1 Non-Hydrostatic RegCM4 (RegCM4-NH): Model description and

2 case studies over multiple domains.

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11 Abstract. We describe the development of a non-hydrostatic version of the regional 12 climate model RegCM4, called RegCM4-NH, for use at convection-permitting resolutions. The non-hydrostatic dynamical core of the Mesoscale Model MM5 is introduced in the 13 14 RegCM4, with some modifications to increase stability and applicability of the model to 15 long-term climate simulations. Newly available explicit microphysics schemes are also 16 described, and three case studies of intense convection events are carried out in order to 17 illustrate the performance of the model. They are all run at convection-permitting grid 18 spacing of 3 km over domains in northern California, Texas and the Lake Victoria region, 19 without the use of parameterized cumulus convection. A substantial improvement is found 20 in several aspects of the simulations compared to corresponding coarser resolution (12) 21 km) runs completed with the hydrostatic version of the model employing parameterized 22 convection. RegCM4-NH is currently being used in different projects for regional climate 23 simulations at convection-permitting resolutions, and is intended to be a resource for 24 users of the RegCM modeling system.

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26 Keywords

27 Regional climate models; RegCM4; km-scale resolution; climate change

28 Introduction

Since the pioneering work of Dickinson et al. (1989) and Giorgi and Bates (1989), 29 30 documenting the first regional climate modeling system (RegCM, version 1) in literature, 31 the dynamical downscaling technique based on limited area Regional Climate Models 32 (RCMs) has been widely used worldwide, and a number of RCM systems have been 33 developed (Giorgi 2019). RegCM1 (Dickinson et al., 1989, Giorgi and Bates, 1989) was 34 originally developed at the National Center for Atmospheric Research (NCAR) based on 35 the Mesoscale Model version 4 (MM4) (Anthes et al, 1987). Then, further model versions 36 followed: RegCM2 (Giorgi et al. 1993a,b), RegCM2.5, (Giorgi and Mearns 1999), 37 RegCM3 (Pal et al. 2007), and lastly RegCM4 (Giorgi et al 2012). Except for the transition 38 from RegCM1 to RegCM2, in which the model dynamical core was updated from that of 39 the MM4 to that of the MM5 (Grell et al. 1994), these model evolutions were mostly based 40 on additions of new and more advanced physics packages. RegCM4 is today used by a 41 large community for numerous projects and applications, from process studies to paleo 42 and future climate projections, including participation in the Coordinated Regional 43 Downscaling EXperiment (CORDEX, Giorgi et al. 2009; Gutowski et al. 2016). The model 44 can also be coupled with ocean, land and chemistry/aerosol modules in a fully interactive 45 way (Sitz et al. 2017).

46 The dynamical core of the standard version of RegCM4 is hydrostatic, with sigma-p 47 vertical coordinates. As a result, the model can be effectively run for grid spacings of ~10 48 km or larger, for which the hydrostatic assumption is valid. However, the RCM community 49 is rapidly moving to higher resolutions of a few km, i.e. "convection-permitting" (Prein et 50 al. 2015; Coppola et al. 2020) and therefore the dynamical core of RegCM4 has been 51 upgraded to include a non-hydrostatic dynamics representation usable for very high 52 resolution applications. This upgrade, which we name RegCM4-NH, is essentially based 53 on the implementation of the MM5 non-hydrostatic dynamical core within the RegCM4 54 framework, which has an entirely different set of sub-grid model physics compared to MM5. 55

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57 RegCM4-NH is already being used in some international projects focusing on climate 58 simulations at convection-permitting km-scales, namely the European Climate Prediction 59 System (EUCP, Hewitt and Lowe 2018) and the CORDEX Flagship Pilot Study dedicated 60 to convection (CORDEX-FPSCONV, Coppola et al. 2020), and it is starting to be used 61 more broadly by the RegCM modeling community.

62 For example, the recent papers by Ban et al. (2021) and Pichelli et al. (2021) document 63 results of the first multi-model experiment of 10-year simulations at the convection-64 permitting scales over the so-called greater Alpine region. Two different simulations with RegCM4-NH for present day conditions have contributed to the evaluation analysis of 65 Ban et al. (2021). They were carried out at the International Centre for Theoretical Physics 66 (ICTP) and the Croatian Meteorological and Hydrological Service (DHMZ) using two 67 68 different physics configurations. The results show that RegCM4-NH largely improves the 69 precipitation simulation as compared to available fine scale observations when going from 70 coarse to high resolution, in particular for higher order statistics, such as precipitation 71 extremes and hourly intensity. Pichelli et al. (2021) then analyse multi-model ensemble 72 simulations driven by selected CMIP5 GCM projections for the decades 1996–2005 and 73 2090-2099 under the RCP8.5 scenario. ICTP contributed to the experiment with 74 simulations using RegCM4-NH driven by the MOCH-HadGEM GCM (r1i1p1) in a two level nest configuration (respectively at 12 and 3 km grid). The paper shows new insights 75 76 into future changes, for example an enhancement of summer and autumn hourly rainfall 77 intensification compared to coarser resolution model experiments, as well as an increase 78 of frequency and intensity of high-impact weather events.

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In this paper we describe the structure of RegCM4-NH and provide some illustrative examples of its performance, so that model users can have a basic reference providing them with background information on the model. In the next section we first describe the new model dynamical core, while the illustrative applications are presented in section 4. Section 5 finally provides some discussion of future developments planned for the RegCM system.

87 Model description

In the development of RegCM4-NH, the RegCM4 as described by Giorgi et al. (2012) was modified to include, the non-hydrostatic dynamical core (*idynamic* = 2 namelist option as described in RegCM-4.7.1/Doc/README.namelist of the source code) of the mesoscale model MM5 (Grell et al. 1994). This dynamical core was selected because RegCM4 already has the same grid and variable structure as MM5 in its hydrostatic core, which substantially facilitated its implementation (Elguindi et al. 2017).

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95 The model equations with complete description of the Coriolis force and a top radiative boundary condition, along with the finite differencing scheme, are given in Grell et al. 96 (1994). Pressure, p, temperature, T, and density, ρ , are first decomposed into a 97 98 prescribed reference vertical profile plus a time varying perturbation. The prognostic 99 equations are then calculated using the pressure perturbation values. Compared to the 100 original MM5 dynamical core, the following modifications were implemented in order to 101 achieve increased stability for long term climate simulations (Elguindi et al. 2017 102 document any modifications which follow the choice of the non-hydrostatic dynamical 103 core through the namelist parameter idynamic = 2; further available user-dependent 104 options, and the corresponding section in the namelist, are explicitly indicated):

105

i) The reference state temperature profile is computed using a latitude dependent
climatological temperature distribution and thus is a function of the specific domain
coordinates (*base_state_pressure*, *logp_lrate* parameters in *&referenceatm*) (Elguindi et
al. 2017). These two parameters were hard-coded in the original MM5 while for the
RegCM are user configurable;

111

ii) The lateral time dependent boundary conditions (*iboudy* in *&physicsparam*) for each
prognostic variable use the same exponential relaxation technique (*iboudy* = 5) described
in Giorgi et al. (1993). The linear MM5 relaxation scheme is also kept as an option (*iboudy*= 1);

iii) The advection term in the model equations, which in the MM5 code is implemented
using a centered finite difference approach, was changed to include a greater upstream
weight factor as a function of the local Courant number (Elguindi et al. 2017). The
maximum value of the weight factor is user configurable (*uoffc* in &*dynparam*). As detailed
in the MM5 model description (Grell et al., 1994), the horizontal advection term for a scalar
variable X contributes to the total tendency as:

123

$$\Delta_{adv} (p^*X)_G = -m^2 |_G \left[\frac{\left(p^*X|_b \frac{u}{m}|_b - p^*X|_a \frac{u}{m}|_a \right)}{dx} + \frac{\left(p^*X|_d \frac{v}{m}|_d - p^*X|_c \frac{v}{m}|_c \right)}{dy} \right]$$

124 125

where the *m* is the projection mapping factor and, with respect to Figure 1, assuming that the computation is to be performed for the gold cross point *G*, the averages are performed in the points *a*, *b*, *c*, *d*. For the u/m and v/m terms, the average value is computed using respectively the values in points AC, BD, CD, AB.

130 In RegCM4 for the term p^*X , the model computes a weighted average value of the field 131 using the value in gold+cyan and gold+green cross points with weights increasing the 132 relative contribution of the upstream point up as a function of the local courant number: 133

134
$$p^*X|_a = 0.5((1 - f_1)p^*X|_G + (1 + f_1)p^*X|_{c_1})$$

135 $p^*X|_b = 0.5((1 - f_1)p^*X|_{c_2} + (1 + f_1)p^*X|_G)$
136 $p^*X|_c = 0.5((1 - f_2)p^*X|_G + (1 + f_2)p^*X|_{g_1})$

- 137 $p^*X|_d = 0.5((1-f_2)p^*X|_{g_2} + (1+f_2)p^*X|_G)$
- where f_1, f_2 are defined as the local Courant number for the 1D advection equations multiplied for a control factor:
- 140

141
$$f_{1} = \mu_{fc} dt \frac{(u|_{a} + u|_{b})}{2dx}$$
$$f_{2} = \mu_{fc} dt \frac{(v|_{c} + v|_{d})}{2dy};$$

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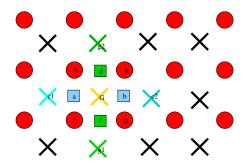


Figure 1 Schematic representation showing the horizontal advection scheme
staggering. Circles are U,V points. X are scalar variable points.

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iv) The water species (cloud, ice,rain, snow) term uses the same advection scheme as
the other variables (Elguindi et al. 2017) and not a complete upstream scheme as in the
MM5 code (Grell et al. 1994);

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v) A local flux limiter reduces the advection terms in order to remove unrealistic strong
gradients and its limits are user configurable (in the *&dynparam* section the maximum
gradient fraction for advection: temperature, *t_extrema*, specific humidity, *q_rel_extrema*,
liquid cloud content, *c_rel_extrema* and for tracers, *t_rel_extrema*). This was hardcoded
in the MM5 code and the limits were not user configurable;

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vi) The diffusion stencil of the Laplace equation uses a nine point approach as in LeVeque
(2006) and a topography dependent environmental diffusion coefficient is added to
reduce spurious diffusion along pressure coordinate slopes (Elguindi et al. 2017) as in
the hydrostatic version of the code (Giorgi et al. 1993b). The change in stencil does not
affect the overall fourth order precision of the model, but reduces the computational
stencil size, thus reducing the communication overhead;

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vii) The top boundary radiative condition (*ifupr* = 1 in *&nonhydroparam*) adopted in the semi-implicit vertical differencing scheme to reduce the reflection of energy waves uses coefficients on a 13x13 matrix which are re-computed every simulation day and not kept constant throughout the whole simulation as in the MM5 code. This allows the model to be run for longer simulation times while not being strongly tied to the initial atmosphericconditions;

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174 viii) The dynamical control parameter β in the semi-implicit vertical differencing scheme 175 (*nhbet* in *&nonhydroparam*) used for acoustic wave damping (Elguindi et al. 2017) is user 176 configurable (Klemp and Dudhia, 2008), while it is hard-coded in the MM5;

177

ix) A Rayleigh damping (*ifrayd* = 1 in *&nonhydroparam*) of the status variables towards
the input GCM boundary conditions can be activated in the top layers (*rayndamp*configuring the number of top levels to apply) with a configurable relaxation time
(*rayalpha0*, Klemp and Lilly, 1978, Durran and Klemp, 1983. This is consistent to what is
implemented in the WRF model);

183

184 x) The water species time filtering uses the Williams (2009) modified filter with $\alpha = 0.53$ 185 instead of the RA filter used by all the other variables. The v factor in the RA filter is user 186 configurable (*gnu1* and *gnu2* in *&dynparam*). This reduces the damping introduced by the 187 Robert-Asselin filter and the computational diffusion introduced by the horizontal 188 advection scheme.

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With these modifications, the model basic equations, under leap-frog integration scheme,are (Elguindi et al. 2017) :

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$$\frac{\partial p^* u}{\partial t} = -m^2 \left[\frac{\partial p^* u u/m}{\partial x} + \frac{\partial p^* v u/m}{\partial y} \right] - \frac{\partial p^* u \dot{\sigma}}{\partial \sigma} + u DIV - \frac{m p^*}{\rho} \left[\frac{\partial p'}{\partial x} - \frac{\sigma}{p^*} \frac{\partial p^*}{\partial x} \frac{\partial p'}{\partial \sigma} \right] + p^* f v - p^* e w \cos \theta + D_u$$
194 (1)

195

$$\frac{\partial p^* v}{\partial t} = -m^2 \left[\frac{\partial p^* u v/m}{\partial x} + \frac{\partial p^* v v/m}{\partial y} \right] - \frac{\partial p^* v \dot{\sigma}}{\partial \sigma} + v DIV - \frac{m p^*}{\rho} \left[\frac{\partial p'}{\partial y} - \frac{\sigma}{p^*} \frac{\partial p^*}{\partial y} \frac{\partial p'}{\partial \sigma} \right] - p^* f u + p^* e w \sin \theta + D_v$$
(2)

$$\frac{\partial p^* w}{\partial t} = -m^2 \left[\frac{\partial p^* u w/m}{\partial x} + \frac{\partial p^* v w/m}{\partial y} \right] - \frac{\partial p^* w \dot{\sigma}}{\partial \sigma} + w DIV + p^* g \frac{\rho_0}{\rho} \left[\frac{1}{p^*} \frac{\partial p'}{\partial \sigma} + \frac{T'_v}{T} - \frac{T_0 p'}{T p_0} \right] - p^* g \left[(q_c + q_r) \right] + p^* e \left(u \cos \theta - v \sin \theta \right) + D_w \quad (3)$$
100

$$\frac{\partial p^* p'}{\partial t} = -m^2 \left[\frac{\partial p^* u p'/m}{\partial x} + \frac{\partial p^* v p'/m}{\partial y} \right] - \frac{\partial p^* p' \dot{\sigma}}{\partial \sigma} + p' DIV - m^2 p^* \gamma p \left[\frac{\partial u/m}{\partial x} - \frac{\sigma}{mp^*} \frac{\partial p^*}{\partial x} \frac{\partial u}{\partial \sigma} + \frac{\partial v/m}{\partial y} - \frac{\sigma}{mp^*} \frac{\partial p^*}{\partial y} \frac{\partial v}{\partial \sigma} \right] + \rho_0 g \gamma p \frac{\partial w}{\partial \sigma} + p^* \rho_0 g \quad (4)$$
201

$$\frac{\partial p^*T}{\partial t} = -m^2 \left[\frac{\partial p^* u T/m}{\partial x} + \frac{\partial p^* v T/m}{\partial y} \right] - \frac{\partial p^* T \dot{\sigma}}{\partial \sigma} + T D I V + \frac{1}{\rho c_p} \left[p^* \frac{D p'}{D t} - \rho_0 g p^* w - D_{p'} \right] + p^* \frac{\dot{Q}}{c_p} + D_T$$
(5)

Where:

$$DIV = m^2 \left[\frac{\partial p^* u/m}{\partial x} + \frac{\partial p^* v/m}{\partial y} \right] + \frac{\partial p^* \dot{\sigma}}{\partial \sigma}$$

$$\dot{\sigma} = -\frac{\rho_0 g}{p^*} w - \frac{m\sigma}{p^*} \frac{\partial p^*}{\partial x} u - \frac{m\sigma}{p^*} \frac{\partial p^*}{\partial y} v$$
$$\tan \theta = -\cos \phi \frac{\partial \lambda / \partial y}{\partial \phi / \partial x}$$

$$p(x, y, z, t) = p_0(z) + p'(x, y, z, t)$$

$$T(x, y, z, t) = T_0(z) + T'(x, y, z, t)$$

$$\rho(x, y, z, t) = \rho_0(z) + \rho'(x, y, z, t)$$

with the vertical sigma coordinate defined as:

$$\sigma = \frac{(p_0 - p_t)}{(p_s - p_t)}$$

 P_s is the surface pressure and P_0 is the reference pressure profile. The total pressure

at each grid point is thus given as:

217
$$p(x, y, z, t) = p^* \sigma(k) + p_t + p'(x, y, z, t)$$

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219 With P_t being the top model pressure assuming a fixed rigid lid.

220 The model physics schemes for boundary layer, radiative transfer, land and ocean 221 surface processes, cloud and precipitation processes are extensively described in Giorgi 222 et al. (2012) and summarized in Table 1. For each physics component a number of 223 parameterization options are available (Table 1), and can be selected using a switch 224 selected by the user. As mentioned, the use of non-hydrostatic dynamics is especially 225 important when going to convection-permitting resolutions of a few km (Prein et al. 2015). 226 At these resolutions the scale separation assumption underlying the use of cumulus 227 convection schemes is not valid any more, and explicit cloud microphysics 228 representations are necessary. The RegCM4 currently includes two newly implemented 229 microphysics schemes, the Nogherotto-Tompkins (Nogherotto et al. 2016) and the WSM5 230 scheme from the Weather Research Forecast (WRF, Skamarok et al. 2008) model, which 231 are briefly described in the next sections for information to model users.

Model physics (Namelist flag)	Options	n. option	Reference
Dynamical core (<i>idynamic</i>)	Hydrostatic	1	Giorgi et al. 1993a,b Giorgi et al. 2012
	Non-Hydrostatic (*)	2	present paper
Radiation	CCSM	0	Kiehl et al. 1996
(irrtm)	RRTM (*)	1	Mlawer et al. 1997
Microphysics	Subex	1	Pal et al 2000

(ipptls)	Nogherotto Thompkins	2	Nogherotto et al. 2016	
	WSM5 (*)	3	Hong et al 2004	
Cumulus	Кио	1	Anthes et al. 1987	
(icup)	Grell	2	Grell 1993	
	Emanuel	4	Emanuel 1991	
	Tiedtke	5	Tiedtke 1989, 1993	
	Kain-Fritsch	6	Kain and Fritsch, 1990; Kain 2004	
	MM5 Shallow cumulus (only mixing) (*)	-1	Grell et al. 1994	
Planetary Boundary Layer	Modified-Holtslag	1	Holtslag et al., 1990	
(ibltyp)	UW	2	Bretherton et al. 2004	
Land Surface (code compiling	BATS	/	Dickinson et al. 1993; Giorgi et al. 2003	
option)	CLM4.5	/	Oleson et al. 2013	
Ocean Fluxes	BATS	1	Dickinson et al. 1993	
(iocnflx)	Zeng	2	Zeng et al. 1998	
	COARE	3	Fairall et al. 1996a,b	

Interactive lake (lakemod)	1D diffusion/convection	1	Hostetler et al. 1993
Tropical band (i_band)	RegT-Band	1	Coppola et al. 2012
Coupled ocean (iocncpl)	RegCM-ES	1	Sitz et al. 2017

- 233
 Table 1 Core and sub-grid physics scheme available in RegCM-NH. New schemes

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 available with this release are starred (*).
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237 Explicit microphysics schemes

238 Nogherotto-Tompkins Scheme

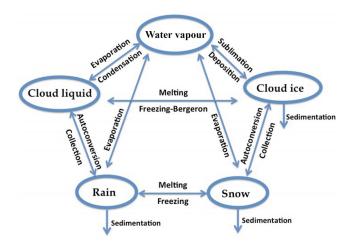
239 A new parameterization for explicit cloud microphysics and precipitation built upon the 240 European Centre for Medium Weather Forecast's Integrated Forecast System (IFS) module (Tiedtke [1993], Tompkins [2007]), was introduced in RegCM4 (ipptls = 2 in 241 242 µparam) by Nogherotto et al. [2016]. In the present configuration, the scheme 243 implicitly solves 5 prognostic equations for water vapor, qv, cloud liquid water, ql, rain, qr, 244 cloud ice, gi, and snow, gs, but it is also easily extendable to a larger number of variables. 245 Water vapor, cloud liquid water, rain, cloud ice and snow are all expressed in terms of the 246 grid-mean mixing ratio.

247 Cloud liquid and ice water content are independent, allowing the existence of supercooled 248 liquid water and mixed-phase clouds. Rain and snow precipitate with a fixed terminal fall 249 speed and can then be advected by the three dimensional winds. A check for the conservation of enthalpy and of total moisture is ensured at the end of each timestep. The 250 251 governing equation for variable each is: 252

$$\frac{\partial q_x}{\partial t} = S_x + \frac{1}{\rho} \frac{\partial}{\partial z} (\rho V_x q_x)$$

255 The local variation of the mixing ratio qx of the variable x is given by the sum of 256 Sx, containing the net sources and sinks of qx through microphysical processes (i.e. 257 condensation, evaporation, auto-conversion, melting, etc.), and the sedimentation term, 258 which is a function of the fall speed Vx. An upstream approach is employed to solve the equations. The sources and sinks contributors are divided in two groups according to the 259 260 duration of the process they describe: processes that are considered to be fast relative to 261 the model time step are treated implicitly while slow processes are treated explicitly. The 262 processes taken into account (shown in Figure 2) are the microphysical pathways across 263 the 5 water variables: condensation, autoconversion, evaporation, cloud water collection 264 (accretion), and autoconversion for warm clouds, and freezing, melting, deposition, 265 sublimation for cold clouds.

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267

Figure 2: Depiction of the new scheme, showing the five prognostic variables and how they are related to each other through microphysical processes

For each microphysical pathway, phase changes are associated with the release or absorption of latent heat, which then impacts the temperature budget. The impact is 272 calculated using the conservation of liquid water temperature TL defined as:

273

274
$$T_L = T - \frac{L_v}{C_p}(q_l + q_r) - \frac{L_s}{C_p}(q_i + q_s).$$

Given that dTL =0, the rate of change of the temperature is given by the following equation:

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$$\frac{\partial T}{\partial t} = \sum_{x=1}^{m} \frac{L(x)}{C_{p}} \left(\frac{dq_{x}}{dt} - D_{q_{x}} - \frac{1}{\rho} \frac{\partial}{\partial z} (\rho V_{x} q_{x}) \right)$$

278 279

where L(x) is the latent heat of fusion or evaporation, depending on the process considered, Dqx is the convective detrainment and the third term in brackets is the sedimentation term.

At the end of each time step a check is carried out of the conservation of total water and

284 moist static energy: $h = C_P T + g_Z + Lq_x$.

The scheme is tunable through parameters in the *µparam* section of the namelist (RegCM-4.7.1/Doc/README.namelist; Elguindi et al. 2017).

287

288 WSM5 Scheme

289 RegCM4-NH also employs the Single-Moment 5-class microphysics scheme of the WRF 290 model (Skamarock et al., 2008). This scheme (ipptls = 3 in µparam) follows Hong 291 et al. (2004) and, similarly to Nogherotto et al. (2016), includes vapor, rain, snow, cloud 292 ice, and cloud water hydrometeors. The scheme separately treats ice and water 293 saturation processes, assuming water hydrometeors for temperatures above freezing, 294 and cloud ice and snow below the freezing level (Dudhia, 1989, Hong et al., 1998). It 295 accounts for supercooled water and a gradual melting of snow below the melting layer 296 (Hong et al., 2004, and Hong and Lim, 2006). Therefore, the WSM5 and Nogherotto-297 Tompkins schemes have similar structures (Figure 2), but also important differences.

298 Differently from the Nogherotto-Tompkins scheme, the WSM5 (as well as the other WSM 299 schemes in WRF) prescribes an inverse exponential continuous distribution of particle 300 size (ex. Marshall and Palmer (1948) for rain, Gunn and Marshall (1958) for snow). It also 301 includes the size distribution of ice particles and, as a major novelty, the definition of the 302 number of ice crystals based on ice mass content rather than temperature. Both the 303 Nogherotto-Tompkins and WSM5 schemes include autoconversion, i.e. sub-time step 304 processes of conversion of cloud water to rain and cloud ice to snow. For rain, Hong et 305 al. (2004) use a Kessler (1969) type algorithm in WSM5, but with a stronger physical basis 306 following Tripoli and Cotton (1980). The Nogherotto-Tompkins scheme also includes the 307 original Kessler (1969) formula as an option, but it makes available other three 308 exponential approaches following Sundqvist et al. (1989), Beheng (1994), and 309 Khairoutdinov and Kogan (2000). For ice autoconversion the Nogherotto-Tompkins 310 scheme uses an exponential approach (Sundqvist, 1989) with a specific coefficient for ice 311 particles (following Lin et al., 1983) depending on temperature, while the WSM5 uses a 312 critical value of ice mixing ratio (depending on air density) and a maximum allowed ice 313 crystal mass (following Rutledge and Hobbs, 1983) that suppresses the process at low 314 temperatures because of the effect of air density. Finally, the WSM5 has no dependency 315 on cloud cover for condensation processes while the Nogherotto-Tompkins scheme uses 316 cloud cover to regulate the condensation rate in the formation of stratiform clouds.

317 Illustrative case studies

319 Three case studies (Table 2) of Heavy Precipitation Events (HPE) have been identified in 320 order to test and illustrate the behavior of the non-hydrostatic core of the RegCM4-NH. 321 with focus on the explicit simulation of convection over different regions of the world. In 322 two of the test cases, California and Lake Victoria, data from the ERA-Interim reanalysis 323 (Dee et al. 2011) are used to provide initial and lateral meteorological boundary conditions 324 (every 6 hours) for an intermediate resolution run (grid spacing of 12 km, with use of 325 convection parameterizations), which then provides driving boundary conditions for the 326 convection-permitting experiments (Figure 3). In the Texas case study, however, we 327 nested the model directly in the ERA-Interim reanalysis given that such configuration 328 was able to accurately reproduce the HPE intensity. In this case the model uses a large 329 LBC relaxation zone which allows the description of realistic fine-scale features driving 330 this weather event (although not fully consistent with the Matte et al. (2017) criteria). All 331 simulations start 24-48 hours before the HPE (Table 2). The analysis focuses on the total 332 accumulated precipitation over the entire model domain at 3 km resolution (Figure 2) for 333 the periods defined in Table 2. In the cases of California and Texas the evaluation also 334 includes the time series of 6 hourly accumulated precipitation averaged on the region of 335 maximum precipitation (black rectangles in Figures 5a and 7a) because high temporal 336 resolution observations (NCEP/CPC) are also available (Table 3). The discussion of the 337 case studies is presented in the next sections; the configuration files (namelists) with full 338 settings for the three test cases are available at https://zenodo.org/record/5106399. 339

340 A key issue concerning the use of CP-RCMs is the availability of very high resolution, 341 high quality observed datasets for the assessment and evaluation of the models, which 342 is lacking for most of the world regions. Precipitation measurements come from 343 essentially three distinct sources: in-situ rain-gauges, ground radar and satellite. In the 344 present study we use 7 observational datasets depending on the case study and the area 345 covered, as described in Table 3. We have used: Precipitation Estimation from Remotely 346 Sensed Information using Artificial Neural Networks - Climate Data Record (PERSIAN-347 CDR), Climate Hazards Group InfraRed Precipitation with Station data (CHIRPS), the 348 Climate Prediction Center morphing method (CMORPH), Tropical Rainfall Measuring 349 Mission (TRMM), NCEP/CPC-Four Kilometer Precipitation Set Gauge and Radar 350 (NCEP/CPC), CPC-Unified gauge-based daily precipitation estimates (CPC) and 351 Parameter-elevation Regressions on Independent Slopes Model (PRISM) (Table 3). 352 NCEP/CPC is a precipitation analysis which merges a rain gauge dataset with radar 353 estimates. CMORPH and PERSIAN-CDR are based on satellite measurements, CHIRPS 354 incorporates satellite imagery with in-situ station data. CPC is a gauge-based analysis of 355 daily precipitation. The PRISM dataset gathers climate observations from a wide range 356 of monitoring networks, applying sophisticated quality control measures and developing 357 spatial climate datasets which incorporate a variety of modeling techniques at multiple 358 spatial and temporal resolutions.

359

Case	ACRONYM	Region of The event	Domains size lon x lat x vertical levels	
1	CAL	California	480 x 440 x 41	15 Feb 2004 00:00 19 Feb 2004 00:00
2	TEX	Texas	480 x 440 x 41	9 June 2010 00:00 12 June 2010 00:00
3	LKV	Lake Victoria	550 x 530 x 41	25 Nov 1999 00:00 1 Dec 1999 00:00

360 Table 2: List of acronyms and description of the test cases with corresponding
 361 3km domain sizes and simulation period.

Dataset name	Region	Spatial Resolution	Temporal Resolution	Data Source	Reference
TRMM	World	0.5°	Daily	Satellite	Huffman et al. (2007)

CHIRPS	World	0.05°	Daily	Station data+Satellit e	Funk et al. (2015)
CMORPH	World	0.25°	Daily	Satellite	Joyce et al. (2004)
NCEP/CPC	USA	0.04°	Hourly	Gauge and Radar	https://doi.or g/10.5065/D 69Z93M3. Accessed: 27/06/2018
CPC	World	0.5°	Daily	Station data	Chen and Xie (2008)
PRISM	USA	0.04°	Daily	Station data	PRISM Climate Group. 2016.
PERSIAN- CDR	World	0.25°	Daily	Satellite	Ashouri et al. (2015)

Table 3: List of observed precipitation datasets used for comparison.

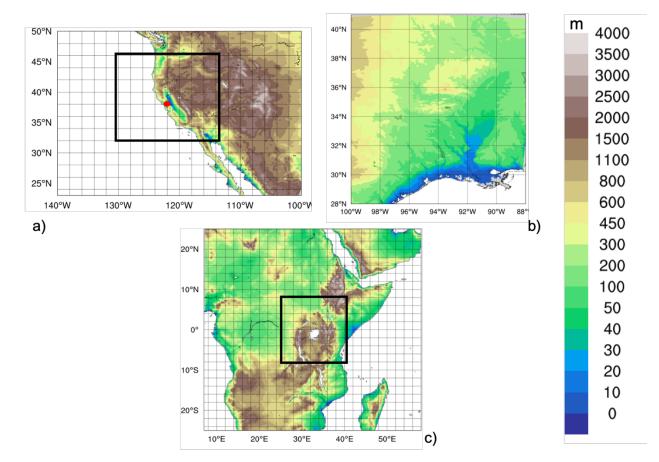


Figure 3: Domains tested , a) California (CAL) , b) Texas (TEX), c) Lake Victoria (LKV) . For CAL (a) and LKV (b) the black square shows the 3 km simulation domains nested in the 12 km domain in figure. For TEX case (b) the 3 km domain simulation is fed directly with the ERA-Interim reanalysis fields.

370

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372 California

The first case, referred to as CAL in Table 2, is a HPE which occurred on February 16-18 2004, producing flooding conditions for the Russian River, a southward-flowing river in the Sonoma and Mendocino counties of northern California (red-dot in Figure 3a). The event is documented in detail by Ralph et al. (2006), who focused their attention on the impact of narrow filament-shaped structures of strong horizontal water vapor transport 378 over the eastern Pacific Ocean and the western U.S. coast, called Atmospheric Rivers 379 (ARs). ARs are typically associated with a low-level jet stream ahead of the cold front of 380 extratropical cyclones (Zhu and Newell 1998; Dacre et al. 2015; Ralph et al. 2018), and 381 can induce heavy precipitation where they make landfall and are forced to rise over 382 mountain chains (Gimeno et al. 2014). The CAL event consists of a slow propagating 383 surface front arching southeastward towards Oregon and then southwestward offshore 384 of California (Figure 4a,c). Rain began over the coastal mountains of the Russian River 385 watershed at 0700 UTC of February 16, as a warm front descended southward, and also 386 coincided with the development of orographically favoured low-level upslope flow (Ralph 387 et al., 2006).

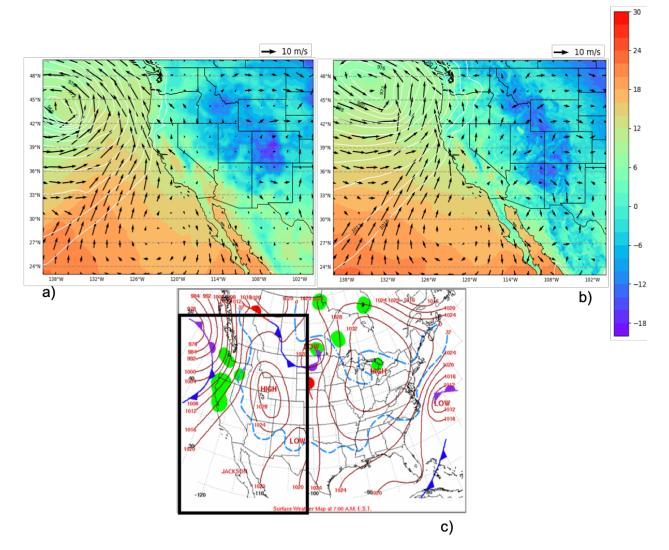


Figure 4: a,b) mean sea level pressure (mslp, hPa, white contour lines), surface
temperature (color shading, °C) and 100-m wind direction (black arrows, m/s) at 0700 UTC,
February 16, 2004 of ERA5 reanalysis and RegCM 12km respectively. c) NCEP-NOA
Surface Analysis of pressure and fronts. The black box in (c) bounded the area represented
in (a) and (b)

394 The intermediate resolution (12 km) domain (Figure 3a) covers a wide area 395 encompassing California and a large portion of the coastal Pacific Ocean, with 23 vertical 396 levels and a parameterization for deep convection based on the Kain-Fritsch scheme 397 (Kain, 2004). The ERA-Interim driven simulation is initialized at 0000 UTC, February 15 2004 (Table 2) and lasts until 0000 UTC February 19 2004. This simulation is used as a 398 399 boundary conditions for a RegCM4-NH run over a smaller area centered over northern 400 California (Fig. 3a) at 3 km horizontal resolution, with 41 vertical levels and boundary 401 conditions updated every 6 hours. In RegCM4-NH only the shallow convection code of 402 the Tiedtke scheme (Tiedtke, 1996) is activated. Simulated precipitation is compared 403 with the CHIRPS, CMORPH, TRMM, PRISM, NCEP/CPC observations (Table 3).

404 As shown in Figure 4 the February 16 synoptic conditions for mean sea level pressure 405 (mslp), surface temperature and wind direction of this case study, are well reproduced by 406 RegCM4 at 12 km (Fig. 4b) when compared to ERA5 reanalysis (Fig. 4a). The surface 407 analysis of pressure and fronts derived from the operational weather maps prepared at 408 the National Centers for Environmental Prediction, Hydrometeorological Prediction 409 National Weather Center, Service 410 (https://www.wpc.ncep.noaa.gov/dailywxmap/index 20040216.html) is also reported in 411 Figure 4c.

The available observed precipitation datasets show similar patterns for the total accumulated precipitation (Figure 5), in particular CHIRPS (Figure 5a), PRISM (Figure 5d) and NCEP (Figure 5e) exhibit similar spatial details and magnitudes of extremes. CHIRPS shows a maximum around 42°N which is not found in the other datasets. CMORPH (Figure 5b) and TRMM (Figure 5c) show lower precipitation maxima and lesser spatial details due to their lower resolution, indicating that the performance of satellite-

418 based products may be insufficient as a stand alone product to validate the model for this419 case.

The largest observed maxima are placed on the terrain peaks, with extreme rainfall greater than 250 mm in 60 hours over the coastal mountains and between 100 – 175 mm elsewhere (Fig. 5). The black box in Fig 5a shows the area of the Russian River watershed where the largest rainfall rates were detected (269 mm and 124 mm in 60-h accumulated rainfall between 0000 UTC February 16 and 1200 UTC February 18, 2004, respectively) (Ralph et al., 2006).

The convection-permitting simulation captures the basic features of the observed precipitation, both in terms of spatial distribution (Fig. 5f) and of temporal evolution of rainfall (Fig. 6a). However, it shows higher precipitation rates than observed over the sea and over the mountain chains, with lower intensities than observed in the south-east part of the mountain chain (Fig. 5). The 12-km simulation instead severely underestimates the magnitude of the event (Fig. 5g).

432 Figure 6a shows the 6-hourly accumulated precipitation averaged over the black box in Figure 5a. The 3 km and 12 km simulations capture the onset of the event, but the peak 433 434 intensity is strongly underestimated by the 12 km run, while it is well simulated by the 3 435 km run, although the secondary maximum is overestimated. These results demonstrate 436 that only the high resolution convection-permitting model is able to captures this extreme 437 event, and that parameterized convection has severe limits in this regard (Done et al. 438 2004; Lean et al. 2008; Weisman et al. 2008; Weusthoff et al. 2010; Schwartz 2014; Clark 439 et al. 2016).

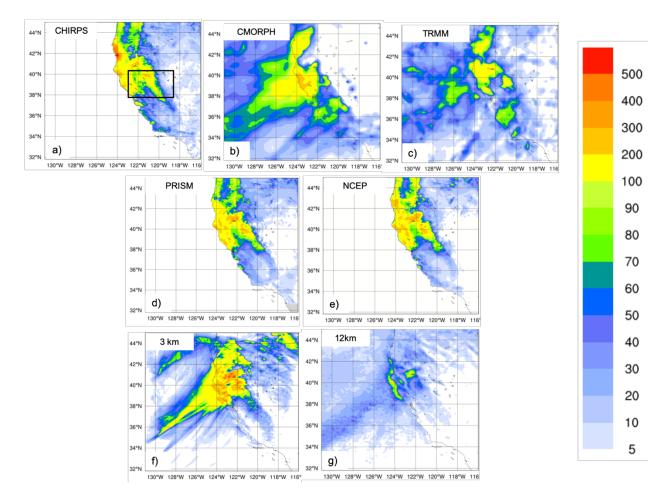




Figure 5 : Total accumulated precipitation (mm) during the California case: CHIRPS (a), CMORPH (b), TRMM (c) observations, PRISM (d) and NCEP Reanalysis (e) and convectionpermitting simulation with RegCM4-NH at 3km (f) and RegCM4 at 12km (g). The black box denotes the area where the spatial average of 6-hourly accumulated precipitation is calculated for Figure 6a.

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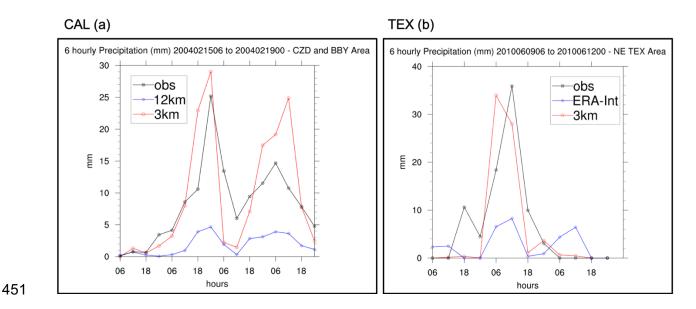


Figure 6: Time series of the 6 hourly accumulated precipitation (in mm on the y-axis) during the CAL event (a) and during the TEX event (b). The blue lines show RegCM4 12 Km and ERA interim 6 hourly accumulated precipitation averaged over the areas indicated by the black squares in Figures 5 and 7 while the red line shows the 6 hourly accumulated precipitation simulated by RegCM4-NH. The observations are shown with a black line.

458 **Texas**

459 Case 2, hereafter referred to as TEX (Table 2), is a convective precipitation episode 460 exhibiting characteristics of the "Maya Express" flood events, linking tropical moisture 461 plumes from the Caribbean and Gulf of Mexico to midlatitude flooding over the central 462 United States (Higgins 2011). During the TEX event, an upper-level cutoff low over 463 northeastern Texas, embedded within a synoptic-scale ridge, moved slowly 464 northeastward. Strong low-level flow and moisture transport from the western Gulf of 465 Mexico progressed northward across eastern Texas. The event was characterized by low-level moisture convergence, weak upper-level flow, weak vertical wind shear, and 466 467 relatively cold air (center of cutoff low), which favored the slow-moving convective storms and nearly stationary thunderstorm outflow boundaries. The main flooding event in 468 469 eastern Texas occurred on June 10, 2010, with a daily maximum rainfall of 216.4 mm for 470 the region in the black box of Figure 7a (Higgins 2011).

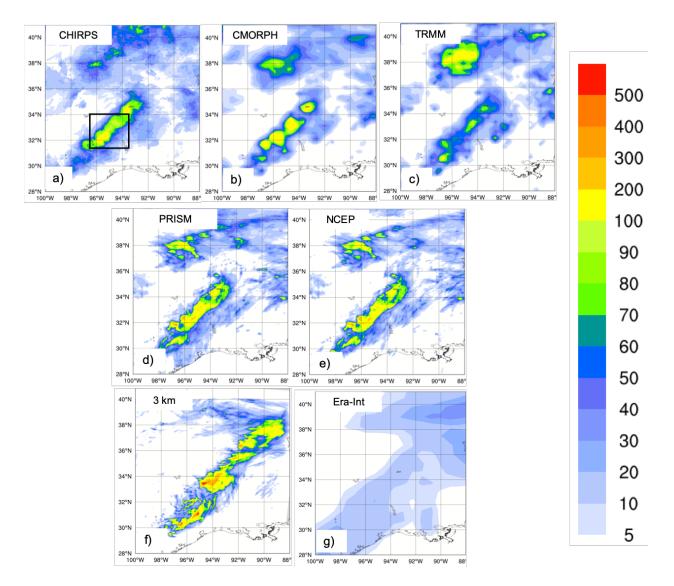


Figure 7: Total accumulated precipitation (mm) during the Texas case: CHIRPS (a), CMORPH (b), TRMM (c), PRISM (d), NCEP Reanalysis (e) and convection-permitting simulation with RegCM4-NH at 3 km grid spacing (f) and ERA-Interim (g). The black box (a) shows the area where the spatial average of 6-hourly accumulated precipitation was calculated for Figure 6b

As for the California case, the observed precipitation datasets show coherent patterns for the total accumulated precipitation (Figure 7), with the highest values related to the mesoscale convective system in eastern Texas (~ 200 mm), and another smaller area of high precipitation more to the north, approximately over Oklahoma. PRISM (Figure 7d)and NCEP (Figure 7e) capture similar spatial details and magnitudes of extremes, 482 CHIRPS (Figure 7a) has lower precipitation extremes in the north compared to the other 483 datasets, while CMORPH (Figure 7b) and TRMM (Figure 7c) show the lowest 484 precipitation extremes and reduced spatial details as already noted for the California 485 case.

486 Figure 7f and Figure 7g present precipitation as produced by the RegCM4-NH and the 487 ERA-Interim reanalysis (driving data) respectively. ERA-Interim reproduces some of the 488 observed features of precipitation, but with a substantial underestimation over the areas 489 of maximum precipitation because of its coarse resolution. By comparison, the RegCM4-490 NH simulation (Fig. 7f) shows an improvement in both pattern and intensity of 491 precipitation, and is substantially closer to observations over eastern Texas. However, 492 the precipitation area is slightly overestimated and the model is not capable of 493 reproducing the small region of maximum precipitation in the north.

494

495 The time series of precipitation over eastern Texas from June 9 to 12, 2010 for 496 observations (black line), ERA-Interim (blue line) and RegCM4-NH (red line) are reported 497 in figure 6b. Precipitation increases over this region from 0000 UTC until it reaches the 498 observed maximum at 1200 UTC, on June 10 (~35 mm), gradually decreasing afterwards 499 until 0600 UTC, on June 11. The RegCM4-NH simulation shows a more realistic temporal 500 evolution than the ERA-Interim, which exhibits an overall underestimation of precipitation. 501 The non-hydrostatic model produces precipitation values closer to the observations, 502 however the simulated maximum is reached 6 hours earlier than observed.

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- 504

505 Lake Victoria

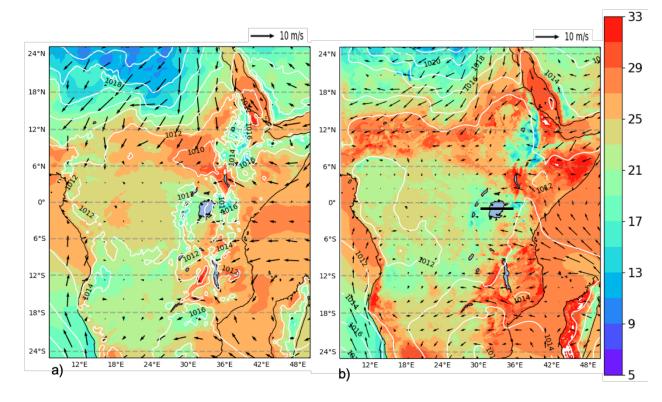
506 Case 3 focuses on Lake Victoria (LKV), with the purpose of testing RegCM4-NH on a 507 complex and challenging region in terms of convective rainfall. It is estimated that each 508 year 3,000-5,000 fishermen perish on the lake due to nightly storms (Red Cross, 2014). 509 In the Lake Victoria basin, the diurnal cycle of convection is strongly influenced by 510 lake/land breezes driven by the thermal gradient between the lake surface and the 511 surrounding land. As the land warms during the course of the day, a lake breeze is 512 generated which flows from the relatively cooler water towards the warmer land surface. 513 The circulation is effectively reversed at night, when the land surface becomes cooler 514 than the lake surface, leading to convergence over the lake and associated thermal 515 instability.

516 In the LKV region, prevailing winds are generally easterly most of the year with some 517 variability due to the movement of the ITCZ. The local diurnal circulation created by the 518 presence of the lake creates two diurnal rainfall maxima. During daylight hours, when the 519 lake breeze begins to advance inland, convergence is maximized on the eastern coast of 520 the lake as the lake breeze interacts with the prevailing easterlies. Studies have also 521 noted the importance of downslope katabatic winds along the mountains to the east of 522 the lake in facilitating convergence along the eastern coastal regions (Anyah et al. 2006). 523 This creates a maximum in rainfall and convection on the eastern coast of LKV. 524 Conversely, during nighttime hours, when the local lake circulation switches to flow from 525 the land towards the lake, the prevailing easterlies create locally strong easterly flow 526 across the lake and an associated maximum in convergence and rainfall on the western 527 side of LKV.

528 The LKV simulation starts on November 25, 1999 and extends to the beginning of 529 December 1999 (Table 2), covering a 5-day period which falls within the short-rain season 530 of East Africa. The choice of 1999, an ENSO neutral year, was made in order to focus the 531 analysis on local effects, such as the diurnal convection cycle in response to the lake/land 532 breeze, with no influence of anomalous large scale conditions. A 1-dimensional lake 533 model (Hostetler et al. 1993; Bennington et al. 2014) interactively coupled to RegCM4-534 NH was utilized to calculate the lake surface temperature (LST), since lake-atmosphere 535 coupling has been shown to be important for LKV (Sun et al. 2015; Song et al. 2004). 536 This coupled lake model has been already used for other lakes, including Lake Malawi in 537 southern Africa (Diallo et al. 2018). As with the other experiments, the boundary 538 conditions are provided by a corresponding 12 km RegCM4 simulation employing the 539 convection scheme of Tiedtke (1996).

540 At the beginning of the simulation, the LST over the lake is uniformly set to 26°C, and is 541 then allowed to evolve according to the lake-atmosphere coupling. This initial LST value 542 is based on previous studies. For example, Talling (1969) finds Lake Victoria surface temperatures ranging from 24.5-26°C during the course of the year. Several studies have used RCMs to investigate the Lake Victoria climate (Anyah et al., 2006; Anyah and Semazzi 2009, Sun et al. 2015), and found a significant relationship between lake temperature and rainfall depending on season. The value of 26°C is typical of the winter season and was chosen based on preliminary sensitivity tests using different values of initial temperature ranging from 24°C to 26°C.

549 The synoptic feature favorable for the production of precipitation over the LKV in this 550 period corresponds to a large area of southeasterly flow from the Indian Ocean (Fig. 8a), 551 which brings low-level warm moist air into the LKV region facilitating the production of 552 convective instability and precipitation. This synoptic situation, with a low-level south-553 easterly jet off the Indian Ocean, is a common feature associated with high precipitation 554 in the area (Anyah et al. 2006), and can be seen in ERA5 data (Figure 8a). Although 555 some bias in terms of magnitude, this is reasonably well reproduce by the 12 km 556 simulation (Figure 8b).



557

558 Figure 8: Mean sea level pressure (mslp) (hPa) (white contour lines), surface temperature 559 (color shading) (°C) and 100-m wind (black arrows) averaged over the period 25 November

560 0000 UTC - 1 December 0000 UTC, by ERA5 reanalysis (a) and RegCM 12km (b). The black 561 line (b) shows the cross-section position represented in Fig. 9

562 The LKV region dynamics are quite distinct between nighttime and daytime and the 563 rainfall in and around the lake has a pronounced diurnal cycle. To understand this strong 564 diurnal cycle, Figure 9 shows a cross-section through the lake (32E to 34E, black line in right panel of Fig. 8b) along 1°S latitude at a period during strong nighttime (Fig. 9b,d; 565 566 0600Z November 30) and daytime convection (Figure 9a.c; 12Z November 29). Wind 567 vectors in Figure 9 show the zonal-wind anomaly across 0°-2°S to highlight the 568 circulations associated with LKV. During the day, surface heating around the lake leads 569 to a temperature difference between the land and lake sufficient to generate a lake 570 breeze, which causes divergence over the lake, while over the highlands to the east the environment is more conducive to convection where convergence is focused (9a,c). 571 572 Conversely, during the night, a land breeze circulation is generated, which induces 573 convergence and convection over the lake (Figure 9b,d). In Figure 10, the evolution of 574 the nighttime land breeze is illustrated with cooler temperature anomalies propagating 575 westward onto the lake during the night.

576 Comparing the 3 km simulation to the 12 km forcing run, we find that the localized 577 circulations created by local forcings (i.e. convection) are much stronger in the convection 578 permitting resolution experiment. We also find stronger and more localized areas of 579 convective updrafts compared to the 12 km simulation (Figure 9c,d; omega is shown 580 instead of vertical velocity here because of the difference in dynamical core). As an 581 example during the nighttime event (Figure 9b,d) there is a broad area of upward motion 582 over the lake and the associated broad convergence in the 12km simulation, while in the 583 convection permitting 3km simulation, convection is much more local and concentrated 584 over the western part of the lake. Indeed, nighttime rainfall tends to be concentrated over 585 the western part of the lake (Sun et al. 2015; Figure 11a-d). Stronger convection 586 simulated in the 3 km experiment could also be tied to stronger temperature anomalies 587 shown over the lake and land and between day and night relative to the 12km simulation 588 (Figure 10). The 3km simulation also shows a more pronounced land breeze propagation 589 at night compared to the 12km simulation.

590 This demonstrates that the 3km simulation is better equipped to simulate the localized 591 circulations associated with this complex land-lake system.

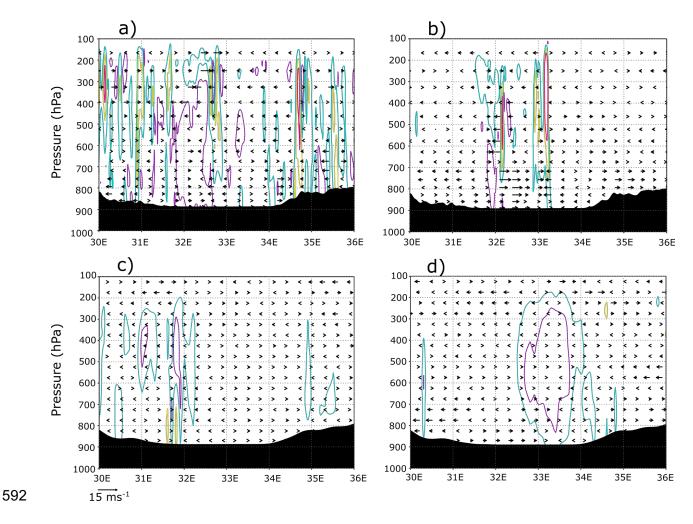




Figure 9. Cross-section through 1°S (black line in Fig. 8b) of the zonal-wind anomaly (0°-594 595 2°S) vectors and the mean contoured vertical velocity (m/s) over 0°-2°S at a) 12Z 29 596 November and b) 6Z 30 November from the 3km simulation. Purple dashed contours 597 indicate -0.1 m/s, light blue contours indicate 0.1 m/s, yellow contours indicate 0.3 m/s, 598 and red contours indicate 0.5 m/s. Lake Victoria encompasses about 32°E to 34°E. The 599 bottom 2 panels show the same as in a) and b) but from the 12km simulation at c) 12Z 29 600 November and d) 6Z 30 November. Purple dashed contours indicate -0.01 hPa/s, light blue 601 dashed contours indicate -0.005 hPa/s, and yellow dashed contours indicate 0.005 hPa/s.

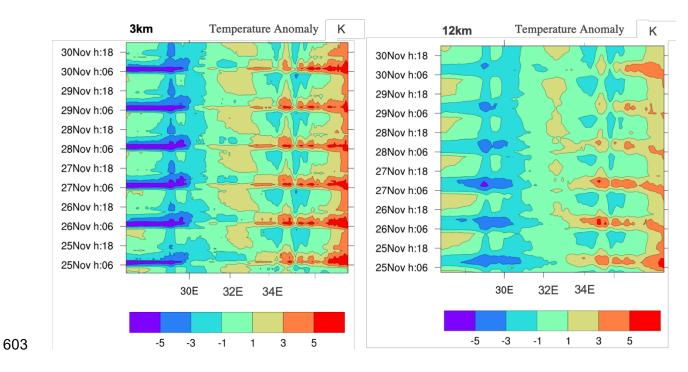


Figure 10 : Longitude-time (hourly) Hovmöller diagram of LKV domain surface temperature
 anomaly (shading, in K). Panels correspond to the 3km simulation (left) and 12km
 simulation (right). The lake Victoria is between 32°E and 34°E longitude

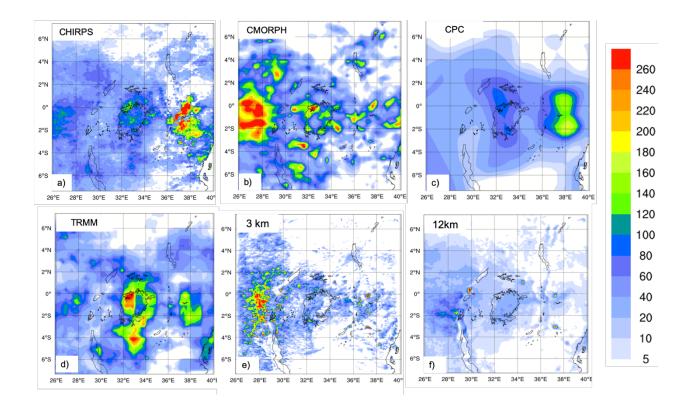


Figure 11: Total event accumulated precipitation (mm) during the LKV case (November 25,
1999-December 1, 1999) measured by CHIRPS (a), CMORPH (b), CPC (d) TRMM (e) and
calculated by RegCM4 at 3 km (e) and 12 km (f).

612

613 Figure 11 reports the total accumulated precipitation observed and simulated for the LKV 614 case. TRMM (Figure 11d) and CPC (Figure 11c) show a similar pattern, with two-rainfall maxima of different intensities over the southeastern and northwestern lake areas. 615 616 CMORPH (Figure 11b) shows a western rainfall maximum similar to TRMM and one large 617 rainfall area almost entirely centered over the highlands to the west of the lake. 618 Conversely in CHIRPS (Figure 11a) a maximum is found to the east of the lake while 619 several localized maxima occur over the lake. The differences among the observed 620 datasets highlight the issue of observational uncertainty and the need to take into 621 consideration shortcomings associated with the types of observational datasets 622 considered. Different datasets can have significantly different climatologies, especially in 623 areas of low data availability. For example, Prein and Gobiet (2017) analyzed two gauge-624 based European-wide datasets, and seven global low-resolution datasets, and found

625 large differences across the observation products, often of similar magnitude as the 626 difference among model simulations. In this case and for this area the observation 627 uncertainty plays a big role especially at high resolution, and highlights the need for an adequate observational network for model validation. However, despite the large 628 629 uncertainty among the different observed datasets (Figure 11 a-d), we find a significant 630 underestimation of the precipitation by the 12 km run over the lake independently of the 631 dataset used as a reference (Figure 11f). In contrast, the 3 km simulation (Figure 11e) 632 shows substantially greater detail, with rainfall patterns more in agreement with the 633 CMORPH data. In particular, the 3 km simulation reproduces well the local rainfall 634 maxima on the western side of the lake, although these appear more localized and with 635 a multi-cell structure compared to CMORPH and TRMM. Additionally, the 12 km simulation underestimates the observed heavy rainfall totals in the highlands to the west 636 637 of the lake region especially when compared to CMORPH, which are instead reproduced 638 by the 3 km simulation.

This last test case demonstrates the ability of RegCM4-NH in simulating realistic convective activity over a such morphologically complex region, which is a significant improvement compared to the hydrostatic-coarse resolution model configuration.

642

643 **Conclusions and future outlook**

644

645 In this paper we have described the development of RegCM4-NH, a non hydrostatic 646 version of the regional model system RegCM4, which was completed in response to the 647 need of moving to simulations at convection-permitting resolutions of a few kilometers. 648 The non-hydrostatic dynamical core of MM5 has been incorporated into the RegCM4 649 system previously based on the MM5 hydrostatic core. Some modifications to the MM5 650 dynamical core were also implemented to increase the model stability for long term runs. 651 RegCM4-NH also includes two explicit cloud microphysics schemes needed to explicitly 652 describe convection and cloud processes in the absence of the use of cumulus 653 convection schemes. Finally, we presented a few case studies of explosive convection to

654 illustrate how the model provides realistic results in different settings and general 655 improvements compared to the coarser resolution hydrostatic version of RegCM4 for 656 such types of events.

657

658 As already mentioned, RegCM4-NH is currently being used for different projects, and 659 within these contests, is being run at grid spacings of a few kilometers for continuous 660 decadal simulations, driven by reanalyses of observations or GCM boundary conditions 661 (with the use of an intermediate resolution domains) over different regions, such as the 662 Alps, the Eastern Mediterranean, Central-Eastern Europe and the Caribbeans. These 663 projects, involving multi-model inter-comparisons, indicate that the performance of 664 RegCM4-NH is generally in line with that of other convection-permitting models, and 665 exhibits similar improvements compared to coarser resolution models, such as a better 666 simulation of the precipitation diurnal cycle and of extremes at hourly to daily time scales. 667 The results obtained within the multi-model context confirm previous results from single-668 model studies (Kendon et al. 2012, 2017, Ban et al. 2014, 2015; Prein et al. 2015, 2017), 669 but also strengthen the robustness of the findings through reduced uncertainty compared 670 to coarse resolution counterpart (Ban et al., 2021, Pichelli et al., 2021). The convection-671 permitting scale can thus open the perspective of more robust projections of future 672 changes of precipitation, especially over sub-daily time scales.

673

674 One of the problems of the RegCM4-NH dynamical core is that, especially for long runs 675 with varied meteorological conditions, a relatively short time step is needed for stability 676 reasons. This makes the model rather computationally demanding, although not more 677 than other convection-permitting modeling systems such as the Weather Research and 678 Forecast model (WRF, Skamarok et al. 2008). For this reason, we are currently 679 incorporating within the RegCM system a very different and more computationally efficient 680 non-hydrostatic dynamical core, which will provide the basis for the next version of the 681 model, RegCM5, to be released in the future.

682

Following the philosophy of the RegCM modeling system, RegCM4-NH is intended to be a public, free, open source community resource for external model users. The nonhydrostatic dynamical core has been implemented in a way that it can be activated in place of the hydrostatic dynamics through a user-set switch, which makes the use of RegCM4-NH particularly simple and flexible. We therefore envision that the model will be increasingly used by a broad community so that a better understanding can be achieved of its behavior, advantages and limitations.

- 690
- 691 Code availability: <u>https://zenodo.org/record/4603556</u>
- 692 Cases study configuration files: https://zenodo.org/record/5106399
- 693 694

Author contribution: CE prepared the manuscript with contributions from all co-authors and coordinated research, SP, TA, GR carried out and analysed the simulations, PE investigated solutions to stabilize/adapt the model at the km-scale and performed preliminary validation tests, GG developed/adapted the model code, FDS contributed to develop the coupled version of the model, NR developed one of the microphysics scheme, GF supervised and coordinated all activities.

- 701
- 702 **Competing interests**: The authors declare that they have no conflict of interest.
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- 704

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