

Dear Editor and Referee #2,

We thank you for the time spent to evaluate and review our manuscript and for the supportive recommendations.

The major and minor remarks brought by Referee #2 were all addressed in the revised manuscript. In our point-by-point answers below, we explain how and where each point raised was addressed in the revised manuscript.

Answers are in blue text and appear after each reviewer comment, followed by

a box showing the associated changes applied to the revised manuscript.

Sincerely,

Georges Kesserwani and James Shaw

Major remarks

“The inclusion of a second-order DG2 discretisation and implementation in the LISFLOOD-FP 8.0 flood forecasting suite, on a regular mesh, of the shallow water equations with friction and rainfall, for flood forecasting, is presented, including analyses of three test cases and comparison with a local inertia solver ACC and a first-order FV1 discretisation and implementation. Parallelisation and GPU results are intercompared in terms of performance and runtime. While the paper is very interesting and the work constitutes a useful addition in a flood forecasting suite for public use, which is very important, there seem to be several loose ends and unclear aspects that need to be resolved before publication is warranted. I therefore recommend to return the manuscript for major revisions.”

Thank you for highlighting the interest and importance of this paper, and for the useful comments and recommendations. They were very helpful to clarify the presentation of the DG2 solver and the discussions of its performance.

“It would be desirable that the following major issues are addressed, either clarified or resolved:”

Explanations on how these issues were resolved or clarified are explained below each comment.

“(i) Since ACC, FV1 and DG2 schemes are tested and compared more clarity on their characteristics is desired in order to understand the differences in the results presented. Why is DG2 much slower than FV1 in the third test? The time stepping schemes of FV1 and DG2 are unclear. Figuring out the split scheme for DG2 given rainfall and damping is too hard or impossible also, since the references are unclear; it is easier to simply state the full scheme; there is sufficient space in (9) to do so and remove any ambiguity. The split scheme for Sf is not implicit, once one looks into the references. Why is the handling of friction (Sf) in DG2 leading to troubles for thin, fast overland flow while it is not troubling for FV1 and ACC? Perhaps use the same time stepping scheme in DG2 as for FV1 or make aspects of Sf semi-implicit. Please clarify and make improvements”

Section 2, mostly subsection 2.1, has been substantially revised to include the technical details of the DG2 and FV1 solvers including their time stepping schemes for the splitting

friction integration and the discretisation of the rainfall source term. In the new subsection 2.1.2, the split friction scheme with DG2/FV1 is explicitly presented. The subsection also includes an explanation on why the handling of friction with DG2 can reduce the time-step size for thin, fast overland flow compared to FV1.

2.1.2 Discretisation of the friction source term

The discretisation of the friction source term is ~~discretised using a~~ based on the split implicit scheme (Liang and Marche, 2009; Kesserwani and Sharifian, 2020), and the of Liang and Marche, 2009. Without numerical stabilisation, the friction function $C_f = gn_M^2/h^{1/3}$ can grow exponentially as the water depth vanishes at a wet-dry front, but the scheme adopted here is designed to ensure numerical stability by limiting the frictional force to prevent unphysical flow reversal.

The implicit friction scheme is solved directly (see Sect. 3.4, Liang and Marche 2009) such that frictional forces are applied to the x -directional discharge component q_x over a time-step Δt , yielding a retarded discharge component q_{fx} :

$$q_{fx}(\mathbf{U}) = q_x + \Delta t \frac{S_{fx}}{\mathcal{D}_x}, \quad (11a)$$

where the denominator \mathcal{D}_x is

$$\mathcal{D}_x = 1 + \left(\frac{\Delta t C_f}{h} \right) \left(\frac{2u^2 + v^2}{\sqrt{u^2 + v^2}} \right). \quad (11b)$$

To update the element-average discharge coefficient $q_{x_{i,j,0}}$, Eqn. (11) is evaluated at the element centre :

$$q_{x_{i,j,0}}^{n+1} = q_{fx}(\mathbf{U}_{i,j,0}^n), \quad (12a)$$

while the slope coefficients $q_{x_{i,j,1x}}$ and $q_{x_{i,j,1y}}$ are updated by calculating the x - and y -gradients using evaluations of Eqn. (11) at Gaussian quadrature points Gx1, Gx2, and Gy1, Gy2 (Fig. 1):

$$q_{x_{i,j,1x}}^{n+1} = \frac{1}{2} [q_{fx}(\mathbf{U}_{i,j}^{\text{Gx2}}) - q_{fx}(\mathbf{U}_{i,j}^{\text{Gx1}})], \quad (12b)$$

$$q_{x_{i,j,1y}}^{n+1} = \frac{1}{2} [q_{fx}(\mathbf{U}_{i,j}^{\text{Gy2}}) - q_{fx}(\mathbf{U}_{i,j}^{\text{Gy1}})]. \quad (12c)$$

Similarly, frictional forces are applied to the y -directional discharge component q_y yielding a retarded discharge q_{fy} :

$$q_{fy}(\mathbf{U}) = q_y + \Delta t \frac{S_{fy}}{\mathcal{D}_y}, \quad (13a)$$

$$\mathcal{D}_y = 1 + \left(\frac{\Delta t C_f}{h} \right) \left(\frac{u^2 + 2v^2}{\sqrt{u^2 + v^2}} \right). \quad (13b)$$

While this friction scheme has been successfully adopted in finite-volume and discontinuous Galerkin settings for modelling dam break flows and urban flood events (Want et al., 2011; Kesserwani and Wang, 2014), it can exhibit spuriously large velocities and correspondingly small time-steps for large-scale, rainfall-induced overland flows, involving widespread, very thin water layers flowing down hill slopes and over steep river banks, as demonstrated by Xia et al., 2017. Due to the involvement of the slope coefficients, water depths at Gaussian quadrature points can be much smaller (and velocities much larger) than the element-average values. Therefore, for overland flow simulations, the LISFLOOD-DG2 time-step size is expected to be substantially reduced compared to LISFLOOD-FV1, which only involves element-average values.

“(ii) In test 3, but also in earlier tests, DG2 has topography at $dx=40$ (or another dx for tests 1 and 2) and seemingly the higher degrees of freedom used for the variables is not used for the topography, i.e. the $dx=10$ topography information can be used to make finer projections onto the topography given that DG2 is second order and the topography should not be planar. Hence, the (at times and in certain simulations) observed simulation underperformance of DG2 seems to be/is more severe than needed.”

It is important to clarify that the topography with the proposed DG2 solver has been, and must be, projected using three degrees of freedom (Kesserwani et al. 2018, Kesserwani and Sharifian 2020, Ayog et al. 2021). This is what we referred to as “piecewise-planar” topography as its projection is expanded based on three degrees of freedom reconstructed from four DEM data evaluation at the four vertices (as opposed to with ACC/FV1 that uses “piecewise-constant” topography, extracted from a single evaluation from the DEM data). To avoid this misunderstanding, Section 2, mostly subsection 2.1 and figure 1, has been revised to describe how the DG2 topography projection was generated for the DG2 solver based on three degrees of freedom.

2.1.1 Initialisation of piecewise-planar topography coefficients from a DEM raster file

The topography coefficients $[z_{i,j,0}, z_{i,j,1x}, z_{i,j,1y}]$ are initialised to ensure the resulting piecewise-planar topography is continuous at face centres, where Riemann fluxes are calculated and the wetting-and-drying treatment is applied under the well-balancedness property (Kesserwani et al., 2018). The topographic elevations at the N, S, E, and W face centres are calculated by averaging the DEM raster values taken at the NW, NE, SW and SE vertices (Fig. 1) such that $z_{i,j}^N = (z_{i,j}^{NW} + z_{i,j}^{NE})/2$ and similarly for $z_{i,j}^E$, $z_{i,j}^S$, and $z_{i,j}^W$. The element-average coefficient $z_{i,j,0}$ is then calculated as:

$$z_{i,j,0} = \frac{1}{4} [z_{i,j}^{NW} + z_{i,j}^{SW} + z_{i,j}^{NE} + z_{i,j}^{SE}], \quad (14a)$$

while the slope coefficients $z_{i,j,1x}$ and $z_{i,j,1y}$ are calculated as the gradients across opposing face centres:

$$z_{i,j,1x} = \frac{1}{2\sqrt{3}} (z_{i,j}^E - z_{i,j}^W), \quad (14b)$$

$$z_{i,j,1y} = \frac{1}{2\sqrt{3}} (z_{i,j}^N - z_{i,j}^S). \quad (14c)$$

LISFLOOD-FP 8.0 includes a utility application, `generateDG2DEM`, that loads an existing DEM raster file and outputs new raster files containing the element-average, x -slope and y -slope topography coefficients, ready to be loaded by the LISFLOOD-DG2 solver.

“(iii) DG2 is used on a regular mesh; DG is most optimal for non-uniform situations but this option or restriction is not discussed; also not discussed are hybrid 1D and 2D schemes or hybrid FV1 and DG2 options, the latter which should be easy to accommodate. At least a discussion is warranted.”

Discussions have been added in the revised paper to highlight potential alternatives to DG2 on uniform grids, including the consideration of a non-uniform DG2 solver, linkage with the existing sub-grid channel model on LISFLOOD-FP and a hybrid FV1/DG2.

3.3.1 River channel free-surface elevation hydrographs

Spatially-adaptive solvers (Kesserwani and Sharifian, 2020; Özgen-Xian et al., 2020) and non-uniform meshing techniques (Kolega and Syme, 2019) offer another alternative to improve flow predictions by selectively capturing fine-scale channel geometries, and such methods are under development for inclusion in a future LISFLOOD-FP release.

3.3.2 Maximum flood extent over Carlisle

The representation of these flood defences could be improved by adopting the recently-developed LISFLOOD-FP levee module^a (Wing et al., 2019; Shustikova et al., 2020), or by implementing a spatially-adaptive multi-resolution method that selectively refines the grid resolution around river channels and other fine-scale features (Kesserwani and Sharifian, 2020).

^aNot yet available with the FV1 or DG2 solvers.

4 Summary and conclusions

However, FV1 and DG2 are the first solvers in LISFLOOD-FP to gain a dynamic rain-on-grid capability, with this capability being added to the optimised ACC solver in a future release. To further improve efficiency and accuracy at coarse resolutions over large catchments, one future direction would be to port the sub-grid channel model—currently integrated with the CPU-optimised ACC solver—to GPU architectures. Another useful direction would be to enable a multi-resolution solver based on Kesserwani and Sharifian, 2020, and introduce a hybrid DG2/FV1 solver that downgraded the DG2 formulation to FV1 in regions of very thin water layer, or in regions of finest grid resolution, to further reduce the computational cost. Both directions are being investigated for inclusion in future LISFLOOD-FP releases.

“(iv) Some of the speed and/or accuracy comparisons seem incomplete or unfair: e.g., in Fig. 6, not the same resolution should (only) be compared but (also) the speed for mixed resolutions with (roughly) the same accuracy”

Test 1 and Test 2 (Subsection 3.1 and 3.2) have been extended to also compare speed-ups among the solvers for mixed resolutions leading to outputs with (roughly) the same accuracy (see revised Fig. 7, revised Fig. 12, revised Table 1, and the changes made listed below).

3.1.3 Solver runtimes for a varying number of elements

As seen earlier in the inset panel of Fig. 7, similar wave-fronts were predicted by DG2 at $\Delta x = 5$ m, ACC at $\Delta x = 2$ m, and FV1 at $\Delta x = 0.5$ m. At these resolutions, DG2-CPU, DG2-GPU and ACC achieved a similar solution quality for a similar runtime cost, with all solvers completing in about 4 minutes (Fig. 8a). Meanwhile, the DG2 solvers on a ten-times coarser grid were $140\times$ faster than FV1-CPU (10 hours 42 minutes) and $28\times$ faster than FV1-GPU (1 hour 47 minutes).

3.1.4 Multi-core CPU scalability

The FV1, ACC and DG2 solvers are at least 5 solvers converged on similar water depth solutions with successive grid refinement. Owing to its first-order accuracy, FV1 requires a very fine resolution grid to match the solution quality of DG2 or ACC, though FV1-GPU enables runtimes up to $6\times$ more costly than ACC at the same grid resolution but, thanks faster than the 16-core FV1-CPU solver. Thanks to its second-order accuracy, DG2 water depth predictions are spatially converged at coarser resolutions (Fig. 7). Hence, DG2 is able to replicate the modelling quality of FV1 at a much coarser resolution, and the multi-core DG2-CPU solver is more efficient on a competitive choice for grids with fewer than 100,000 (10^5) elements, while DG2-GPU is more efficient on grids with a higher number of elements. Both DG2-CPU and FV1-CPU solvers scale efficiently up to 16 CPU cores. elements.

Table 1

Solver runtimes at the standard resolution grid spacings of $\Delta x = 50$ m, $\Delta x = 20$ m, and the finest resolution of $\Delta x = 10$ m. ACC, FV1-CPU and DG2-CPU solvers are run on a 16-core CPU; FV1-GPU and DG2-GPU solvers are run on a single GPU. ACC solver runtimes were obtained for the ACC implementation of Neal et al., 2012a.

	$\Delta x = 50$ m 57000 57 000 elements	$\Delta x = 20$ m 850 000 elements	$\Delta x = 10$ m 1.7 million elements
ACC	20 s	466 s (8 mins)	1779 s (30 mins)
FV1-CPU	22 s	739 s (12 mins)	2188 s (36 mins)
FV1-GPU	19 s	145 s (2 mins)	965 s (16 mins)
DG2-CPU	788 s (13 mins)	4133 s (69 mins)	33009 s (9 hours)
DG2-GPU	448 s (7 mins)	2304 s (38 mins)	13606 s (4 hours)

3.2.3 Runtime cost

DG2-CPU and DG2-GPU at $\Delta x = 50$ m outperform ACC, FV1-CPU and FV1-GPU at ~~the standard resolution, becoming $2.5\Delta x = 10$ m,~~ while still achieving similarly accurate flood map predictions at a $5\times$ ~~faster~~ coarser resolution (Fig. 12). DG2-CPU at ~~the finest resolution $\Delta x = 50$ m~~ is $2\times$ faster than ACC at $\Delta x = 10$ m, while DG2-GPU is twice as fast again. DG2-GPU flood maps at an improved resolution of $\Delta x = 20$ m are obtained at a runtime cost of 38 mins, which is still competitive with ACC at $\Delta x = 10$ m (with a runtime cost of 30 mins).

In summary, all solvers predicted similar water depth and velocity hydrographs, though ACC experienced a short period of numerical instability in a localised region where the Froude number exceeded the limit of the local inertia equations. The shock-capturing FV1 and DG2 shallow water solvers yield robust predictions throughout the entire simulation, with FV1-GPU being consistently faster than ACC on a 16-core CPU. As found earlier in Sect. 3.1.3, ~~GPU parallelisation became more efficient as the total number of elements was increased~~ DG2 at a $2-5\times$ coarser resolution is a competitive alternative to ACC and FV1, with the GPU implementation being preferable when running DG2 on a grid with more than 100,000 elements.

“(v) There is a large accumulation of minor remarks, see below, which should be refuted and addressed. The above major remarks are also reflected, sometimes on several occasions, in several of the minor/detailed remarks given below.”

Description of how the minor/detailed remarks have been addressed is provided below.

Minor/detailed remarks

“Line 23: simplify (plural)”

Revised to resolve any ambiguity in ACC and ATS simplifications.

1 Introduction

LISFLOOD-FP already includes a ~~diffusive wave~~ local inertia (or ‘~~zero-inertia~~ gravity wave’) solver, ~~LISFLOOD-ATS~~ LISFLOOD-ACC, and a ~~local-inertia~~ diffusive wave (or ‘~~gravity wave~~ zero-inertia’) solver, LISFLOOD-ATS. The LISFLOOD-ACC ~~, that solver~~ simplifies the full shallow water equations by neglecting convective acceleration, while LISFLOOD-ATS neglects both convective and inertial acceleration.

“Line 42: But there are plenty of FV2 solvers, which may be faster; has a comparison with those been made or at least discussed? Please discuss.”

A discussion has been added to the introduction section.

1 Introduction

Second-order finite volume (FV2) methods offer an alternative approach to obtain second-order accuracy, with many FV2 models adopting the Monotonic Upstream-centred Scheme for Conservation Laws (MUSCL) method. While FV2-MUSCL solvers can achieve second-order convergence (Kesserwani and Wang, 2014), the MUSCL method relies on global slope limiting and non-local, linear reconstructions across neighbouring elements that can affect energy conservation properties (Ayog et al., 2021) and affect wave arrival times when the grid is too coarse (Kesserwani and Wang, 2014). Hence, although FV2-MUSCL is typically 2–10× faster than DG2 per element (Ayog et al., 2021), DG2 can improve accuracy and conservation properties on coarse grids, which is particularly desirable for efficient, long-duration continental- or global-scale simulations that rely on DEM products derived from satellite data (Bates 2012, Yamazaki et al., 2019).

“Line 56: What about mixed 1D and 2D solvers, see the literature?”

On this aspect, a sentence has been added to only discuss the literature specific to LISFLOOD-FP.

1 Introduction

LISFLOOD-FP includes extension modules to provide efficient rainfall routing (Sampson et al., 2013), modelling of hydraulic structures (Wing et al., 2019; Shustikova et al., 2020), and coupling between two-dimensional flood-plain solvers and a one-dimensional sub-grid channel model (Neal et al., 2021a)

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“Line 58: This paper [use of ”remainder” is a bit much on page 2].” Amended.

1 Introduction

The ~~remainder of this~~ paper is structured as follows: ...

“Line 72: What is new or simplified relative to K2018, and why?” Addressed by adding the following introductory text at the beginning of Sect. 2.1.

2.1 The new LISFLOOD-DG2 solver

The LISFLOOD-DG2 solver implements the DG2 formulation of Kesserwani et al., 2018 that adopts a simplified ‘slope-decoupled’ stencil compatible with raster-based Godunov-type finite volume solvers. Piecewise-planar topography, water depth and discharge fields are modelled by an element-average coefficient and dimensionally-independent x -slope and y -slope coefficients. This DG2 formulation achieves well-balancedness for all discharge coefficients in the presence of irregular, piecewise-planar topography with wetting-and-drying (Kesserwani et al., 2018). A piecewise-planar treatment of the friction term is applied to all discharge coefficients prior to each time-step, based on the split implicit friction scheme of Liang and Marche, 2009. Informed by the findings of Ayog et al., 2021, the automatic local slope limiter option in LISFLOOD-DG2 is deactivated for the flood-like test cases presented in Sect. 3. This slope-decoupled, no-limiter approach can achieve a $5\times$ speed-up over a standard tensor-product stencil with local slope limiting (Kesserwani et al., 2018; Ayog et al., 2021), meaning this DG2 formulation is expected to be particularly efficient for flood modelling applications.

“Please specify Eqn (2): define x , y , t .”

These are now defined after they first appear in the partial derivatives in Eqn (1):

2.1 The new LISFLOOD-DG2 solver

The DG2 formulation (Kesserwani et al., 2018) discretises the two-dimensional shallow water equations, written in conservative vectorial form as

$$\partial_t \mathbf{U} + \partial_x \mathbf{F}(\mathbf{U}) + \partial_y \mathbf{G}(\mathbf{U}) = \mathbf{S}_b(\mathbf{U}) + \mathbf{S}_f(\mathbf{U}) + \mathbf{R}, \quad (15)$$

where ∂_t , ∂_x and ∂_y denote partial derivatives in the horizontal spatial dimensions x and y , and temporal dimension t .

“Line 80: define L and T .”

Line 83: I don’t like calling C_f a coefficient since it includes a dynamic variable, i.e. h ; rephrase; it is a friction function. I would take out $h^{1/3}$ (of C_f) and even write it is h^α since α is a parameterization anyway.”

L and T were defined in the revised paper, and C_f is referred to as “friction function”. Apart from these requested revisions, the friction source term is written as it appears in Liang and Marche, 2009.

2.1 The new LISFLOOD-DG2 solver

Units are notated in square brackets [·], where L denotes unit length and T denotes unit time. The two-dimensional topographic elevation data is denoted z [L] and g is the gravitational acceleration [L/T²]. The frictional forces in the x - and y -directions are $S_{fx} = -C_f u \sqrt{u^2 + v^2}$ and $S_{fy} = -C_f v \sqrt{u^2 + v^2}$ respectively, where the ~~roughness coefficient is $C_f = gn_M^2/h^{1/3}$~~ , friction function is $C_f = gn_M^2/h^{1/3}$ and $n_M(x, y)$ is Manning's coefficient [T/L^{1/3}].

“Line 90: Why is this a dot product? No dot; it is a matrix-vector product.”

This ~~was removed~~.

“Line 119: there should be a flag doing that (application of a slope limiter) automatically for smooth solutions, not a hand-switch. See Krivodonova or extensions for improvements. In any real flooding case, this should be done via an automatic switch. Please clarify and add a sensible automatised switch. That the test cases do not involve shock wave propagation, is that not an omission? Tests 2 and 3 must have seen some shock-wave propagation prior to settling into a quasi-steady state, given the $F > 1$, $F < 1$ transition observed and discussed. Please clarify”

The local limiter is actually automatic, when activated, but can still entail around double the runtime costs as it needs to search where to apply or not apply slope limiting (Ayog et al. 2021). Informed by the study Ayog et al. (2021), flooding cases over a highly rough and dry topography do not need local slope limiting, and switching it off completely greatly accelerates DG2 solver's runtimes.

2.1 The new LISFLOOD-DG2 solver

While LISFLOOD-DG2 is equipped with a generalised minmod slope limiter (Cockburn and Shu, 2001) localised by the Krivodonova shock detector (Krivodonova 2004), the automatic local slope limiter was deactivated for the sake of efficiency: none of the test cases presented in Sect. 3 ~~as none~~ involve shock wave propagation —

~~The friction source term S_f and rainfall source term R are applied separately at the beginning of each time-step: the frictional since all waves propagate over an initially dry bed and are rapidly retarded by frictional forces (Néelz and Pender, 2013; Xia et al., 2019). The lack of shock wave propagation means that all LISFLOOD-FP solvers—DG2, FV1 and ACC—are capable of realistically simulating all test cases presented in Sect. 3.~~

“Line 121: In KS2020, I do not see thus split implicit scheme explained? Rather KS2020 refers to KL2010, where matters are also not explained, so reference is inappropriate. LM2009 do show a split scheme but it is not really implicit. I find (8) to (9) vague. Too convoluted. Simply give the entire scheme in 9 (there is ample space) then we all know what is done rather than having to piece somehow together what is factually done (which I failed to do). Comment whether the schemes for S_f and rain are 2nd or 1st order also per the comment on

eqn (8) below. Given that Sf is kind of quadratic in u (or hu) a semi-implicit scheme is almost instantly made up. That would be straightforward for (9a) and even easy for (9b). One can check the formal time accuracy for the case with only Sf and R on the RHS. And perhaps even for a constant bottom slope river kinematic limit (which is an exact limit within the SWE).”

As mentioned above in our reply to major remark (i), subsection 2.1.2 has been revised to explain how the split friction discretisation of Liang and Marche, 2009 was integrated in our DG2, and in what sense is implicit. In addition, a new subsection 2.1.3 was added to explain how the rain term was integrated.

2.1.3 Discretisation of the rainfall source term

The discretisation of rainfall source term ~~is discretised explicitly to evolve~~ evolves the water depth element-average coefficients $h_{i,j,0}$:

$$h_{i,j,0}^{n+1} = h_{i,j,0}^n + \Delta t R_{i,j}^n, \quad (14)$$

where $R_{i,j}^n$ denotes the prescribed rainfall rate at element (i, j) ~~at and~~ time level n , ~~and Δt is the time step.~~ ~~The original water depth slope coefficients are preserved by.~~ Eqn. (14) ~~is first-order-accurate in space and time, which is deemed sufficient since rainfall data is typically available at far coarser spatial and temporal resolutions than the computation grid, leading to zero element-wise slope coefficients for the rainfall source term in order to preserve the existing local water surface gradient.~~ After applying friction and rainfall source terms, flow coefficients $U_{i,j}$ are evolved from time level n to $n + 1$ using an explicit two-stage Runge-Kutta scheme (Kesserwani et al., 2010): ~~where element indices (i, j) are omitted for clarity of presentation.~~ The time step Δt is calculated according to the CFL condition using the maximum stable Courant number of 0.33 (Cockburn and Shu, 2001). ~~The~~ Recall that the rainfall source term, friction source term, and remaining flux and bed slope terms are treated separately such that, at each timestep, the flow variables updated by Eqn. (14) are subsequently updated by Eqn. (12), and finally by Eqns. (8)–(9). The complete DG2 model workflow is summarised by the flowchart in Fig. 2, wherein each operation is parallelised using the CPU and GPU parallelisation strategies discussed next.

“Line 121: Why is DG2 slower; due to the CFL=1/3 criterion relative to CFL \sim 0.9 for FV1? Eqn (8): why is this scheme for rainfall 2nd order? Why is rain not directly included into (9). $R=R(t)$ so non-autonomous RK would be fine? Please explain your choices better. Please provide the entire scheme for the entire system in (9); there is space and currently the time stepping scheme is unclear. Please clarify.” [Specification for the CFL number with DG2 and FV1:](#)

2.1 The new LISFLOOD-DG2 solver

The remaining terms are the spatial fluxes and topographic slope terms, which are discretised by an explicit second-order two-stage Runge-Kutta scheme (Kesserwani et al., 2010) to evolve the flow coefficients via the $U_{i,j}$ from time level n to $n + 1$:

$$\mathbf{U}^{\text{int}} = \mathbf{U}^n + \Delta t \mathbf{L}(\mathbf{U}^n), \quad (15a)$$

$$\mathbf{U}^{n+1} = \frac{1}{2} [\mathbf{U}^n + \mathbf{U}^{\text{int}} + \Delta t \mathbf{L}(\mathbf{U}^{\text{int}})], \quad (15b)$$

where element indices (i, j) are omitted for clarity of presentation. The initial time-step Δt is a fixed value specified by the user, and the time-step is updated thereafter according to the CFL condition using the maximum stable Courant number of 0.33 (Cockburn and Shu, 2001).

Clarification on why the rain terms is not directly integrated in equation 9:

2.1 The new LISFLOOD-DG2 solver

A standard splitting approach is adopted such that the friction source term \mathbf{S}_f and rainfall source term \mathbf{R} in Eqn. (15) are applied separately at the beginning of each time-step. By adopting a splitting approach, friction or rainfall source terms are only applied as required by the particular test case, for better runtime efficiency. The discretisation of the friction source term is described later in Sect. 2.1.2, and the rainfall source term in Sect. 2.1.3.

“Figure 2: How can one apply rainfall if dt is not know yet, since dt is determined later in the diagram? Order seems off?”

The revised manuscript clarifies that, “The initial time-step is a fixed value specified by the user and the time-step is updated thereafter according to the CFL condition”. Figure 2 has also been clarified by using ‘Update Δt ’ instead of ‘Calculate Δt ’.

“Line 157: typo; L missing (only subscript y seen).” “Line 173: add OUP ‘and,’ at end of third option.

Typo fixed and an ‘and’ was added.

“Line 183: Can it be clarified earlier whether or not (I think it is ”or not”) DG2 can be truly C^0 ? Since the $x*y$ term is not include, therefore it cannot represent C^0 solutions or topography. While speed is gained this drops formal C^0 smoothness. Say so explicitly, probably earlier, also why this option with only polynomials 1, x and y is chosen, I assume to enhance computational speed.”

The introductory text added at the start of subsection 2.1 summarises the rationale for adopting the proposed DG2 formulation. Also, it is stated in the added subsection 2.1.1 the continuity property for the DG2 topography projection is only satisfied at the face-centers where wetting-and-drying treatment and Riemann fluxes are evaluated (i.e. at the location

that matters the most!). A theoretical proof can be found in (Kesserwani et al., 2018). (Revised text shown in reply to an earlier comment.)

“Line 186: Nothing is said about the time steps for FV1 and ACC? Why does FV1 do better for thin overland flows? How are Sf and R dealt with in FV1? This information is relevant to understand the test results provided and analysed later. Please clarify (earlier on).”

The CFL numbers used with FV1 and ACC were added in Section 2; and, the added subsection 2.1.2 offers an explanation on why FV1 can better handle thin overland flows.

“Eqn (14): Perhaps add some spacing around Δt .”

Space added.

“Line 207: Second-order wrt to what since it cannot be 2nd order wrt solutions requiring the advective velocity terms? Please define clearly what is meant here? When advection is important, it cannot be second order, of course?”

Clarification was made to explain in what sense ACC offers a second-order discretisation in space (see also the revised Appendix B). As the set of PDEs used for the ACC solver neglect the advective velocity terms, we agree that its discretization cannot be valid for advection dominated flow, irrespective of the order-of-accuracy of its spatial discretisation.

“Fig. 6: Since FV1 is order one that use of the same grid spacing in the different models is a priori unfair. What happens if FV1 is adjusted (and maybe ACC) such that the expected errors are the same? Please add that situation. The left bottom panel could be done for finer resolution as well, as in Fig 7”

Answered in our reply to major remark (iv).

“Fig. 7: Figure a bit large. Otherwise nice”

The figure size is based on the line length, which is full-width in the proof version, but the figure will only occupy a single column in the final publication.

“Line 270 paragraph: But this comparison is unfair; one needs to compare run times for the same or at least similar accuracy; either focussing on overall accuracy or the accuracy at the front since the latter may be most relevant. This needs an extension.

Line 283: Statement by authors themselves underscores my previous remark. So why is this case not made in the paper with the same convergence/accuracy compared with runtime?

Fig. 8b,c: Same remark; what is the right comparison here? Please address. Caption: page number and caption overlap.”

Answered in our reply to major remark (iv).

“Fig. 12: Can you show simulations for DG2 with $dx=50m$, say, as well since that runs in less time than the other models but the solution may still be better?

Table 1: So DG2 $dx=50m$ should be compared with runtimes of the other models or even DG2-GPU at $dx=25m$ should be compared with the other models, in line with the previous five remarks.”

Addressed in our reply to major remark (iv).

“Line 352: Why is the domain initially dry? How long does it take for the memory of this odd initial condition to disappear?”

This model spin-up procedure is originally specified by Xia et al., 2019, and the memory of the initially dry condition disappears after 1.5 simulated days.

3.3 Catchment-scale rain-on-grid simulation

As specified by Xia et al., 2019, ~~Manning's coefficient n_M is 0.035 for river channels and 0.075 elsewhere.~~ ~~The simulation~~ the simulation comprises a spin-up phase and subsequent analysis phase. The spin-up phase starts at 00:00 3 December 2015 ~~with from~~ an initially dry domain. Water is introduced into the domain via the rainfall source term (Eqn. 14), using Met Office rainfall radar data at a 1 km resolution updated every 5 minutes (Met Office, 2013). ~~Heavy rainfall fills all river channels by~~ The spin-up phase ends and the analysis phase begins at 12:00 4 December 2015, ~~when analysis of results begins~~ once the memory of the dry initial condition has disappeared, and water depths and discharges in all river channels have reached a physically realistic initial state (Xia et al., 2019). The simulation ends at 12:00 8 December 2015 after a total of 5.5 simulated days.

“Line 356: I do not understand this remark ‘No Data’; translate this computer remark.”
The explanation has been revised to avoid using computer jargon:

3.3 Catchment-scale rain-on-grid simulation

An open boundary condition is imposed along the irregular-shaped catchment perimeter by adjusting the ~~DEM so that ‘NoData’ values outside the perimeter of the irregular-shaped catchment are~~ terrain elevation of elements lying outside the catchment such that their elevation is below mean sea level, thereby allowing water to drain out of the River Eden into the Solway Firth. At each time-step, water flowing ~~into ‘NoData’ elements out of the Solway Firth~~ is removed by zeroing the water depth in elements lying outside the catchment.

“Line 364: But that seems unfair. Finer data need to be projected on the topography that DG2 can model, which is not $dx=40m$ but the finer one projected? In fact, DG2 would then do a bit better I suspect. Address this please; perhaps really redo the DG2 case; it should have 3 dofs per cell, also for the topography.”

Answered in our reply to major remark (ii).

“Line 367: Why is repositioning required since across the river channel the river level is within 0.2m, say, the same.” Approximate gauging station coordinates are provided by the UK Environment Agency, but these are often positioned near the river bank and not in the channel itself.

“Is the level meant the river level height above sea level?” Correct. For each model run, hydrographs of free-surface elevation above mean sea level are measured in river channels at sixteen gauging stations.

“What is recorded?” Observed free-surface elevation hydrographs are calculated from Environment Agency measurements of water depth and river bed elevation above mean sea level.

“So why is this repositioning needed, since variation across the river channel will be minimal.” [Approximate gauging station coordinates are provided by Environment Agency, 2020, but these are often positioned near the river bank and not in the channel itself.](#)

“It is free-surface height that matters. Please clarify.” [Yes, this is a more specific term, so we now use “free-surface elevation” consistently across Section 3.3.](#)

“Fig. 14: I am not convinced by the DG2 results due to not taking into account slope information in the bathymetry. So DG2 may be underperforming while this can partially be repaired. Please update and clarify. Line 408: As stated above this anomaly of using only a planar representation in DG2 should be fixed: you do not use DG2 optimally with only 40m planar resolution on bathymetry, while you are allowed to include higher resolution and project on the DG2 basis you use. This leads to an unfair comparison which seems not aligned with the DG2 capabilities.”

[Answered in our reply to major remark \(ii\).](#)

“Line 418: What is the error in the observations? What is the variation in surface level across a (flooded) river channel? What is consequently the result of shifting the measurement position? I would guess the error is circa 0.2m (from watching river levels bob about ± 0.1 m for a river in flood next to a river gauge). What is the variation within 10m along stream of the surface level in the simulation and across the channel, in order to get an error estimate? Please clarify and add a discussion on errors, in both the observations and simulations.”

[As these high-frequency temporal variations average out in our long-running simulation, we wouldn't count them as observation errors. A discussion on the expected observation errors has been added:](#)

3.3.1 River channel free-surface elevation hydrographs

~~River levels~~ [Free-surface elevation hydrographs](#) at the sixteen [river](#) gauging stations are shown in Fig. 14. [Observed free-surface elevation hydrographs are calculated from Environment Agency measurements of water depth and river bed elevation above mean sea level \(Environment Agency, 2020\). While water depths can be measured to an accuracy of \$\sim 0.01\$ m \(Bates et al., 2014\), discrepancies between in-situ, point-wise river bed elevation measurements and the remotely-sensed, two-dimensional DEM can result in systematically biased free-surface elevations, as reported by Xia et al., 2019.](#)

“Line 425: Please add the word ‘maximum’ before the last word ‘flood’.”

[Added.](#)

“Line 428: ‘More notable differences’ Where? Clarify please, e.g. by using arrows. Fig. 16: Graphs in Fig. 16 look all the same; use arrows to highlight areas you want the reader to focus on.”

[Arrows were added, see revised Fig. 16.](#)

“Line 428: On flood defence walls not being captured on coarse grids: unless, of course, you put those in by hand, which should be done (regardless). Here is where variable/non-uniform mesh capabilities of DG2 or FV1 would have come in handy. These are then essentially

(vertical) walls with a finite vertical extent. Also on a coarse mesh one can add vertical walls at cell edges as approximation”

Answered in our reply to major remark (iii).

“Line 432: So, what then is right or wrong, ACC or FV1 and how do we know that? Please clarify”

The text has been revised to clarify that ACC at the finest resolution, of 10m, was used as the reference solution and appropriate metrics to measure floodplain prediction capability were used.

3.3.2 Maximum flood extent over Carlisle

The DG2 and FV1 predictions of maximum flood extent can be quantified against the ACC prediction at $\Delta x = 10$ m, which is treated as the reference solution. The hit rate measures flood extent underprediction as the proportion of wet elements in the reference solution that were also predicted as wet. The false alarm ratio measures flood extent overprediction as the proportion of predicted wet elements that were dry in the reference solution. The critical success index measures both over- and underprediction. All three metrics range between 0 and 1, and further details are provided by Wing et al., 2017.

“Fig. 17: I still disagree with a DG0 bottom topography at $dx=40$ for DG2. That is a false comparison, cf. previous remarks made.”

Answered in our reply to major remark (ii).

“Further: what flood accuracy do we actually care about and what are the error bars given errors in rainfall? Please discuss.” The following text has been added to clarify the level of accuracy for the considered case study.

3.3 Catchment-scale rain-on-grid simulation

While rainfall data errors can influence model outputs, (Ming et al., 2020) found that a prescribed 10% rainfall error lead to only 5% relative mean error in predicted water depth hydrographs. As such, modelling uncertainties due to rainfall errors are not quantified in these deterministic model runs.

“Line 491: But as said before, your topography is incorrect, in not matching DG2 accuracy that can be used. Also, I am not convinced by your time stepping in DG2 as compared to what is done in ACC and FV1 (the latter which is unclear as well and makes it difficult to understand the comparison). Please clarify”

Answered in our reply to major remark (ii).

“Line 501: The conclusion here is a bit of a downer for DG2; no discussion is made on using hybrid approaches, including variables time and spatial resolution for which DG should be good and more promising. Please comment and update.”

Answered in our reply to major remark (vi).

“Appendix B: I am a bit confused; why do we need to linearise to establish whether ACC is second order? Please clarify.” Linearisation is necessary in order to perform a Taylor series expansion of the discrete local inertia equations. The manuscript now justifies this linearisation more clearly:

Appendix B: LISFLOOD-ACC order-of-accuracy

The formal order-of-accuracy of LISFLOOD-ACC is determined by ~~analysing a numerical analysis of~~ the discrete local inertia equations (de Almeida et al., 2012). ~~The numerical analysis starts by linearising the frictionless one-dimensional~~ To begin, the local inertia equations ~~:- in which H is~~ are linearised by assuming small perturbations in free-surface elevation η about a constant reference depth H [L]., leading to the linearised frictionless one-dimensional local inertia equations:

$$\frac{\partial \eta}{\partial t} + \frac{\partial q}{\partial x} = 0, \quad (\text{B1a})$$

$$\frac{\partial q}{\partial t} + gH \frac{\partial \eta}{\partial x} = 0, \quad (\text{B1b})$$

This linear assumption is valid for gradually-varying, quasi-steady flows (de Almeida et al., 2012), and ensures that the remainder of the analysis is tractable.