec{u}}{\partial z^2}$ 179 $\frac{\partial S}{\partial t} = (k_h + v_h) \frac{\partial^2 S}{\partial x^2}$ 180 where R_{sn} is the net solar radiation at the surface (W m⁻²), F(z) is the fraction 181 182 (dimensionless) of R_{sn} that penetrates to the depth z, and k_h and k_m are eddy diffusion coefficients for heat and momentum (m^2 -s⁻¹), respectively. The value of k_h within the 183 cool skin layer and that of k_m within the viscous layer were set to zero. Molecular 184 185

transport is the only mechanism for the vertical diffusion of heat and momentum inthe cool skin and viscous layer, respectively (Hasse, 1971; Grassl, 1976; Wu, 1985). The parameters v_m and v_h are the molecular diffusion coefficients for momentum and temperature, respectively, ρ_{w0} is the density (kg m⁻³) of water, and e_{w} is the specific heat capacity at constant pressure (J kg⁻¹-K⁻¹). S is salinity (‰), \vec{u}

190 is the current velocity (m s⁻¹), f is the Coriolis parameter (dimensionless), and \hat{k} is the vertical unit vector (m s⁻¹).

The eddy diffusivity for momentum k_m is simulated using an eddy kinetic energy

193 approach based on the Prandtl–Kolmogorov hypothesis as follows:

$$k_m = c_k l_k \sqrt{E} \tag{3}$$

where $c_k = 0.1$ (Gaspar et al., 1990), l_k is the mixing length (m), and

196 $E = 0.5(u'^2 + v'^2 + w'^2)$ is turbulent kinetic energy. The turbulent kinetic energy (E)

197 is determined using a 1-D equation (Mellor and Yamada, 1982) as follows:

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$$\frac{\partial E}{\partial t} = \frac{\partial}{\partial z} k_m \frac{\partial E}{\partial z} + k_m \left(\frac{\partial \overline{u}}{\partial z}\right)^2 + k_h \frac{g}{\rho_W} \frac{\partial \rho_W}{\partial z} - c_{\varepsilon} \frac{E^{3/2}}{l_F}$$
(4)

199 where $c_{\varepsilon} = 0.7$ (Gaspar et al., 1990), g is the gravity (m s⁻²), ρ_{w} is the density of

water (kg m⁼³), and l_z is the characteristic dissipation length (m). The mixing length

 (l_k) and dissipation length (l_{ε}) were determined following the approach reported by

Gaspar et al. (1990). This approach is valid for determining the eddy diffusivity of

both the ocean mixed layer and surface layer.

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In the SIT model setting, the specific heat of sea water is a constant (4186.84 J

205 kg⁻¹-K⁻¹), and the Prandtl number in water is defined as the ratio of momentum

diffusivity to thermal diffusivity, which is a dimensionless number set as a constant

207 (1.0). The kinematic viscosity is a constant (1.14 \times 10⁻⁶ m² s⁻¹; Paulson and

Simpson, 1981), and the downward solar radiative flux into water with nine

209 wavelength bands was determined following the approach reported by Paulson and

Simpson (1981). The minimum turbulent kinetic energy is set to 10⁻⁶ m² s⁻², and the

211 zero displacement is set to 0.03 m.

The SIT model determines the vertical profiles of the temperature and

213	momentum of a water column from the surface down to the seabed, except in the
214	fixed ocean model bottom experiment. The default setting of vertical discretization
215	(e.g., in the control coupled experiment) is 41 layers with 12 layers in the first 10.5 m,
216	6 layers between 10.5 m and 107.8 m (Supplementary Information I). In the 1-D TKE
217	ocean model, temperature and salinity below 107.8 m, where vertical mixing is
218	greatly weakened, are nudged toward the climatological values of GODAS data until
219	4607 m. The extra high vertical resolution is needed to catch detailed temporal
220	variation of upper ocean temperature characterized by the warm layer and cool skin
221	(Tu and Tsuang, 2005). To account for the neglected horizontal advection heat flux,
222	the ocean is weakly nudged (by using a 30-day time scale) between 10 <u>.5 m</u> and
223	100107.8 m and strongly nudged (by using a 1-day time scale) below 100107.8 m
224	according to the NCEP GODAS climatological ocean temperature; no. No nudging is
225	performed for depths under 10 m. Considerably fine 41-layer vertical discretization is
226	applied, with 12 layers in within the upper-most 10.5 m. The resolution in the upper 10
27	m is considerably fine to capture the upper-ocean warm layer, and the thickness of the
228	first layer below sea surface is 0.05 mm to reproduce the ocean surface cool skin. The
229	41 levels are at the surface and at the depths of 0.05 mm, 1.0 cm, 2.0 cm, 3.0 cm, 4.0
230	cm, 5.0 cm, 6.0 cm, 7.0 cm, 8.0 cm, 9.0 cm, 10.0 cm, 16.8 cm, 29.5 cm, 43.6 cm, 59.2
231	cm, 76.9 cm, 96.8 m, 119.4 cm, 145.3 cm, 174.9 cm, 208.9 m, 248.3 cm, 293.8 cm,
232	346.8 cm, 408.4 cm, 480.2 cm, 564.3 cm, 662.6 cm, 777.9 cm, 913.1 cm, 1072.0 cm,
233	1258.8 cm, 1478.6 cm, 1737.3 cm, 2042.0 cm, 2401.1 cm, 2824.4 cm, 3323.6 cm,
234	3912.4 cm, and 4607.1 cm. The SIT model calculates data two timestwice for each
235	CAM5 time step (30 min; i.e., coupling 48 times per day).
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237	2.3. Experimental setupdesign
238	Five sets A series of 30-year numerical experiments (Table 1) were conducted to

39	investigate the effect of the air–sea interaction on the MJO simulation. In all-
40	simulations, The HadSST1 used to force the CAM5.3 coupled and uncoupled model
41	was forced by observed the climatological monthly SST except mean SST averaged
42	over 1982-2001. The monthly SST was linearly interpolated to daily SST fluctuation
43	that forced the model. The SST in theair-sea coupling region-where the SIT model
44	determined the upper ocean temperature. The was recalculated by the SIT during the
45	simulation, while the prescribed annual cycle of SST was used in the areas outside the
46	coupling region. Ocean bathymetry of the SIT was derived from the NOAA ETOPO1
47	data (Amante and Eakins, 2009) and interpolated into 1.9° × 2.5° horizontal
48	resolution.
49	All simulations were driven by the prescribed annual cycle of SST repeatedly for
50	30 years. The strategy is to evaluate the simulation capacity of climate models under
51	the same condition without considering interannual variation induced by SST. This
52	approach has been widely adopted in many studies (Delworth et al., 2006; Haertel et
53	al., 2020; Subramanian et al., 2011; Tseng et al., 2014; Wang et al., 2005).
54	Atmospheric initial conditions and external forcing such as CO ₂ , ozone, and
55	aerosol in near-equilibrium climate state around the year 2000 were taken from
56	F_2000_CAM5 component set based on CESM1.2.2 framework development. The
57	data has been commonly used in present-day simulations using CAM5 (e.g., He et al.,
58	<u>2017).</u>
59	The setup of five sets of experiment sets were conducted in this study are
60	described as follows.
61	(1) aA standalone CAM5.3 simulation forced by observed climatological monthly
62	SSTHadISST1 (A-CTL) and athe control experiment of coupled CAM5-SIT-
63	v1.0 simulation (C-30NS; 41 vertical levels, coupling in the entire tropics
64	between 30°SN and 30°NS with a diurnal cycle);).

265	(2) an upper Upper-ocean vertical resolution experiment (C LR12m and C LR34m):
266	two coarse vertical resolution simulations with a thickness of 11.8 and 34.2 m,
267	respectively, at the third layer; (3) a lower ocean boundary experiment: three Two
268	simulations with the lower boundary of the SIT model first layer centering at 12
269	m (C-LR12m) and 34 m (C-LR34m). Further details of the experimental design
270	are shown in supplementary Fig. S1.
271	(3) Shallow ocean bottom experiment: Three simulations with the ocean model
272	bottom at 10 m (C-HR1mB10m), 30 m (C-HR1mB30m), and 60 m (C-
273	HR1mB60m);) (supplementary Fig. S2).
274	(4) a regional Regional coupling experiment: Four simulations with fourthe coupling
275	domains, namely the latitudinal effect [region in 0°N-30°N (C-0_30N) and 0°S-
276	30°S (C-0_30S)] and the longitudinal) for latitudinal effect-[, and 30°E-180°E
277	(C-30_180E) and 30°E-75°W (C-30E_75W)] (see the) for longitudinal effect.
278	The coupling domain domains are shown in Fig. 1); and .
279	(5) a-A non-diurnal coupling experiment: a nondiurnal simulation (C-30NS-nD) that
280	considers the air–sea interaction by only once a day, namely, calculating ocean
281	surface fluxes SHF and LHF based on daily mean atmospheric variables and SST
282	(C 30NS nD), with. To prevent the inconsistent local time in different regions,
283	the coupling frequency maintained at each grid point remained 48 times per day
284	to prevent the local time in different regions from being inconsistent when
285	coupling once a day. Greenhouse gas concentrations were fixed at the using the
286	same daily means of atmospheric variables and SST at that particular point. In
287	contrast, the control simulation calculates air-sea fluxes 48 times a day based on
288	instantons values observed in the year 2000. A comparison between the non-
289	diurnal simulation and the control simulation reveals the effect of diurnal cycle
290	in air0sea coupling.

291	The main codes of the SIT model in Fortran 90 are packaging in independent
292	and original subprograms, with data and interface blocks in modules, that creates
293	explicit interfaces between the CAM5.3 and the SIT model without a coupler. In
294	addition, these modules contain dynamically allocable arrays and the independent I/O
295	procedures of the SIT model. The coupler in the CAM5 SIT only brokers
296	communication interchanges between the simulated SST and calculated oceanic
297	surface fluxes.
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300	3. Results and Discussion
301	The realistic simulation of the MJO has always been a major bottleneck in the
302	development of climate models. In this section, we demonstrate howthe sensitivity of
303	air-sea coupling experiments using a 1-D high-resolution ocean-mixed-layer model
304	significantly improves the MJO simulation by the CAM5.3. The period between
305	November and April when the MJO is the most prominent was the targeted season in
306	this study.
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308	43.1 Improvement of MJO simulation through air-sea coupling
309	This subsection compares the MJO simulation of the <u>control</u> coupled
310	modelexperiment (C-30NS) with that of the uncoupled AGCM (A-CTL) forced by
311	climatological monthly SST <u>of HadISST1</u> to demonstrate the effect of air–sea
312	coupling on the MJO simulation by coupling the SIT model to the CAM5.3 in the
313	tropical belt (30°N–30°S).
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315	43.1.1 Wavenumber-frequency spectra and eastward propagation characteristics
316	A wavenumber-frequency spectrum (W-FS) analysis was conducted to quantify

propagation characteristics simulated in different experiments. The spectra
of unfiltered U850 in observation ERA-I reanalysis, C-30NS, and A-CTL are shown
in Fig. 2a-c, respectively. The coupled C-30NS effectively simulated considering the
observed coupling in 30°N-30°S realistically simulates eastward-propagating signals
at zonal wavenumber 1 and 30-80-day periods (Fig. 2a-b), although with a slightly
larger amplitude: compared with ERA-I. By contrast, the uncoupled A-CTL diddoes
not effectively simulate the observed characteristics yield realistic simulation; instead,
it simulated simulates both eastward (wavenumber 1)- and westward (wavenumber 2)-
propagating signals with an unrealistic spectral shift to time scales longer than the
observed 30-80-day period.
The major features of the simulated MJO propagation were examined. Figure
2d-f show the time evolution of intraseasonal-precipitation and U850 anomalies in
Hovmöller diagrams; specifically, which represent lagged correlation coefficients
between the precipitation at averaged over 10°S-5°N, 75-100°E with and the average
precipitation at and U850 averaged over 10°N-10°S and U850 anomalies along the
equatoron intraseasonal timescales. Figure 2d indicates eastward propagation for both
precipitation and U850 from the eastern IO to the dateline, with precipitation leading
U850 by approximately a quarter of a cycle. The Hovmöller diagram derived from the
C-30NS (Fig. 2e) exhibits the key characteristics of eastward propagation for both
precipitation and U850 and the relative phases between the two, although the
simulated correlation wasis slightly weaker than that observed derived from GPCP
and ERA-I. By contrast, the uncoupled A-CTL simulated intraseasonal
signals that propagated propagate westward over the IO and simulated weak and much
slower eastward propagation crossing the MC and WP (Fig. 2f). The contrast between
Fig. 2e and 2f demonstrated demonstrate that coupling a 1-D-ocean TKE ocean model
alone could lead to a significant improvement in an AGCM in simulating the major

characteristics (e.g., amplitude, propagation direction and speed, and phase relationship between precipitation and circulation) of the MJO.

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43.1.2 Coherence of the simulated MJO

Cross-spectral analysis was performed conducted to examine the coherence and phase lag between tropical circulation and convection, which were plotted over the tropical wave spectra. Figure 2g-i show the symmetric part (e.g., Wheeler and Kiladis, 1999) of OLR and U850 in observation ERA-I/NOAA data, C-30NS, and A-CTL, respectively. We present only a magnified display of the spectra between the frequency of 0 to 0.35 day⁻¹ to highlight the MJO and equatorial Kelvin waves. The most prominent characteristic observed was characteristics seen in ERA-I/NOAA data are the peak coherence at wavenumbers 1–3 and a phase lag of approximately 90° in the 30-80-day band for the symmetric component associated with the MJO (Ren et al., 2019; Wheeler and Kiladis 1999). The coupled experiment C-30NS simulated simulates strong coherence in this low-frequency band (wavenumber 1) and exhibited exhibits a realistic phase lag relationship between U850 and OLR perturbations. However, the coherence at wavenumbers 2–3 for the 30–80-day period simulated by C-30NS wasis weaker than that observed. In addition, this in ERA-I/NOAA data. This undersimulation was also noted in CCSM4 (Subramanian et al., 2011), the uncoupled and coupled CAM4 and CAM5 (Li et al., 2016), and NorESM1-M (Bentsen et al., 2013), which had a version of the CAM as an AGCM. In summary, C-30NS produced considering the coupling between 30°N-30°S produces coherent and energetic patterns in the eastward-propagating intraseasonal fluctuations of U850 and OLR in the tropical IO and WP that are generally consistent with the MJO characteristics. By contrast, the MJO characteristics in A-CTL were are considerably weaker than those in C-30NS and that observed in ERA-I/NOAA data.

370	43.1.3 Horizontal and vertical structures of the MJO across the MC
371	Figure 2j-o show the horizontal and vertical structures of the MJO when deep
372	convection is the strongest over the MC (i.e., phase 5). Figure 2j-1 present the 20-
373	100-day filtered OLR (W m ⁻² , shaded) and 850-hPa wind (m s ⁻¹ , vector). C-30NS
374	realistically simulated the enhanced tropical convection over the eastern IO and the
375	Kelvin-wave-like easterly anomalies over the tropical WP despite undersimulating
376	the convection over the MC (Fig. 2j and 2k). By contrast, A-CTL failed to simulate
377	the enhanced convection over the eastern IO and MC; instead, it simulated
378	considerably weaker convection and easterly winds over the MC and WP,
379	respectively, than that observed in ERA-I/NOAA data (Fig. 2j and 2l).
380	Figure 2m-o show the vertical-longitudinal profiles of 20-100-day filtered
381	15°N-15°S averaged vertical velocity (OMEGA; Pa s ⁻¹ , shaded) and moist static
382	energy (MSE) anomalies (W m ⁻² , contour) at phase 5. The spatial distribution of
383	negative OMEGA (ascending motion) anomalies generally agreed with OLR
384	anomalies in C-30NS simulation and observation NOAA data over the Indo-Pacific
385	region (Fig. 2m and 2n). The observed relative relatively spatial relationship between
386	the ascending motion and MSE wasseen in ERA-I is well simulated in the coupled
387	experiment C-30NS. For example, positive MSE anomalies on the eastern side of the
388	anomalous ascent demonstrated demonstrate that the energy recharge process occurs in
389	advance of the MJO convection over the lower-tropospheric easterlies (Fig. 2j2m and
390	2k2n), whereas negative MSE anomalies on the western side revealed reveal that the
391	discharge process occurs during and after convection over the lower-tropospheric
392	westerlies. By contrast, this phase relationship, considered to be an essential feature
393	leading to the eastward propagation of an MJO (Hannah and Maloney 2014; Heath et
394	al., 2021), wasis not properly simulated in the uncoupled experiment A-CTL (Fig. 33

20), in which the simulated weak negative OMEGA <u>wasis</u> located between negative and positive MSE anomalies over weak lower-tropospheric wind anomalies and associated with weak convection over the MC (Fig. 21).

The observed temporal evolution of NOAA OLR and ERA-I U850 (Fig. 3a) indicated indicates that convection originating in the western IO wasis enhanced during its eastward propagation to the MC where it reached reaches the peak amplitude and then gradually weakened when continuing moving eastward to the dateline. In In the coupled experiment C-30NS, this evolution of convectively coupled circulation wasis realistically simulated, although it wasis weaker than the observed strength seen in NOAA OLR (Fig. 3b). Moreover, the split of convection into two cells off the equator in phase 6 wasis appropriately simulated in C-30NS (P6 in Fig. 3a and 3b). This split was caused by the topographic and land-sea contrast effects of the MC (Tseng et al., 2017). Associated with the split wasis the southward detouring of the anomalous convection during the passage of the MJO through the MC (Kim et al. 2017, Tseng et al., 2017; Wu and Hsu, 2009). After the passage of the MJO through the MC, the anomalous convection stayedstays south of the equator and continued moving eastward to the dateline. In the uncoupled A-CTL, the systematic eastward propagation of convectively coupled MJO circulation from the IO into the MC wasis not simulated. Instead, the convection over the MC developed develops in situ at a later stage than that observed (e.g., P6 in Fig. 3c) and dissipated rapidly. The A-CTL simulated simulates a pair of off-equator convection anomalies in the eastern IO during phase 2 (P2 in Fig. 3c) that movedmoves westward toward the central IO and were amplified at later stages (e.g., P4 in Fig. 3c). This unrealistic evolution explains the westward propagation tendency observed in the Hovmöller diagram (Fig. 2f).

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43.1.4 Characteristics of air—sea interaction

Figure 4a-c show the longitude-phase diagram in which the 20-100-day filtered precipitation (shaded) and SST (contour) anomalies were averaged over 10°S-10°N to determine the relationship between precipitation and SST fluctuations and to establish a link between air—sea coupling and convection. The propagation of the enhanced convection with positive SST anomalies to the east could be clearly seen in observation GPCP/OISST and the coupled experiment C-30NS (Fig. 4a and 4b). The highest SST anomaly (SSTA) ledleads the maximum precipitation anomaly by approximately 2–3 phases, and the SSTA began begins to decrease following the onset of enhanced precipitation. The observation revealed ERA-I and OISST data reveal the following relationship between net surface flux and SST: the decreased (increased) latent/sensible heat fluxesLHF/SHF and increased (decreased) downward radiation flux leading (lagging) the positive (negative) SSTA east (west) of anomalous deep convection. This well-known lead-lag relationship reflecting the active air-sea interaction in an MJO wasis realistically simulated in the coupled experiment C-30NS (not shown). The contrast between C-30NS and A-CTL confirms the key role of the air-sea interaction in contributing to the eastward propagation and demonstrates that the eastward propagation simulation can be markedly improved by incorporating the air-

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43.1.5 Vertically tilting structure

The warm SST was the key forcing that contributed to the boundary layer convergence before the onset of deep convection (Li et al., 2020; Tseng et al., 2014). Hence, the warmer upper ocean enhances the low-level atmospheric convergence and

sea interaction process in the model, even when using a simple 1-D ocean model such

then leads to enhanced low-level moisture and preconditioned deep convection and eastward propagation. This moistening process associated with warm ocean surface temperature wasis well simulated in the coupled experiment C-30NS but is not shown here. Instead, we present the coupling of moisture divergence (MD) and atmospheric circulation.

MD and zonal wind anomalies from the surface to the upper troposphere

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averaged over the 10°S-10°N and 120-150°E region are shown in Fig. 4d-f to depict the relationship between the vertically tilting structure of MD and zonal wind anomalies. Note that the active convection occurred around phase 5. The coupled experiment C-30NS (Fig. 4e) realistically simulated simulates the observed deepening of coupled MD and zonal wind anomalies with time (Fig. 4d). An evolution from the right to left seen in each panel of Fig. 4d-f wasis equivalent to the eastward movement of vertically tilting circulation from the eastern IO into the MC because of the eastward-propagating nature of the MJO. Figure 4d and 4e show that in both observation and ERA-I reanalysis and the coupled experiment C-30NS, the near-surface convergence (negative MD) occurring in the easterly anomalies ledlead the convection and continued deepening up to 500 hPa from phase 2 to phase 6 when the easterly anomalies switchedswitch to westerly anomalies. By contrast, this observed evolution of coupled MD-zonal wind anomalies were are not appropriately simulated in the uncoupled experiment (Fig. 4f). For example, a slow deepening with time wasis observed in the MD anomaly but not in the zonal wind anomaly that exhibited exhibits a vertically decayed structure, suggesting that MD and wind anomalies were are not well coupled, as noted in observation the ERA-I/NOAA data and the control coupled experiment.

In <u>observationthe ERA-I reanalysis data</u>, the negative near-surface MD anomalies <u>appeared appear</u> first under the easterly anomaly and <u>continued continued</u>

deepening between the easterly and westerly anomalies. This development in the phase relationship between MD and zonal wind anomalies in both observation ERA-I reanalysis data and the coupled simulation is consistent with the well-known structure embedded in the MJO, namely the near-surface convergence in the easterly phase (i.e., a boundary-layer moistening process; Kiranmayi and Maloney 2011; Li et al., 2020; Tseng et al., 2014), followed by the deep convection when transitioning to the westerly phase. This close phase relationship that is key to the eastward propagation wasis appropriately simulated in the coupled experiment but not in the uncoupled experiment.

43.1.6 Intraseasonal variance of precipitation

Figure 4g–i present the spatial distribution of intraseasonal variance of precipitation. In observation the GPCP data, the maximum variance wasis noted over the tropical eastern IO, MC, and tropical WP. The maximum variance south of the island in the MC and the equator in the tropical WP reflects the southward shift of the MJO deep convection when passing through the MC, partly due to the blocking effect of mountainous islands and the higher moisture content over high SST south of the equator in the region during boreal winter (Kim et al., 2017; Ling et al., 2019; Sobel et al., 2008; Tseng et al., 2017; Wu and Hsu, 2009). Although the control coupled experiment failedfails to simulate the variance maximum in the tropical eastern IO, it appropriately simulated imulates the maximum variance over the tropical WP, reflecting its ability to simulate the eastward propagation of the MJO through the MC. By contrast, the uncoupled A–CTL experiment simulated simulates considerably weaker intraseasonal variance in both the tropical eastern IO and the tropical WP. Figure 4j–l are the 20–100-day filtered SST (K, shaded) and 850-hPa wind (m s⁻¹, vector) during MJO phase 7 when deep convection is the strongest over the dateline.

The <u>coupled experiment C-30NS</u> realistically <u>simulated simulates</u> the negative SST
anomaly over the MC and WP when enhanced tropical convection passed through
the MC to the dateline, indicating the capability of the SIT model to reproduce the
observed SST anomaly by exchanging surface fluxes LHF/SHF between the
atmosphere and ocean. In A-CTL, no SST anomaly wasis evident because the model
was forced by prescribed climatological SST. The contrast seen in Fig. 4j-l
demonstrates the essential role of atmosphere-ocean coupling in shaping the MJO.
A delayed air-sea interaction mechanism was noted, where a preceding active-phase
MJO may trigger an inactive-phase MJO through the delayed effect of the induced
SST anomaly. In addition, the westerly winds at 850 hPa moving southward between
MC and WP wereare captured by the control experiment C-30NS and wereare
similar to the observed ERA-I reanalysis winds (Fig. 4j and 4k). By contrast, A-CTL
forced by climatological monthly SST (<0.05 K phase ⁻¹ anomaly) failed fails to
simulate the southward westerly wind of the region extending from the MC to the
dateline (Fig. 41).
4.23.2 Effect of upper-ocean vertical resolution
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in eastward-propagating wavenumber 1 at three timescales (e.g., longer than 80 days, 30-80 days, and approximately 30 days; Fig. 5a). In C-LR34m, both eastward and westward signals were are simulated with the dominant W-FS timescale that was longer than 80 days (Fig. 5b). The appearance of both eastward and westward signals at a lower frequency implied a stronger stationary tendency or weaker eastwardpropagating tendency. This result is consistent with that reported by Tseng et al. (2014) that the <u>scientific reproducibility of coarser the</u> resolution is, thecauses a longer intraseasonal periodicity and slower is the eastward propagation of the MJO. The effect of vertical resolution on the MJO simulation can be seen in the Hovmöller diagram. The eastward propagation simulated in C-LR12m (Fig. 5c) markedly weakened after crossing the MC compared with that simulated in the control experiment C-30NS- (Fig. 2e). In C-LR34m, the quasi-stationary fluctuation and westward propagation were are simulated over the IO (Fig. 5d), appearing similar to those in A-CTL- (Fig. 2f). The observed lead-lag relationship between precipitation (zonal wind) and SST wasis poorly simulated in C-LR12m (Fig. 5e) and even more poorly simulated in C–LR34m (Fig. 5f). This result confirms the finding reported by Tseng et al. (2014) that a higher vertical resolution in the firstupper few meters below the surface allows for a faster air—sea interaction, thus resulting in a more realistic simulation of the MJO.

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43.3 Effect of the lowest boundary of the SIT model

The ocean is a vital energy source for the MJO. Although vertical resolution is crucial for the efficiency of air—sea interaction, the thickness of the upper ocean that interacts with the atmosphere represents the <u>ocean</u> heat content to substantiate the MJO. A key question is how thick <u>an oceanica vertically-gridded ocean</u> mixed layer should be for a realistic simulation. To explore this issue, three experiments with a

model ocean with a 1-m vertical resolution and the ocean model (SIT) bottom at 10,
30, and 60 m, which included the top 11, 1312, 14, and 1516 levels, respectively, as
listedshown in Section 2supplementary Fig. S2 and Table 1, were conducted. The
spectra and the Hovmöller diagrams shown in Fig. 6a-c and Fig. 6d-f, respectively,
demonstrate that the thicker ocean model ocean simulated simulates a stronger MJO
with a frequency closer to those in the observation and an coupled experiment C-
30NS and ERA-I/NOAA data, and more realistic eastward propagation similar to that
in C 30NS and observations In addition, the lead-lag relationship between
precipitation (wind) and SST wasis more realistically simulated with increasing
thickness of the ocean model (Fig. 6g-i).
This result suggests that the thickness of the upper oceanoceanic mixed layer that
interacts with the atmosphere strongly affects the frequency of the simulated MJO. A
thinner (thicker) oceanic mixed layer is more quickly (slowly) recharged and
discharged through heatSHF and LHF exchange between the atmosphere and ocean
and therefore likely fluctuates at a faster (slower) tempo. The simulated periodicity is
therefore affected by the thickness of oceanic mixed layer (or ocean heat content).
Although this studythe result suggests 60 m is an appropriate thickness to realistically
simulate the periodicity of the MJO, we did not intend to suggest the exact thickness
required for a proper simulation because it might depend on the model. The oceanic
mixed layer should be adequately thick to contain a certain amount of heat to generate
appropriate periodicity that is close to that observed. However, the reason for the
intraseasonal time scale (i.e., 20-100 days) should be determined in future studies.
This finding does not suggest a constant periodicity because periodicity might be
affected by the time-varying structure of the atmosphere and ocean in the real world.

The MJO is a planetary-scale phenomenon. Given its large-scale circulation, the
air-sea interaction affecting the MJO likely occurs in a much larger area than the
region near the major convection anomalies. In this section, we discuss whether and
how the effect of coupling domain affects aon model's ability to simulate the eastward
propagation speed and periodicity of the MJO. Four experiments considering the
coupling in various domains (C-0_30N, C-0_30S, C-30_180E, and C-30E_75W,_
Fig. 1) were conducted to investigate for the effect of the coupling domain on the
eastward propagation speed and periodicity of the MJO in the simulation.purpose.
The results are shown in Fig. 7. The domains of the four experiments are shown in
Fig. 1. The C-0_30N that considered the coupling in the tropics between the equator
and 30°N simulatedsimulates the least realistic MJO propagation in terms of W-FS
(Fig. 7a), zonal wind-precipitation coupling (Fig. 7e), and SST-precipitation (Fig.
7i) of among the four regional coupling experiments. By contrast, coupling only the
tropics between the equator and 30°S simulated simulates a more realistic MJO in all
three aspects (i.e., spectrum in Fig. 7b, temporal evolution of precipitation/wind, and
precipitation/SST coupling in Fig. 7f and 7j). Figure. 8a indicates that the positive
precipitationnegative OLR anomalies at phase 5 simulated in C-0_30N stayedstays
mainly north of the equator and diddoes not shift southward in the MC as
observedrevealed in ERA-I reanalysis and in NOAA OLR and in the control
experiment C-30NS, and the convection over the IO wasis unrealistically weak. By
contrast, the southward detouring in the MC <u>wasis</u> realistically simulated in C-0_30S
that coupled only the tropical ocean between the equator and 30°S. This result
indicates that air-sea coupling occurring south of the equator is the key to producing
appropriate eastward propagation and detouring of the MJO through the MC. Without
this coupling, the C-0_30N experiment failed fails to realistically simulate the
eastward propagation of the MJO ₇ (Fig. 7e). This contrast can be attributed to the—

603	observed warmer ocean surface and higher moisture content found south of the
604	equator in boreal winter, which comprise a more favorable environmental condition
605	for air-sea coupling and convection-circulation coupling and the occurrence of the
606	MJO.
607	MJO simulations can be affected by air-sea coupling in the longitudinal domain.
608	Tseng et al. (2014) examined this effect by allowing coupling in different regions
609	(e.g., the IO, WP, and IO + WP) and found that the IO + WP coupling experiment
610	yielded the most satisfactory MJO simulation in terms of the zonal W-FS and
611	eastward propagation characteristics. In this study, we conducted sensitivity
612	experiments in which we allowed coupling in the tropics in two longitudinal domains,
613	namely $30^{\circ}E-180^{\circ}E$ (C-30_180E) and $30^{\circ}E-75^{\circ}W$ (C-30E_75W). The $30^{\circ}E-180^{\circ}E$
614	region covered the IO and WP, and the 30°E-75°W region covered the IO and the
615	entire tropical Pacific. As shown in Fig. 7, the C-30E_75W experiment simulated the
616	MJO, yielding results simulates more similar to the observation and those in C
617	30NS realistic MJO than tothe C-30_180E experiment, with stronger eastward
618	propagation and larger amplitudes in the spectrum (Fig. 7c and 7d) and Hovmöller
619	diagrams of precipitation/wind (Fig. 7g and 7h) and precipitation/SST (Fig. 7k and
620	7l). The simulated MJO in C-30E_75W propagated further farther east than that in C-
621	30_180E, particularly evident in Fig. 7k and 7l. The spatial distributions of circulation
622	and precipitation OLR shown in Fig. 8c and 8d indicated indicate the presence of a
623	stronger convective-coupled circulation system over the MC and WP in C-30E_75W.
624	These results suggest that coupling over the entire tropical IO and Pacific could
625	enhance the strength and eastward propagation of the MJO and encourage
626	further farther propagation to the central Pacific.
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The Previous studies showed that the diurnal cycle in the MC can weaken the
MJO and its eastward propagation (Hagos et al., 2016; Oh et al., 2013). We conducted
an experiment to determine whether the computing surface heat fluxes using daily
mean value values, instead of instantaneous values, of atmospheric variables and SST
with the same coupling frequency would affect the MJO simulation. The coupling in
the model was performed conducted through heat flux the SHF and LHF exchange
between the atmosphere and ocean, that were calculated based on simulated winds,
moisture, and temperature. As mentioned in Section 2.3, air-sea fluxes were
calculated twice for every time step (coupling 48 times per day,) in the control
coupled experiment (C-30NS) based on the instantaneous values of atmospheric and
oceanic variables. In the experiment in which the diurnal cycle was removed (C-
30NS-nD), air-sea fluxes were calculated as in C-30NS but were based on daily
mean data. means of both atmospheric variables and SST. Doing this removed certain
diurnal effects of air-sea coupling. The results shown in Fig. 9 reveal the enhancement
of the eastward-propagating signals in the MJO (e.g., a larger amplitude in spectrum;
Fig. 9a) and further eastward and faster propagation (Fig. 9b) as well stronger
coupling between precipitation and SST (Fig. 9c)-) in C-30NS-nD. The overall
results are consistent with previous finding that the diurnal cycle tends to reduce the
amplitude and propagation of the MJO, indicating that the weakening effect occurs
through air-sea coupling in addition to those processes in the atmosphere. Previous
studies have hypothesized that rapid interaction processes in the diurnal time scale
tend to extract energy from the MJO, thus reducing both the strength and propagation
tendency of the MJO. However, a comparison between the spectra of C-30NS and C-
30NS-nD indicated indicates that the experiment in which the diurnal cycle wasis
removed appeared to oversimulate the MJO with unrealistic strength, suggesting that
the effect of the diurnal cycle should be considered in the model to simulate a more

realistic MJO. However, whether this is a common result in different models remain to be examined.

Air-sea coupling is a key mechanism for the successful simulation of the MJO

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5 4. Discussion and conclusions

(Chang et al., 2019; DeMott et al., 2015; Jiang et al., 2015, 2020; Kim et al., 2010; Li et al., 2016; Li et al., 2020; Newman et al., 2009; Tseng et al., 2014). This study, following the study conducted by of Tseng et al. (2014), demonstrated that coupling a high-resolution 1-D TKE ocean model (namely the SIT model) to the CAM5, namely the CAM5–SIT, significantly improved the MJO simulation over the standalone CAM5. The CAM5 SIT realistically simulated By coupling SIT model to an AGCM different from Tseng et al. (2014), this study confirms the scientific reproducibility for the improvement of MJO simulation in modeling science. The CAM5–SIT realistically simulates the MJO characteristics in many aspects (e.g., intraseasonal periodicity, eastward propagation, coherence in the low-frequency band, detouring propagation across the MC, tilting vertical structure, and intraseasonal variance in the WP). Systematic sensitivity experiments were conducted to investigate the effects of the vertical resolution and the thickness of the 1-D ocean model, coupling domains, and the absence of the diurnal cycle. The results of all the sensitivity experiments are summarized in Fig. 10a and 10b, which show four common metrics for MJO evaluation. The four metrics are the propagation speed of the MJO (estimated from the U850 Hovmöller diagram as Fig. 2d-f) versus the power ratio of eastward- and westward-propagating 30–80-day signals (E/W ratio, derived from the zonal W–FS) in Fig. 10a— and the eastward propagation speed of the 30–80-day filtered

681 precipitation anomaly (estimated from the precipitation Hovmöller diagram) versus 682 the variance explained by RMM1 and RMM2 (i.e., the sum of the variance explained 683 by EOF1 and EOF2 based on Wheeler and Hendon, 2004) in Fig. 10b. Based on the 684 maximum precipitation anomaly and zero values of U850 (indicating deep convection 685 region), propagation speeds of precipitation and U850 were calculated from 686 Hovmöller diagrams between 60°E and 150°W. Overall, the control experiment C-687 30NS simulates the most realistic MJO among all sensitivity experiments. 688 As for vertical resolution, we determined that the MJO simulation efficiency 689 decreased when the vertical resolution of the SIT model wasis decreased from 1 m to 690 12 or 34 m, as observed in the C-LR12m and C-LR34m experiments, 691 respectively. This finding, consistent with that reported by Tseng et al. (2014), 692 suggests that a finer vertical resolution more effectively resolves temperature 693 variations in the ocean warm layer and enhances atmospheric—ocean coupling, thus 694 enabling the upper ocean to more efficiently respond to atmospheric forcing by 695 providing sensible and latent heat fluxes; this results in superior synchronization 696 between the lower atmosphere and the upper ocean. 697 We observed that the thinnershallower ocean mixed layermodel bottom could 698 speed up the eastward propagation of the MJO by producing more perturbations of 699 shorter periodicity (Fig. 6) and resultedresults in a weaker MJO. The shallower 700 oceanic mixed layer likely respondedresponds more quickly to atmospheric forcing 701 but provided provides less sensible and latent heat fluxes to the atmosphere. Thus, the 702 MJO propagated propagates too fast with a weaker amplitude. 703 In the coupling domain sensitivity experiments, we investigated the essential 704 coupling domain required to simulate the realistic MJO and the effect of the domain 705 on the MJO simulation. Coupling only the northern tropics failedfails to simulate the 706 eastward propagation, whereas coupling only the southern tropics vielded yields a

more realistic MJO simulation, although this simulation wasis inferior to coupling the entire tropics. This contrast reveals the importance of the southern tropical ocean, especially in the MC where high SST and moisture content are noted. Coupling in the southern tropics is therefore essential for providing the energy required to maintain the MJO and its eastward propagation. By contrast, the northern tropics are relatively dry and cool. Coupling in this region is therefore less effective in improving MJO simulation.

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In the longitudinal domain sensitivity experiments, we found that the MJO amplitude and the eastward extend of its eastward propagation were enhanced by extending the eastern boundary of the coupling domain from the tropical eastern IO to the tropical WP and further to the tropical eastern Pacific (Fig. 1). Further extension of the domain to cover the tropical Atlantic diddoes not exhibit further enhancement (not shown). This result indicates that coupling in the tropical central and eastern Pacific, although not the major MJO signal regions (i.e., from the tropical IO to the tropical WP), still played a marked role in sustaining the MJO. We propose the following to explain this effect. Because of the planetary scale of the MJO, the nearsurface easterly circulation to the east of the convection core often extended to the tropical central and eastern Pacific where the climatological easterly prevailed. The coupling beyond the WP increased low-level moisture transport and convergence to the east of the convection and establish an environment suitable for the further eastward propagation of the MJO. This effect was likely terminated by the landmass of Central America when the tropical Atlantic was further included. Thus, a further eastward extension of the coupling domain exerted little effect on further enhancing the MJO. A diagnostic study on the effect of the longitudinal coupling domain is being conducted, and the results will be reported in a following paper.

The diurnal versus nondiurnal cycle experiment indicated indicates that

733	nondiurnal coupling tended to enhance eastward-propagating signals but slow down
734	the <u>eastward</u> propagation- (Fig. 10a-b). This result is consistent with the finding of
735	previous studies that the diurnal cycle in the atmosphere extracts energy from the
736	MJO, thus weakening it.
737	In this study, we demonstrated how air-sea coupling can improve the MJO
738	simulation in a GCM. The findings are as follows.
739	(1) Better resolving the fine structure of the upper-ocean temperature and therefore
740	the air-sea interaction ledleads to more realistic intraseasonal variability in both
741	SST and atmospheric circulation.
742	(2) An adequate thickness of the oceanic mixed layer is required to simulate a delayed
743	response of the upper ocean to atmospheric forcing and lower-frequency
744	fluctuation.
745	(3) Coupling the tropical eastern Pacific, in addition to the tropical IO and the tropical
746	WP, can enhance the MJO and facilitate the further eastward propagation of the
747	MJO to the dateline.
748	(4) Coupling the southern tropical ocean, instead of the norther tropical ocean, is
749	essential for simulating a realistic MJO.
750	(5) Stronger MJO variability can be obtained without considering the diurnal cycle in
751	coupling.
752	Our study confirmed the effectiveness of air-sea coupling for improving MJO
753	simulation in a climate model and demonstrated how and where to couple. The
754	findings enhance our understanding of the physical processes that shape the
755	characteristics of the MJO.
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757	Code and data availability. The model code of CAM5-SIT is available at
758	https://doi.org/10.5281/zenodo.5510795. Input data of CAM5-SIT using the

- 759 climatological Hadley Centre Sea Ice and Sea Surface Temperature dataset and
- GODAS data forcing, including 30-year numerical experiments, are available at
- 761 https://doi.org/10.5281/zenodo.5510795.

- Author contributions. HHH is the initiator and the primary investigator of the
- 764 Taiwan Earth System Model project. YYL is the CAM5–SIT model developer and
- writes the majority part of the paper. WLT and LCJ assist in MJO analysis.

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767 *Competing interests.* The authors declare that they have no conflict of interest.

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- 769 Acknowledgements. The contribution from YYL, HHH, WLT, and LCJ to this study is
- supported by Ministry of Science and Technology of Taiwan under contracts MOST
- 771 110-2123-M-001-003, MOST 110-2811-M-001-603, MOST 109-2811-M-001-624
- and MOST108-2811-M-001-643. Our deepest gratitude goes to the editors and
- anonymous reviewers for their careful work and thoughtful suggestions that have
- helped improve this paper substantially. We sincerely thank the National Center for
- 775 Atmospheric Research and their Atmosphere Model Working Group (AMWG) for
- release CESM1.2.2. We thank the computational support from National Center for
- High530 performance Computing of Taiwan. This manuscript was edited by Wallace
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- 780 **Reference**
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Table 1. List of experiments 1095

Section	Category	Experiments	Description
4 <u>3</u> .1	Coupled or	A-CTL	Standalone CAM5.3 forced by observed forced
	uncoupled		by the monthly mean Hadley Centre SST dataset
			version 1 climatology
		C-30NS_(the_	CAM5.3 coupled with SIT over the tropical
		control coupled	domain (30°SN-30°NS), with 41 layers of finest
		experiment)	vertical resolution (up to submarine
			topographythe seabed) and diurnal cycle; the
			frequency of CAM5 being exchanged with CPL
			is 48 times per day
4 <u>3</u> .2	Upper- ocean vertical resolution	C-LR12m	The first ocean vertical level starts at 11.85 m_
			with 31 layers (beside SST and cool skin layer_
			are 11.5 m, 29.5 m and 43.6 m up to the seabed)
		C–LR34m	The first ocean vertical level starts at 34.2 m33.9
			m with 28 layers (beside SST and cool skin layer)
			are 33.9 m, 76.9 m and 96.8 m up to the seabed)
4 <u>3</u> .3	Lowest boundary of SIT	C–HR1mB10m	The lowest boundary of SIT has a depth of 10 m
			(middle gridmodel depth between 0 m and 10 m)
		C–HR1mB30m	The lowest boundary of SIT has a depth of 30 m
			(middle gridmodel depth between 0 m and 30 m)
		C–HR1mB60m	The lowest boundary of SIT has a depth of 60 m
			(middle gridmodel depth between 0 m and 60 m)
4 <u>3</u> .4	Regional coupling domain in	C-0_30N	Coupled in the tropical northern hemisphere
			(0°N–30°N, 0°E–360°E)
		C-0_30S	Coupled in the tropical southern hemisphere
	latitude		(0°S–30°S, 0°E–360°E)
	Regional coupling domain in	C-30_180E	Coupled in the Indo-Pacific (30°SN-30°NS,
			30°E–180°E)
		C-30E_75W	Coupled over the Indian Ocean and Pacific
	longitude		Ocean (30°SN-30°NS, 30°E-75°W)
4 <u>3</u> .5	Absence of the diurnal cycle	C-30NS-nD	Absence of the diurnal cycle in C–30NS; the
			CAM5.3 daily atmospheric mean of surface
			wind, temperature, total precipitation, net
			surface heat flux, u-stress and v-stress over
			water trigger the SIT and daily mean SST
			feedback to atmosphere; the frequency of CAM5
			is exchanged with CPL 48 times per day

- 1097 Experiment abbreviations: "A" means standalone AGCM simulation. "C" means the
- 1098 CAM5.3 coupled to the SIT model.

1099	Figure List
1100	Figure 1. Schematics of coupled and uncoupled domains in the regional coupling
1101	experiment: (a) C-30NS, (b) C-0_30N, (c) C-0_30S, (d) C-30_180E, and (e) C-
1102	30E_75W. The background is the climatological mean SST in December–February (DJF).
1103	
1104	Figure 2. (a)–(c) Zonal wavenumber–frequency spectra for 850-hPa zonal wind averaged
1105	over 10°S-10°N in boreal winter after removing the climatological mean seasonal cycle.
1106	Vertical dashed lines represent periods at 80 and 30 days, respectively. (d)-(f) Hovmöller
1107	diagrams of the correlation between the precipitation averaged over 10°S-5°N, 75-100°E
1108	and the intraseasonally filtered precipitation (color) and 850-hPa zonal wind (contour)
1109	averaged over 10°N-10°S. (g)-(i) Zonal wavenumber-frequency power spectra of
1110	anomalous OLR (colors) and phase lag with U850 (vectors) for the symmetric component
1111	of tropical waves, with the vertically upward vector representing a phase lag of 0° with
1112	phase lag increasing clockwise. Three dispersion straight lines with increasing slopes
1113	represent the equatorial Kelvin waves (derived from the shallow water equations)
1114	corresponding to three equivalent depths, 12, 25, and 50 m, respectively. (j)-(l)
1115	Composites of 20-100-day filtered OLR (W m ⁻² , shaded) and 850-hPa wind (m s ⁻¹ ,
1116	vector) for MJO phase 5 when deep convection is the strongest over the MC and 850-hPa
1117	wind, with the reference vector (1 m s ⁻¹) shown at the top right of each panel, and (m)-
1118	(o) 15°N-15°S averaged p-vertical velocity anomaly (Pa s ⁻¹ , shaded) and moist static
1119	energy anomaly (W m ⁻² , contour, interval 0.003); solid, dashed, and thick-black lines
1120	represent positive, negative, and zero values, respectively. The number of days used to
1121	generate the composite is shown at the bottom right corner of each panel. (a), (d), (g), (j),
1122	and (m) are from observations; the ERA-Interim and NOAA post-processed data (abbr.
1 123	ERA-I/NOAA); (b), (e), (h), (k), and (n) are from the control experiment C-30NS; and
1124	(c), (f), (i), (l), and (o) are from the A–CTL.
1125	
1126	Figure 3. Evolution of the filtered OLR anomaly (W m ⁻² , shaded) and 850-hPa wind (m
1127	s ⁻¹ , vector) at phase 2, 4, 6, and 8: (a) observation the ERA-I/NOAA data, (b) the control
1 128	coupled experiment C-30NS, and (c) the uncoupled experiment A-CTL. The unit of the
1129	reference vector shown at the top right corner of each panel is m s ⁻¹ , and the number of
1130	days used for the composite is shown at the bottom right corner of each panel.
1131	
1132	Figure 4. (a)–(c) Phase-longitude Hovmöller diagrams of 20–100-day filtered
1133	precipitation (mm day ⁻¹ , shaded) and SST anomaly (K, contour) averaged over 10°N-
1134	10°S from phase 1 to 8. Contour interval is 0.03; solid, dashed, and thick-black lines
1135	represent positive, negative, and zero values, respectively. (d)-(f) Phase-vertical
1136	Hovmöller diagrams of 20–100-day moisture divergence (shading, $10^{-6}~g~kg^{-1}~s^{-1}$) and

- zonal wind (contoured, m s⁻¹) averaged over 10°N–10°S, 120–150°E; solid, dashed, and
- thick-black curves are positive, negative, and zero values, respectively. (g)–(i) Variation
- of 30–60-day filtered precipitation in the eastern IO and the WP in observation (color
- shading), and the ratio between intraseasonal and total variance (contoured) and (j)–(l)
- 1141 composites 20–100-day filtered SST (K, shaded) and 850-hPa winds (m s⁻¹, vector) at
- phase 7 when deep convection was the strongest over the dateline. Reference vector
- shown at the top right corner of each panel. (a), (d), (g), and (j) are from the
- 1 45 <u>experiment C-30NS</u>; and (c), (f), (i), and (l) are from the <u>uncoupled experiment A-CTL</u>.
- 1 | 147 **Figure 5.** (a)–(b) Same as in Fig. 2(a) but for the C-LR12m and C-LR34m. (c)–(d)
- 1 | 148 Same as in Fig. 2(d) but for the C—LR12m and C—LR34m. (e)—(f) Same as in Fig. 4(a)
- 1 but for the C_LR12m and C_LR34m.
- Figure 6. Same as in Fig. 5 but for the C-HR1mB10m, C-HR1mB30m, and C-
- 1152 HR1mB60m.

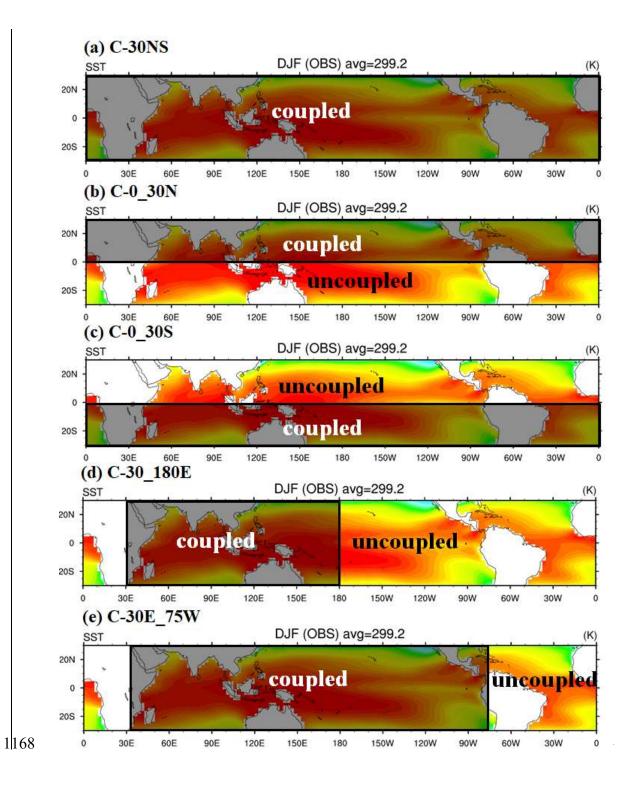
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- 1154 **Figure 7.** Same as in Fig. 5 but for the C-0 30N, C-0 30S, C-30 180E, and C-
- 1155 30E 75W.
- 1157 **Figure 8.** Same as in Fig. 3 but for phase 5 in the C-0 30N, C-0 30S, C-30 180E, and
- 1158 C-30E 75W.
- Figure 9. Similar as in Fig. 5 but for the C–30NS–nD.
- Figure 10. Scattered plots of various MJO indices in observation the ERA-I/NOAA data
- and 12 experiments: (a) power ratio of east/west propagating waves of wavenumber 1–3
- of 850-hPa zonal winds (X-axis) with a 30–80-day period and eastward propagation speed
- of U850 anomaly (Y-axis) from the Hovmöller diagram and (b) RMM1 and RMM2
- variance and eastward propagation speed of the filtered precipitation anomaly derived
- from the Hovmöller diagram.



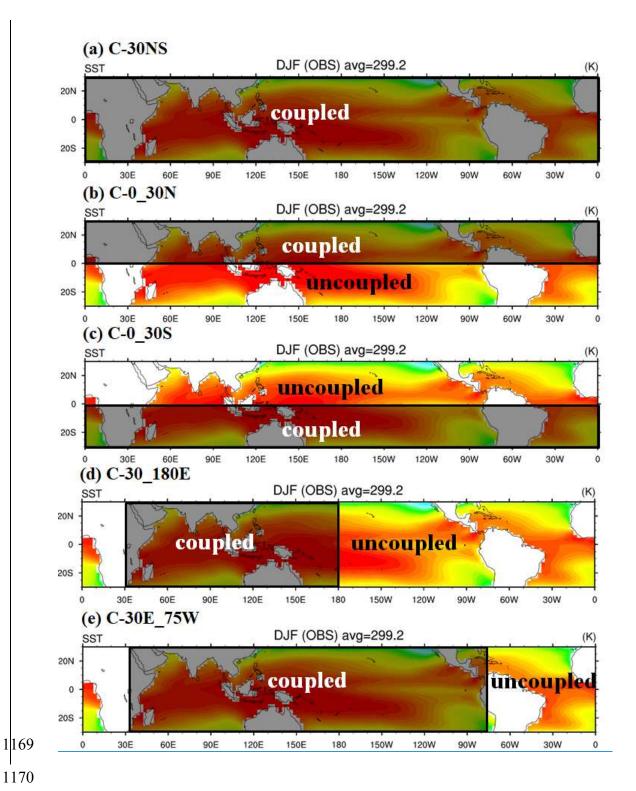
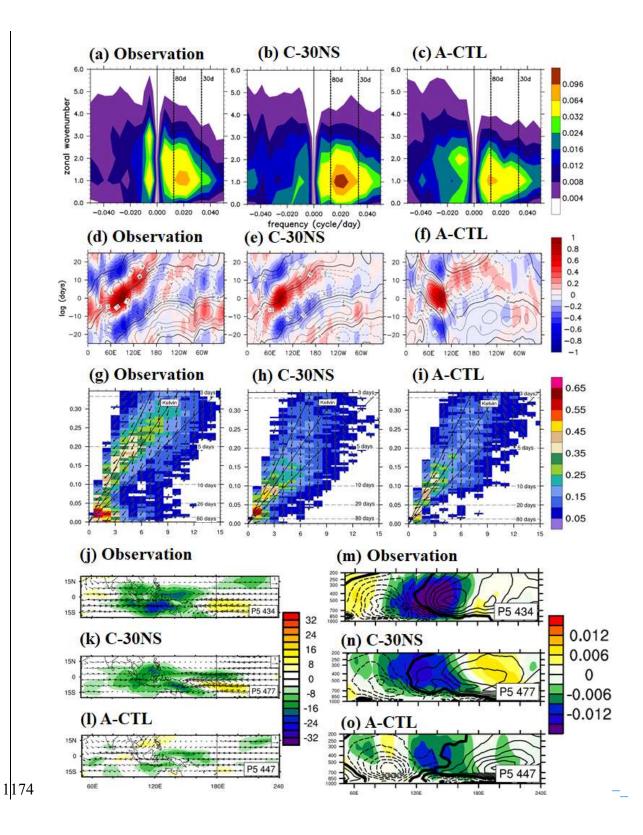


Figure 1. Schematics of coupled and uncoupled domains in the regional coupling experiment: (a) C-30NS, (b) C-0_30N, (c) C-0_30S, (d) C-30_180E, and (e) C-30E_75W. The background is the climatological mean SST in December–February (DJF).



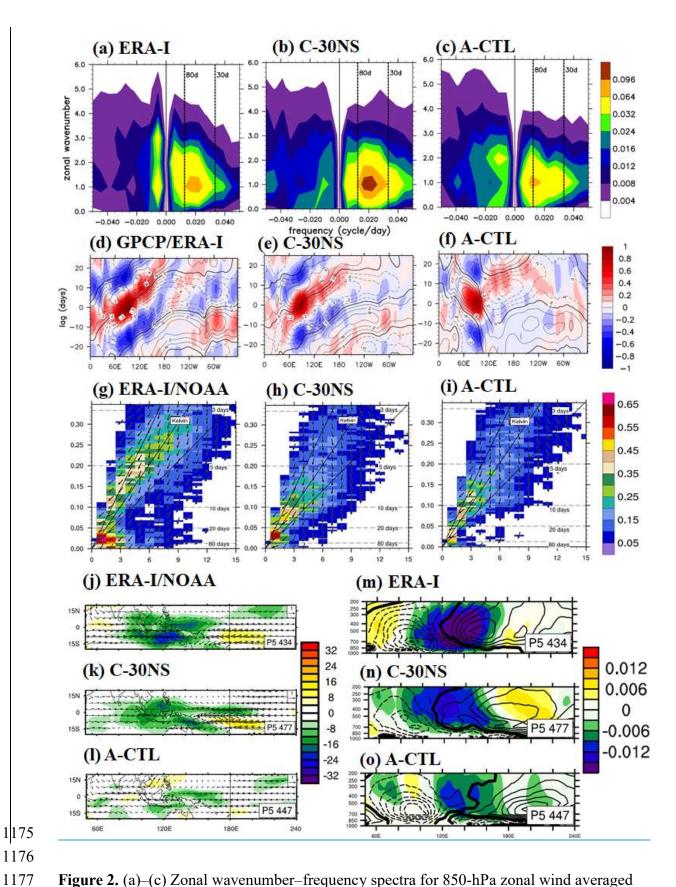


Figure 2. (a)—(c) Zonal wavenumber—frequency spectra for 850-hPa zonal wind averaged over 10°S–10°N in boreal winter after removing the climatological mean seasonal cycle. Vertical dashed lines represent periods at 80 and 30 days, respectively. (d)—(f) Hovmöller

diagrams of the correlation between the precipitation averaged over 10°S-5°N, 75-100°E 1180 and the intraseasonally filtered precipitation (color) and 850-hPa zonal wind (contour) 1181 averaged over 10°N-10°S. (g)-(i) Zonal wavenumber-frequency power spectra of 1182 anomalous OLR (colors) and phase lag with U850 (vectors) for the symmetric component 1183 of tropical waves, with the vertically upward vector representing a phase lag of 0° with 1184 phase lag increasing clockwise. Three dispersion straight lines with increasing slopes 1185 represent the equatorial Kelvin waves (derived from the shallow water equations) 1186 1187 corresponding to three equivalent depths, 12, 25, and 50 m, respectively. (i)–(1) Composites of 20–100-day filtered OLR (W m⁻², shaded) and 850-hPa wind (m s⁻¹, 1188 vector) for MJO phase 5 when deep convection is the strongest over the MC and 850 hPa 1189 wind, with the reference vector (1 m s⁻¹) shown at the top right of each panel, and (m)-1190 1191 (o) 15°N-15°S averaged p-vertical velocity anomaly (Pa s⁻¹, shaded) and moist static energy anomaly (W m⁻², contour, interval 0.003); solid, dashed, and thick-black lines 1192 represent positive, negative, and zero values, respectively. The number of days used to 1193 1194 generate the composite is shown at the bottom right corner of each panel. (a), (d), (g), (j), 1195 and (m) are from observations; the ERA-Interim and NOAA post-processed data (abbr. 1196 ERA-I/NOAA); (b), (e), (h), (k), and (n) are from the control experiment C-30NS; and (c), (f), (i), (l), and (o) are from the A-CTL. 1197

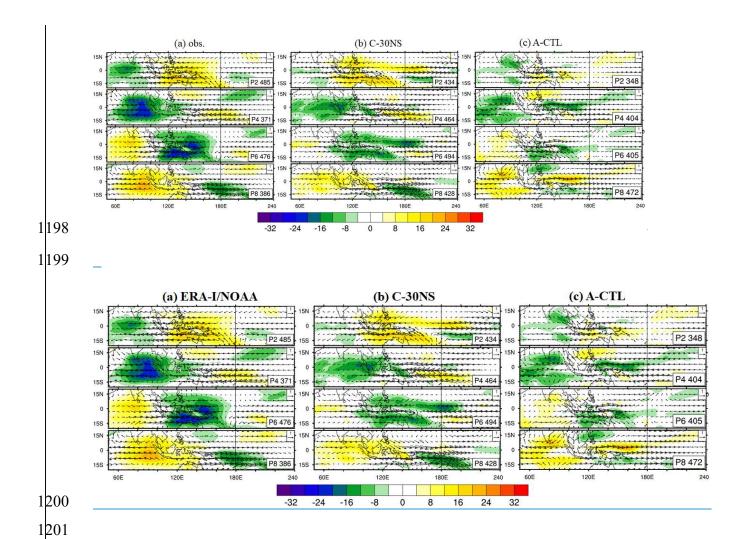
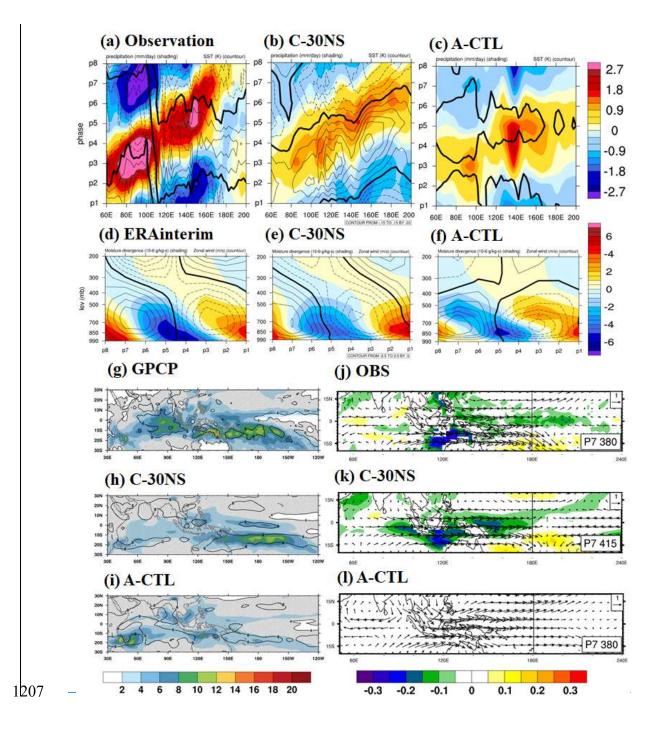


Figure 3. Evolution of the filtered OLR anomaly (W m⁻², shaded) and 850-hPa wind (m s⁻¹, vector) at phase 2, 4, 6, and 8: (a) observation the ERA-I/NOAA data, (b) the control coupled experiment C-30NS, and (c) the uncoupled experiment A-CTL. The unit of the reference vector shown at the top right corner of each panel is m s⁻¹, and the number of days used for the composite is shown at the bottom right corner of each panel.



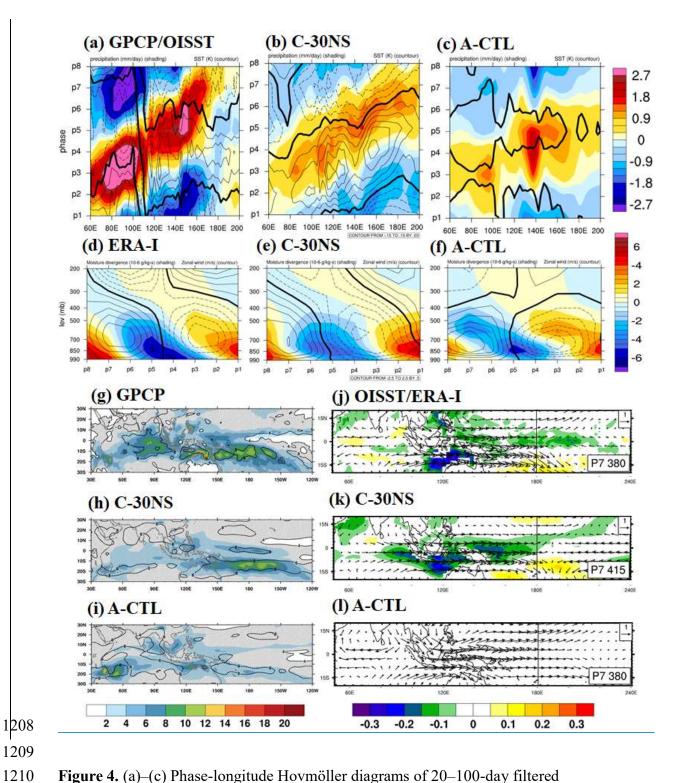
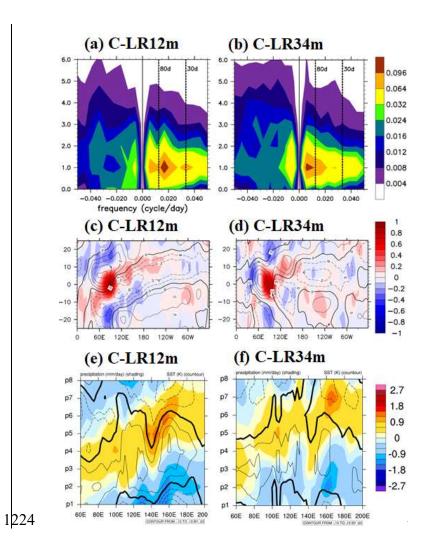


Figure 4. (a)–(c) Phase-longitude Hovmöller diagrams of 20–100-day filtered precipitation (mm day⁻¹, shaded) and SST anomaly (K, contour) averaged over 10°N–10°S from phase 1 to 8. Contour interval is 0.03; solid, dashed, and thick-black lines represent positive, negative, and zero values, respectively. (d)–(f) Phase-vertical Hovmöller diagrams of 20–100-day moisture divergence (shading, 10^{-6} g kg⁻¹ s⁻¹) and zonal wind (contoured, m s⁻¹) averaged over 10° N– 10° S, 120– 150° E; solid, dashed, and thick-black curves are positive, negative, and zero values, respectively. (g)–(i) Variation

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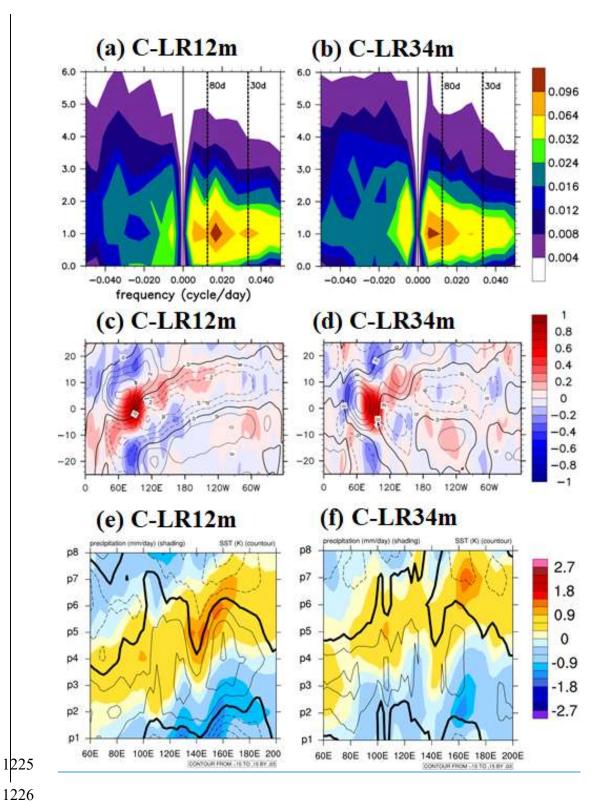
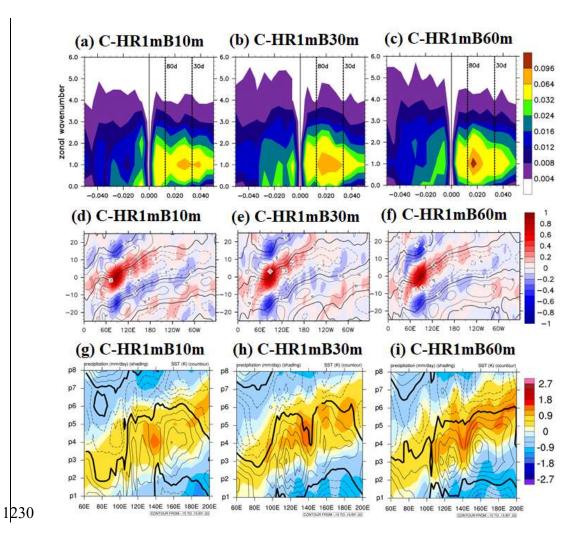


Figure 5. (a)–(b) Same as in Fig. 2(a) but for the C—LR12m and C—LR34m. (c)–(d) Same as in Fig. 2(d) but for the C—LR12m and C—LR34m. (e)–(f) Same as in Fig. 4(a) but for the C—LR12m and C—LR34m.



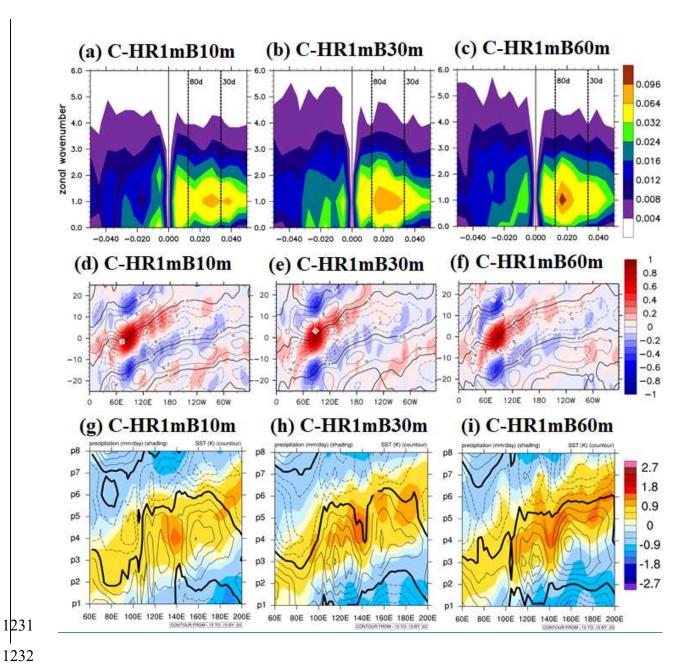
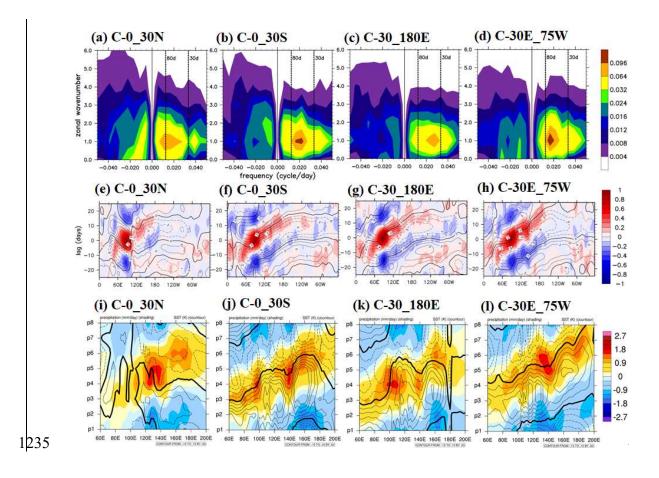


Figure 6. Same as in Fig. 5 but for the C–HR1mB10m, C–HR1mB30m, and C–HR1mB60m.



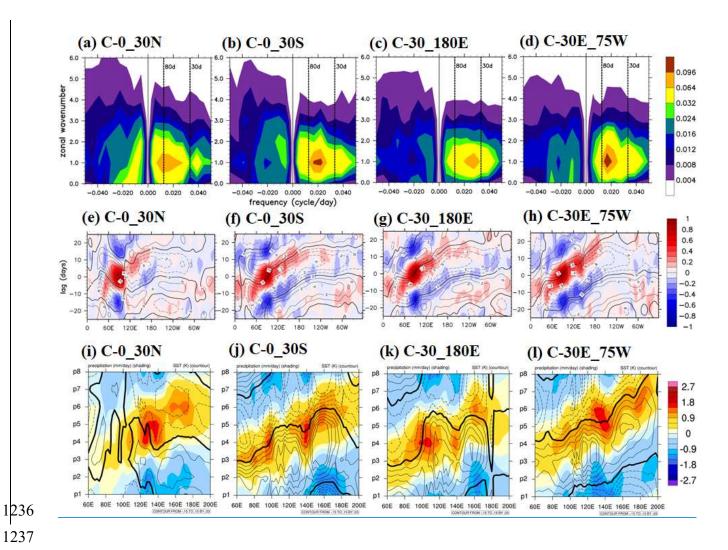
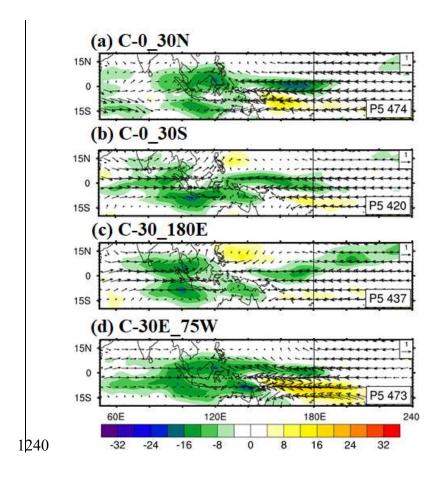


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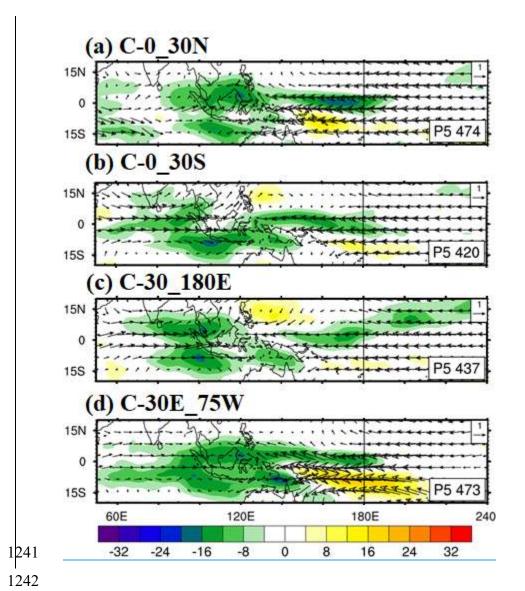


Figure 8. Same as in Fig. 3 but for phase 5 in the C-0_30N, C-0_30S, C-30_180E, and C-30E_75W.

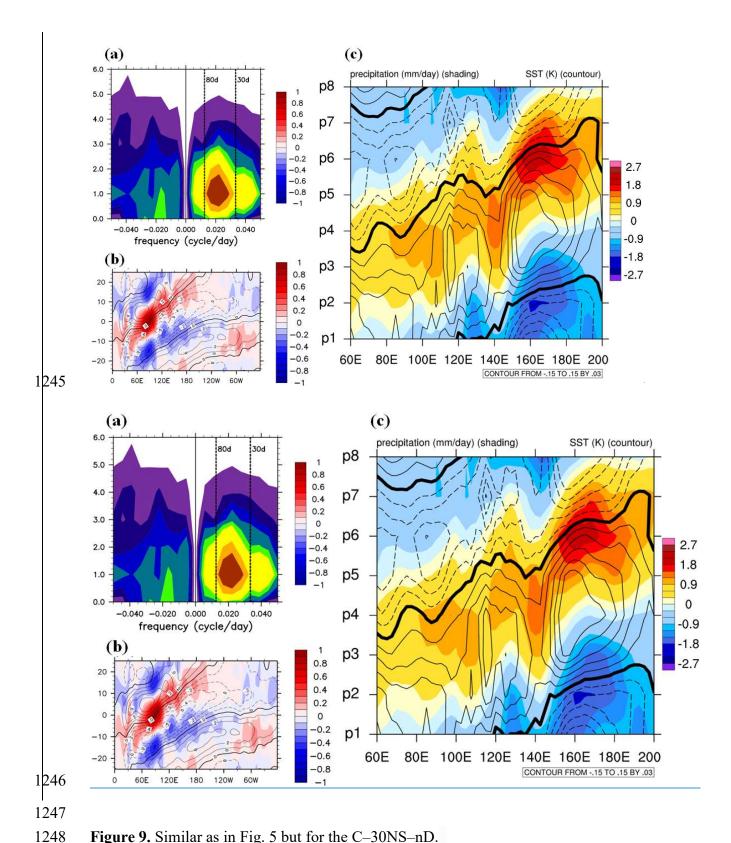


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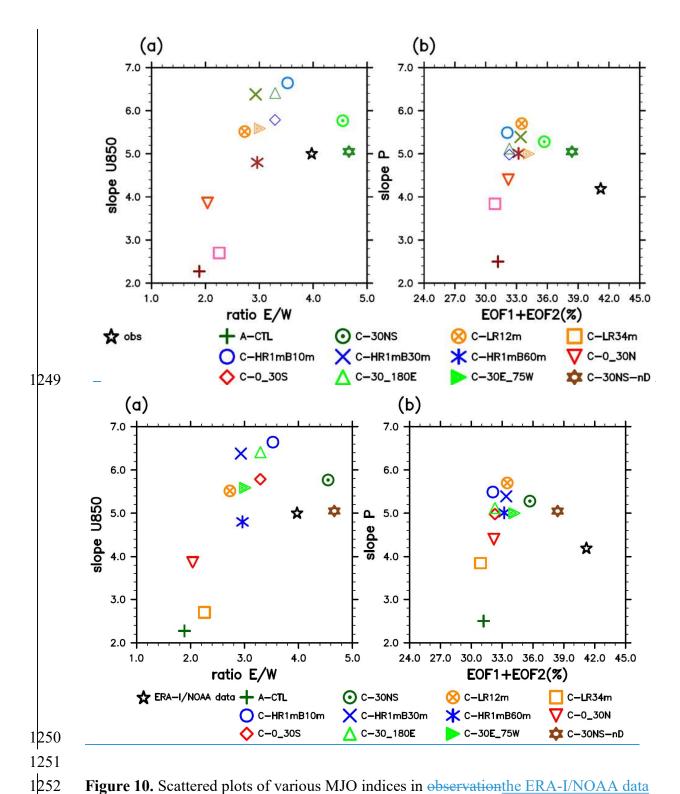


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Supplementary Information 1 2 Embedding a One-column Ocean Model (SIT 1.06) in the 3 Community Atmosphere Model 5.3 (CAM5.3; CAM5–SIT 4 v1.0) to Improve Madden-Julian Oscillation Simulation in 5 **Boreal Winter** 6 7 Yung-Yao Lan, Huang-Hsiung Hsu*, Wan-Ling Tseng, and Li-Chiang Jiang 8 9 10 Research Center for Environmental Changes, Academia Sinica, Taipei 11529, Taiwan 11 *Correspondence to: Huang-Hsiung Hsu (hhhsu@gate.sinica.edu.tw) 12

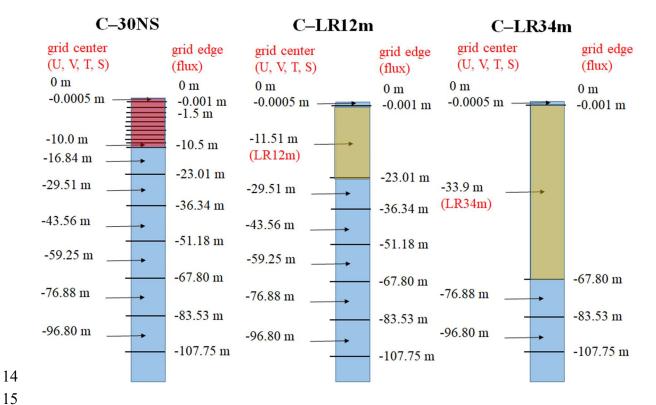


Figure S1. Diagram showing the vertical grid within 107.8 m in C–30NS, C–LR12m and C–LR34m, the model is as thick as 107.8 meters and with several layers between surface and model bottom. C–LR12m (31 vertical layers) and C–LR34m (28 vertical layers) have a first layer with grid center of 12 m and 34 m, respectively, but have the same vertical discretization as in the control experiment (C–30NS, 41 vertical layers) below the first layer.

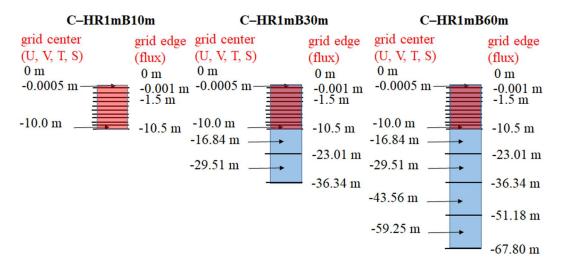


Figure S2. Diagrams showing the vertical grids in C–HR1mB10m, C–HR1mB30m and C–HR1mB60m. The model bottoms are 10, 30, and 60 m, respectively, unless the seabed is shallower than the above depth.