

1 **The Effects of Ocean Surface Waves on Global**  
2 **Intraseasonal Prediction: Case Studies with a Coupled**  
3 **CFSv2.0-WW3 System**

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16 **Abstract.** This article describes the implementation of a coupling between a global forecast model  
17 (CFSv2.0) and a wave model (WW3), and investigates the effects of ocean surface waves on [the](#) air-sea  
18 interface in the new framework. Several major wave-related processes, including the Langmuir mixing,  
19 Stokes-Coriolis force with entrainment, air-sea fluxes modified by Stokes drift and momentum roughness  
20 length, are evaluated in two groups of 56-day experiments, one for boreal winter and the other for boreal  
21 summer. Comparisons are made against in-situ buoys, satellite measurements and reanalysis data, to  
22 evaluate the influence of waves on intraseasonal prediction of sea surface temperature (SST), 2-m air  
23 temperature (T02), mixed layer depth (MLD), 10-m wind speed (WSP10) and significant wave height  
24 (SWH). The wave-coupled experiments show that overestimated SSTs and T02s, as well as  
25 underestimated MLDs at mid-high latitudes in summer from original CFSv2.0 are significantly improved  
26 due to enhanced vertical mixing generated by Stokes drift. For WSP10s and SWHs, the wave-related  
27 processes generally ~~reduce lead to reduction of~~ biases in regions where WSP10s and SWHs are  
28 overestimated. On one hand, the decreased SSTs stabilize [the](#) marine atmospheric boundary layer,  
29 weaken WSP10s and then SWHs. On the other hand, the increased roughness length due to waves  
30 ~~reduces leads to reduction in~~ the originally overestimated WSP10s and SWHs. In addition, the effects of  
31 Stokes drift and current on air-sea fluxes also rectify WSP10s and SWHs. These cases are helpful for the  
32 future development of the two-way CFSv2.0-wave coupled system.

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## 34 1 Introduction

35 Ocean surface gravity waves play an important role in modifying physical processes at the  
36 atmosphere–ocean interface, which can influence momentum, heat and freshwater fluxes across the air-  
37 sea interface (Li and Garrett, 1997; Taylor and Yelland, 2001; Moon et al., 2004; Janssen 2004; Belcher  
38 et al., 2012; Moum and Smyth, 2019). For instance, ocean surface waves modify ocean surface roughness  
39 to influence the marine atmospheric boundary layer and thus change the momentum, latent heat, and  
40 sensible heat transfer (Janssen 1989, 1991; Taylor and Yelland, 2001; Moon et al., 2004; Drennan et al.,  
41 2003, 2005). The breaking waves inject turbulent kinetic energy in the upper ocean, which enhances the  
42 mixing process (Terray et al. 1996). Nonbreaking surface waves also affect mixing in the upper ocean  
43 by adding a wave-related Reynolds stress (Qiao et al., 2004; Ghantous and Babanin, 2014). The wave-  
44 related Stokes drift interacts with Coriolis force and produces the Coriolis-Stokes force (Hasselmann  
45 1970). The shear of Stokes drift is critical for generation of Langmuir circulation, which significantly  
46 deepens the mixed layer by strong vertical mixing both at climate scales (Li and Garrett 1997; Belcher  
47 et al., 2012) and at weather scales (Kukulka et al., 2009).

48 Various wave-related parameterizations have been proposed and ~~applied~~used in ~~modeling~~modelling.  
49 The wave-related Charnock parameter ( $C_{ch}$ ) defines sea surface roughness and affects wind stress  
50 estimates (Pineau-Guillou et al. 2018; Sauvage et al. 2020). There are primarily three methods for  
51 ~~defining~~  $C_{ch}$ , assessed from the wave-induced kinematic stress (Janssen 1989, 1991), the wave age  
52 (Drennan et al., 2003, 2005; Moon et al., 2004), or the steepness (Taylor and Yelland, 2001). The former  
53 two are based on the wind-sea conditions, whereas the latter includes both swells and wind-sea waves.  
54 Modifications to these Charnock parameterizations were suggested in recent studies for the leveling off  
55 roughness under high winds (~~e.g.~~e.g., Fan et al., 2012; Bidlot et al., 2020; ECMWF, 2020; Li et al., 2021).

56 In the oceanic boundary layer, waves influence upper ocean mixing via wave dissipation and Stokes  
57 drift-induced processes. In Breivik et al. (2015), the wave dissipation-related turbulent kinetic energy  
58 flux is found to yield the largest sea surface temperature (SST) differences in the extratropics. The Stokes  
59 drift-induced Langmuir turbulence can improve temperature simulation over most of the world oceans,  
60 particularly in the Southern Ocean (Belcher et al., 2012; Li et al., 2016). Polonichko (1997), Van Roekel  
61 et al. (2012) and Li et al. (2017) indicated that the Langmuir cell intensity strongly depends on the  
62 alignment of winds and waves, reaching maximum when they are aligned. Li et al. (2016) found the  
63 effect of Langmuir cell can be further enhanced by entrainment. In Couvelard et al. (2020), the Stokes  
64 drift-related forces can also contribute [modestly](#) to [the](#) deepening of the mixed layer depth (MLD)  
65 [modestly](#). In the First Institute of Oceanography Earth System Model, Bao et al. (2019) indicated that  
66 the non-breaking wave-induced mixing, Stokes drift-affected air-sea fluxes, as well as sea spray are all  
67 important for climate estimates.

68 The wave-related processes at the air-sea interface are complex and important in global coupled  
69 systems (e.g., Breivik et al. 2015; Law-Chune and Aouf, 2018; Bao et al. 2019; Couvelard et al. 2020).  
70 Most of the coupled models with a wave component at global scale were developed for climate research  
71 (e.g., Law-Chune and Aouf, 2018; Bao et al. 2019; Couvelard et al. 2020). Exceptionally, an Integrated  
72 Forecasting System (IFS) with fully coupled atmosphere, ocean and wave components, developed by  
73 European Centre for Medium-Range Weather Forecasts (ECMWF) (Janssen 2004; Bidlot et al. 2019,  
74 2020), has been released with great flexibility for global forecasts from medium-range weather scales to  
75 seasonal scales (Breivik et al. 2015).

76 The effects of wave-related processes are worth further evaluation in different global coupled  
77 modelling systems. Since it takes sufficient periods for the wave energy to develop (Janssen 2004), we

78 investigate the impact of individual wave process at intraseasonal timescale in a new global atmosphere-  
79 ocean-wave system. To achieve this, we coupled the WAVEWATCH III (WW3) to the Climate Forecast  
80 System model version 2.0 (CFSv2.0) and then conducted sensitivity experiments in boreal winter and  
81 summer for comparison. The effects of upper ocean mixing modified by Langmuir cell, Stokes-Coriolis  
82 force and entrainment, air-sea fluxes modified by surface current and Stokes drift, and momentum  
83 roughness length are evaluated. The CFSv2.0 is a coupled system with [the](#) main application for  
84 intraseasonal and seasonal prediction ([e.g.e.g.](#), Saha et al. 2014). The National Centers for Environmental  
85 Prediction (NCEP) is establishing its [own](#)-atmosphere-ocean-wave system, in which the Global Forecast  
86 System (GFS; the atmosphere module in CFSv2.0 system) is one-way coupled with WW3. Our work can  
87 provide insights for two-way wave coupling of CFSv2.0, and is helpful for the future development of the  
88 CFSv2.0-wave coupling system. Two groups of 56-day predictions were conducted for boreal winter and  
89 boreal summer, respectively. [Then,](#) the predictions were [then](#)-compared with observations and reanalysis  
90 data. For each group, sensitivity experiments with different wave parameterizations were carried out to  
91 evaluate the effects of individual wave-related process.

92 The rest of the paper is structured as follows: methods and numerical experiments with different  
93 parameterizations are described in Section 2; the observations and reanalysis data are introduced in  
94 Section 3, and the results of experiments are evaluated and compared in Section 4. Finally, a summary  
95 and discussion are given in Section 5.

## 96 **2 Methods and Experiments**

### 97 **2.1 Coupling WAVEWATCH III with CFSv2.0**

98 The version 5.16 of WW3 (WAVEWATCH III Development Group, 2016) developed by the National

99 Oceanic and Atmospheric Administration (NOAA)/NCEP has been incorporated into the CFSv2.0 (Saha  
100 et al., 2014) as a new model component. The latitude range of WW3 is 78°S--78°N with a spatial  
101 resolution of 1/3°; the frequency range is 0.04118-0.4056Hz and the total number of frequencies is 25;  
102 the number of wave directions is 24 with a resolution of 15°; the maximum global time step and the  
103 minimum source term time step are both 180 s.

104 The CFSv2.0 contains two components, the GFS (details are available at  
105 <http://www.emc.ncep.noaa.gov/GFS/doc.php>) as the atmosphere component and the Modular Ocean  
106 Model version 4 (MOM4; Griffies et al., 2004) as the ocean component. The MOM4 is integrated on a  
107 nominal 0.5° horizontal grid with [the](#) enhanced horizontal resolution to 0.25° in the tropics, and has 40  
108 vertical levels; the vertical spacing is 10 m in the upper 225 m, and then increases in unequal intervals to  
109 the bottom at 4478.5 m. A three-layer sea ice model is included in MOM4 (Wu et al. 2005). The GFS  
110 uses a spectral triangular truncation of 382 waves (T382) in the horizontal, which is equivalent to a grid  
111 resolution of nearly 35 km, and 64 sigma-pressure hybrid layers in the vertical. The time steps of both  
112 MOM4 and GFS are 180 s. The ocean and atmosphere components are then coupled at the same rate. In  
113 the original two-way coupled system, the GFS receives SST from MOM4 and sends fluxes of heat,  
114 momentum, freshwater to MOM4 (black arrows in Fig. 1).

115 The Chinese Community Coupler version 2.0 (C-Coupler2; Liu et al., 2018) is [applied](#) used to  
116 interpolate and pass variables between atmosphere and wave components as well as ocean and wave  
117 components. Each component receives inputs and supplies outputs on its own grids. The C-Coupler2 is  
118 a common, flexible and user-friendly coupler, which contains a dynamic 3-D coupling system and  
119 enables variables to remain conserved after interpolation.

120 A schematic diagram of the coupled atmosphere-ocean-wave system is shown in Fig. 1. As illustrated,

121 WW3 is two-way coupled with MOM4 and GFS, through the C-Coupler2. WW3 is forced by 10-m wind  
122 from GFS (green arrows) and surface current from MOM4 (blue arrows), and then generates and evolves  
123 the wave action density spectrum. Meanwhile, the surface Stokes drift velocity, the Stokes transport and  
124 the turbulent Langmuir number are passed to MOM4 (red arrows; see Section 2.3) from WW3, and the  
125 surface Stokes drift velocity and the Charnock parameter are passed to GFS (red arrows; see Section 2.4  
126 and 2.5). The high frequency tail assumption for Stokes drift in WW3 is used with a spectral level  
127 decaying as  $f^{-5}$  (frequency). Additionally, the regular ocean surface current velocities from MOM4 are  
128 also passed to GFS, to calculate the relative wind velocity for the turbulent fluxes together with surface  
129 Stokes drift (blue arrows; see Section 2.4).

130 Both the CFSv2.0 and WW3 use warm starts; the initial fields at 00:00 UTC of the first day in each  
131 experiment for CFSv2.0 were generated by the real time operational Climate Data Assimilation System  
132 (Kalnay et al., 1996), downloaded from the CFSv2.0 official website  
133 (<http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod>). To get initial conditions for WW3, a stand-  
134 alone WW3 model is set up synchronously (see Section 2.2). Since the interactions between waves and  
135 sea ice are complicated and beyond the scope of the study, we turn off the coupling between WW3 and  
136 CFSv2.0 in areas with sea ice.

137 In addition, to properly select the coupling frequency between CFSv2.0 and WW3, the root mean  
138 square errors (RMSEs) of SST, significant wave height (SWH) and 10-m wind speed (WSP10) with  
139 different coupling steps for the fully coupled experiment (ALL; details in Section 2.6) are calculated and  
140 compared (Table S1 of the [supplementary supplementary material](#)). [The three components are coupled](#)  
141 [every time step \(180 s\) in 1\\_STEP\\_ALL experiment, every 5 steps \(900 s\) in 5\\_STEP\\_ALL experiment](#)  
142 [and every 10 steps \(1800 s\) in 10\\_STEP\\_ALL experiment. In 10\\_STEP\\_WW3, only the WW3 is coupled](#)

143 every 10 time steps, whereas the GFS and the MOM4 remain the one time step (180 s) coupling frequency  
144 as the original settings in CFSv2.0. From Table S1, the 10\_STEP\_WW3 experiment has a relatively short  
145 runtime and small RMSEs. Therefore, the time steps of the 10\_STEP\_WW3 are selected to compromise  
146 computing time consumption and model RMSEs.

## 147 **2.2 Initialization of WAVEWATCH III**

148 In WW3, input of momentum and energy by wind, and dissipation for wave-ocean interaction are two  
149 important terms (combined as input-dissipation source term) in the energy balance equation  
150 (WAVEWATCH III Development Group, 2016), which include the estimation of the Charnock  
151 parameter ~~related estimation~~. Several different packages to calculate the input-dissipation source term  
152 (ST) are available offered in the WW3 version 5.16, including ST2 (Tolman and Chalikov, 1996), ST3  
153 (Janssen, 2004; Bidlot, 2012), ST4 (Ardhuin et al., 2010), and ST6 (Zieger et al., 2015).

154 The initial wave fields were generated from 10-day simulation starting from rest in a stand-alone WW3  
155 model. To minimize the biases of initial wave fields, we tested simulations with ST2, ST3, ST4, and ST6  
156 schemes respectively, and compared the results with Janson-3 observations. Two 10-m wind datasets,  
157 the Cross-Calibrated Multi-Platform (CCMP; Atlas et al., 2011) data and the fifth generation European  
158 Centre for Medium-Range Weather Forecasts (ECMWF) Reanalysis (ERA5; Hersbach et al., 2020) data,  
159 were used to drive the wave model respectively. Comparing ~~ingee~~ all results, the ST4 scheme with ERA5  
160 wind forcing generates the minimum SWH bias (Table S2 in the supplementary supplementary material),  
161 consistent with findings in Stopa et al. (2016). Thus, the ST4 scheme was chosen to calculate the input  
162 and dissipation term, and generate initial wave fields with ERA5 wind forcing for experiments listed in  
163 Table 1. The parameters used for ST4 scheme followed TEST471f from WAVEWATCH III

164 Development Group (2016), which is the CFSR (CFS Reanalysis) tuned setup and is commonly-used at  
 165 global scale.

### 166 2.3 Parameterizations of Stokes Drift-Related Ocean Mixing

167 The full Stokes drift profile used in MOM4 is obtained by the method of Couvelard et al. (2020),  
 168 which is based on the work of Breivik et al. (2014; 2016). Breivik et al. (2016) derived the full Stokes  
 169 drift profile as

$$u_s B16(z) = u_s(0) [\exp(2k_p z) - \sqrt{-2\pi k_p z} \operatorname{erfc}(\sqrt{-2k_p z})], \quad (1)$$

170 ~~where~~The  $u_s(0)$  is the surface Stokes drift velocity,  $k_p = \frac{u_s(0)}{6V_s}$ ,  $V_s$  is the Stokes transport, and  $\operatorname{erfc}$   
 171 is the complementary error function. Eqn. 1 is depth-averaged within each vertical grid interval as

$$u_s(z) = \frac{u_s(0)}{(th)_k} [I(z_{k+1/2}, k_p) - I(z_{k-1/2}, k_p)], \quad (2)$$

$$I(z, k_p) = \frac{1}{6k_p} [e^{2k_p z} + 4k_p z \frac{u_s B16(z)}{u_s(0)}], \quad (3)$$

172 where  $th$  is the thickness of layer  $k$ , following Li et al., (2017), Wu et al., (2019) and Couvelard et al.,  
 173 (2020).

#### 174 2.3.1 Mixing of Langmuir Turbulence

175 McWilliams and Sullivan (2000) modified the turbulent velocity scale  $W$  in K-Profile Parameterization  
 176 (KPP) for vertical mixing by introducing an enhancement factor  $\varepsilon$ , to account for both boundary layer  
 177 depth changes and nonlocal mixing by Langmuir turbulence. Based on their work, Van Roekel et al.  
 178 (2012) improved the enhancement factor corresponding to alignment and misalignment of winds and  
 179 waves. Li et al. (2016) evaluated these parameterizations in a coupled global climate model, and found  
 180 that the difference between parameterizations with alignment and with misalignment was not significant,

181 owing to the relatively coarse resolution which cannot accurately represent the refraction by coasts and  
 182 current features. We ~~applied~~ used the parameterization from Van Roekel et al. (2012) as well. Because  
 183 the resolution in our model is relatively coarse too, and the angles between winds and waves are less than  
 184  $30^\circ$  in most areas (Fig. S1i&j in the ~~supplementary material~~ ~~supplementary~~), we didn't consider  
 185 misalignment in the study.

186  $W$  ( $W=ku_*/\phi$ , where  $u_*$  is the surface friction velocity,  $\phi$  is the dimensionless flux profile, and  
 187  $k=0.4$  is the von Kármán constant) ~~depends on~~ ~~varies in proportion to~~ the turbulent Langmuir number,  
 188 that is,

$$W = \frac{ku_*}{\phi} \varepsilon, \quad (4)$$

$$\varepsilon = \sqrt{1 + (3.1La_t)^{-2} + (5.4La_t)^{-4}}, \quad (5)$$

189 where  $La_t$  is the turbulent Langmuir number, defined as

$$La_t = \sqrt{\frac{u_*}{|u_s(0)|}}, \quad (6)$$

190 with  $u_s(0)$  ~~is~~ the surface Stokes drift velocity.

191 Furthermore, the enhanced  $W$  will influence the calculation of boundary layer depth. In KPP the  
 192 boundary layer depth is determined as the smallest depth at which the bulk Richardson number equals  
 193 the critical value  $Ri_{cr} = 0.3$ , that is,

$$Ri_b(h) = \frac{gh[\rho_r - \rho(h)]}{\rho_0[|u_r - u(h)|^2 + W^2]} = Ri_{cr}, \quad (7)$$

194 where  $g$  is acceleration of gravity,  $\rho$  is density,  $u$  is velocity,  $\rho_r$  is surface density,  $u_r$  is surface  
 195 velocity,  $\rho_0$  is ~~the an~~ average value of the density and  $h$  is the boundary layer depth. Hence, when  $W$   
 196 is enhanced, the boundary layer depth  $h$  is deepened accordingly.

### 197 2.3.2 Stokes–Coriolis Force and Associated Entrainment

198 Because the Stokes drift velocity is an increment superimposed on the original current velocity, the  
199 Coriolis force and the Stoke drift together produce an additional so-called Stokes–Coriolis (SC) force  
200 (Hasselmann 1970), that is,

$$SC\ Force = \overline{u_s(\mathbf{z})} \times f\hat{z}. \quad (8)$$

201 Here  $\overline{u_s(\mathbf{z})}$  is the Stokes drift velocity vector,  $f$  is the Coriolis frequency, and  $\hat{z}$  is the vertical unity  
202 vector. For consistency, the Stokes drift velocity is also included in advection terms of tracers ([e.g.,](#)  
203 temperature, salinity) and convergence terms (Law-Chune and Aouf, 2018; Couvelard et al., 2020). And  
204 the free surface condition for barotropic mode is correspondingly modified to

$$\frac{\partial \eta}{\partial t} = -\nabla M_{curr} - \nabla M_{st}, \quad (9)$$

205 where  $\eta$  is surface elevation,  $M_{curr}$  and  $M_{st}$  are the total vertical integral of regular Eulerian current  
206 and Stokes drift, respectively.

207 To depict the entrainment below the ocean surface boundary layer induced by Stokes drift, Li et al.  
208 (2016) suggested ~~adding to add~~ the square of surface Stokes drift velocity ( $|u_s(0)|^2$ ) to the denominator  
209 of Eqn. 7, that is,

$$Ri_b(h) = \frac{gh[\rho_r - \rho(h)]}{\rho_0[|u_r - u(h)|^2 + W^2 + |u_s(0)|^2]} = Ri_{cr}. \quad (10)$$

210 The boundary layer depth  $h$  in KPP from Eqn. 10 is then enhanced due to Stokes drift velocity.

### 211 2.4 Stokes Drift and Sea Surface Current on Air–Sea Fluxes

212 At [the](#) air-sea boundary layer, the momentum flux ( $\tau$ ), sensible heat flux (SH) and freshwater flux (E)  
213 are calculated as

$$\tau = \rho_a C_d |\Delta \vec{V}| \Delta \vec{V}, \quad (11)$$

$$SH = \rho_a C_h |\Delta \vec{V}| \Delta \theta, \quad (12)$$

$$E = \rho_a C_e |\Delta \vec{V}| \Delta q, \quad (13)$$

214 where  $C_d$ ,  $C_h$ ,  $C_e$  are surface exchange coefficients for momentum, sensible heat and freshwater.  $\rho_a$   
 215 is air density.  $\Delta \theta$ ,  $\Delta q$  are potential temperature and humidity differences between air and sea, and  $\Delta \vec{V}$   
 216 is velocity of air relative to water flow.

217 In CFSv2.0,  $\Delta \vec{V}$  is set to be wind speed ( $\overrightarrow{U_{wind}}$ ). However, the effect of ocean surface current should  
 218 not be ignored. Luo et al. (2005) first indicated that including ocean surface current ( $\overrightarrow{U_{surf}}$ ) improves  
 219 estimates of  $\tau$  and subsequent ocean response. Renault et al. (2016) further indicated that the  
 220 improvements of  $\tau$  by  $\overrightarrow{U_{surf}}$  also feed back into atmosphere. At present,  $\Delta \vec{V} = \overrightarrow{U_{wind}} - \overrightarrow{U_{surf}}$  is  
 221 widely used in coupled ocean-atmosphere models (e.g., Hersbach and Bidlot, 2008; Takatama et al.,  
 222 2017; Renault et al., 2021). Furthermore, Bao et al. (2019) indicated that as a part of the sea surface water  
 223 movement with speed magnitude comparable to surface current in mid-high latitudes, the surface Stokes  
 224 drift ( $\overrightarrow{u_s(0)}$ ) should also be included, that is,

$$\Delta \vec{V} = \overrightarrow{U_{wind}} - \overrightarrow{U_{surf}} - \overrightarrow{u_s(0)}. \quad (14)$$

225 ~~To account for the effects of the surface currents and of the Stokes drift, Eqn. 14 was used~~  
 226 ~~the effects of Stokes drift velocity, the Eqn. 14 was applied~~ in the coupled experiments (Table 1). To  
 227 complete the coupling, the corresponding modification of the tridiagonal matrix (Lemarié 2015) has been  
 228 implemented ~~is also conducted~~ in CFSv2.0. Note that the direction of Stokes drift is generally consistent  
 229 with 10-m wind (Fig. S1i&j in supplementary material), but the directions of surface  
 230 current and 10-m wind are usually different ~~with an angle~~ due to Coriolis effect (Fig. S1g&h).  
 231 Consequently, the effects of  $\overrightarrow{U_{surf}}$  and  $\overrightarrow{u_s(0)}$  on  $\Delta \vec{V}$  depend on the angles between them and  $\overrightarrow{U_{wind}}$ .

## 232 2.5 Parameterizations of Momentum Roughness

233 In CFSv2.0, the fluxes of momentum, heat, and freshwater are passed from atmosphere to ocean, and  
234 the estimates of them are critically important. The fluxes are in part determined by surface roughness  
235 length, which can be converted to surface exchange coefficients ~~in Eqn. 11-13~~ based on the Monin-  
236 Obukhov similarity theory (Monin and Obukhov, 1954).

### 237 2.5.1 The Momentum Roughness Length in GFS

238 In GFS, the momentum roughness length  $z_0$  has two terms. The first term  $z_{ch}$  is parameterized by  
239 the Charnock relationship (Charnock, 1955) representing wave-resulted sea surface roughness, and the  
240 second term  $z_{vis}$  is the viscous contribution (Beljaars, 1994) for low winds and smooth surface, that is,

$$z_0 = z_{ch} + z_{vis} = \frac{C_{ch}u_*^2}{g} + \frac{0.11\nu}{u_*}. \quad (15)$$

241 Here  $C_{ch} = 0.014$  is the constant Charnock parameter,  $\nu$  is the air kinematic viscosity. The relation of  
242  $z_0$  in GFS versus 10-m wind speed is shown in Fig.2 (black line).

### 243 2.5.2 The Charnock Relationship Related to Wave State

244 When ocean surface waves are explicitly considered, the Charnock parameter  $C_{ch}$  is not a constant  
245 (Janssen 1989, 1991; Taylor and Yelland, 2001; Moon et al., 2004; Drennan et al., 2003, 2005). In the  
246 study, we adopted a method developed by Moon et al. (2004), which considered the surface roughness  
247 leveling off under extremely high wind speed (Powell et al., 2003; Donelan et al., 2004). Based on  
248 observations, Moon et al. (2004) proposed Eqn. 16 to estimate the Charnock parameter by the wave age  
249  $\frac{c_{pi}}{u_*}$  ( $c_{pi}$  is the peak phase speed of the dominant wind-forced waves) with constant values of  $a$  and  $b$   
250 changing with 10-m wind speed every 5 m/s in the range of 10 m/s to 50 m/s.

$$C_{ch} = a \left( \frac{c_{pi}}{u_*} \right)^b, \quad (16)$$

251 To obtain continuous values of  $a$  and  $b$ , we derive a new relationship (Eqn. 17) to estimate  $a$  and  $b$  from  
 252 10-m wind speed  $U_{10}$  by fitting the values in Table 1 of Moon et al. (2004),

$$a = \frac{1}{0.1477U_{10}^2 - 0.7395U_{10} - 10.9995}, \quad (17)$$

$$b = 1.5661E^{-5}U_{10}^3 - 0.002U_{10}^2 + 0.1017U_{10} - 1.6182.$$

253 Because the observations in Moon et al. (2004) were obtained under tropical cyclones, Eqn. 17 is used  
 254 for  $U_{10} > 15$  m/s, whereas the original Charnock relationship of WW3 ST4 scheme (Janssen 1989, 1991)  
 255 is used for  $U_{10} \leq 15$  m/s. The revised parameterization is called ST4-M04. Figure S2 in  
 256 [supplementary supplementary material](#) shows the  $C_{ch}$  distribution obtained by Eqn. 16-17. In general  
 257 small wind direction variations at low latitudes lead to large wave age and thus low  $C_{ch}$ . The situation is  
 258 opposite at mid-high latitudes.

259 The relationships between  $z_0$  and  $U_{10}$  in GFS, WW3 ST4 scheme (Janssen 1989, 1991) and ST4-  
 260 M04 scheme were compared in Fig.2. The  $z_0$  in GFS increases relatively slowly with increasing wind  
 261 speed (black). The value of  $z_0$  from ST4 scheme (purple) increases rapidly with wind speed at high  
 262 winds. In comparison, in ST4-M04 scheme (blue) the rapid increase of  $z_0$  at high wind speed is  
 263 obviously restrained, although the mean  $z_0$  is slightly higher than that in GFS at wind speed  $>10$  m/s  
 264 due to larger  $C_{ch}$  ( $>0.014$  in Fig. S2). Furthermore, since the Charnock number is constant in GFS, the  
 265 standard deviation (STD) of  $z_0$  at a given wind speed is near zero. Since the  $z_0$  is determined only by  
 266 wind-sea conditions in ST4 and ST4-M04 scheme, the STD at a given wind speed is mainly owing to  
 267 variations in wind fetch and development stage of sea state. The reduced STDs in ST4-M04 scheme,  
 268 compared to ST4, imply less sensitivity of  $z_0$  to fetch and sea state. Note that the ST4-M04 is used in

269 GFS, while the  $z_0$  in WW3 is still calculated by the ST4 source term to avoid affecting the balance of  
270 adjusted wind input and dissipation.

## 271 **2.6 Set of Experiments**

272 A series of numerical experiments was conducted to evaluate the effects of aforementioned wave-  
273 related processes on ocean and atmosphere in two 56-day periods, from January 3 to February 28, 2017  
274 and from August 3 to September 28, 2018 for boreal winter and boreal summer, respectively.

275 The reference experiment (CTRL) is a one-way coupled experiment, in which CFSv2.0 provides 10-  
276 m wind and surface current to WW3, whereas no variable is transferred from WW3 to CFSv2.0. The  
277 results of CFSv2.0 in CTRL are consistent with the corresponding CFS Reanalysis data (Saha et al.,  
278 2010). For each period, four sensitivity experiments were carried out (Table 1). Based on CTRL, the first  
279 is [the](#) VR12-AL-SC-EN experiment, in which the Langmuir mixing parameterization is ~~applied~~[used](#) with  
280 Stokes–Coriolis force and entrainment in MOM4. The second is [the](#) Z0-M04 experiment, in which the  
281 constant  $C_{ch}$  in GFS is replaced by  $C_{ch}$  from WW3 ST4-M04 scheme. The effect of fluxes in GFS  
282 generated by  $\Delta\vec{V}$  (Eqn. 14) is tested in the FLUX experiment. The last experiment is the ALL, which  
283 includes all three parameterizations.

## 284 **3 Data**

285 Due to the availability of in situ and reanalysis data in the simulation periods, only sea surface  
286 temperature (SST), ocean subsurface temperature and salinity (T/S), 2-m air temperature (T02), 10-m  
287 wind speed (WSP10), and significant wave height (SWH) were used to evaluate the simulation results.

288 The daily average satellite Optimum Interpolation SST (OISST) data were obtained from NOAA, with

289 0.25° × 0.25° resolution (Reynolds et al., 2007; <https://www.ncdc.noaa.gov/oisst>). The global Argo  
290 observational profiles of T/S (Li et al., 2019) were from China Argo Real-time Data Center  
291 ([www.argo.org.cn](http://www.argo.org.cn)). The ERA5 datasets of T02, WSP10 and SWH with a spatial resolution of 0.5° were  
292 also used (Hersbach et al., 2020; [https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-](https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels)  
293 [single-levels](https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels)), which assimilated huge amounts of historical data and thus provided reliable hourly  
294 estimates. Additionally, the WSP10 and SWH observations from the available National Data Buoy  
295 Center (NDBC) buoy data (<https://www.ndbc.noaa.gov>) were ~~used~~ **applied** for comparison.

## 296 4 Experimental Results

297 In this section, an evaluation of simulation results was presented. Comparisons were made between  
298 model results and observations/reanalysis data. ~~The results in the first three days were excluded in the~~  
299 ~~evaluation, since the wave influences were weak at the beginning~~~~The results in the first three days were~~  
300 ~~excluded in the evaluation, since the initial wave influences were too weak.~~ Compared with observations  
301 or ERA5, the general increase of the biases in all experiments is likely a drift from the initial conditions  
302 since no data are assimilated.–

### 303 4.1 Sea Surface Temperature (SST) and 2-m Air Temperature (T02)

304 Figure 3a shows the spatial distribution of 53-day (day 4 to day 56) averaged SST biases in CTRL in  
305 boreal winter, defined as SST in CTRL minus OISST. The global mean SST bias is approximately 0.32°C,  
306 and the average RMSE is about 1.09°C from day 4 to day 56 in CTRL (Fig. 3a). The simulated SSTs are  
307 generally overestimated, and the large biases (>1.0°C) are mainly distributed in the Southern Ocean. In  
308 Fig. 3b, the global-averaged RMSEs of CTRL (black) increase with time in the first month and then

309 gradually level off. Compared with CTRL, the RMSEs are reduced continuously in VR12-AL-SC-EN  
310 and ALL (yellow and red), but not in Z0-M04 and FLUX (purple and blue).

311 To understand the critical key-process responsible for the bias reduction in ALL, the SST differences  
312 are compared across all four experiments (Fig.3c-f). Clearly, the difference in experiment VR12-AL-SC-  
313 EN is similar to that in ALL (Fig. 3c&3f). The spatial correlation coefficient between the SST differences  
314 with CTRL ~~of~~ SST of the two experiments (Fig. 3c&3f) is 0.67, significant at 99% confidence level, and  
315 the RMSEs of SST are not different significantly (red and yellow lines in Fig. 3b), indicating the Stokes  
316 drift-related parameterizations in VR12-AL-SC-EN mainly contribute to the SST positive bias reduction.  
317 ~~This contrasts with~~~~This is different with~~ Couvelard et al. (2020), where SST overestimations and MLD  
318 underestimations are reduced mainly due to the directly modified turbulence kinetic energy scheme. The  
319 global mean SST bias in ALL is 0.02 °C with RMSE of 1.03°C, and in most areas the SST differences  
320 compared with CTRL are significant ( $P \leq 0.05$ ) (dotted areas in Fig. 3f). Large SST improvements  
321 mainly appear in the Southern Ocean, with a regional RMSE decreases from 1.27 to 1.04 °C south of  
322 45°S (Fig. 3f and red line in Fig. 5a). The reduction of overestimated SSTs in CTRL (red in Fig 3a) is  
323 because the Stokes drift-related parameterizations in MOM4 inject turbulent kinetic energy into the  
324 ocean, which enhance vertical mixing, and subsequently cool the surface waters (Belcher et al., 2012; Li  
325 et al. 2016). The modified roughness and relative velocity in Z0-M04 and FLUX also influence upper  
326 ocean mixing (Fig. 3d&e) via changing momentum flux, and lead to generally warmer SSTs (purple and  
327 blue lines in Fig. 3b&5a). The effect from Stokes drift-related ocean mixing parameterizations dominates  
328 SST changes in ALL.

329 In boreal summer, the global mean SST bias in CTRL is overestimated approximately 0.29°C, and the  
330 averaged RMSE from day 4 to day 56 is about 1.19°C. The overestimated SSTs ( $>1.0^\circ\text{C}$ ) mainly occur

331 in the Northern Hemisphere (Fig. 4a). The global-averaged RMSEs are also generally lower in VR12-  
332 AL-SC-EN and ALL than in CTRL (Fig. 4b). The cooling effects in VR12-AL-SC-EN lead to a global  
333 mean bias of 0.06°C, and the large SST improvements mainly occur north of 50°N (Fig. 4c and yellow  
334 line in Fig. 5b). The changes of SST in Z0-M04 and FLUX (Fig. 4d&e; purple and blue lines in Fig.  
335 4b&5b) are relatively small. The global mean bias in ALL is 0.04°C with an RMSE of 1.14°C (Fig. 4f).

336 As aforementioned, large improvements of overestimated SST mainly occur at mid-high latitudes in  
337 local summer. The time series of RMSEs and correlation coefficients of SST between model and  
338 observation in the region (0-360°E, 45°-78°S in boreal winter and 0-360°E, 50°-78°N in boreal summer)  
339 are shown in Fig. 5c-f. The RMSEs in CTRL (blue in Fig. 5c&d) increase in the first few weeks and then  
340 gradually decrease afterward. Compared with CTRL, RMSEs in VR12-AL-SC-EN (yellow) and ALL  
341 (red) are significantly ( $P \leq 0.01$ ) reduced by about 0.3°C. The spatial correlation coefficients decrease  
342 with time but remain high ( $>0.90$ ) for all experiments (Fig. 5e&f) with higher values in experiment  
343 VR12-AL-SC-EN (yellow).

344 We also compared T02 from experiments with ERA5 (Fig. 6). Warm biases of T02 appear in both  
345 winter and summer in CTRL (Fig. 6a&b). The changes of T02 in sensitivity experiments (Fig. 6c-j) are  
346 generally consistent with the changes of SST in the same experiments (Fig.3&4). The correlation  
347 coefficients between the SST and the T02 changes for the ALL experiment in boreal winter and summer  
348 (Fig. 3f&6f and Fig. 4f&6j) are 0.61 and 0.53 respectively, significant at 99% confidence level. In boreal  
349 winter, all wave-coupled experiments except FLUX reduce the T02 mean bias (Fig.6c-f). VR12-AL-SC-  
350 EC has the largest T02 bias reduction compared with CTRL, from 0.55°C to 0.17°C (Fig.6c). In boreal  
351 summer, both VR12-AL-SC-EC and ALL have the largest T02 bias reduction, from 0.29°C to 0.08°C  
352 (Fig.6g&j). Noticeably, the improvements in RMSEs are not large for all experiments, because the

353 improvements mainly occur in areas with overestimated temperature.

#### 354 **4.2 Mixed Layer Depth (MLD)**

355 To further evaluate the direct effect of the wave-related processes on [the](#) upper ocean, we compared  
356 the MLD of all experiments with that estimated from Argo profiles in summer. The simulated T/S were  
357 interpolated onto the positions of Argo profiles at the nearest time. The MLD was estimated as the depth  
358 where the change of potential density reaches the value corresponding to a 0.2°C decrease of potential  
359 temperature with unchanged salinity from surface (de Boyer Montégut et al., 2004; Wang and Xu, 2018).

360 The time series of MLDs from numerical experiments and Argo south of 45°S in boreal winter (north  
361 of 45°N in boreal summer) are compared in Fig. 7a (7b). The simulated MLDs are generally within the  
362 STD of Argo MLDs (shading in Fig. 7). In CTRL, the mean bias (CTRL minus Argo) with STD is -  
363 13.15±7.82 m (-6.75±5.29 m) in boreal winter (summer). The correlation coefficient of MLDs in CTRL  
364 with Argo MLDs is 0.55 (0.68) with  $P \leq 0.01$ , and the mean RMSE is 15.30 m (8.55 m) in boreal winter  
365 (summer). In ALL, the mean bias (ALL minus Argo) with STD is 7.70±10.42 m (3.30±7.78 m) in boreal  
366 winter (summer), and the correlation coefficient of MLDs enhances to 0.63 (0.78). The RMSE south of  
367 45°S decreases from 15.30 m in CTRL to 12.96 m in ALL. The RMSE north of 45°N decreases from  
368 6.71 m in CTRL to 5.55 m in ALL in the first six weeks but the value increases in the last two weeks due  
369 to overestimation of MLDs. Compared with CTRL (orange in Fig. 7), VR12-AL-SC-EN (yellow) and  
370 ALL (dark blue) show significant improvements ( $P \leq 0.01$ ) on the underestimated MLDs time series,  
371 whereas the MLDs difference between CTRL and Z0-M04 (purple)/FLUX (blue) is non-significant.

### 372 4.3 Wind Speed at 10 m (WSP10) and Significant Wave Height (SWH)

373 Compared with ERA5, the WSP10s in CTRL are generally overestimated in both winter and summer  
374 (Fig.8a&9a). The global averaged RMSEs of WSP10s in CTRL are 4.25 m/s (4.26 m/s) in boreal winter  
375 (summer). The global averaged RMSEs of WSP10s in all experiments increase with time in the first two  
376 weeks and then gradually level off ~~with perturbations~~ (Fig. 8b&9b). The differences of RMSEs between  
377 CTRL and other experiments are tiny in the first 10 days, and afterwards the RMSEs in Z0-M04 and  
378 ALL (purple and red) become clearly smaller than in CTRL over most of the time.

379 The comparisons of the simulated SWHs in CTRL with the ERA5 also show that the SWHs are  
380 overestimated in both winter and summer (Fig. 10a&11a). In boreal winter, the global mean SWH bias  
381 in CTRL is approximately 0.20 m with overestimates ( $> 0.30$  m) in the Pacific, the North Atlantic and  
382 the Southern Ocean (Fig.10a), and the average RMSE is about 1.29 m. In boreal summer, the global  
383 mean bias in CTRL is approximately 0.17 m with 1.22 m RMSE (Fig. 11a). Similar to WSP10s, the  
384 RMSEs of SWHs also increase in the first two weeks and then gradually level off ~~with perturbations~~ (Fig.  
385 10b&11b). The RMSEs in Z0-M04 and ALL (purple and red) are smaller than in CTRL over most of the  
386 time, consistent with changes of WSP10s. The correlation coefficients between changes of WSP10s and  
387 changes of SWHs in ALL are 0.77 and 0.73 in boreal winter and summer respectively (Fig. 8f&10f and  
388 Fig. 9f&11f), significant at 99% confidence level, indicating that the SWHs changes are closely related  
389 to changes of wind speeds.

390 In VR12-AL-SC-EN, the reduction of SST warm biases affects air temperature and stabilizes [the](#)  
391 marine atmospheric boundary layer (Sweet et al. 1981; O'Neill et al. 2003), and subsequently reduces  
392 WSP10s and SWHs with decreased global bias in boreal winter (Fig.8c&10c). In Z0-M04, the  
393 overestimated WSP10s and SWHs are also reduced (Fig. 8d&10d) due to the larger  $z_0$  with the ST4-

394 M04 scheme at wind speed  $> 10$  m/s (Fig. 2). The increase of  $z_0$  enhances wind stress and momentum  
395 transferred into the ocean, and therefore reduces surface winds (Pineau-Guillou et al. 2018; Sauvage et  
396 al. 2020) and consequently reduces SWHs. In FLUX (Fig. 8e&10e),  $\overline{U}_{\text{surf}}$  and  $\overline{u_s(0)}$  decrease wind  
397 stress and momentum transfer when their directions are consistent with wind directions, and vice versa  
398 (Hersbach and Bidlot, 2008; Renault et al., 2016). For instance, the angles between wind and current are  
399 relatively small ( $<90^\circ$ ) in the northeastern Pacific, reducing the wind stress and thus enhancing WSP10s  
400 (Fig. 8e). In contrast, the large angles ( $>90^\circ$ ) between the northwesterlies and the Kuroshio in the  
401 northwestern Pacific enhance wind stress, and decrease WSP10s (Fig. 8e). Consequently, ~~the~~  
402 improvements occur in areas with misalignment of winds and currents. With all combined effects, the  
403 biases of WSP10s and SWHs in ALL in most regions are decreased (Fig. 8f&10f), with the reduced  
404 global RMSEs of 4.17 m/s and 1.18 m respectively. In boreal summer, the improvements of WSP10s  
405 and SWHs are relatively small in terms of global averaged RMSEs, because of smaller positive biases in  
406 CTRL (Fig. 9a&11a). In ALL, the global averaged bias of WSP10s (SWHs) is -0.01 m/s (0.03 m). The  
407 largest reduction primarily appears in the Southern Ocean (Fig. 9f&11f) to improve the overestimated  
408 westerlies and SWHs in CTRL (Fig. 9a&11a).

409 Previous studies indicated that ocean surface winds in ERA5 are underestimated in some regions  
410 (Belmonte Rivas and Stoffelen 2019; Kalverla et al. 2020; Sharmar and Markina 2020). To better  
411 demonstrate the effects of waves on WSP10s and SWHs, comparisons of WSP10s and SWHs with the  
412 NDBC buoy data are made (Table 2 and Fig. 12). The differences between sensitivity experiments and  
413 CTRL are all statistically significant at 95% confidence level. Buoys are mainly located in the  
414 northeastern Pacific, the tropical Pacific and the northwestern Atlantic oceans (Fig. S3), and buoy  
415 identifiers with total numbers, longitudes and latitudes are listed in Table S3. The method from Hsu et al.

416 (1994) was used to adjust wind speeds from buoy data to the reference height of 10 m.

417 Compared to the NDBC data, the WSP10s and the SWHs in CTRL are generally overestimated in both  
418 winter and summer with positive mean biases (Table 2 and Fig. 12). The reduction of mean biases appears  
419 in all experiments except FLUX in boreal winter. The wave-related processes are most effective in areas  
420 with positive biases, consistent with previous comparisons with ERA5. In boreal winter, the angles  
421 between winds and currents are small. The wind stresses are then reduced in FLUX, and the WSP10s are  
422 enhanced. So the positive bias is further enhanced. The improvements in ALL are generally the largest  
423 (Table 2), with the WSP10s RMSE of 1.04 m/s (1.15 m/s) and the SWHs RMSE of 0.36 m (0.24 m) in  
424 boreal winter (summer). As shown in Fig. 12, with the increase of WSP10s and SWHs, the reduction of  
425 overestimation in ALL compared with CTRL is more prominent.

## 426 **5 Summary and Discussion**

427 To investigate the individual role played by wave-related processes on atmosphere and ocean interface  
428 in a coupled global atmosphere-ocean-wave ~~modeling~~modelling system on intraseasonal scale, we  
429 implemented the version 5.16 of WW3 into CFSv2.0 for global oceans from 78°S-78°N, using the C-  
430 Coupler2. In this coupled system, the WW3 was forced by 10-m wind and surface current generated in  
431 CFSv2.0. Stokes drift-related Langmuir mixing, Stokes-Coriolis force and entrainment in ocean, air-sea  
432 fluxes modified by surface current and Stokes drift, and momentum roughness length ( $z_0$ ) were  
433 considered separately, and the results of sensitivity experiments were compared against in-situ buoys,  
434 satellite measurements and ERA5 reanalysis. The effects of waves on intraseasonal prediction were  
435 examined in two 56-day cases, one for boreal winter and the other one for boreal summer.

436 The following key results were found:

- 437 1. Overestimated SST, T02 and underestimated MLD in [the](#) mid-high latitudes in CFSv2.0 are  
438 significantly improved, particularly in local summer. Because enhanced vertical mixing  
439 generated by Langmuir turbulence, Stokes-Coriolis force and entrainment in VR12-AL-SC-  
440 EN changes temperature structure in the upper ocean, and further affects air temperature. In  
441 boreal winter, the regional RMSE of SST (T02) in the Southern Ocean decreases from 1.27  
442 (1.93) in CTRL experiment to 1.04 (1.67) °C in ALL experiment. In boreal summer, the effect  
443 is weaker because of the smaller ocean areas in [the](#) mid-high latitudes of the Northern  
444 Hemisphere.
- 445 2. In general, all wave-related processes ~~reduce~~[lead to reduction of](#) biases for WSP10s and SWHs,  
446 particularly in regions where WSP10s and SWHs are overestimated. The decreased SSTs in  
447 VR12-AL-SC-EN stabilize [the](#) marine atmospheric boundary layer, and lead to weakened  
448 WSP10s and SWHs. The modified roughness in Z0-M04 generally enhances momentum  
449 transfer into the ocean, and so decreases WSP10s and SWHs. The relative wind-wave-current  
450 speed in FLUX also affects wind stress, and further influences WSP10s and SWHs. Compared  
451 with NDBC buoy observations and ERA5, the ALL experiment shows significant  
452 improvements.

453 In addition to the variables aforementioned, the changes of simulated enthalpy fluxes were also  
454 compared, which mainly depend on the WSP10s changes. However, the wave-related effects on enthalpy  
455 fluxes are non-significant for the 2-month simulation, so the results are not shown.

456 The wave-related parameterizations used in the study mainly improve model biases at mid-high  
457 latitudes, and SST biases in tropical oceans are only slightly improved (Fig. 3&4). Breivik et al. (2015)  
458 improved SST as well as subsurface temperature simulations in Nucleus for European Modelling of the

459 Ocean (NEMO) with parametrizations including the wave-related Charnock parameter, modification of  
460 water-side stress with wind input and wave dissipation, wave dissipation-related turbulent kinetic energy  
461 flux and the Stokes-Coriolis force. Based on a global NEMO-WW3 coupled framework, Couvelard et al.  
462 (2020) modified the Charnock parameter, the Stokes drift-related forces and the Langmuir cell with  
463 misalignment of winds and waves, the oceanic surface momentum flux and the turbulence kinetic energy  
464 to reduce SST and MLD biases. In addition, sea sprays can enhance air-sea heat fluxes in the tropics  
465 (Andreas et al. 2008; Andreas et al. 2015). We will consider more processes in ~~the~~ future [studiesstudy](#).

466 Different parameterizations for the same wave-related process also deserve discussion. For ocean  
467 surface roughness, the most classic parametrizations are those developed by Janssen (1989, 1991), Taylor  
468 and Yelland (2001) and Drennan et al. (2003). The method of Taylor and Yelland (2001) requires the  
469 peak wavelength for the total spectrum, whereas that of Drennan et al. (2003) only requires the peak of  
470 wind-sea waves. This difference leads to the fact that the former is more suitable for a mixed sea state,  
471 while the latter is more suitable for a young sea state (Drennan et al., 2005). And the effect of Janssen's  
472 parameterization (1989, 1991) is similar to that of Drennan et al. (2003), since it is also based on the  
473 wind-sea conditions (Shimura et al., 2017).

474 The case studies indicate that there remain significant biases in the coupled system, probably owing  
475 to inaccuracy of coarse resolution, absence of a coupled wave-ice modular, and deficiency of initial fields.  
476 In addition, to further improve the model and eliminate the biases, as Breivik et al. (2015) proposed,  
477 extra adjusting of the individual model components in the coupled systems is also necessary. All of these  
478 require further efforts to investigate efficient methods to improve fully coupled systems.

479 **Code and data availability**

480 The code developed for the coupled system can be found under <https://doi.org/10.5281/zenodo.5811002>  
481 (Shi et al., 2021), including the coupling, preprocessing, run control and postprocessing scripts. The  
482 initial fields for CFSv2.0 are generated by the real time operational Climate Data Assimilation System,  
483 downloaded from the CFSv2.0 official website  
484 (<http://nomads.ncep.noaa.gov/pub/data/nccf/com/cfs/prod>). The daily average satellite Optimum  
485 Interpolation SST (OISST) data are obtained from NOAA (<https://www.ncdc.noaa.gov/oisst>), and the  
486 National Data Buoy Center (NDBC) buoy data are also obtained from NOAA  
487 (<https://www.ndbc.noaa.gov>). The Argo observational profiles of T/S are available at China Argo Real-  
488 time Data Center ([www.argo.org.cn](http://www.argo.org.cn)). The ERA5 reanalysis are available at the Copernicus Climate  
489 Change Service (C3S) Climate Data Store  
490 (<https://cds.climate.copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels>).

491 **Author contribution**

492 FX and RS designed the experiments and RS carried them out. RS developed the code of coupling  
493 parametrizations and produced the figures. ZF contributed to the installation and operation of CFSv2.0.  
494 LL and HY contributed to the application of C-Coupler2. XL and YZ provided the original code of  
495 CFSv2.0. RS prepared the manuscript with contributions from all co-authors. FX and HL contributed to  
496 review and editing.

497 **Acknowledgments**

498 The authors would like to extend thanks to all developers of the CFSv2.0 model  
499 (<https://cfs.ncep.noaa.gov/cfsv2/downloads.html>). The work was supported by National Key Research

500 and Development Program of China (no. 2016YFC1401408), and Tsinghua University Initiative  
501 Scientific Research Program (2019Z07L01001). [We thank Professor Wei Xue and Professor Yao Liu for](#)  
502 [helping us test the coupling frequency. We also thank three anonymous reviewers and the handling editor](#)  
503 [for their constructive comments.](#)~~We also thank two anonymous reviewers for their constructive comments.~~

## 504 **References**

505 Andreas, E. L., Mahrt, L., and Vickers, D.: An improved bulk air-sea surface flux algorithm, including  
506 spray-mediated transfer, Quarterly Journal of the Royal Meteorological Society, 141, 642-654, 2015.

507 Andreas, E. L., Persson, P. O. G., and Hare, J. E.: A bulk turbulent air-sea flux algorithm for high-wind,  
508 spray conditions, Journal of Physical Oceanography, 38, 1581-1596, 2008.

509 Ardhuin, F., Rogers, E., Babanin, A. V., Filipot, J., Magne, R., Roland, A., Der Westhuysen, A. V.,  
510 Queffelec, P., Lefevre, J. M., and Aouf, L.: Semiempirical Dissipation Source Functions for Ocean  
511 Waves. Part I: Definition, Calibration, and Validation, Journal of Physical Oceanography, 40, 1917-1941,  
512 <http://dx.doi.org/10.1175/2010JPO4324.1>, 2010.

513 Atlas, R., Hoffman, R. N., Ardizzone, J., Leidner, S. M., Jusem, J. C., Smith, D. K., and Gombos, D.: A  
514 Cross-calibrated, Multiplatform Ocean Surface Wind Velocity Product for Meteorological and  
515 Oceanographic Applications, Bulletin of the American Meteorological Society, 92, 157-174,  
516 <http://dx.doi.org/10.1175/2010BAMS2946.1>, 2011.

517 Bao, Y., Song, Z., & Qiao, F.: FIO-ESM version 2.0: Model description and evaluation. Journal of  
518 Geophysical Research: Oceans, 125, e2019JC016036. <https://doi.org/10.1029/2019JC016036>, 2019.

519 Belcher, S. E., Grant, A. L. M., Hanley, K., Foxkemper, B., Van Roekel, L., Sullivan, P. P., Large, W. G.,  
520 Brown, A. R., Hines, A., and Calvert, D.: A global perspective on Langmuir turbulence in the ocean

521 surface boundary layer, *Geophysical Research Letters*, 39, <http://dx.doi.org/10.1029/2012GL052932>,  
522 2012.

523 Beljaars, A. C. M.: The parametrization of surface fluxes in large-scale models under free convection. *Q.*  
524 *J. R. Meteorol. Soc.*, 121, 255-270. <https://doi.org/10.1002/qj.49712152203>, 1994.

525 Belmonte Rivas, M., and Stoffelen, A.: Characterizing ERA-Interim and ERA5 surface wind biases using  
526 ASCAT, *Ocean Science*, 15, 831-852, [10.5194/os-15-831-2019](https://doi.org/10.5194/os-15-831-2019), 2019.

527 Bidlot, J.-R., Prates, F., Ribas, R., et al.: Enhancing tropical cyclone wind forecasts, ECMWF newsletter,  
528 164, 33-37. [https://www.ecmwf.int/en/newsletter/164/meteorology/enhancing-tropical-cyclonewind-](https://www.ecmwf.int/en/newsletter/164/meteorology/enhancing-tropical-cyclonewind-forecasts)  
529 [forecasts](https://www.ecmwf.int/en/newsletter/164/meteorology/enhancing-tropical-cyclonewind-forecasts), 2020.

530 Bidlot, J.-R.: Model upgrade improves ocean wave forecasts. ECMWF newsletter, 159, 10-10.  
531 <https://www.ecmwf.int/en/newsletter/159/news/modelupgrade-improves-ocean-wave-forecasts>, 2019.

532 Bidlot, J.-R.: Present status of wave forecasting at ECMWF, Workshop on ocean waves, 25-27, 2012.

533 Breivik, Ø., Bidlot, J.-R., and Janssen, P. A.: A Stokes drift approximation based on the Phillips spectrum,  
534 *Ocean Modelling*, 100, 49-56, 2016.

535 Breivik, Ø., Janssen, P. A., and Bidlot, J.-R.: Approximate Stokes drift profiles in deep water, *Journal of*  
536 *Physical Oceanography*, 44, 2433-2445, 2014.

537 Breivik, Ø., Mogensen, K., Bidlot, J., Balmaseda, M., and Janssen, P. A. E. M.: Surface wave effects in  
538 the NEMO ocean model: Forced and coupled experiments, *Journal of Geophysical Research*, 120, 2973-  
539 2992, <http://dx.doi.org/10.1002/2014JC010565>, 2015.

540 Charnock, H.: Wind stress on a water surface, *Quarterly Journal of the Royal Meteorological Society*, 81,  
541 639-640, <http://dx.doi.org/10.1002/qj.49708135027>, 1955.

542 Couvelard, X., Lemarié, F., Samson, G., Redelsperger, J.-L., Ardhuin, F., Benshila, R., and Madec, G.:  
543 Development of a two-way-coupled ocean-wave model: assessment on a global NEMO(v3.6)-  
544 WW3(v6.02) coupled configuration, *Geosci. Model Dev.*, 13, 3067-3090, [https://doi.org/10.5194/gmd-](https://doi.org/10.5194/gmd-13-3067-2020)  
545 13-3067-2020, 2020.

546 de Boyer Montégut, C., Madec, G., Fischer, A. S., Lazar, A., and Iudicone, D.: Mixed layer depth over  
547 the global ocean: An examination of profile data and a profile - based climatology, *Journal of*  
548 *Geophysical Research*, 109, <http://dx.doi.org/10.1029/2004JC002378>, 2004.

549 Donelan, M. A., Haus, B. K., Reul, N., Plant, W. J., Stiassnie, M., Graber, H. C., Brown, O. B., and  
550 Saltzman, E. S.: On the limiting aerodynamic roughness of the ocean in very strong winds, *Geophysical*  
551 *Research Letters*, 31, [10.1029/2004gl019460](http://dx.doi.org/10.1029/2004gl019460), 2004.

552 Drennan, W. M., H. C. Graber, D. Hauser, and C. Quentin: On the wave age dependence of wind stress  
553 over pure wind seas, *J. Geophys. Res.*, 108(C3), 8062, [doi:10.1029/2000JC000715](https://doi.org/10.1029/2000JC000715), 2003.

554 Drennan, W. M., P. K. Taylor, and M. J. Yelland: Parameterizing the sea surface roughness, *J. Phys.*  
555 *Oceanogr.*, 35(5), 835-848, 2005.

556 ECMWF: Official IFS Documentation CY47R1. In chap. PART VII: ECMWF wave model. Reading,  
557 UK: ECMWF, [10.21957/31drbygag](https://www.ecmwf.int/en/forecasts/ifs-documentation/part-vii-wave-model), 2020.

558 Fan, Y., Lin, S., Held, I. M., Yu, Z., and Tolman, H. L.: Global Ocean Surface Wave Simulation Using a  
559 Coupled Atmosphere–Wave Model, *Journal of Climate*, 25, 6233-6252, [http://dx.doi.org/10.1175/JCLI-](http://dx.doi.org/10.1175/JCLI-D-11-00621.1)  
560 D-11-00621.1, 2012.

561 Ghantous, M., and Babanin, A. V.: One-dimensional modelling of upper ocean mixing by turbulence due  
562 to wave orbital motion, *Nonlinear Processes in Geophysics*, 21, 325-338, [10.5194/npg-21-325-2014](https://doi.org/10.5194/npg-21-325-2014),  
563 2014.

564 Griffies, S. M., Harrison, M. J., Pacanowski, R. C., and Rosati, A.: A technical guide to MOM4, GFDL  
565 Ocean Group Tech. Rep, 5, 342, 2004.

566 Hasselmann, K.: Wave-driven inertial oscillations, *Geophysical and Astrophysical Fluid Dynamics*, 1,  
567 463-502, <http://dx.doi.org/10.1080/03091927009365783>, 1970.

568 Hersbach H. and Bidlot, J.-R.: The relevance of ocean surface current in the ECMWF analysis and  
569 forecast system. Proceeding from the ECMWF Workshop on Atmosphere-Ocean Interaction, 10-12,  
570 [https://www.ecmwf.int/sites/default/files/elibrary/2009/9866-relevance-ocean-surface-current-ecmwf-](https://www.ecmwf.int/sites/default/files/elibrary/2009/9866-relevance-ocean-surface-current-ecmwf-analysis-and-forecast-system.pdf)  
571 [analysis-and-forecast-system.pdf](https://www.ecmwf.int/sites/default/files/elibrary/2009/9866-relevance-ocean-surface-current-ecmwf-analysis-and-forecast-system.pdf), 2008.

572 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., et al.: The ERA5 Global Reanalysis, *Quarterly Journal*  
573 *of the Royal Meteorological Society*, Soc. 00: 363 (2020), <https://doi.org/10.1002/qj.3803>, 2020.

574 Hsu, S. A., Eric A. Meindl, and David B. Gilhousen: Determining the Power-Law Wind-Profile Exponent  
575 under Near-Neutral Stability Conditions at Sea, *Applied Meteorology*, Vol. 33, No. 6,  
576 [https://doi.org/10.1175/1520-0450\(1994\)033<0757:DTPLWP>2.0.CO;2](https://doi.org/10.1175/1520-0450(1994)033<0757:DTPLWP>2.0.CO;2), 1994.

577 Janssen, P. A. E. M.: The interaction of ocean waves and wind, Cambridge University Press, pp 312,  
578 2004.

579 Janssen, P. A. E. M.: The Quasi-linear theory of wind wave generation applied to wave forecasting. *J.*  
580 *Phys. Oceanogr.* 21, 1631-1642. [https://doi.org/10.1175/1520-](https://doi.org/10.1175/1520-0485(1991)021<1631:QLTOWW>2.0.CO;2)  
581 [0485\(1991\)021<1631:QLTOWW>2.0.CO;2](https://doi.org/10.1175/1520-0485(1991)021<1631:QLTOWW>2.0.CO;2), 1991.

582 Janssen, P. A. E. M.: Wave-induced stress and the drag of air flow over sea waves. *J. Phys. Oceanogr.*  
583 19(6), 745-754. [https://doi.org/10.1175/1520-1310\\_0485\(1989\)019<0745:WISATD>2.0.CO;2](https://doi.org/10.1175/1520-1310_0485(1989)019<0745:WISATD>2.0.CO;2), 1989.

584 Kalnay, E., Kanamitsu, M., Kistler, R., Collins, W. D., Deaven, D. G., Gandin, L. S., Iredell, M. D., Saha,  
585 S., White, G. H., and Woollen, J.: The NCEP/NCAR 40-Year Reanalysis Project, *Bulletin of the*

586 American Meteorological Society, 77, 437-471, <http://dx.doi.org/10.1175/1520->  
587 0477(1996)077%3C0437:TNYRP%3E2.0.CO;2, 1996.

588 Kalverla, P. C., Holtslag, A. A. M., Ronda, R. J., and Steeneveld, G.-J.: Quality of wind characteristics  
589 in recent wind atlases over the North Sea, *Quarterly Journal of the Royal Meteorological Society*, 146,  
590 1498-1515, 10.1002/qj.3748, 2020.

591 Kukulka, T., Plueddemann, A. J., Trowbridge, J. H., and Sullivan, P. P.: Significance of Langmuir  
592 circulation in upper ocean mixing: comparison of observations and simulations, *Geophysical Research*  
593 *Letters*, 36, <http://dx.doi.org/10.1029/2009GL037620>, 2009.

594 Law-Chune, S., and Aouf, L.: Wave effects in global ocean modeling: parametrizations vs. forcing from  
595 a wave model, *Ocean Dynamics*, 68, 1739-1758, <http://dx.doi.org/10.1007/s10236-018-1220-2>, 2018.

596 Lemarié, F.: Numerical modification of atmospheric models to include the feedback of oceanic currents  
597 on air-sea fluxes in ocean-atmosphere coupled models, INRIA Grenoble-Rhône-Alpes, Laboratoire Jean  
598 Kuntzmann, <https://hal.inria.fr/hal-01184711/document>, 2015.

599 Li, D., Staneva, J., Bidlot, J.-R., Grayek, S., Zhu, Y., and Yin, B.: Improving Regional Model Skills  
600 During Typhoon Events: A Case Study for Super Typhoon Lingling Over the Northwest Pacific Ocean,  
601 *Frontiers in Marine Science*, 8, 10.3389/fmars.2021.613913, 2021.

602 Li, M., and Garrett, C.: Mixed Layer Deepening Due to Langmuir Circulation, *Journal of Physical*  
603 *Oceanography*, 27, 121-132, <http://dx.doi.org/10.1175/1520->  
604 0485(1997)027%3C0121:MLDDTL%3E2.0.CO;2, 1997.

605 Li, Q., Foxkemper, B., Breivik, O., and Webb, A.: Statistical models of global Langmuir mixing, *Ocean*  
606 *Modelling*, 113, 95-114, <http://dx.doi.org/10.1016/j.ocemod.2017.03.016>, 2017.

607 Li, Q., Webb, A., Foxkemper, B., Craig, A., Danabasoglu, G., Large, W. G., and Vertenstein, M.:  
608 Langmuir mixing effects on global climate: WAVEWATCH III in CESM, *Ocean Modelling*, 103, 145-  
609 160, <http://dx.doi.org/10.1016/j.ocemod.2015.07.020>, 2016.

610 Li, Z., Liu, Z., & Xing, X.: User Manual for Global Argo Observational data set (V3.0) (1997-2019),  
611 China Argo Real-time Data Center, Hangzhou, 37pp, 2019.

612 Liu, L., Zhang, C., Li, R., and Wang, B.: C-Coupler2: a flexible and user-friendly community coupler for  
613 model coupling and nesting, *Geoscientific Model Development Discussions*, 11, 1-63,  
614 <http://dx.doi.org/10.5194/gmd-11-3557-2018>, 2018.

615 Luo, J.J., Masson, S., Roeckner, E., Madec, G., & Yamagata, T.: Reducing climatology Bias in an ocean-  
616 atmosphere CGCM with improved coupling physics. *Journal of Climate*, 18(13), 2344-2360.  
617 <https://doi.org/10.1175/JCLI3404.1>, 2005.

618 McWilliams, J. C., and Sullivan, P. P.: Vertical Mixing by Langmuir Circulations, *Spill Science &*  
619 *Technology Bulletin*, 6, 225-237, [http://dx.doi.org/10.1016/S1353-2561\(01\)00041-X](http://dx.doi.org/10.1016/S1353-2561(01)00041-X), 2000.

620 [Monin, A. S., and Obukhov, A. M.: Basic laws of turbulent mixing in the surface layer of the atmosphere,](#)  
621 [Contrib. Geophys. Inst. Acad. Sci. USSR, 151\(163\), e187, 1954.](#)

622 Moon, I., Ginis, I., and Hara, T.: Effect of surface waves on Charnock coefficient under tropical cyclones,  
623 *Geophysical Research Letters*, 31, <http://dx.doi.org/10.1029/2004GL020988>, 2004.

624 Moum J.N., and Smyth W.D.: Upper Ocean Mixing. In Cochran, J. Kirk; Bokuniewicz, J. Henry; Yager,  
625 L. Patricia (Eds.) *Encyclopedia of Ocean Sciences*, 3rd Edition. vol. 1, pp. 71-79, Elsevier. ISBN: 978-  
626 0-12-813081-0, 2019.

627 O'Neill, L. W., Chelton, D. B., and Esbensen, S. K.: Observations of SST-induced perturbations of the  
628 wind stress field over the Southern Ocean on seasonal timescales, *Journal of Climate*, 16, 2340-2354,  
629 10.1175/2780.1, 2003.

630 Pineau-Guillou L., Arduin,F., Bouin,M.N.,Redelsperger,J.L.,Chapron,B., Bidlot, J. R., & Quilfen, Y.:  
631 Strong winds in a coupled wave-atmosphere model during a North Atlantic storm event: evaluation  
632 against observations. *Quarterly Journal of the Royal Meteorological Society*, 144(711), 317-332, 2018.

633 Polonichko, V.: Generation of Langmuir circulation for nonaligned wind stress and the Stokes drift, *J.*  
634 *Geophys. Res.*, 102, 15773–15780, <https://doi.org/10.1029/97JC00460>, 1997.

635 Powell, M. D., Vickery, P. J., and Reinhold, T. A.: Reduced drag coefficient for high wind speeds in  
636 tropical cyclones, *Nature*, 422, 279-283, 10.1038/nature01481, 2003.

637 Qiao, F., Yuan, Y., Yang, Y., Zheng, Q., Xia, C., and Ma, J.: Wave - induced mixing in the upper ocean:  
638 Distribution and application to a global ocean circulation model, *Geophysical Research Letters*, 31,  
639 <http://dx.doi.org/10.1029/2004GL019824>, 2004.

640 Renault, L., Arsouze, T., and Ballabrera-Poy, J.: On the Influence of the Current Feedback to the  
641 Atmosphere on the Western Mediterranean Sea Dynamics, *Journal of Geophysical Research-Oceans*, 126,  
642 10.1029/2020jc016664, 2021.

643 Renault, L., Molemaker, M. J., McWilliams, J. C., Shchepetkin, A. F., Lemarie, F., Chelton, D., Illig, S.,  
644 and Hall, A.: Modulation of Wind Work by Oceanic Current Interaction with the Atmosphere, *Journal of*  
645 *Physical Oceanography*, 46, 1685-1704, 10.1175/jpo-d-15-0232.1, 2016.

646 Reynolds, R. W., Smith, T. M., Liu, C., Chelton, D. B., Casey, K. S., and Schlax, M. G.: Daily High-  
647 Resolution-Blended Analyses for Sea Surface Temperature, *Journal of Climate*, 20, 5473-5496,  
648 <http://dx.doi.org/10.1175/2007JCLI1824.1>, 2007.

649 Saha, S., Moorthi, S., Pan, H., Wu, X., Wang, J., Nadiga, S., Tripp, P., Kistler, R., Woollen, J., and  
650 Behringer, D.: The NCEP Climate Forecast System Reanalysis, *Bulletin of the American Meteorological*  
651 *Society*, 91, 1015-1057, <http://dx.doi.org/10.1175/2010BAMS3001.1>, 2010.

652 Saha, S., Moorthi, S., Wu, X., Wang, J., Nadiga, S., Tripp, P., Behringer, D., Hou, Y., Chuang, H., and  
653 Iredell, M. D.: The NCEP Climate Forecast System Version 2, *Journal of Climate*, 27, 2185-2208,  
654 <http://dx.doi.org/10.1175/JCLI-D-12-00823.1>, 2014.

655 Sauvage, C., Lebeau-pin Brossier, C., Bouin, M. N., & Ducrocq, V.: Characterization of the air-sea  
656 exchange mechanisms during a Mediterranean heavy precipitation event using realistic sea state  
657 modelling. *Atmospheric Chemistry & Physics*, 20(3), 2020.

658 Sharmar, V., and Markina, M.: Validation of global wind wave hindcasts using ERA5, MERRA2, ERA-  
659 Interim and CFSRv2 reanalyses, *IOP Conference Series: Earth and Environmental Science*, 606, 012056  
660 (012059 pp.)-012056 (012059 pp.), 10.1088/1755-1315/606/1/012056, 2020.

661 Shimura, T., Mori, N., Takemi, T., et al.: Long term impacts of ocean wave-dependent roughness on  
662 global climate systems. *Journal of Geophysical Research: Oceans*, 122(3),  
663 <https://doi.org/10.1002/2016JC012621>, 2017.

664 Stopa, J. E., Ardhuin, F., Babanin, A. V., and Zieger, S.: Comparison and validation of physical wave  
665 parameterizations in spectral wave models, *Ocean Modelling*, 103, 2-17,  
666 <http://dx.doi.org/10.1016/j.ocemod.2015.09.003>, 2016.

667 Sweet, W., Fett, R., Kerling, J., and Laviolette, P.: Air-sea interaction effects in the lower troposphere  
668 across the north wall of the Gulf Stream, *Mon. Weather Rev.*, 109, 1042-1052, 10.1175/1520-  
669 0493(1981)109<1042:Asieit>2.0.Co;2, 1981.

670 Takatama, K., and Schneider, N.: The Role of Back Pressure in the Atmospheric Response to Surface  
671 Stress Induced by the Kuroshio, *Journal of the Atmospheric Sciences*, 74, 597-615, 10.1175/jas-d-16-  
672 0149.1, 2017.

673 Taylor, P. K., and Yelland, M. J.: The Dependence of Sea Surface Roughness on the Height and Steepness  
674 of the Waves, *Journal of Physical Oceanography*, 31, 572-590, [http://dx.doi.org/10.1175/1520-0485\(2001\)031%3C0572:TDOSSR%3E2.0.CO;2](http://dx.doi.org/10.1175/1520-0485(2001)031%3C0572:TDOSSR%3E2.0.CO;2), 2001.

676 Terray, E. A., Donelan, M. A., Agrawal, Y. C., Drennan, W. M., Kahma, K. K., Williams, A. J., Hwang,  
677 P. A., and Kitaigorodskii, S. A.: Estimates of Kinetic Energy Dissipation under Breaking Waves, *Journal*  
678 *of Physical Oceanography*, 26, 792-807, [http://dx.doi.org/10.1175/1520-0485\(1996\)026%3C0792:EOKEDU%3E2.0.CO;2](http://dx.doi.org/10.1175/1520-0485(1996)026%3C0792:EOKEDU%3E2.0.CO;2), 1996.

680 Tolman, H. L., and Chalikov, D. V.: Source Terms in a Third-Generation Wind Wave Model, *Journal of*  
681 *Physical Oceanography*, 26, 2497-2518, [http://dx.doi.org/10.1175/1520-0485\(1996\)026%3C2497:STIATG%3E2.0.CO;2](http://dx.doi.org/10.1175/1520-0485(1996)026%3C2497:STIATG%3E2.0.CO;2), 1996.

683 Van Roekel, L. P., Foxkemper, B., Sullivan, P. P., Hamlington, P. E., and Haney, S.: The form and  
684 orientation of Langmuir cells for misaligned winds and waves, *Journal of Geophysical Research*, 117,  
685 <http://dx.doi.org/10.1029/2011JC007516>, 2012.

686 Wang, L., and Xu, F.: Decadal variability and trends of oceanic barrier layers in tropical Pacific, *Ocean*  
687 *Dynamics*, 68, 1155-1168, 10.1007/s10236-018-1191-3, 2018.

688 WAVEWATCH III Development Group.: User manual and system documentation of WAVEWATCH III  
689 version 5.16, NOAA/NWS/NCEP/MMAB Technical Note 329, 326, 2016.

690 Wu, L., Staneva, J., Breivik, Ø., Rutgersson, A., Nurser, A. G., Clementi, E., and Madec, G.: Wave effects  
691 on coastal upwelling and water level, *Ocean Modelling*, 140, 101405, 2019.

692 Wu, X., K. S. Moorthi, K. Okamoto, and H. L. Pan: Sea ice impacts on GFS forecasts at high latitudes,  
693 Proceedings of the 85th AMS Annual Meeting, 8th Conference on Polar Meteorology and Oceanography,  
694 San Diego, CA, 2005.

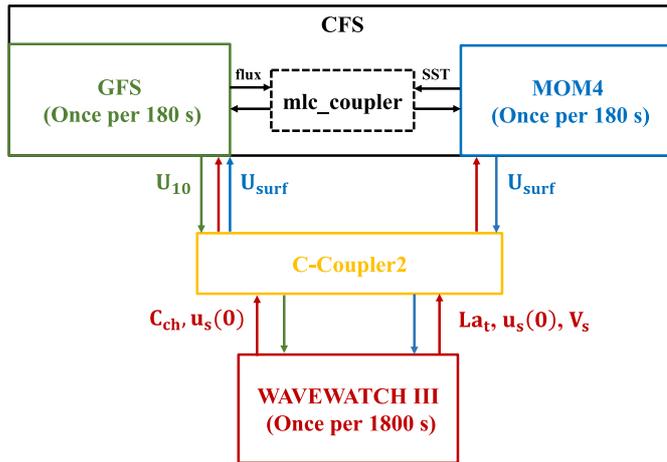
695 Zieger, S., Babanin, A. V., Rogers, W. E., and Young, I. R.: Observation-based source terms in the third-  
696 generation wave model WAVEWATCH, Ocean Modelling, 96, 2-25,  
697 <http://dx.doi.org/10.1016/j.ocemod.2015.07.014>, 2015.

**Table 1.** List of numerical experiments: setups different from CTRL are marked with bold

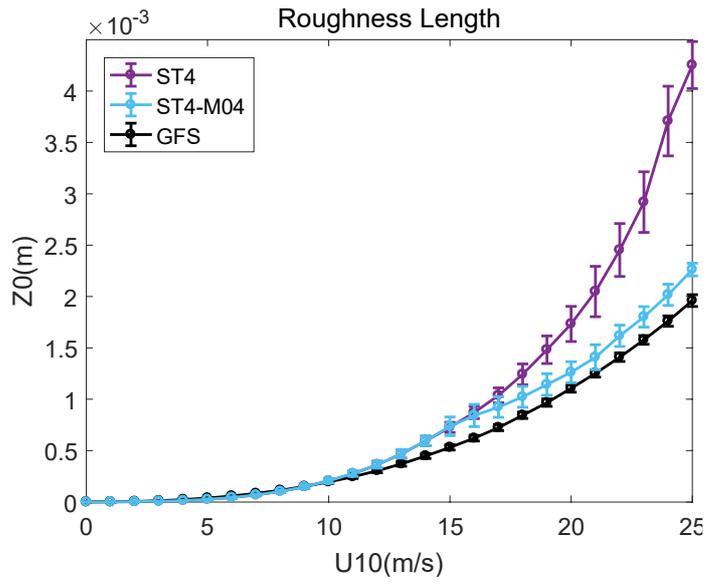
Experiments	Physical Process/Parameterization		
	Langmuir Cell with Stokes-Coriolis Force and Entrainment	Roughness (Charnock Parameter)	Relative Velocity in Flux
CTRL	Off	Off	Off
VR12-AL-SC-EN	<b>Eqn. 1-6, 8-10</b>	Off	Off
Z0-M04	Off	<b><math>C_{ch}</math> from Eqn. 16, 17</b>	Off
FLUX	Off	Off	<b><math>\Delta\vec{V}</math> from Eqn. 14</b>
ALL	<b>Eqn. 1-6, 8-10</b>	<b><math>C_{ch}</math> from Eqn. 16, 17</b>	<b><math>\Delta\vec{V}</math> from Eqn. 14</b>

**Table 2.** The 53-day mean bias with standard deviation (STD) and RMSE for WSP10 and SWH compared with NDBC buoy observation: the bias is calculated as simulation minus NDBC.

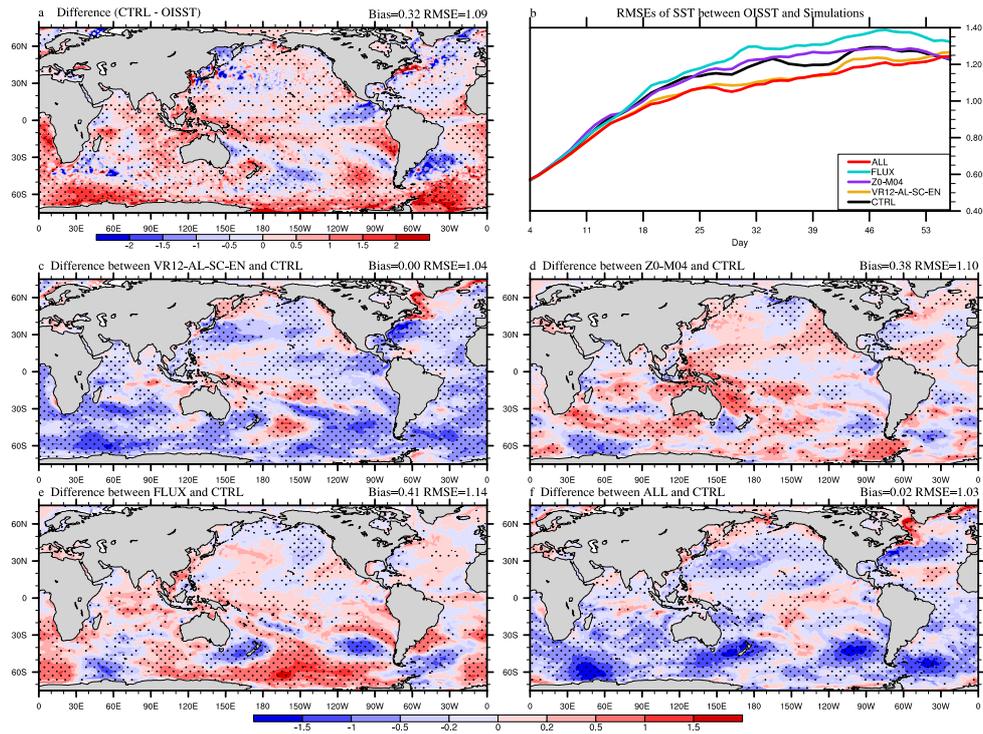
<b>Boreal Winter WSP10</b>	<b>Bias with STD</b>	<b>RMSE</b>
CTRL	0.16±1.23	1.24
VR12-AL-SC-EN	0.01±1.12	1.12
Z0-M04	-0.01±1.07	1.07
FLUX	0.39±1.20	1.26
ALL	0.07±1.04	1.04
<b>Boreal Winter SWH</b>	<b>Bias with STD</b>	<b>RMSE</b>
CTRL	0.21±0.38	0.44
VR12-AL-SC-EN	0.14±0.35	0.37
Z0-M04	0.10±0.30	0.32
FLUX	0.24±0.34	0.42
ALL	0.12±0.34	0.36
<b>Boreal Summer WSP10</b>	<b>Bias with STD</b>	<b>RMSE</b>
CTRL	0.15±1.23	1.24
VR12-AL-SC-EN	-0.03±1.22	1.22
Z0-M04	-0.04±1.21	1.21
FLUX	-0.22±1.18	1.20
ALL	-0.17±1.14	1.15
<b>Boreal Summer SWH</b>	<b>Bias with STD</b>	<b>RMSE</b>
CTRL	0.28±0.25	0.38
VR12-AL-SC-EN	0.19±0.24	0.30
Z0-M04	0.22±0.26	0.34
FLUX	0.14±0.25	0.29
ALL	0.12±0.21	0.24



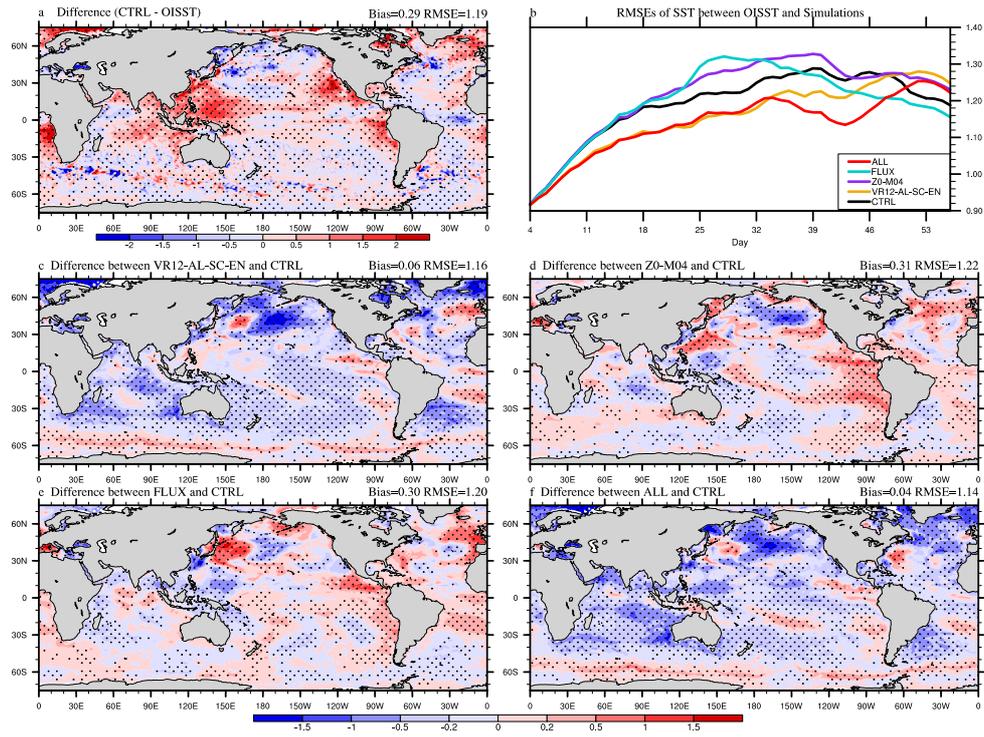
**Figure 1.** A schematic diagram of the atmosphere-ocean-wave coupled modeling system. The arrows indicate the coupled variables that are passed between the model components. In the diagram,  $C_{ch}$ ,  $La_t$ ,  $u_s(0)$ ,  $V_s$ ,  $U_{10}$ , and  $U_{surf}$  are Charnock parameter (red arrows), turbulent Langmuir number (red arrows), surface Stokes drift velocity (red arrows), Stokes transport (red arrows), 10-m wind (green arrows) and surface current (blue arrows), respectively.



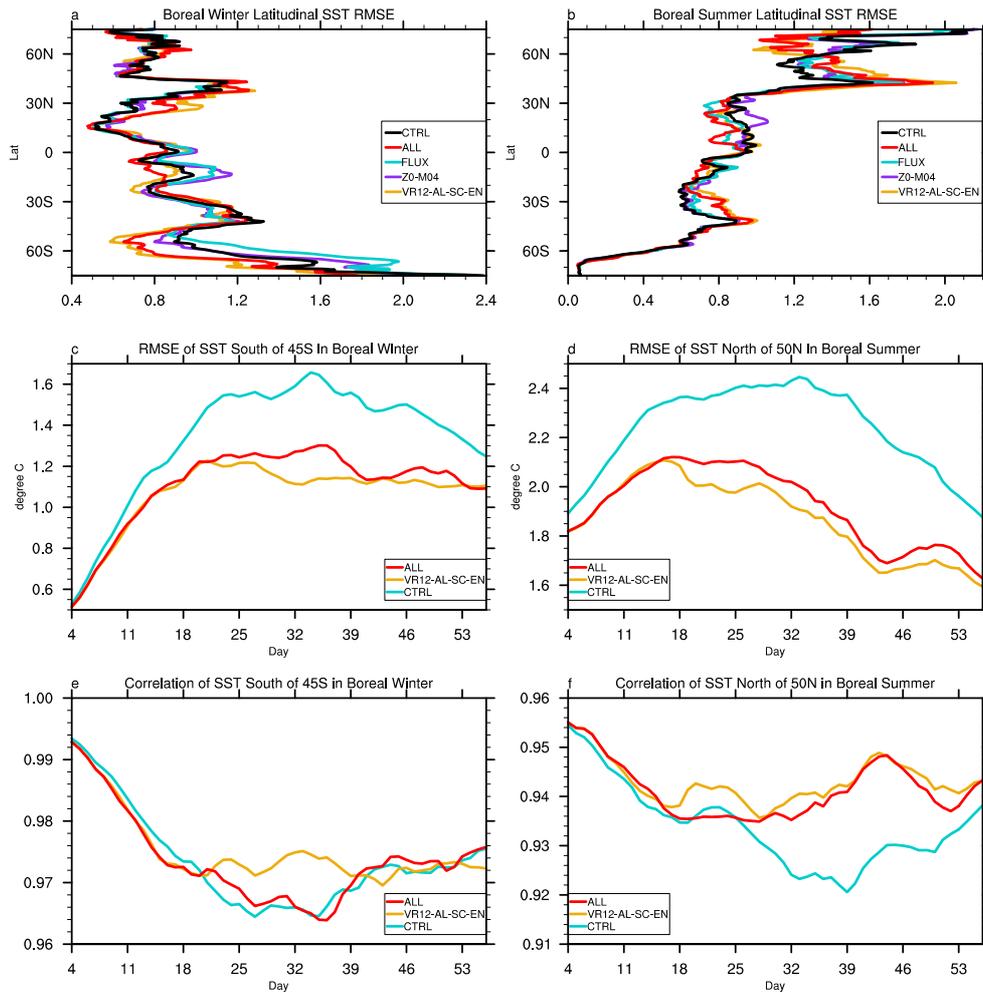
**Figure 2.** Relationships between momentum roughness length  $z_0$  (m) in the coupled system and 10-m wind speed (m/s); error bars indicate twice the standard deviations for each point.



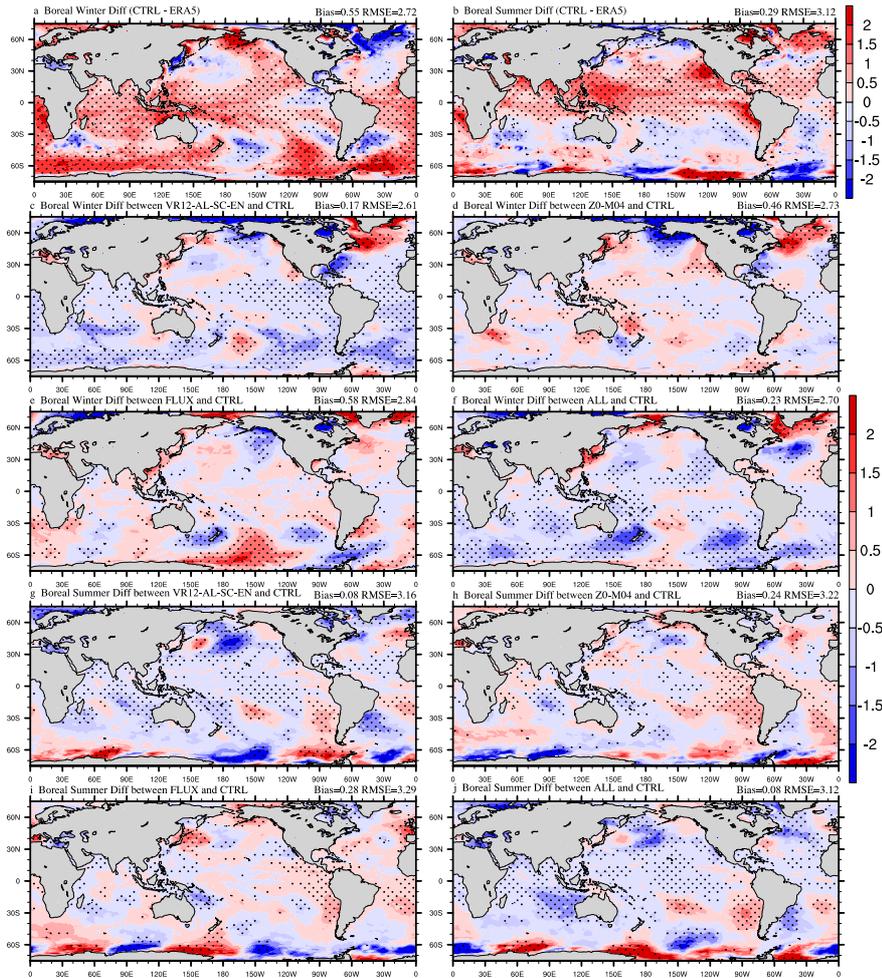
**Figure 3.** The 53-day average SST ( $^{\circ}\text{C}$ ) bias in CTRL (a; CTRL minus OISST), the time series of global-averaged RMSE (b), and the differences between VR12-AL-SC-EN (c)/Z0-M04 (d)/ FLUX (e)/ ALL (f) and CTRL in Jan-Feb, 2017 (VR12-AL-SC-EN/Z0-M04/FLUX/ALL minus CTRL). The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.



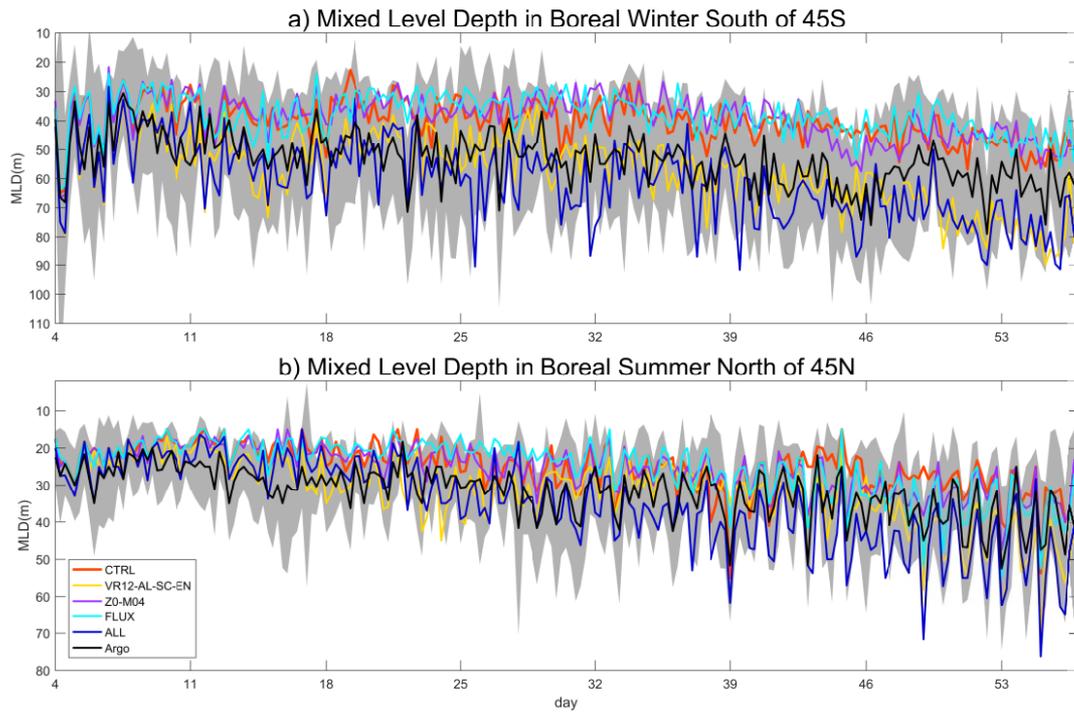
**Figure 4.** The same as Figure 3 but for Aug-Sep, 2018.



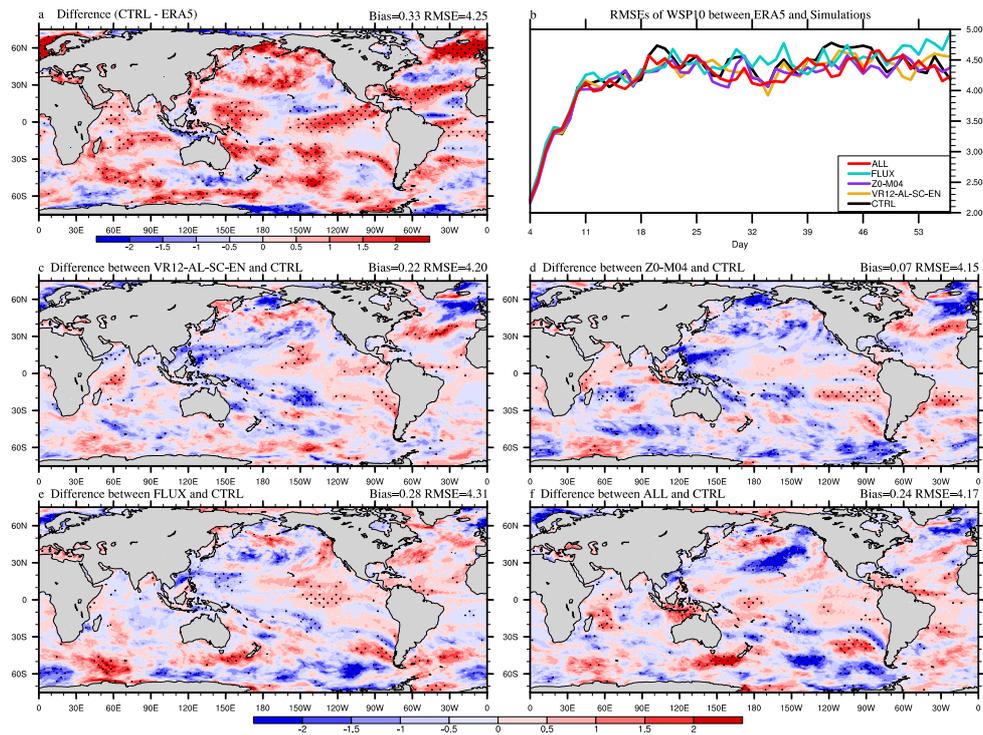
**Figure 5.** The 53-day averaged latitudinal distribution of SST root mean square errors (RMSE), time series of domain-averaged SST RMSE and correlation coefficient: **a/b** the latitudinal RMSE in boreal winter/summer compared with OISST, **c/d** the time series of domain-averaged (0-360°E, 45-78°S/50-78°N) SST RMSE in boreal winter/summer, **e/f** the time series of domain-averaged (0-360°E, 45-78°S/50-78°N) SST correlation coefficient in boreal winter/summer; differences of RMSE and correlation coefficient time series between VR12-AL-SC-EN/ALL and CTRL are statistically significant at 99% confidence level, except those in Fig. e.



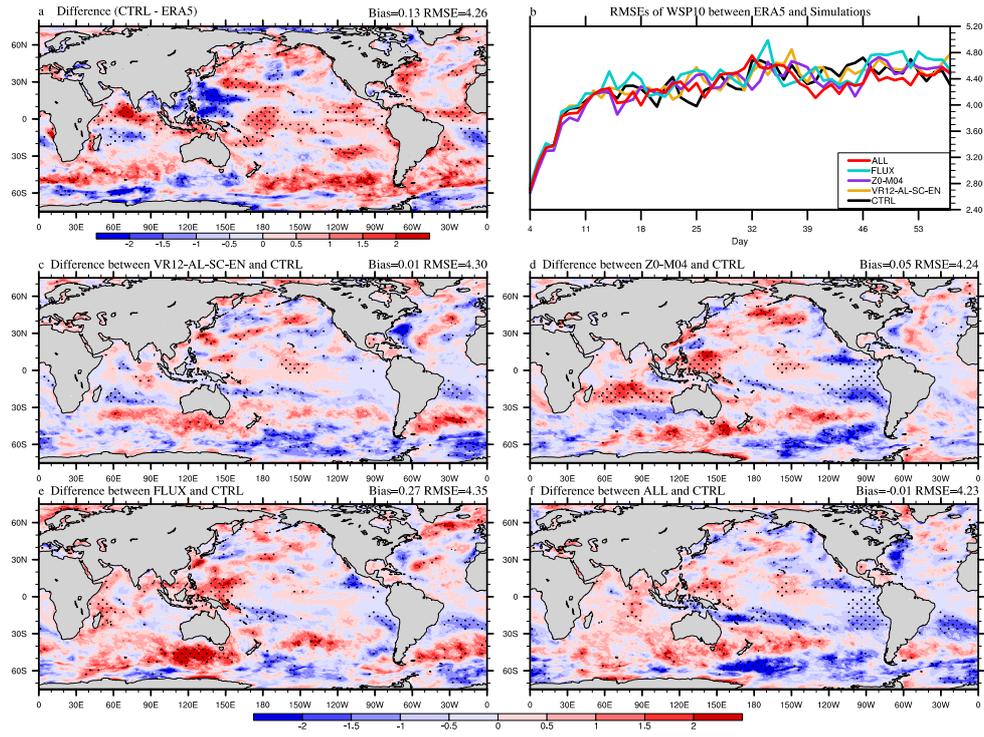
**Figure 6.** The 53-day average T02 (°C) bias in CTRL (a&b; CTRL minus ERA5), and the differences between VR12-AL-SC-EN (c&g)/Z0-M04 (d&h)/ FLUX (e&i)/ ALL (f&j) and CTRL (VR12-AL-SC-EN/Z0-M04/FLUX/ALL minus CTRL). The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level. a/c/d/e/f are for Jan-Feb, 2017, and b/g/h/i/j are for Aug-Sep, 2018.



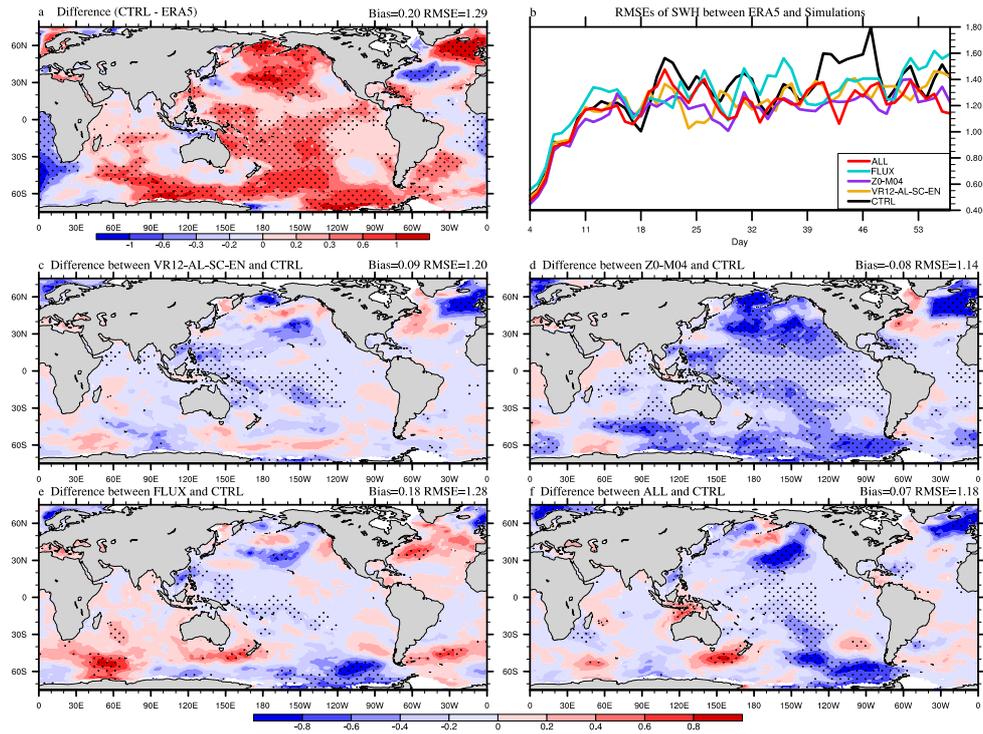
**Figure 7.** The 53-day time series of domain-averaged (0-360°E, 45-78°S/N) mixed layer depth (MLD; m) in boreal winter/summer: the difference between CTRL and VR12-AL-SC-EN/ALL passes the student's t-test at 99% confidence level; the time intervals are 6 hours; shaded areas indicate twice the standard deviations for Argo.



**Figure 8.** The 53-day average WSP10 (m/s) bias in CTRL (a; CTRL minus ERA5), the time series of global-averaged RMSE (b), and the differences between VR12-AL-SC-EN (c)/Z0-M04 (d)/ FLUX (e)/ ALL (f) and CTRL in Jan-Feb, 2017 (VR12-AL-SC-EN/Z0-M04/FLUX/ALL minus CTRL). The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.



**Figure 9.** The same as Figure 8 but for Aug-Sep, 2018.



**Figure 10.** The 53-day average SWH (m) bias in CTRL (a; CTRL minus ERA5), the time series of global-averaged RMSE (b), and the differences between VR12-AL-SC-EN (c)/Z0-M04 (d)/ FLUX (e)/ ALL (f) and CTRL in Jan-Feb, 2017 (VR12-AL-SC-EN/Z0-M04/FLUX/ALL minus CTRL). The first 3-day simulation is discarded. The dotted areas are statistically significant at 95% confidence level.

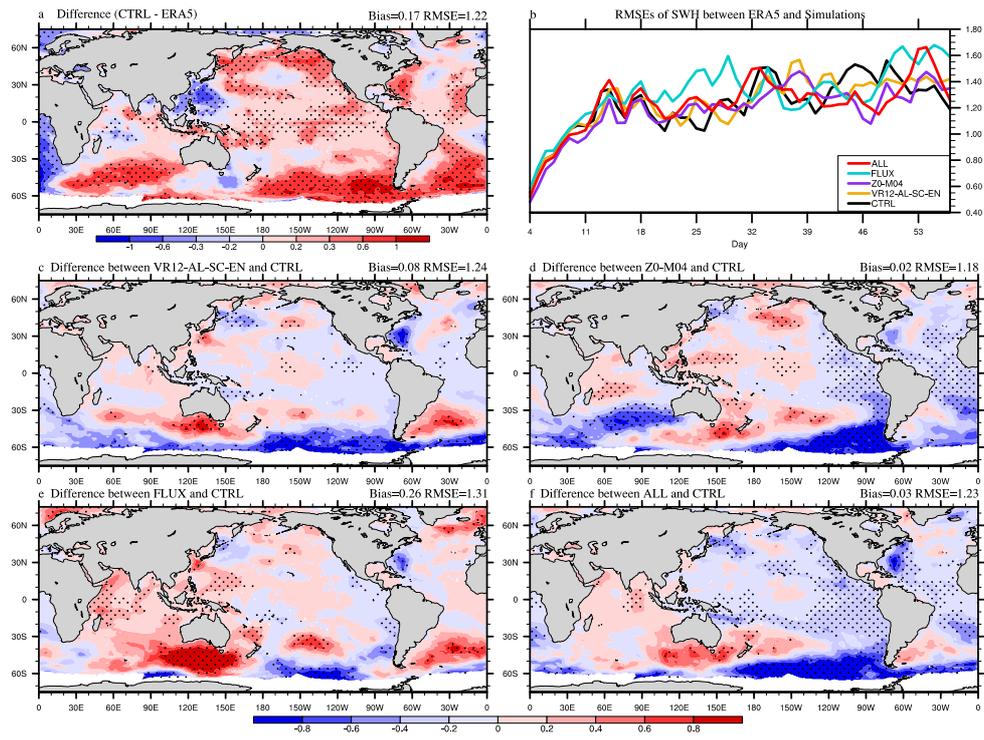
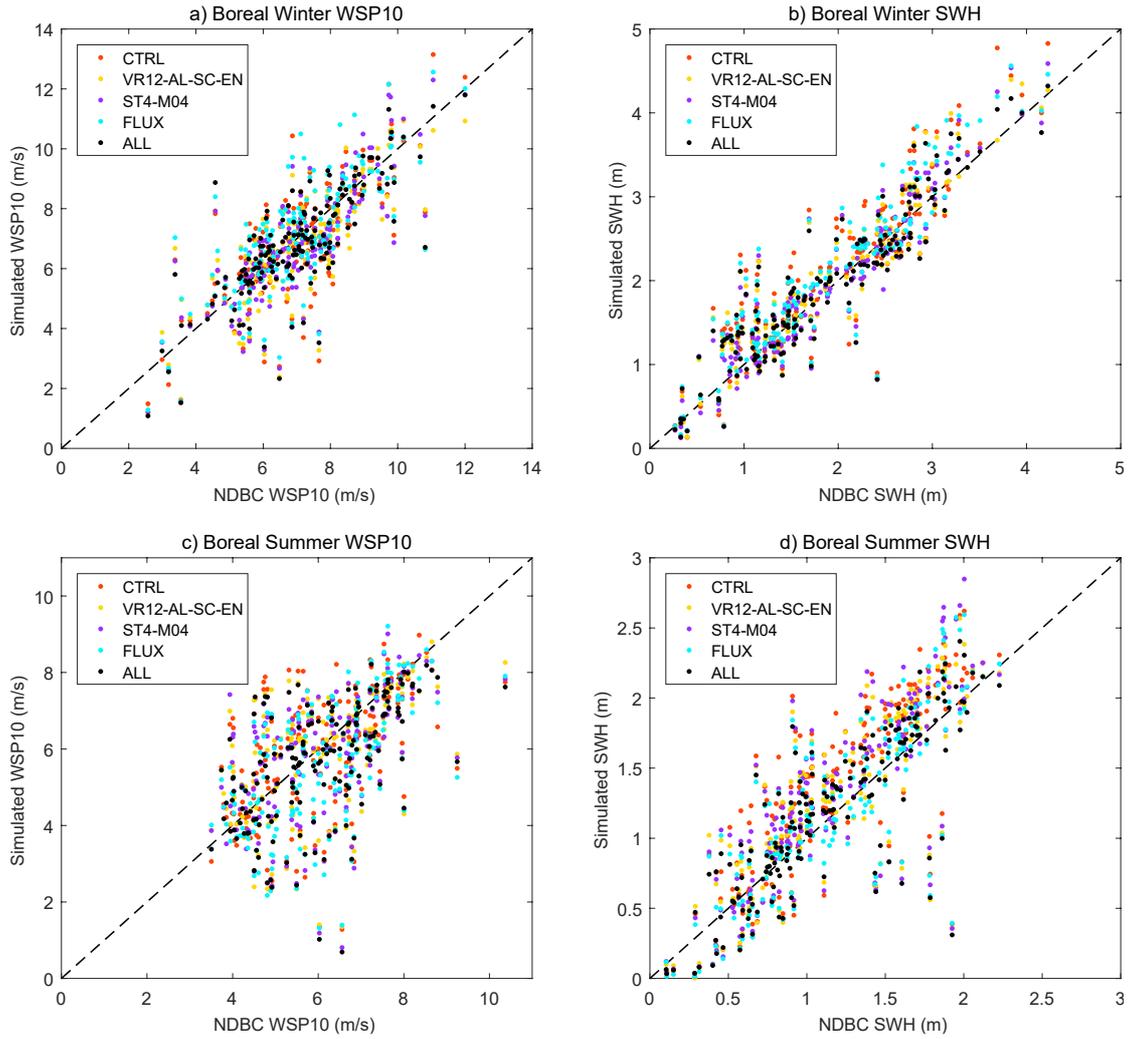


Figure 11. The same as Figure 10 but for Aug-Sep, 2018.



**Figure 12.** Scatter plots of simulated WSP10/SWH (y-axis) vs buoy WSP10/SWH (x-axis): (a) the WSP10 in Jan-Feb, 2017, (b) the SWH in Jan-Feb, 2017, (c) the WSP10 in Aug-Sep, 2018, and (d) the SWH in Aug-Sep, 2018. The dotted line is  $y=x$ . The corresponding mean biases with standard deviations and RMSEs for every experiment are shown in Table 2.