



# Mitigation of the double ITCZ syndrome in BCC-CSM2-MR through improving parameterizations of boundary-layer turbulence and shallow convection

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**Abstract.** The spurious double intertropical convergence zone (ITCZ) is one of the most prominent systematic biases in coupled atmosphere-ocean general circulation models (CGCMs), and the underestimated marine stratus over eastern subtropical oceans has been recognized as a possible contributor. Rather than modifying the cloud scheme itself, this study significantly promotes the marine stratus simulation through improving parameterizations of boundary-layer turbulence and shallow convection in the medium-resolution Beijing Climate Center Climate System Model version 2 (BCC-CSM2-MR). The University of Washington moist turbulence scheme is implemented in BCC-CSM2-MR to better represent the stratocumulus, and a decoupling criterion is also introduced to the shallow convection scheme for improving the simulation of the stratocumulus-to-cumulus transition. Results show that the simulated precipitation in the eastern Pacific south of the equator is largely reduced, alleviating the double ITCZ problem. The tropical precipitation asymmetry index increases from  
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20  $-0.024$  in the original BCC-CSM2-MR to  $0.147$  in the revised BCC-CSM2-MR, which is much closer to the observation. The study suggests that improving parameterizations of boundary-layer turbulence and shallow convection is effective for mitigating the double ITCZ syndrome in CGCMs.

## 1 Introduction

The coupled atmosphere-ocean general circulation models (CGCMs) have been widely used in the studies of climate variability, change and prediction. Despite decades of development, there are still many systematic biases in CGCMs, hampering the reliability of model results and limiting their utility. A prominent tropical bias in generations of CGCMs is the double intertropical convergence zone (ITCZ) syndrome, which is characterized by two parallel zonal bands of annual precipitation straddling the equator over the central and eastern Pacific, while it is absent in the observations (Mechoso et al., 1995; Lin, 2007; Zuidema et al., 2016; Zhang et al., 2019). The observed convergence zone south of the equator extends  
30 southeastward from the western Pacific, whereas most CGCMs simulate a southern zonal rainfall band extending too far



eastward. This bias is often associated with excessive warm sea surface temperatures (SSTs) in the southeastern Pacific (SEP). Similarly, in the tropical Atlantic basin, most CGCMs present a fake southern zonal convergence zone, mirroring the actual zonal convergence zone north of the equator. The spurious double ITCZ not only affects the intensity of the Hadley circulation and the distribution of the trade winds directly related to the simulation of El Niño events, but also creates biases  
35 of latent heating in the tropics that can impact midlatitude weather and climate through atmospheric teleconnections (Schneider et al., 2009; Manganello and Huang, 2009).

The double ITCZ syndrome in CGCMs has been a long-standing problem, and it remains a serious impediment. In earlier phases of Coupled Model Intercomparison Project (CMIP), from Phase 3 (CMIP3) to Phase 5 (CMIP5), most CGCMs suffered from the double ITCZ problem to various degrees (Lin, 2007). Compared with CMIP3 models, there is no  
40 evidence of improvements in reducing the excessive precipitation and warmer SST in the SEP by the CMIP5 models, with results from CMIP5 somewhat worse than those from CMIP3 (Zhang et al., 2015). In the latest CMIP6, the biases persist in some of the current state-of-the-art CGCMs (Williams et al., 2018; Wu et al., 2019), indicating that it remains a tough challenge to alleviate the double ITCZ bias in CGCMs.

Many efforts have been devoted to identifying the possible contributors to the double ITCZ problem, and both oceanic  
45 and atmospheric causes are suggested (Zhang et al., 2019). For instance, by restoring model ocean temperature and salinity to observations in the upper ocean, the effects of ocean eastern boundary biases along the Peruvian-Chilean coast was investigated, and it was found that the coastal SST and salinity biases exert significant influences on the SEP precipitation (Large and Danabasoglu, 2006). Modelling studies show that the double ITCZ bias is very sensitive to atmospheric processes, such as convections and cloud. As ITCZ precipitation originates from convections, convection parameterization  
50 schemes are often blamed for the double ITCZ problem. By modifications in closure, trigger function, and lateral entrainment of the convection scheme, the double ITCZ can be mitigated to varying degrees (Song and Zhang, 2009; Zhang and Song, 2010; Queslatti and Bellon, 2013; Song and Zhang, 2018). Deficiencies in the extratropical cloud simulation are also suggested to be possible causes of the tropical double ITCZ (Hwang and Frierson, 2013; Li and Xie, 2014). It is argued that the negative cloud amount biases over the Southern Ocean and the associated warming can induce anomalous northward  
55 cross-equatorial atmospheric energy transport, resulting in a meridional shift of ITCZ. However, based on this atmospheric teleconnection argument and by reducing the shortwave radiation bias over the Southern Ocean in the fully coupled National Center for Atmospheric Research (NCAR) Community Earth System Model version 1 (CESM1), Kay et al. (2016) found that the double ITCZ is not improved.

The negative cloud amount biases off the west coast of South America, which are common in CGCMs, are regarded as  
60 another cause for the double ITCZ problem. It is believed that the underestimated cloud cover leads to more net heat flux into the ocean and the warm SST biases in the SEP, which is associated with stronger convections and precipitation. Previous studies to ameliorate the double ITCZ bias mostly focused on the role of the cloud fraction parameterization, either through prescribed increases or through physical parameterization changes (Ma et al., 1996; Yu and Mechoso, 1999; Dai et al., 2003, 2005; Qin and Lin, 2018). Furthermore, increased cloud fraction and the associated shortwave cloud forcing over



65 the SEP can be driven by changes in representation for cloud microphysics (Woelfle et al., 2019). Modifications on the cloud scheme itself are helpful to mitigate the double ITCZ to some extent; however, it still plagues most CGCMs.

Rather than directly modifying the cloud macro/microphysics schemes in CGCMs, this study changes parameterizations of the boundary-layer turbulence and shallow convection to improve the low-level cloud simulation indirectly. The low-level cloud near the South American west coast is the steadiest and most persistent stratocumulus regime in the world, which is  
70 closely related to multiple processes, such as boundary layer (BL) mixing, surface sensible and latent heat fluxes, cloud-top radiative cooling, and entrainment (Wood 2012). Furthermore, a prominent feature of the low-level cloud in the SEP is that the stratocumulus regime progressively transforms into the trade cumulus regime moving downstream off the coast. In the transition, interactions between BL turbulence and shallow convection play a critical role. When multiple separate BL and shallow convection schemes are combined to complement each other in a CGCM, inconsistencies between parameterizations  
75 are likely to be introduced. This study aims to find out whether improving parameterizations of BL turbulence and shallow convection accounts for the double ITCZ bias in CGCMs. This is done through a two-step process, using the medium resolution Beijing Climate Center Climate System Model version 2 (BCC-CSM2-MR). The first step is to determine if the stratocumulus simulation is improved with prescribed SSTs through modifying parameterizations of BL turbulence and shallow convection in the atmospheric component of BCC-CSM2-MR. The second step, which indicates the atmosphere-  
80 ocean feedback, is to perform a pair of ocean-atmosphere coupled simulations to demonstrate the impact of improved stratocumulus representation on the simulation of the SSTs and precipitation in the SEP.

The paper is organized as follows. Section 2 describes the BCC-CSM2-MR, focusing on the new configurations of the BL and shallow convection schemes, as well as the model setups for the control and sensitive experiments. A brief review of the observational data used to evaluate the model results is also included. Section 3 compares the cloud simulation results  
85 from the atmosphere-only runs, and section 4 discusses the impact of improved cloud simulation on the SSTs and precipitation in the coupled runs. Further discussion on the effectiveness of the new configuration of the BL turbulence and shallow convection parameterizations is shown in section 5. At last, section 6 presents the summary and conclusions.

## 2 Model description, experimental design, and observational data

### 2.1 Brief description of the BCC-CSM2-MR

90 In the present study, we use the version of BCC-CSM2-MR participating in CMIP6, which is a fully coupled model with atmosphere, ocean, land surface, and sea ice components (Wu et al., 2019). The atmospheric component in BCC-CSM2-MR is based on the Beijing Climate Center atmospheric general circulation model (BCC-AGCM; Wu et al., 2010). The spectral dynamical core of BCC-AGCM is featured by introducing a reference stratified atmospheric temperature and a reference surface pressure into the governing equations (Wu et al., 2008). The Beijing Climate Center Atmosphere-Vegetation  
95 Interaction Model (BCC-AVIM; Li et al., 2019) serves as the land component of BCC-CSM2-MR. It includes major land surface biophysical and plant physiological processes. The oceanic component is based on the Modular Ocean Model



version 4 (MOM4; Griffies et al., 2005) and the sea ice component is the Sea Ice Simulator (SIS; Winton, 2000). Over the sea ice, a new bulk aerodynamic algorithm is formulated for computing surface exchange fluxes (Lu et al., 2013). The above four components are physically coupled through fluxes of momentum, energy, and water at their interfaces, which is realized using the NCAR flux coupler version 5.

The atmospheric component of BCC-CSM2-MR has a horizontal resolution of T106 (approximately 1.125° latitude by 1.125° longitude) with 46 hybrid vertical levels. The deep convection scheme used in BCC-CSM2-MR is based on a mass-flux bulk cloud model approach, in which the mass change for the adiabatic ascent cloud parcel with altitude is derived from a total energy conservation equation of the whole adiabatic system involving the updraft cloud parcel and the environment (Wu 2012). Compared with its previous version in CMIP5, BCC-CSM2-MR has notably improved the simulation skills of atmospheric variability in the tropics, such as the Madden-Julian oscillation and the stratospheric quasi-biennial oscillation (Wu et al., 2019; Lu et al., 2020).

However, the version of BCC-CSM2-MR participating in CMIP6 still suffers from the double ITCZ syndrome. The mean precipitation errors are dominated by systematic errors along the ITCZ. Here we show that a revised version of BCC-CSM2-MR, configured with new parameterizations of BL turbulence and shallow convection, significantly reduces the double ITCZ bias.

## 2.2 Parameterization of BL processes

The CMIP6 configuration of the BCC-CSM2-MR employs the Holtslag and Boville (1993) parameterization (hereafter called the HB scheme), which is optimized for the simulation of dry convective BLs over land. The HB scheme is based on the eddy diffusivity approach. And the eddy diffusivity of variables  $\chi$  is given by

$$K_{\chi} = kw_{\chi}z\left(1 - \frac{z}{h}\right)^2, \quad (1)$$

where  $k$  is the von Karman constant;  $z$  is the height;  $w_{\chi}$  is a turbulent velocity; and  $h$  is the boundary layer height. For neutral and stable conditions,  $w_{\chi}$  is proportional to the friction velocity, while for unstable conditions,  $w_{\chi}$  is proportional to the convective velocity scale  $w_*$

$$w_* = \left[ \frac{g}{\theta_{vs}} h \overline{(w'\theta'_v)_s} \right]. \quad (2)$$

Here,  $g$  is the gravitational acceleration;  $\theta_{vs}$  is the virtual potential temperature at surface; and  $\overline{(w'\theta'_v)_s}$  represents the buoyancy heat flux at surface. The above formula suggests that the HB scheme assumes the BL turbulence to be forced exclusively from the surface heating and friction velocity.

In the marine stratocumulus-capped BLs, the turbulence structure depends strongly on the dominant turbulence generating mechanism resulting from both evaporative and radiative cooling at cloud top. To provide a more physically



realistic treatment of stratocumulus-topped BLs, the University of Washington moist turbulence (UWMT) scheme from Bretherton and Park (2009) is implemented in the revised BCC-CSM2-MR to replace the HB scheme. The UWMT scheme also uses first-order  $K$  diffusion to represent all the turbulence, in which the eddy diffusivity is calculated based on the turbulent kinetic energy (TKE,  $e$ ) and proportional to the stability-corrected length scale  $lS_\chi$ , given by

$$130 \quad K_\chi = lS_\chi \sqrt{e}. \quad (3)$$

In the case of an inversion layer at the top of convective BLs, the diffusivities are parameterized as following

$$K_\chi = w_e \Delta z_e, \quad (4)$$

where  $w_e$  is the entrainment rate, and  $\Delta z_e$  is the thickness of the entrainment layer. The UWMT scheme uses the  $w_e$  entrainment closure raised by Nicholls and Turton (1986):

$$135 \quad w_e = A \frac{w_*^3}{(g \Delta^E s_{vl} / s_{vl})(z_t - z_b)}. \quad (5)$$

Here,  $z_t$  and  $z_b$  are the top and bottom heights of the entrainment layer, respectively;  $\Delta^E$  denotes a jump across the entrainment layer; and  $s_{vl}$  is the liquid virtual static energy.  $A$  is a nondimensional entrainment efficiency, which is affected by evaporative cooling of mixtures of cloud-top and above-inversion air.

### 2.3 Shallow cumulus Parameterization

140 To treat shallow cumulus, the original version of the BCC-CSM2-MR adopts a stability-dependent mass-flux representation of moist convective processes with the use of a simple bulk three-level cloud model, as in Hack (1994). Specifically, in a vertically discrete model atmosphere where the level index  $k$  decreases upward and layers  $k$  and  $k+1$  are moist adiabatically unstable, the Hack scheme assumes that there is a non-entraining convective element with roots in level  $k+1$ , condensation and rain out processes in level  $k$ , and limited detrainment in level  $k-1$ . By repeated application of this procedure from the  
145 bottom of the model to the top, the thermodynamic structure is locally stabilized.

The Hack shallow cumulus scheme can be also active in moist turbulent mixing, such as stratocumulus entrainment, which has different physical characteristics than cumulus convection. Shallow cumulus is usually regarded as a decoupled BL regime, in which the vertical mixing processes do not achieve a single well-mixed layer, while the stratocumulus regime represents a well-mixed BL up to the cloud top. The decoupling criterion to distinguish the two regimes is of great  
150 importance for simulating the stratocumulus-to-cumulus transition (Bretherton and Wyant, 1997; Wood and Bretherton, 2004). A number of these decoupling criteria have been explored, such as static stability (Klein and Hartmann, 1993) and the buoyancy flux integral ratio (Turton and Nicholls, 1987). In the light of its robustness, the stability criterion with a threshold of 15 K is introduced into the Hack scheme. The lower tropospheric stability ( $LTS$ ) is defined as



$$LTS = \theta_{700hPa} - \theta_{sfc}, \quad (6)$$

155 where  $\theta_{700hPa}$  and  $\theta_{sfc}$  are potential temperatures at 700 hPa and surface, respectively. In the revised BCC-CSM2-MR,  
the Hack scheme is activated only in the decoupled BL regimes with  $LTS < 15$  K below 700 hPa. Above 700 hPa, the  
Hack scheme is retained to remove any local instability as long as the two adjacent model layers are moist adiabatically  
unstable.

#### 2.4 Cloud fraction parameterization

160 In the original BCC-CSM2-MR, the cloud amount is evaluated via a diagnostic method, depending on relative humidity,  
convective mass fluxes and atmospheric stability (Wu et al., 2019). Three types of clouds are diagnosed: layered stratus  
cloud, convective cloud, and low-level marine stratocumulus. The marine stratocumulus cloud fraction is particularly  
parameterized to compensate the HB dry turbulence scheme. It is calculated by using an empirical relationship between the  
observed stratocumulus cloud amount and the boundary layer stratification, which is evaluated with potential temperatures at  
165 700 hPa and surface, as in Klein and Hartmann (1993).

Considering the advantage of the UWMT scheme to simulate marine stratocumulus, the revised BCC-CSM2-MR  
excludes the empirical calculation of stratocumulus cloud amount and only uses the relative humidity and convective cloud  
fraction to deduce the overall cloud amount. The stratocumulus cloud fraction is assumed to be diagnosed from relative  
humidities. The critical relative humidities  $RH_c$  are used to tune the global annual mean top-of-atmosphere (TOA) shortwave  
170 and longwave radiative energy fluxes close to observations. The original BCC-CSM2-MR uses the low-level (below 750 hPa)  
 $RH_c$  of 0.945, while the revised BCC-CSM2-MR uses the low-level  $RH_c$  of 0.97. Given the exclusion of special  
stratocumulus cloud fraction parameterization and increment in the low-level  $RH_c$  in the revised BCC-CSM2-MR, better  
simulation skills of low-level cloud are attributed to the improved schemes of BL turbulence and shallow convection.

#### 2.5 Experimental design

175 Table 1 summarizes all the experimental setups. Two sets of 11-yr Atmospheric Model Intercomparison Project (AMIP)-  
type simulations are conducted with the same prescribed sea surface boundary conditions using the Hadley Centre SSTs and  
sea ice concentrations. The one using the default configurations of HB and Hack schemes is referred to as REF\_amip, and  
the other using the new configurations of UWMT and modified Hack schemes is referred to as NEW\_amip. For each  
integration, the first year is treated as the spinup, and the last 10 years are used for analysis. Both runs start from identical  
180 initial conditions. Note that the Cloud Feedback Model Intercomparison Project Observation Simulator Package (COSP;  
Bodas-Salcedo et al., 2011) is turned on to better compare with satellite observations.

For the fully coupled experiments, two sets of 12-yr Coupled Model Intercomparison Project (CMIP)-type simulations  
are carried out under pre-industry conditions, using greenhouse gases, ozone, aerosol emission, and others fixed at the level



of 1850. The CMIP simulation with the default HB and Hack schemes is referred to as REF\_cmip, whereas the simulation  
185 using the UWMT and modified Hack schemes is referred to as NEW\_cmip. Both runs are initialized with the data sets at the  
345th year of a pre-industry control run provided by the default BCC-CSM2-MR. These initialization data sets are close to  
the equilibrium state. Note that previous studies have shown that the formation and mitigation of the double ITCZ biases can  
be completed within the first 2 years of simulation (Liu et al., 2012). Therefore, a 12-yr integration for each coupled  
experimental setup is performed, of which the first 2 years are treated as the spinup, and the last 10 years are used for  
190 analysis.

To highlight the relative roles of the UWMT scheme and the modified Hack scheme in the low-level cloud simulation  
in NEW\_amip, two extra AMIP-type sensitivity experiments are designed. The modifications of the boundary-layer  
turbulence scheme and the shallow convection scheme are added, respectively in the two sensitivity experiments. The  
simulation with changes only in the boundary-layer turbulence scheme is referred to as UWMT\_amip, whereas the  
195 simulation with changes only in the shallow convection scheme is referred to as mHack\_amip. The difference between  
UWMT\_amip and REF\_amip will show the isolate impact of change in the boundary-layer turbulence parameterization,  
while the difference between mHack\_amip and REF\_amip will show the isolate impact of change in the shallow convection  
parameterization.

## 2.6 Observational data

200 To evaluate the simulated cloud amount and shortwave cloud radiative forcing (SWCRF), the distribution of cloud fraction  
from the CALIPSO GOCCP data set (GCM-Oriented CALIPSO Cloud Product; Chepfer et al., 2010) and TOA SWCRF  
from the CERES-EBAF data set (Clouds and the Earth's Radiant Energy System Energy Balanced and Filled data product;  
Loeb et al., 2018) are used as references. The simulated precipitation is compared with the Global Precipitation Climatology  
Project (GPCP) monthly precipitation analysis (Adler et al., 2003) for years 1981 to 2010. Considering the uncertainty in  
205 precipitation observations, the Climate Prediction Center (CPC) Merged Analysis of Precipitation (CMAP; Xie and Arkin,  
1997) from 1981 to 2010 is also used. Additional observations include SSTs from the Hadley Centre Sea Ice and Sea Surface  
Temperature data set (HadISST; Rayner, 2003), and wind velocities from the Japanese Meteorological Agency 55-year  
Reanalysis (JRA-55; Kobayashi et al., 2015).

## 3 Improved cloud representation in the atmosphere-only simulations

### 210 3.1 TOA SWCRF

TOA SWCRF is one of the most important metrics to examine for improving the representation of low-level clouds. Figure 1  
compares annual mean TOA SWCRF simulated by REF\_amip with CERES-EBAF observations. The global mean SWCRF  
is  $-45.83 \text{ W m}^{-2}$  in CERES-EBAF and  $-49.51 \text{ W m}^{-2}$  in REF\_amip. REF\_amip simulates the SWCRF distributions relatively  
well, but overestimates the magnitude of SWCRF over regions of Northern Hemisphere storm tracks and underestimates the



215 magnitude of SWCRF over the subtropical marine stratocumulus regions. Compared with REF\_ami (Figure 1d),  
NEW\_ami shows a considerably increased simulated magnitude of SWCRF over the eastern subtropical ocean regions,  
suggesting that the low-level cloud representation is improved with new configurations of BL turbulence and shallow  
convection schemes.

### 3.2 Low-level cloud cover

220 TOA SWCRF strongly depends on the low-level cloud fraction. As shown in Figure 2a, the observed low cloud is prevalent  
over midlatitude storm-track regions, eastern tropical and subtropical oceans and the Southern Ocean. Figure 2b and 2c  
present the simulated annual mean low-level cloud amount from REF\_ami and NEW\_ami, respectively. Generally, both  
simulations reproduce the observed patterns of low-cloud amount with maxima in the midlatitude storm tracks and minima  
in the trade cumulus regime over the ocean. One prominent discrepancy between REF\_ami and observations is a  
225 remarkable underestimation of low-cloud fraction over the eastern subtropical oceans, which contributes to the weak bias in  
the magnitude of TOA SWCRF over these regions in REF\_ami. The observed maxima in the subtropical stratocumulus  
decks are better simulated by NEW\_ami. The global mean low-level cloud fraction increases from 24.64% in REF\_ami to  
28.74% in NEW\_ami, closer to 37.06% in the GOCCP observations. With new configurations of BL turbulence and  
shallow convection parameterizations, the low-cloud fraction is significantly increased by up to 40% over the eastern  
230 subtropical oceans (Figure 2d).

### 3.3 The subtropical Stratocumulus-to-Cumulus (Sc-to-Cu) transition

In order to further illustrate the consistency between the UWMT turbulence scheme and modified Hack shallow convection  
scheme, Figure 3 shows vertical cross sections of cloud fraction depicting the subtropical stratocumulus to trade cumulus  
transition both in the observations and simulations. We focus on the five main regions of the subtropical Sc-to-Cu transition,  
235 with the locations shown in Figure 2a. Observational guidance is provided from GOCCP. GOCCP generally shows a clear  
transition from stratocumulus near the coast to trade cumulus well offshore, marked by a gradual reduction of cloud cover  
along with a rising cloud-top height and thickening cloud depth moving downstream off the coasts (Figure 3, right).  
However, REF\_ami (Figure 3, left) tends to overall underestimate the cloud amount through the cross sections. Specifically,  
REF\_ami produces the cloud fraction much less near the coast and fails to simulate the elevated maxima prevalent in the  
240 observed cross sections. This reflects the limitations of the HB turbulence scheme in simulating stratocumulus. In addition,  
REF\_ami is prone to abruptly reduce the cloud amount out of the coastal stratocumulus regions, resulting in an almost  
cloudless state offshore. No clear transitional regime is simulated, which can be attributed to the excessive vertical mixing  
from the original Hack shallow convection scheme over the western edge of the stratocumulus regime. In contrast to  
REF\_ami, the cross sections in NEW\_ami depict a very different picture. With new configurations of boundary-layer  
245 turbulence and shallow cumulus schemes, NEW\_ami presents more realistic characteristics of subtropical Sc-to-Cu  
transition (Figure 3, middle). The qualitative aspects of the transition are better captured for all regions: most notably the



gradual reduction of cloud cover as it moves offshore and the placement of the maximum cloudiness which generally occurs somewhat offshore, as in GOCCP. A possible reason for the over-extension of Sc in NEW\_amip may be that the decoupling criterion added to the Hack shallow convection scheme is too strong, leading to the weak vertical mixing over the eastern edge of the shallow cumulus regime.

#### 4 Alleuviated double ITCZ in the coupled simulations

The underestimated stratus cloud cover in the SEP has long been recognized as a possible contributor to the double ITCZ bias in the eastern Pacific. The above-mentioned analysis from the atmosphere-only runs demonstrates that the simulated low-cloud amount in the SEP can be increased obviously by modifying parameterizations of BL turbulence and shallow convection. Is it possible that the improved cloud simulation reduces the warm bias of SST in the SEP and hence helps to eliminate the double ITCZ bias? Below will clarify the effects of the increased low-cloud amount in the coupled simulations.

#### 4.1 Precipitation

##### 4.1.1 Precipitation patterns

Figure 4 shows simulated annual-mean precipitation distributions compared with the GPCP and CMAP observational estimates. Two sets of observational estimates are presented here to illustrate the uncertainty in observations. In observations (Figure 4a and 4b), massive precipitation can be found in regions of Asian monsoon and midlatitude storm tracks over the northwest Pacific and Atlantic Oceans. In the tropics, the primary peaks are located in the eastern Indian Ocean and Maritime Continent regions. Furthermore, two zonal precipitation bands are located at  $0^{\circ} - 10^{\circ}\text{N}$  in the equatorial Pacific and Atlantic oceans, respectively, constituting the northern ITCZ. The southern South Pacific convergence zone (SPCZ) is mainly located around  $5^{\circ}\text{S} - 10^{\circ}\text{S}$  near the western Pacific warm pool region and experiences a southeast tilt as it extends eastward into the central Pacific. In the southeast Pacific, the SPCZ, the west coast of South America, and the northern ITCZ bound a triangular-shaded dry region, which is dominated by stratocumulus and trade cumulus. While the main spatial patterns of observed precipitation climatology are properly reproduced, prominent double ITCZ biases develop in the simulation of tropical precipitation in REF\_cmip (Figure 4c). Specifically, the simulated SPCZ does not tilt southeastward as that in the observation but overly extends eastward, leading to the excessive precipitation in the central and eastern equatorial Pacific. The precipitation band of  $> 3 \text{ mm day}^{-1}$  between  $10^{\circ}\text{S}$  and  $5^{\circ}\text{S}$  extends as far east as  $90^{\circ}\text{W}$ . Also, a spurious precipitation band appears in the southern Atlantic at  $10^{\circ}\text{S} - 0^{\circ}$ . In contrast, NEW\_cmip significantly reduces the SPCZ bias in the central and eastern equatorial Pacific and produces no rain belt in the tropical southern Atlantic (Figure 4d). The precipitation of  $> 3 \text{ mm day}^{-1}$  between  $10^{\circ}\text{S}$  and  $5^{\circ}\text{S}$  is confined to west of  $130^{\circ}\text{W}$  in the Pacific Ocean.



#### 275 4.1.2 Asymmetry of the tropical precipitation

The asymmetry in the meridional distribution of tropical precipitation is a characteristic in the observed ITCZ. Figure 5 shows the zonally averaged annual mean precipitation from the GPCP and CMAP observations, and the REF\_cmip and NEW\_cmip simulations. The CMAP observations generally show very similar latitudinal variations to that of GPCP, but with slightly stronger precipitation intensity. These two sets of observational estimates both show a pronounced asymmetry about the equator with strong northern ITCZ peak at 5° – 10°N and weak southern ITCZ peak at 10° – 5°S. The precipitation produced by REF\_cmip is about 50% larger than the observed amounts in the southern ITCZ, which comes largely from the spurious rainfall in the central and eastern Pacific and Atlantic south of the equator. The precipitation rate in the southern ITCZ is comparable to that in the northern ITCZ, and the meridional distribution of the tropical precipitation tends to be symmetrical. In contrast, NEW\_cmip produces much weaker precipitation in the southern ITCZ, which is in good agreement with the observed asymmetry.

To quantitatively examine the ITCZ fidelity in the coupled simulations, the tropical precipitation asymmetry index ( $A_p$ ), which is defined as the precipitation difference between the northern (0° – 20°N) and southern (20°S – 0°) tropics normalized by the tropical mean, is calculated as follows (Hwang and Frierson, 2003; Xiang et al., 2017; Adam et al., 2018)

$$A_p = \left( \overline{P}_{0-20^\circ\text{N}} - \overline{P}_{20^\circ\text{S}-0} \right) / \overline{P}_{20^\circ\text{S}-20^\circ\text{N}}, \quad (7)$$

where the overbar represents a zonal mean of precipitation. The observed annual mean  $A_p$  is 0.194 in GPCP and 0.214 in CMAP, respectively, because the ITCZ is predominantly north of the equator. In the coupled simulations, the  $A_p$  index increases from –0.024 in REF\_cmip to 0.147 in NEW\_cmip, which is much closer to the observed values, primarily because the southern ITCZ in the Pacific and Atlantic sectors is reduced.

#### 4.1.3 Seasonal cycle of precipitation over eastern Pacific

Figure 6 presents the seasonal evolution of monthly precipitation averaged between 90°W and 160°W from the GPCP observations and coupled simulations. In the observation (Figure 6a), the precipitation greater than 3 mm day<sup>-1</sup> occurs only between March and April south of the equator, whereas the precipitation greater than 3 mm day<sup>-1</sup> persists throughout the year in the northern ITCZ region with higher precipitation rates occurring between April and November. The double ITCZ bias in REF\_cmip comes mainly from overestimated precipitation as high as 12 mm day<sup>-1</sup> between 5°S and 10°S in boreal winter and spring (Figure 6b). In contrast, the excessive precipitation south of the equator in boreal winter and spring is reduced in NEW\_cmip and closer to the observation although the maximum rainfall is slightly larger in April (Figure 6c and 6d), indicating the alleviated double ITCZ bias. In addition, the northern ITCZ dry bias in the boreal winter and spring is also alleviated and the northern precipitation band in boreal summer and autumn moves closer to the equator in NEW\_cmip.

A southern ITCZ ( $SI$ ) index, simply defined as the annual mean precipitation rate over southeastern Pacific (20°S – 0°, 90° – 160°W; Bellucci et al., 2010), is used to quantify the coupled model biases in a more objective way. The simulated  $SI$  index decreases from 3.46 mm day<sup>-1</sup> in REF\_cmip to 2.51 mm day<sup>-1</sup> in NEW\_cmip, closer to the observational values (1.36



mm day<sup>-1</sup> in GPCP and 1.66 mm day<sup>-1</sup> in CMAP). The following analysis will focus on the Pacific ocean to understand the impact of improved boundary-layer turbulence and shallow convection schemes on the simulated ITCZ.

#### 4.2 SST and the heat flux into the ocean

310 The annual mean precipitation change is closely related to the SST change (Song and Zhang, 2016). Figure 7 shows the annual mean SST from the HadISST observation, the REF\_cmip and NEW\_cmip simulations, and their differences in the Pacific. The 30-year average of HadISST from 1971 to 2000 is used as the observed climatology, in which the warm SST corresponds to the northern ITCZ and SPCZ precipitation (Figure 7a). The SST from REF\_cmip, featured by a stronger cold tongue extending excessively westward along the equator, is warmer than the HadISST between 5°S and 20°S in the central and eastern Pacific. A band of cold SST bias down to -3 K between 5°S and 5°N across the Pacific can be clearly seen from the difference between them (Figure 7c). There is a conspicuous region of warm bias up to 2 K between 5°S and 20°S east of 150°W, which has been attributed to the underestimation of clouds in this region (e.g., Ma et al. 1996). It should be noted that this study compares the SST from pre-industrial simulations to the present-climate observations. Considering the global warming trend, the pre-industrial SST should be colder than the 20th century observations, suggesting that the cold bias in cold tongue region may be overestimated and the warm bias in the central and eastern Pacific may be underestimated. When the new boundary-layer turbulence and shallow convection schemes are used, the simulated warm water is cooled down by up to 4 K in the SEP relative to the REF\_cmip (Figure 7d).

The SST differences between the NEW\_cmip and REF\_cmip simulations are determined by both the net surface heat flux difference and ocean dynamic heat transport difference. The differences between NEW\_cmip and REF\_cmip runs for the atmospheric forcing on SST via surface heat flux  $\Delta Q_{atm}$  is composed of differences of shortwave radiation  $\Delta Q_{SW}$ , longwave radiation  $\Delta Q_{LW}$ , latent heat flux  $\Delta Q_{LH}$ , and sensible heat flux  $\Delta Q_{SH}$ :

$$\Delta Q_{atm} = \Delta Q_{SW} + \Delta Q_{LW} + \Delta Q_{LH} + \Delta Q_{SH} . \quad (8)$$

The contributions of these four components are presented in Figure 8. In the SEP, the net atmospheric heat flux at sea surface is reduced in NEW\_cmip, contributing to the cooled water in these regions. Among the atmospheric radiative heat flux components, the shortwave radiation is dominant, while the longwave radiation exerts a smaller and opposite effect. The latent heat flux is similar to the longwave radiative flux, and the sensible heat flux is negligible. Overall, the shortwave response explains most of net surface heat flux response in the SEP.

#### 4.3 The Walker circulation and the surface wind stress

Changes in the simulated SSTs lead to changes in the simulated large-scale atmospheric circulation. The JRA-55 reanalysis data shows descending motion east of 150°W, which corresponds to the rainless state in the eastern Pacific between 5°S and 10°S, and intense ascending motion west of the dateline, which corresponds to the notable precipitation in the SPCZ (Figure



9a). In REF\_cmip (Figure 9b), the weaker downward motion is limited to the middle and upper troposphere east of 125°W. The stronger upward motion occurs in the middle and upper troposphere between 150°W and 135°W, and is extended eastward to 100°W in the lower troposphere. Both the strengthened large-scale ascending motion and the weakened descending motion in the eastern part of the Walker circulation contribute to the excessive precipitation in the southern ITCZ region, resulting in the double ITCZ bias in REF\_cmip. When the new boundary-layer turbulence and shallow convection schemes are used, the descending branch of the Walker circulation is better simulated in NEW\_cmip (Figure 9c). Compared with REF\_cmip, the upward motion in the low troposphere shrinks to 120°W in NEW\_cmip. Furthermore, the downward motion in the middle and upper troposphere is enhanced and expanded westward to 140°W, making it in better agreement with the JRA-55 reanalysis. The improved descending branch of the Walker circulation is consistent with the decrease of precipitation relative to REF\_cmip.

Figure 10 shows the annual mean differences of surface wind stress vectors and magnitudes between the NEW\_cmip and REF\_cmip simulations. The modified boundary-layer turbulence and shallow convection schemes result in increased southeasterly winds off the west coast of South America. In addition, the difference between NEW\_cmip and REF\_cmip clearly shows the strengthened southeasterly trade winds in the eastern Pacific between 5°S and 10°S, corresponding to the stronger descending branch of the Walker circulation in NEW\_cmip.

#### 4.4 Enhanced cold advection in the upper ocean

The enhanced Walker circulation and surface wind stress subsequently induce strengthened cold advection in the upper ocean. Figure 11 shows the longitude-depth cross section of zonal ocean current and temperature averaged over 5°S – 10°S for the difference between NEW\_cmip and REF\_cmip. The enhanced westward ocean current over the whole zonal band helps transport cooler water from east to west in NEW\_cmip.

## 5 Discussion

### 5.1 The Sc-to-Cu transition

The relative contribution of the UWMT boundary-layer scheme and modified Hack shallow convection scheme in portraying the Sc-to-Cu transition is further examined. Figure 12 presents the vertical cross sections of cloud fraction through the five main subtropical stratocumulus regions in the sensitivity experiments of UWMT\_amip and mHack\_amip. UWMT\_amip reproduces the elevated cloudiness maxima, which can be expected from the applicability of the UWMT turbulence scheme in simulating stratocumulus, but still suffers from the rapid dissipating of cloud cover downstream off the coasts (Figure 12, left). On the contrary, mHack\_amip produces the gradual reduction of cloud amount offshore but misses the elevated cloudiness maxima, and constrains the cloud to a very low level (Figure 12, right). Neither UWMT\_amip nor mHack\_amip captures all the significant transition characteristics, indicating that the coupling between boundary-layer turbulence and shallow convection is important for an accurate simulation of the Sc-to-Cu transition.



In this study, introducing the decoupling criterion in the Hack scheme aims to better treat the transition between the stratocumulus regime and shallow cumulus regime. Regime-dependent parameterizations have limitations when representing various cloud regimes and their transitions. Two classes of unified parameterizations of boundary-layer turbulence and shallow convection have been documented in the literature, known as the eddy diffusion mass flux (EDMF) approach (Siebesma et al., 2007; Pergaud et al., 2009; Hourdin et al., 2013), and the higher-order turbulence closure (HOC) approach (Bogenschutz et al., 2013; Guo et al., 2010, 2014). Effects of these greater unification of cloud parameterizations on the double ITCZ bias in CGCMs will be further explored in future work.

### 5.2 Robustness of the alleviated double ITCZ

More investigations on the role of boundary-layer turbulence and shallow convection schemes in the double ITCZ formation in different CGCMs are desired. For simplicity, the UWMT and modified Hack schemes are employed in the high-resolution Beijing Climate Center Climate System Model version 2 (BCC-CSM2-HR), which is largely different from BCC-CSM2-MR with respect to model physics and dynamics, to examine the robustness of the alleviated double ITCZ through improving parameterizations of boundary-layer turbulence and shallow convection.

During the transition from BCC-CSM2-MR to BCC-CSM2-HR, the atmospheric component increased its horizontal resolution to T266 ( $\sim 0.45^\circ$ ) with a higher model top, and the physics package was essentially updated, especially the deep convection scheme. Furthermore, the oceanic component was upgraded to the Modular Ocean Model version 5 (MOM5). However, BCC-CSM2-HR suffered from the double ITCZ syndrome until the UWMT and modified Hack schemes were introduced. This suggests that the boundary-layer and shallow convection schemes contribute primarily to the double ITCZ bias in BCC-CSM2-HR. The tropical precipitation patterns simulated in BCC-CSM2-HR barely manifest a double ITCZ, as shown in Figure 13. The triangular-shaded dry region in the SEP reproduced by BCC-CSM2-HR resembles the observed much better than that simulated in BCC-CSM2-MR, probably due to the improved interactions among the boundary-layer turbulence, shallow convection, and other processes. Anyway, improving parameterizations of boundary-layer turbulence and shallow convection shows robustness in mitigating the double ITCZ syndrome in different models of BCC-CSM2-MR and BCC-CSM2-HR.

## 6 Summary and conclusions

It is a challenge to eliminate the double ITCZ problem, one of the most prominent systematic biases in CGCMs. This study investigates the roles of BL turbulence and shallow convection parameterizations in alleviating the double ITCZ bias using the BCC-CSM2-MR. The original BCC-CSM2-MR presents a serious double ITCZ problem in precipitation, with two significant zonal rain bands straddling the equator across the Pacific. In contrast, the revised BCC-CSM2-MR with new configurations of BL turbulence and shallow convection schemes remarkably alleviates the spurious southern ITCZ bias associated with SST warm bias in the SEP, owing to the reduced net surface shortwave radiation associated with the



increased low cloud fraction. Correspondingly, the cooler water in the SEP induces stronger and wider subsiding motion of  
400 the Walker circulation. The stronger Walker circulation and enhanced eastward surface wind stress in turn lead to increased  
oceanic zonal cold advection from east to west in the southern equatorial Pacific.

Consistent with previous studies (Ma et al., 1996; Yu and Mechoso, 1999; Dai et al., 2003, 2005; Qin and Lin, 2018),  
this study emphasizes the importance of stratus clouds and SSTs in the SEP in alleviating the double ITCZ bias via changes  
in the tropical circulation and ocean dynamics. However, different with studies with modifications of cloud scheme itself,  
405 this study focuses on the roles of BL turbulence and shallow convection in the low-level cloud simulation, and the  
subsequent impacts on the mitigation of the double ITCZ. It is found that the simulated low clouds are sensitive to both BL  
turbulence and shallow convection parameterizations. Better consistency between the BL turbulence scheme and the shallow  
convection scheme results in better simulation of the Sc-to-Cu transition. In conclusion, this study indicates that improving  
parameterizations of BL turbulence and shallow convection is an effective way to reduce the double ITCZ in CGCMs.

410 In the present study, the alleviation of the double ITCZ problem is accompanied by an amplification of the cold tongue  
bias, as found by many early efforts focusing on the role of low clouds. The inverse responses of the double ITCZ and cold  
tongue biases to southeast Pacific low clouds suggest that other parameterized processes, e.g., deep convection, may play an  
important role in the accurate simulation of the ITCZ-cold tongue complex. Song and Zhang (2018) demonstrated that the  
double ITCZ bias is largely eliminated and the cold bias in equatorial cold tongue is also significantly reduced through  
415 modifying the deep convection scheme in the Community Earth System Model version 1.2.1 (CESM1.2.1). The effect of  
improving the representation of deep convection on the double ITCZ bias in BCC-CSM2-MR is beyond the scope of this  
paper and will be explored in future work.

*Code and data availability.* Please contact the corresponding author if readers want to validate the model modifications and  
420 to conduct replication experiments. The source codes and required input data of the BCC models are freely available upon  
request from Tongwen Wu (twwu@cma.gov.cn).

*Author contributions.* YL modified the source codes, designed and performed all the experiments presented in the paper, and  
wrote a large part of the paper. TW supervised the BCC-CSM development and provided critical comments on the paper. All  
425 the authors continuously discussed the model development and the results.

*Competing interests.* The authors declare that they have no conflict of interest.

*Acknowledgements.* This work was supported by The National Key Research and Development Program of China (2016YFA  
430 0602100). We gratefully acknowledge the groups of GPCP (<https://www.esrl.noaa.gov/psd/data/gridded/data.gpcp.html>), C  
MAP (<https://www.esrl.noaa.gov/psd/data/gridded/data.cmap.html>), CALIPSO-GOCCP ([http://climserv.ipsl.polytechnique.f](http://climserv.ipsl.polytechnique.fr/cfmpip-obs)  
[r/cfmpip-obs](http://climserv.ipsl.polytechnique.fr/cfmpip-obs)), CERES-EBAF (<https://ceres.larc.nasa.gov/products.php?product=EBAF-Product>), HadISST (<https://www.met>



office.gov.uk/hadobs/hadisst), and the JRA-55 reanalysis (<https://jra.kishou.go.jp/JRA-55>) for providing public access to various observational data sets. All the graphics in this study are created by the NCAR Command Language (NCL; doi:10.5065/435 D6WD3XH5).

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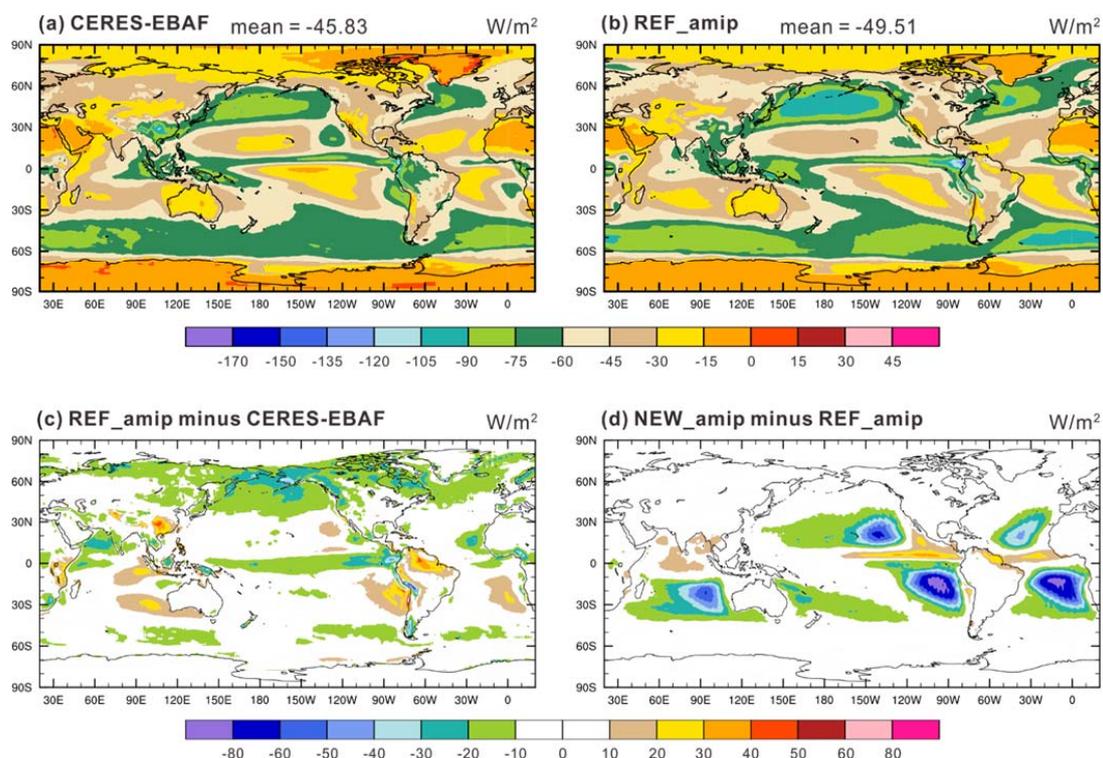
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**Table 1.** Summary of experimental setups.

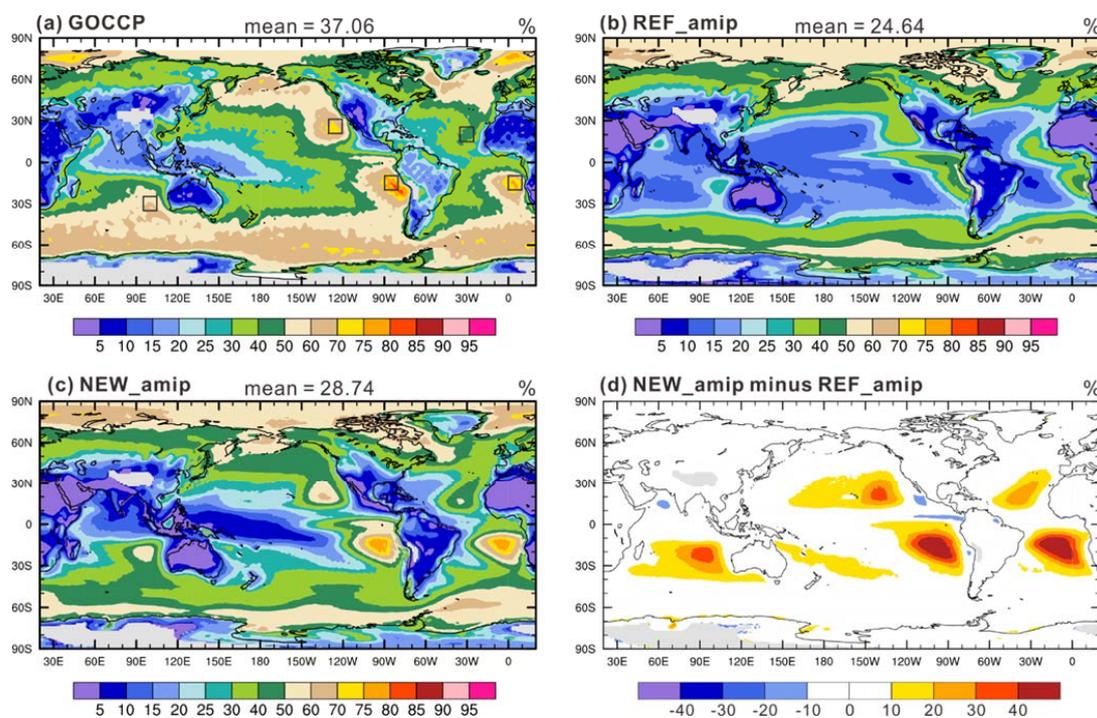
Run name	Boundary layer scheme	Shallow convection scheme
REF_amip	HB	Hack
NEW_amip	UWMT	modified Hack
REF_cmip	HB	Hack
NEW_cmip	UWMT	modified Hack
UWMT_amip	UWMT	Hack
mHack_amip	HB	modified Hack

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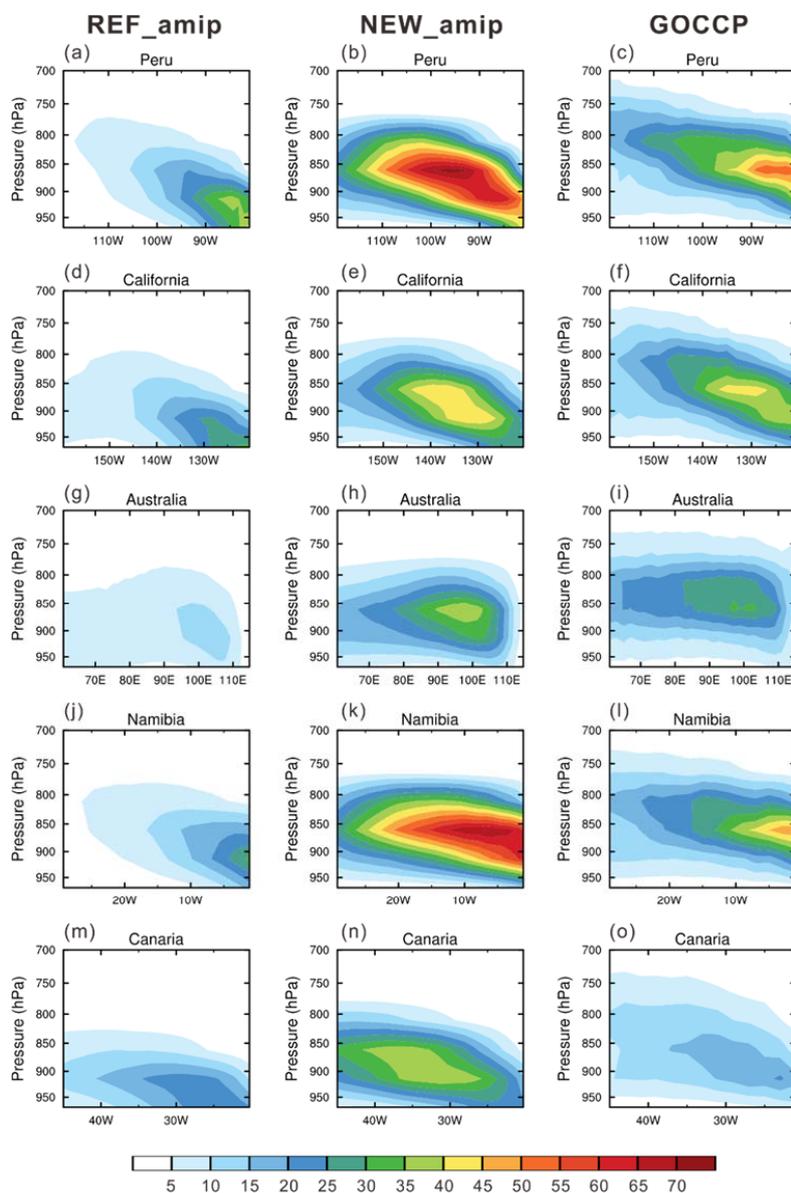


**Figure 1.** Annual mean climatologies of TOA SWCRF ( $\text{W m}^{-2}$ ) from (a) CERES-EBAF observations (from 2001 to 2010), (b) REF\_ampie, and the differences between (c) REF\_ampie and observations, (d) NEW\_ampie and REF\_ampie. Shown atop panels (a) and (b) are global mean values.

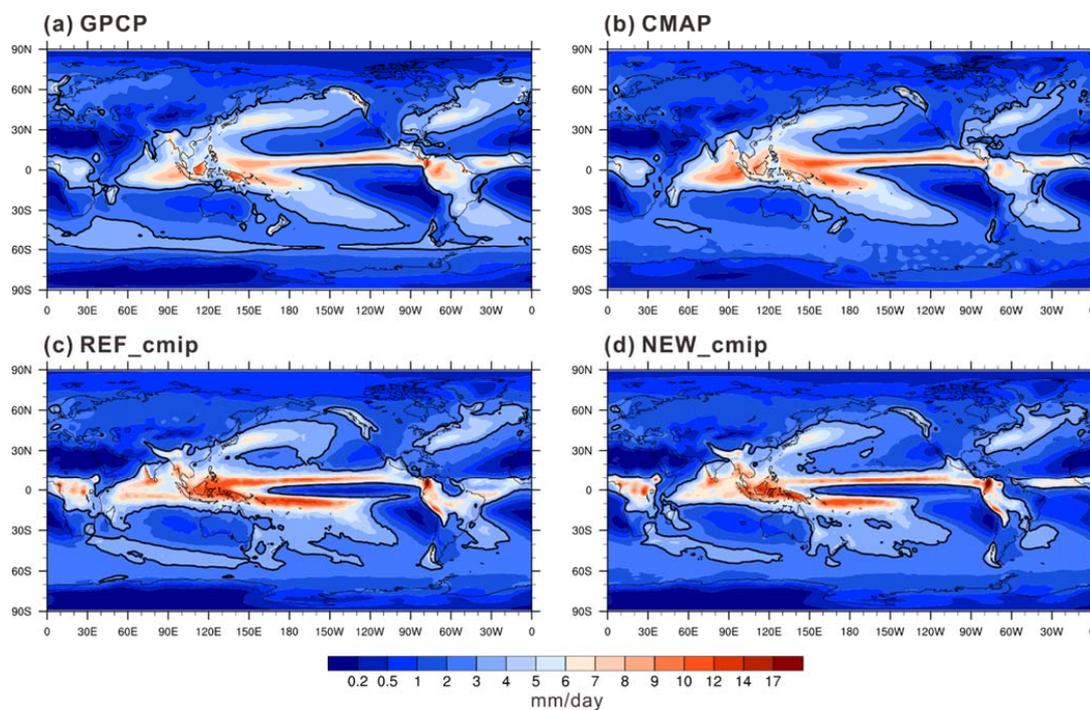
580



585 **Figure 2.** Annual mean low-level cloud fraction (%) from (a) GOCCP (from 2007 to 2012), (b) REF\_amip, (c) NEW\_amip, and (d) the difference between NEW\_amip and REF\_amip. Shown atop panels (a), (b) and (c) are global mean values. Boxes in panel (a) indicate the five main subtropical marine stratocumulus regions, namely, Peru ( $10^{\circ}$ - $20^{\circ}$ S,  $80^{\circ}$ - $90^{\circ}$ W), California ( $20^{\circ}$ - $30^{\circ}$ N,  $120^{\circ}$ - $130^{\circ}$ W), Australia ( $25^{\circ}$ - $35^{\circ}$ S,  $95^{\circ}$ - $105^{\circ}$ E), Namibia ( $10^{\circ}$ - $20^{\circ}$ S,  $0^{\circ}$ - $10^{\circ}$ E), and Canaria ( $15^{\circ}$ - $25^{\circ}$ N,  $25^{\circ}$ - $35^{\circ}$ W).

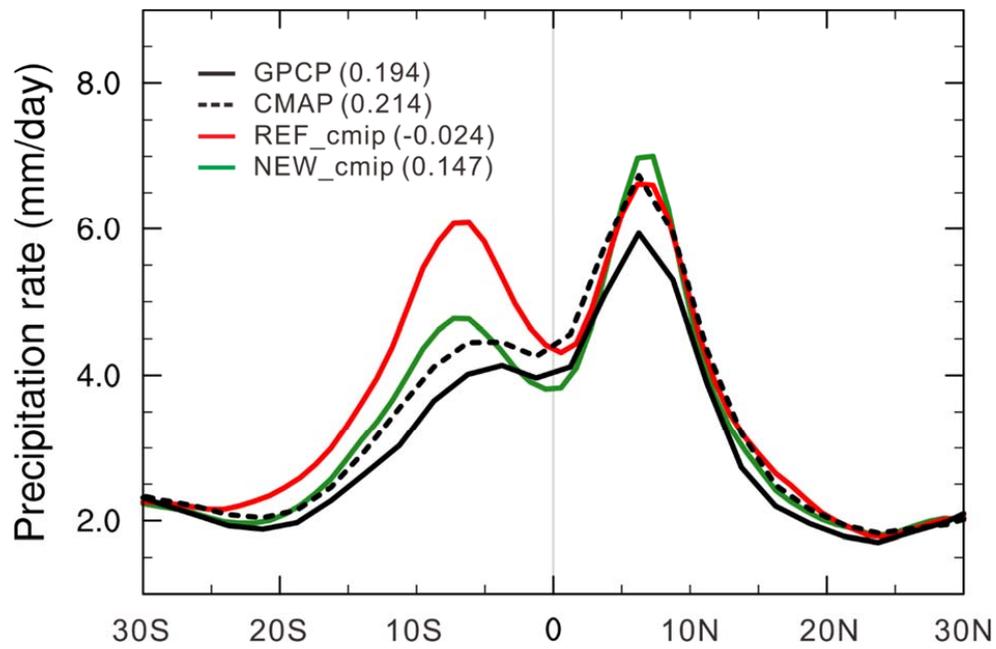


590 **Figure 3.** Cross sections of cloud fraction from five locations for (left) REF\_amip, (center) NEW\_amip, and (right) GOCCP. Refer to the boxes in Fig. 2 for the locations of the cross sections.

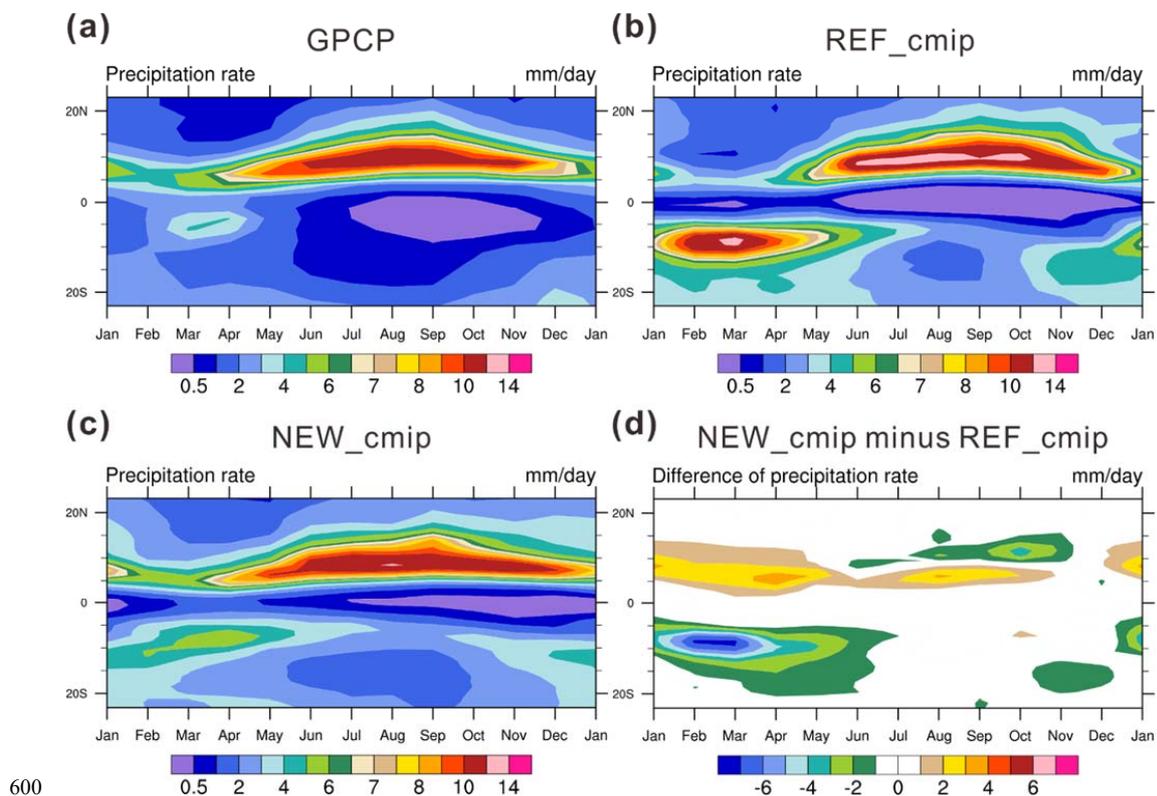


**Figure 4.** Annual mean precipitation rate ( $\text{mm day}^{-1}$ ) from (a) GPCP, (b) CMAP, (c) REF\_cmip, and (d) NEW\_cmip. The 3  $\text{mm day}^{-1}$  contour is included in bold for reference.

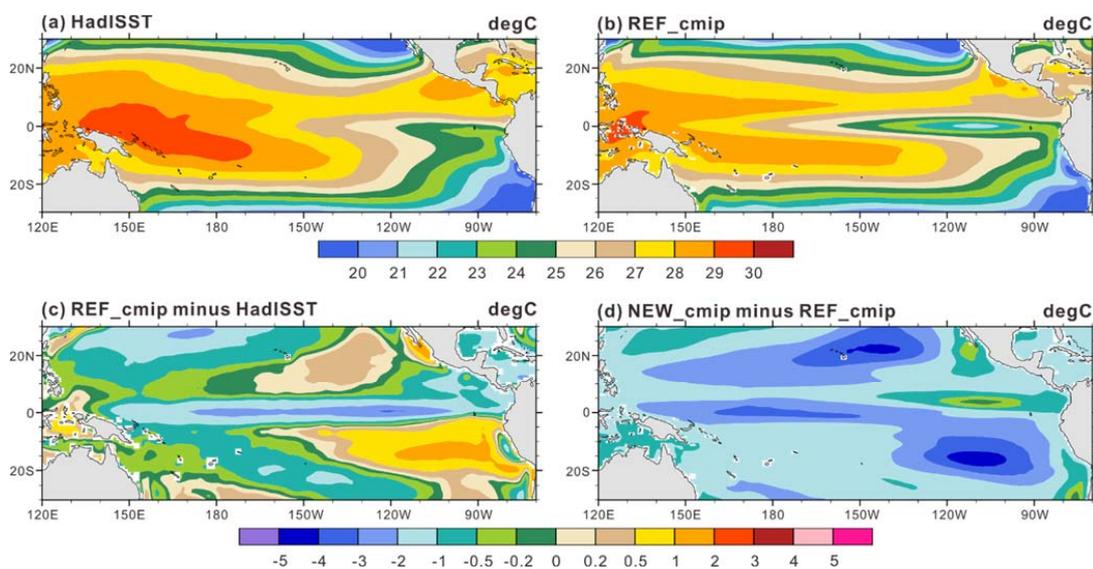
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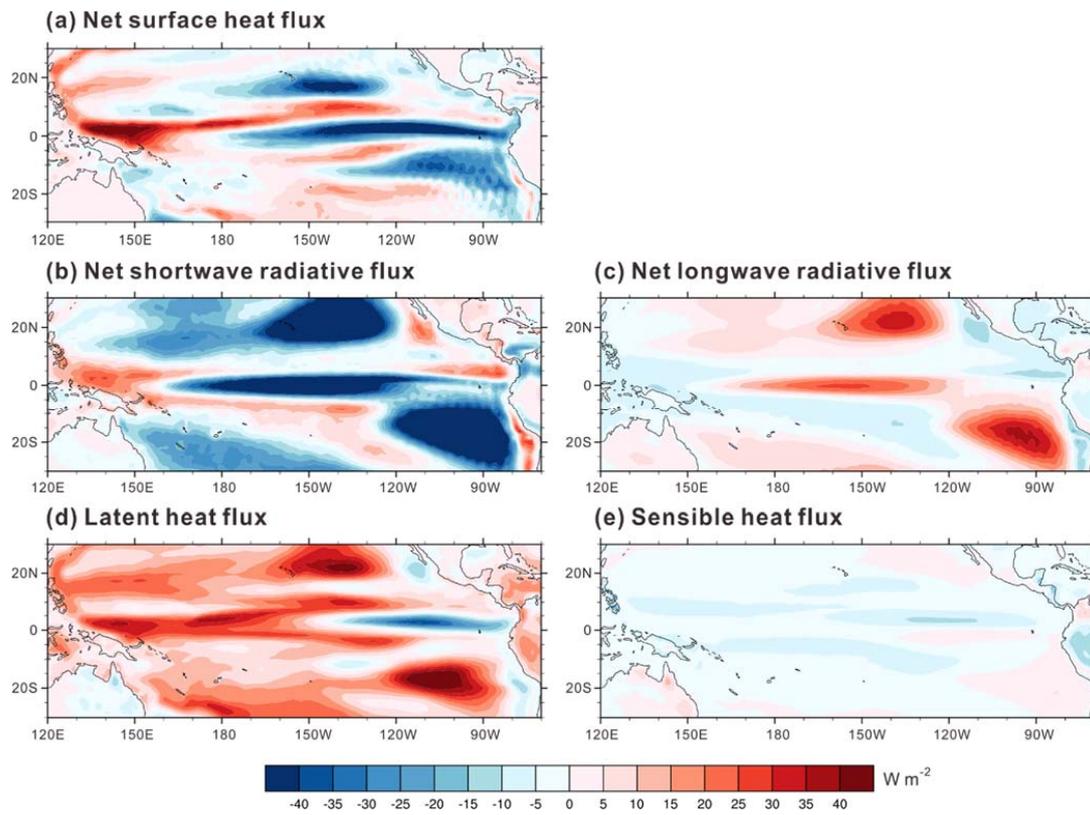
**Figure 5.** Zonal mean precipitation rate ( $\text{mm day}^{-1}$ ) from GPCP, CMAP, REF\_cmip and NEW\_cmip in the tropics. Values of the tropical precipitation asymmetry index are indicated in parentheses.



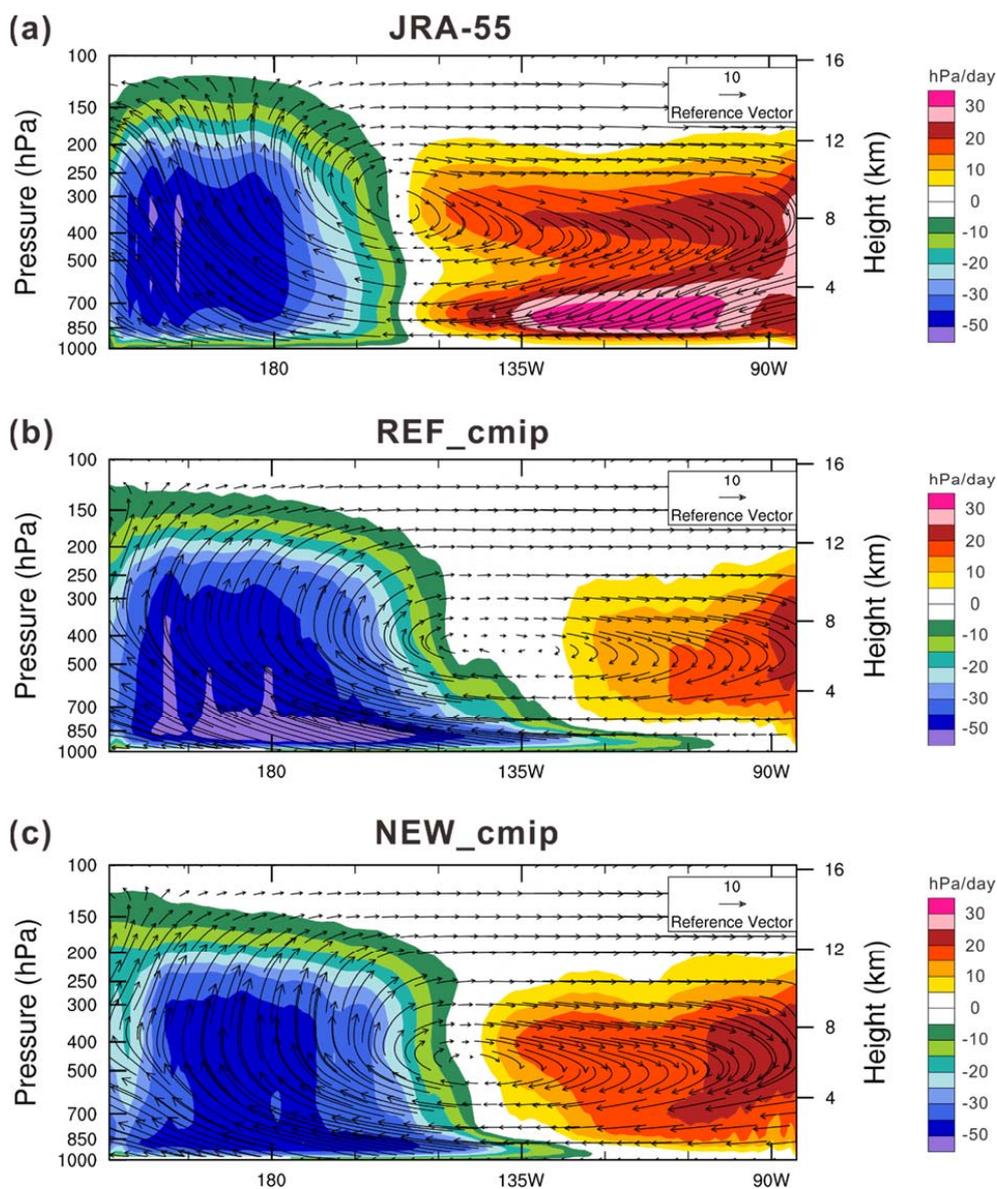
**Figure 6.** Seasonal cycle of precipitation rate ( $\text{mm day}^{-1}$ ) averaged over eastern Pacific ( $90^\circ - 160^\circ\text{W}$ ) for (a) GPCP, (b) REF\_cmip, (c) NEW\_cmip, and (d) the difference between NEW\_cmip and REF\_cmip.



605 **Figure 7.** Annual mean sea surface temperature (°C) from (a) HadISST, and (b) REF\_cmip, and the difference between (c) REF\_cmip and HadISST, and (d) NEW\_cmip and REF\_cmip.

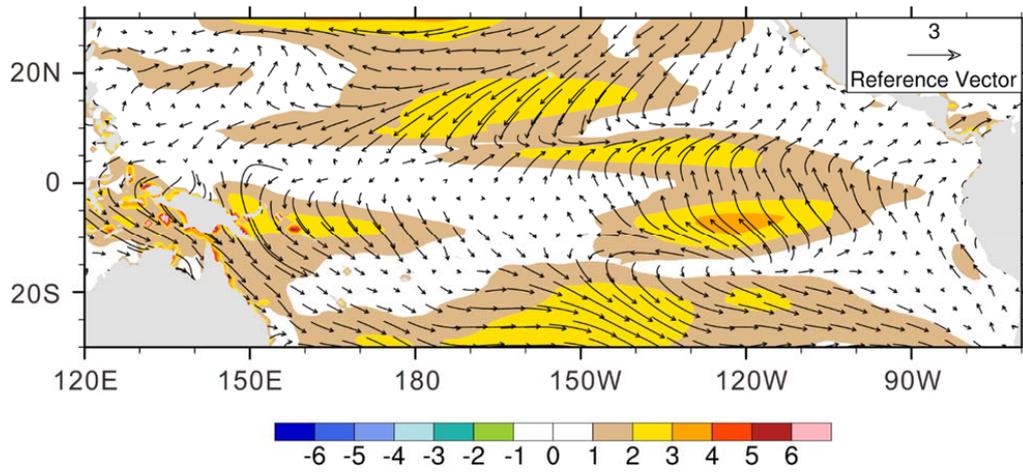


**Figure 8.** The differences between the NEW\_cmip and REF\_cmip simulations (W m<sup>-2</sup>) for (a) net surface heat flux  $\Delta Q_{atm}$ , (b) net  
610 shortwave radiative flux  $\Delta Q_{SW}$ , (c) net longwave radiative flux  $\Delta Q_{LW}$ , (d) latent heat flux  $\Delta Q_{LH}$ , and (e) sensible heat flux  $\Delta Q_{SH}$ .

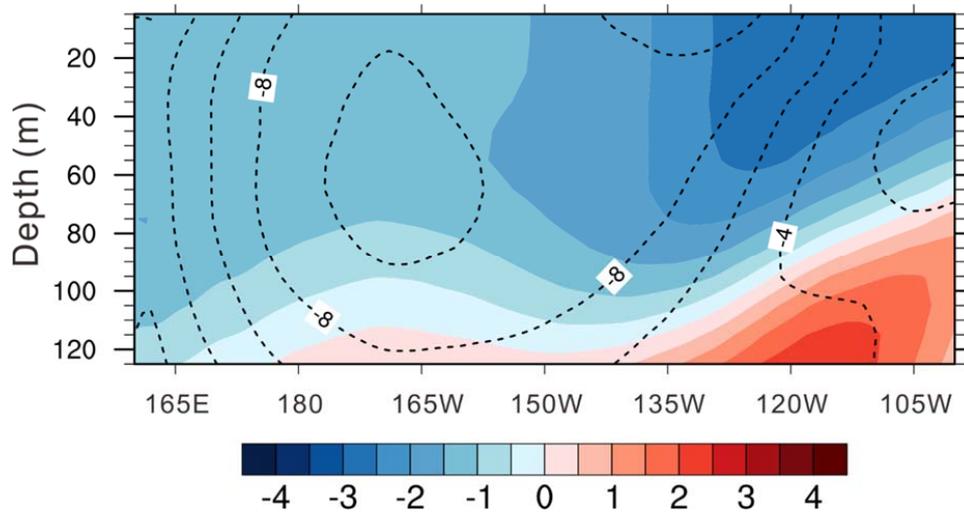


**Figure 9.** Annual mean vertical pressure velocity ( $\text{hPa day}^{-1}$ ; shaded) and wind vectors (arrows) in the longitude-height cross section averaged over  $5^{\circ}\text{S} - 10^{\circ}\text{S}$  for (a) the JRA-55 reanalysis, (b) REF\_cmip, and (c) NEW\_cmip.

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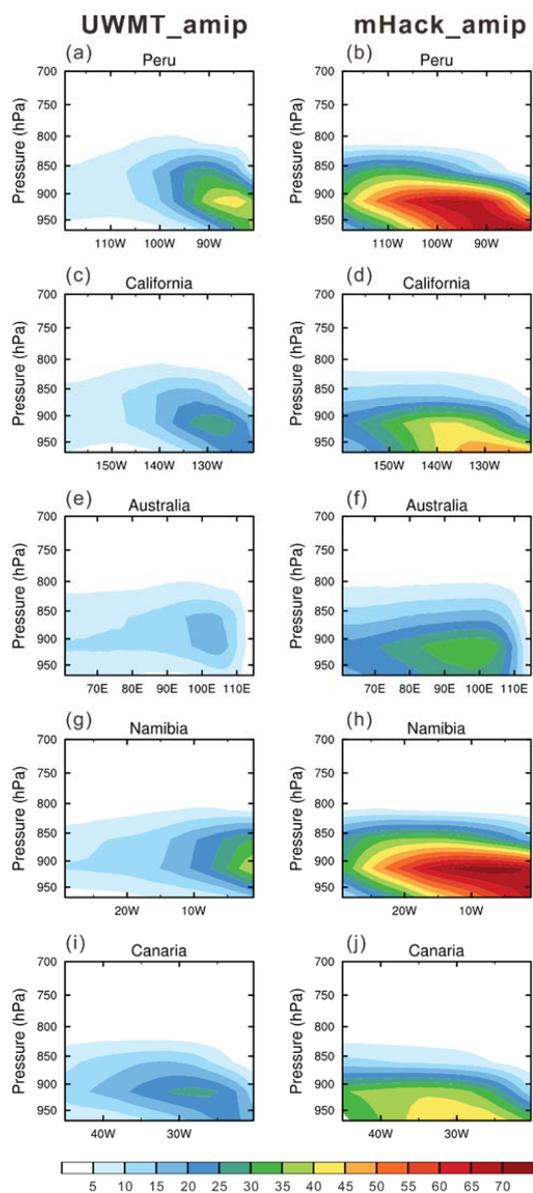


**Figure 10.** The difference of annual mean wind stress vector and wind stress magnitude (shaded,  $10^{-2} \text{ N m}^{-2}$ ) between NEW\_cmp and REF\_cmp.

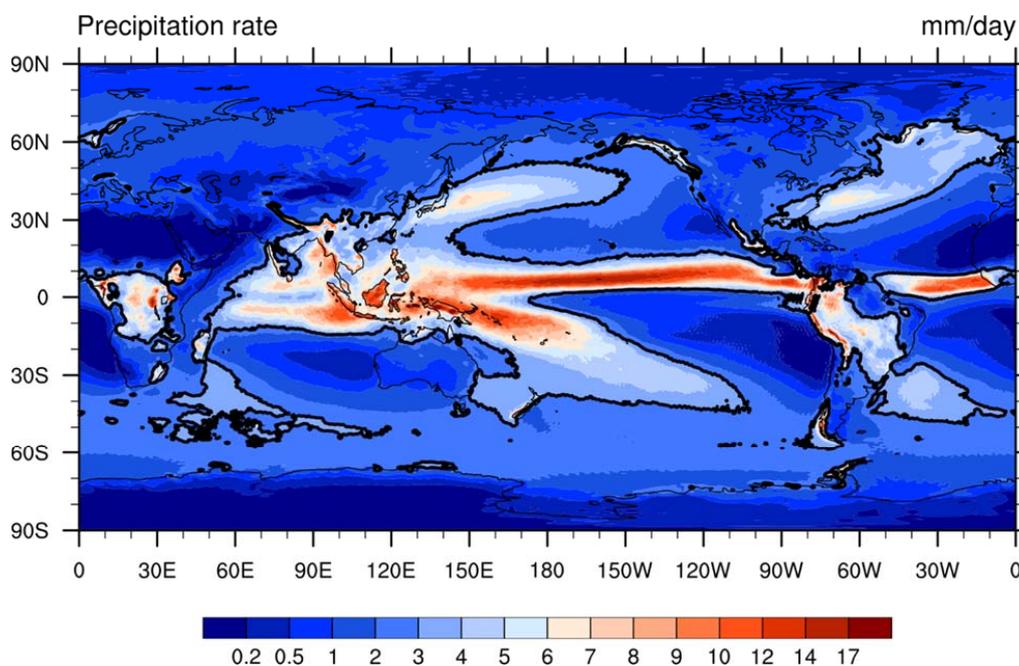


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**Figure 11.** Longitude-depth cross section of zonal ocean current (contours;  $\text{cm s}^{-1}$ ) and temperature (shaded;  $^{\circ}\text{C}$ ) averaged over over  $5^{\circ}\text{S} - 10^{\circ}\text{S}$  for the difference between NEW\_cmip and REF\_cmip.



625 **Figure 12.** Cross sections of cloud fraction from five locations for (left) UWMT\_ami, and (right) mHack\_ami. Refer to the boxes in Fig. 2 for the locations of the cross sections.



630 **Figure 13.** Annual mean precipitation rate ( $\text{mm day}^{-1}$ ) from BCC-CSM2-HR. The  $3 \text{ mm day}^{-1}$  contour is included in bold for reference.