

# Modeling of land-surface interactions in the PALM model system

## 6.0: Land surface model description, first evaluation, and sensitivity to model parameters

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**Abstract.** In this paper the land-surface model embedded in the PALM model system is described and evaluated against in-situ ~~measurement data in Cabauw. For this, two measurements at Cabauw, Netherlands. Two~~ consecutive clear-sky days are simulated and the components of surface energy balance, ~~as well as near-surface~~ potential temperature, humidity and horizontal wind speed are compared ~~against observation data to observations.~~ For the simulated period, components of the energy balance ~~agree well during are consistent with~~ day- and nighttime ~~, and also observations,~~ the daytime Bowen ratio ~~also~~ agrees fairly well ~~compared to the observations. Although the with observations. The~~ model simulates a ~~significantly~~ more stably-stratified nocturnal boundary layer ~~compared to the observation than the observations,~~ near-surface potential temperature and humidity agree fairly well during day. Moreover, ~~we performed a sensitivity study in order to investigate how much the model results depend a sensitivity analysis is performed to investigate dependence of the model~~ on land-surface and soil specifications, as well as atmospheric initial conditions. ~~By this, we find that a false estimation of the, because they represent a major source of uncertainty in the simulation setup. It is found that an inaccurate estimation of~~ leaf area index, ~~the~~ albedo, or ~~the~~ initial humidity causes a ~~serious-significant~~ misrepresentation of the daytime turbulent sensible and latent heat fluxes. During ~~the~~ night, the boundary-layer characteristics are ~~mostly affected by grid size, surface roughness, primarily affected by surface roughness~~ and the applied radiation schemes.

### 1 Introduction

The land surface influences atmospheric dynamics significantly through the exchange of energy, mass, and momentum. Therefore, an accurate representation of surface-atmosphere interactions is essential for any numerical modeling of the atmospheric boundary layer (Garratt, 1993; Betts et al., 1996). More specifically, surface roughness as well as sensible and latent heat fluxes at the surface act as the lower boundary condition for the momentum, temperature and humidity equations in the atmosphere, respectively. If this information is not available, it is necessary to parameterize the land-surface processes with a Land-Surface Model (LSM). In simple terms, the input to an LSM is the type of surface, vegetation and soil, as well as the radiative forcing. Based on that, LSMs solve the surface energy budget equation by means of a set of prognostic equations for the surface tem-

perature and compute soil moisture and temperature in a multi-layer soil model. Nevertheless, the results strongly depend on the input data (e.g. Avissar and Pielke, 1989).

25 LSMs are required in various ~~situations: Most often, numerical setups, e.g., when~~ observational data of the sensible and latent heat fluxes are unavailable or do not adequately reflect the complexity of the landscape ~~-E.g. in case of forecasts or, in weather forecasting,~~ it is ~~trivial that natural to use~~ a prognostic approach ~~has to be chosen~~. Moreover, the use of an LSM allows ~~for an interaction between atmosphere and surface on shortest timescales and at every point in space, whereas the prescription of fluxes cannot represent this interaction. This plays an important role in simulating to study surface-atmosphere interactions on~~ 30 ~~turbulent timescales across the entire domain. The temporal and spacial scales affect~~ feedback effects, e.g., between mesoscale circulations and underlying heterogeneous surfaces (e.g. Patton et al., 2005; Huang and Margulis, 2010), between clouds and radiation (Lohou and Patton, 2014; Horn et al., 2015), or in cities, where shadows of buildings are another reason for highly heterogeneous heat fluxes.

The PALM model system (~~Maronga et al., 2015, 2020~~) (~~formerly an abbreviation for Parallelized Large-eddy Simulation Model and now~~ 35 has been applied for studying a variety of atmospheric boundary-layer flows for about 20 years. Since 2015, PALM comes with a fully interactive LSM. Originally, it was designed similar to the LSM in the ~~LES Large-Eddy Simulation (LES)~~ model DALES (Heus et al., 2010), which contains parameterizations ~~of the Tiled ECMWF Scheme for Surface Exchanges over Land (TESSEL/H TESSEL, Viterbo and Beljaars, 1995; van den Hurk et al., 2000; Balsamo et al., 2009, 2011)~~ from the ~~TESSEL/H TESSEL scheme (Viterbo and Beljaars, 1995; van den Hurk et al., 2000; Balsamo et al., 2009, 2011)~~ but also some found in the ~~Interaction~~ 40 ~~Sol-Biosphère-Atmosphère model (ISBA, Noilhan and Mahfouf, 1996)~~ ~~ISBA model (Noilhan and Mahfouf, 1996)~~ and own extensions. The LSM in PALM has first been described by Maronga and Bosveld (2017) and has been used with respect to radiation fog by Maronga and Bosveld (2017) and Schwenkel and Maronga (2019), also including Cabauw (Netherlands) data for evaluation. Recently, it has also been employed in urban environments (see first results in Maronga et al., 2020).

Other coupled LES-LSM implementations are found in, e.g., UCLA-LES (Huang et al., 2011), ICON-LEM (Dipankar et al., 45 2015), ~~or and~~ the WRF model (Skamarock et al., 2019) in LES mode coupled to NOAH-LSM (Chen and Dudhia, 2001). In early literature, the focus of coupled LES-LSM studies was mainly to analyze the feedback effect that creates heterogeneity at the surface (~~e.g. Patton et al., 2005; Huang et al., 2009; Brunsell et al., 2011~~) (~~e.g. Patton et al., 2005; Huang et al., 2009; Huang and Margulis, 20~~ 50 . Today, ~~more and more most~~ studies require the coupled LES-LSM approach to simulated realistic cases, e.g., in the urban environment or for wind turbine applications. As the methodology gains foothold in engineering and industry, it becomes increasingly important ~~that for~~ the embedded land-surface representation in PALM ~~reflects to reflect~~ reality (Maronga et al., 2020).

This paper is part of a series featuring different parts of the PALM model system 6.0 in this special issue. For the user, a systematic sensitivity test of relevant land surface parameters with the LSM in PALM is of particular interest. In the present study, we evaluate the LSM embedded into PALM against Cabauw data for a selected period with clear-sky conditions and 55 ~~small limited~~ large-scale advection. By this we ensure as far as possible that the developing boundary-layer is not affected by non-local processes, which we neglect in the simulations of this study. ~~Moreover, we will~~ ~~Furthermore, we~~ determine key parameters that influence the diurnal cycle in a coupled LES-LSM framework. The model sensitivity ~~will be is~~ analyzed by

means of a comprehensive set of simulations varying land-surface and soil parameters individually. Therewith, the present study complements the earlier work of Maronga and Bosveld (2017), who focused on the nocturnal boundary layer with developing radiation fog.

Section 2 of the paper describes physical and technical aspects of the LSM in PALM. Section 3 provides information about the Cabauw Experimental Site for Atmospheric Research (~~CESAR Monna and Bosveld, 2013~~) (CESAR, Monna and Bosveld, 2013) and the observations used. ~~In Sect. 4~~ Section 4 describes the model setup and initialization ~~is described and~~ and gives a complete list of ~~the conducted simulations is given~~ simulations conducted. Section 5 outlines the results of the sensitivity study ~~The and discusses the~~ validity and limitations of the LSM ~~is discussed~~. Finally, ~~the summary and conclusions are drawn in~~ Sect. 6 ~~and an outlook is given~~ draws the summary and conclusions and gives an outlook.

## 2 Description of the land-surface model (LSM) in PALM

The ~~LSM implemented in PALM~~ specific implementation of PALM's LSM, which is derived from the HTESSEL scheme is described below. It consists of a solver for the energy balance of the Earth's surface in combination with a multi-layer soil scheme. The ~~original scheme was~~ scheme was initially designed for vegetated surfaces and bare soil only, but it has since been adapted for paved surface materials like asphalt and concrete, ~~and~~; a simplified version for inland and sea water surfaces has ~~been added, too~~ also been added. A tile approach is available for ~~vegetation~~ vegetated surfaces, in which the surface ~~can comprise of~~ is a fraction of bare soil and a fraction covered with vegetation. Furthermore, the LSM has a liquid water reservoir on plants and soil to store and evaporate liquid water from precipitation interception and dew formation. A liquid water reservoir is also available when the surface type is set to pavement, representing the ability of impervious surfaces to store a limited amount of ~~precipitation~~ precipitated water on the surface. ~~The specific implementation of PALM's LSM, which is derived from the HTESSEL scheme is described below~~. For further details, see also Viterbo and Beljaars (1995); van den Hurk et al. (2000); Balsamo et al. (2009, 2011) and literature referenced therein.

### 2.1 Energy balance solver

~~Within the LSM, the~~ The energy balance is calculated as

$$\frac{dT_0}{dt}C_0 = R_n - H - LE - G, \quad (1)$$

where  $C_0$  and  $T_0$  are the heat capacity and radiative temperature of the surface, respectively. Note that  $C_0$  is zero by default in case of surfaces covered by vegetation or water surfaces, where it is assumed that  $T_0$  is the temperature of a skin layer covering the surface that does not have a heat capacity, ~~which we think is a valid approach for low vegetation like~~ grass. The heat capacity is fully customizable by the user and can be adjusted to account for the heat stored in e.g. forests (Lindroth et al., 2010; Swenson et al., 2019). In all other cases (i.e., pavements and bare soils), no skin layer is assumed (see below).  $R_n$ ,  $H$ ,  $LE$ , and  $G$  are the net radiation, sensible heat flux, latent heat flux, and ground (soil) heat flux at the surface, respectively.  $R_n$  is defined positive downwards whereas  $H$ ,  $LE$ , and  $G$  are defined positive away from the surface.  $R_n$  is

defined through the sum of the radiative fluxes:

$$R_n = SW_{\downarrow} - SW_{\uparrow} + LW_{\downarrow} - LW_{\uparrow}, \quad (2)$$

where  $SW_{\downarrow}$ ,  $SW_{\uparrow}$ ,  $LW_{\downarrow}$ , and  $LW_{\uparrow}$  are the shortwave incoming (downward), shortwave outgoing (upward), longwave incoming (downward), and longwave outgoing (upward) flux, respectively. The radiation components are defined positive according to their direction ( $SW_{\downarrow}$  and  $LW_{\downarrow}$  positive downwards;  $SW_{\uparrow}$  and  $LW_{\uparrow}$  positive upwards). The radiative fluxes are provided by one of the available radiation schemes in PALM (for details, see Maronga et al., 2020).

### 95 2.1.1 Parameterization of fluxes

The turbulent heat fluxes both are parameterized using a resistance parameterization.  $H$  is calculated as

$$H = -\rho c_p \frac{1}{r_a} (\theta_{mo} - \theta_0), \quad (3)$$

where  $\rho$  is the density of the air,  $c_p = 1005 \text{ J kg}^{-1} \text{ K}^{-1}$  is the specific heat at constant pressure, and  $r_a$  is the aerodynamic resistance.  $\theta_0$  and  $\theta_{mo}$  are the potential temperature at the surface and at a fixed height within the atmospheric surface layer  
 100 (at height  $z_{mo}$ , usually at height of the first atmospheric grid level, i.e.,  $z_{mo} = 0.5\Delta_z$  where  $\Delta_z$  is the vertical grid spacing), respectively. The potential temperature is linked to the actual temperature via the Exner function, viz.

$$\Pi = \left( \frac{p}{1000 \text{ hPa}} \right)^{R_d/c_p}, \quad (4)$$

~~with pressure where  $p$  and is the pressure and  $R_d$  is~~ the gas constant for dry air  $R_d$ .  $r_a$  is calculated via Monin-Obukhov similarity theory (MOST) as

$$105 \quad r_a = \frac{\theta_{mo} - \theta_0}{u_* \theta_*}, \quad (5)$$

where  $u_*$  and  $\theta_*$  are the friction velocity and characteristic temperature scale, respectively, ~~and which~~. These values are calculated locally ~~based on MOST. Note that the values for  $u_*$  and  $\theta_*$  from the previous time step are used in Eq. (5) using MOST.~~ The roughness lengths are individually set for momentum, heat, and moisture (see Table 1). Note that ~~for water surfaces, a Charnock parameterization can be switched on for taking into account the effect of subgrid-scale wave motions through the roughness lengths. For details on the particular implementation of MOST and the Charnock parameterization in PALM, see Maronga et al. (2020). Note that  $r_a$  in Eq. (5) is calculated based on  $u_*$  and  $\theta_*$  values from the current time step to calculate  $H$  at the prognostic time step.~~ For details on the particular implementation of MOST in PALM, see Maronga et al. (2020).

The ground heat flux,  $G$ , is parameterized after ~~(Duynerkerke, 1999)~~ Duynerkerke (1999) as

$$G = \Lambda(T_0 - T_{soil,1}), \quad (6)$$

115 with  $\Lambda$  being the total thermal conductivity between skin layer and the uppermost soil layer.  $T_{soil,1}$  is the temperature of the uppermost soil layer (calculated at the center of the layer).  $\Lambda$  is calculated via a resistance approach as a combination of the

conductivity between the canopy and the soil-top ( $\Lambda_{\text{skin}}$ , constant value) and the conductivity of the top half of the uppermost soil layer ( $\Lambda_{\text{soil}}$ ):

$$\Lambda = \frac{\Lambda_{\text{skin}} \Lambda_{\text{soil}}}{\Lambda_{\text{skin}} + \Lambda_{\text{soil}}} . \quad (7)$$

120 When no skin layer is used (i.e. ~~in case of~~ for pavements and bare soils),  $\Lambda$  ~~reduces~~ simplifies to the heat conductivity of the uppermost soil layer, viz.

$$\Lambda = \frac{\lambda_{\text{T,pave}}}{\Delta z_{\text{soil},1}} , \quad (8)$$

with  $\lambda_{\text{T,pave}}$  being the thermal conductivity of the pavement and  $\Delta z_{\text{soil},1}$  being the ~~depth~~ thickness of the uppermost soil layer. In this case, it is assumed that the soil temperature is constant within the uppermost 25 % of the top soil layer and equals  
125 the radiative temperature at the surface.  $C_0$  is then set to a non-zero value according to the material properties and the layer thickness.

The total latent heat flux,  $LE$ , is parameterized as

$$LE = -\rho l_v \frac{1}{r_a + r_s} (q_{v,\text{mo}} - q_{v,\text{sat}}(T_0)) . \quad (9)$$

Here,  $l_v = 2.5 \times 10^6 \text{ J kg}^{-1}$  is the latent heat of vaporization,  $r_s$  is the total surface resistance,  $q_{\text{mo}}$  is the water vapor mixing  
130 ratio at height  $z_{\text{mo}}$ , and  $q_{\text{sat}}$  is the water vapor mixing ratio at saturation at the surface, which is a function of  $T_0$ . In practice, up to three individual components are calculated for vegetated surfaces. Transpiration of the the vegetated fraction ( $LE_{\text{veg}}$ ) is parameterized as

$$LE_{\text{veg}} = -\rho l_v \frac{1}{r_a + r_c} (q_{v,\text{mo}} - q_{v,\text{sat}}(T_0)) , \quad (10)$$

where  $r_c$  is the canopy resistance. Analogous the bare soil fraction evaporation ( $LE_{\text{soil}}$ ) is calculated via

$$135 \quad LE_{\text{soil}} = -\rho l_v \frac{1}{r_a + r_{\text{soil}}} (q_{v,\text{mo}} - q_{v,\text{sat}}(T_0)) , \quad (11)$$

with  $r_{\text{soil}}$  being the soil resistance. The liquid water reservoir evaporation ( $LE_{\text{liq}}$ ) is given by

$$LE_{\text{liq}} = -\rho l_v \frac{1}{r_a} (q_{v,\text{mo}} - q_{v,\text{sat}}(T_0)) , \quad (12)$$

i.e., only the aerodynamic resistance exists for liquid water. The total evapotranspiration is then given by a combination of the three individual components by (see Viterbo and Beljaars, 1995)

$$140 \quad LE = c_{\text{veg}}(1 - c_{\text{liq}})LE_{\text{veg}} + c_{\text{liq}}LE_{\text{liq}} + (1 - c_{\text{veg}})(1 - c_{\text{liq}})LE_{\text{soil}} . \quad (13)$$

Here,  $c_{\text{veg}}$  and  $c_{\text{liq}}$  are the fractions of the surface covered with vegetation and liquid water, respectively. Liquid water from precipitation can be stored on the vegetation and bare soil. Note that for paved and water surfaces both  $LE_{\text{veg}}$  and  $LE_{\text{soil}}$  are set to zero and the only possible source of evaporation is the liquid water reservoir.

All equations above are solved locally for each surface element of the model grid.

## 145 2.1.2 Liquid water reservoir

~~In order to~~ To account for the evaporation of liquid water on plants ~~on and~~ on impervious surfaces, an additional equation is solved for the liquid water reservoir:

$$\frac{d m_{\text{liq}}}{d t} = \frac{L E_{\text{liq}}}{\rho_l l_v}, \quad (14)$$

where  $m_{\text{liq}}$  and  $\rho_l$  are the water column on the surface and the density of water, respectively. The maximum amount of water  
150 that can be stored on plants is calculated via

$$m_{\text{liq},\text{max}} = \min \left( 1, m_{\text{liq},\text{ref}} \cdot (c_{\text{veg}} \cdot LAI + (1 - c_{\text{veg}})) \right), \quad (15)$$

where  $m_{\text{liq},\text{ref}} = 0.2\text{mm}$  is the reference liquid water column on a single leaf or bare soil and  $LAI$  the leaf area index. Exceeding liquid water is directly removed from the surface and infiltrated in the underlying soil. For paved surfaces,  $m_{\text{liq},\text{max}}$  is set to 1 mm. Exceeding liquid water is assumed to be drained off. Note that  $m_{\text{liq}}$  enters the calculation of  $L E_{\text{liq}}$  indirectly via  
155  $c_{\text{liq}}$ , which is given either as the ratio  $m_{\text{liq}}/m_{\text{liq},\text{max}}$  for vegetation (following the HTESSEL scheme) or  $(m_{\text{liq}}/m_{\text{liq},\text{max}})^{0.67}$  for pavement following Masson (2000) (based on Noilhan and Planton, 1989).

## 2.1.3 Calculation of resistances

The resistances are calculated separately for bare soil and vegetation following Jarvis (1976). The canopy resistance,  $r_c$ , is calculated as

$$160 \quad r_c = \frac{r_{c,\text{min}}}{LAI} f_1(SW_{\downarrow}) f_2(\tilde{m}) f_3(e_{\text{def}}), \quad (16)$$

with  $r_{c,\text{min}}$  being a minimum canopy resistance.  $f_1 - f_3$  are correction functions ~~depending that depend~~ on  $LAI$ , the incoming shortwave radiation ( $SW_{\downarrow}$ ) and the water-vapor pressure deficit ( $e_{\text{def}} = e_{\text{sat}} - e$ , with  $e_{\text{sat}}$  and  $e$  being the water-vapor pressure at saturation and the current water-vapor pressure, respectively). The layer-averaged volumetric soil moisture content ( $\tilde{m}$ ) is given by

$$165 \quad \tilde{m} = \sum_{k=1}^N R_{\text{fr},k} \max(m_{\text{soil},k}, m_{\text{wilt}}), \quad (17)$$

where  $N$  is the number of soil layers,  $R_{\text{fr},k}$  is the root fraction in layer  $k$ ,  $m_{\text{soil},k}$  is the volumetric soil moisture content in layer  $k$ , and  $m_{\text{wilt}}$  is the permanent wilting point.

The correction functions  $f_1$  ~~and  $f_2$  read is~~

$$\frac{1}{f_1(SW_{\downarrow})} = \min \left( 1, \frac{0.004 SW_{\downarrow}}{0.81(0.004 SW_{\downarrow} + 1)} \right), \quad (18)$$

170 which accounts for the reaction of plants to sunlight (opening/closing stomatas); the reaction of plants to water availability in the soil is ~~considered via~~ given by the correction function  $f_2$ , as

$$\frac{1}{f_2(\tilde{m})} = \begin{cases} 0 & \tilde{m} < m_{\text{wilt}} \\ \frac{\tilde{m} - m_{\text{wilt}}}{m_{\text{fc}} - m_{\text{wilt}}} & m_{\text{wilt}} \leq \tilde{m} \leq m_{\text{fc}} \\ 1 & \tilde{m} > m_{\text{fc}}, \end{cases} \quad (19)$$

with  $m_{\text{fc}}$  being the soil moisture at field capacity. Furthermore, a correction for the water vapor pressure deficit is given by

$$\frac{1}{f_3(e_{\text{def}})} = \exp(g_{\text{D}} e_{\text{def}}), \quad (20)$$

175 where  $g_{\text{D}}$  is a correction factor that is used for ~~high-vegetation~~ tall vegetation such as trees.

The soil resistance ( $r_{\text{soil}}$ ) is calculated as

$$r_{\text{soil}} = r_{\text{soil,min}} \cdot f_4(m_{\text{soil},1}), \quad (21)$$

where  $r_{\text{soil,min}}$  is the minimum soil resistance. The correction function  $f_4$  is given by

$$f_4 = \max\left(\frac{m_{\text{soil},1} - m_{\text{min}}}{m_{\text{fc}} - m_{\text{min}}}, 1\right), \quad (22)$$

180 with  $m_{\text{min}}$  being a minimum soil moisture for the soil matrix based on the wilting point and the residual moisture  $m_{\text{res}}$ , calculated as

$$m_{\text{min}} = c_{\text{veg}} m_{\text{wilt}} + (1 - c_{\text{veg}}) m_{\text{res}}. \quad (23)$$

Note that the total surface resistance ( $r_{\text{s}}$ , ~~cf.~~ see Eq. (9)) is calculated as a diagnostic quantity from  $LE$  after the energy balance is solved.

## 185 2.2 Soil model

~~The soil model consists of prognostic~~ Prognostic equations for the soil temperature and the volumetric soil moisture ~~which are solved for multiple layers~~. ~~The soil model only takes into account vertical transport~~ are solved in multiple layers in the soil model. Transport is restricted to only the vertical direction within the soil and no ice phase is considered ~~at the moment~~ currently. By default, the soil model ~~consists is constructed~~ of eight layers, with default layer ~~depths~~ thicknesses of 0.01, 0.02, 0.04, 190 0.06, 0.14, 0.26, 0.54, and 1.86 m, ~~but the number of layers as well as their depths can be modified~~ however the number and thicknesses of layers are fully customizable. The vertical heat and water transport is modeled using the Fourier law of diffusion and Richards' equation, respectively. For vegetated surface elements, root fractions ~~can be~~ are assigned to each soil layer to account for the explicit ~~water withdrawal of plants~~ withdrawal of water by plants, used for transpiration, from the respective soil layer. Viterbo and Beljaars (1995) and Balsamo et al. (2009) give more details.

**Table 1.** Look-up table for vegetation parameters of 18 predefined vegetation types in the style of the ECMWF-IFS classification, adapted for PALM.  $\Lambda_{\text{skin},s}$  and  $\Lambda_{\text{skin},u}$  are the total thermal conductivities between the skin layer and the surface for near-surface stable and unstable stratification, respectively.  $\epsilon$  is the surface emissivity. All other symbols are used as defined in the main text.

type	description	$r_{c,\min}$ ( $\text{m s}^{-1}$ )	$LAI$	$c_{\text{veg}}$	$g_D$ ( $\text{hPa}^{-1}$ )	$z_0$ (m)	$z_{0,h}$ (m)	$\Lambda_{\text{skin},s}$ ( $\text{W m}^{-2} \text{K}^{-1}$ )	$\Lambda_{\text{skin},u}$ ( $\text{W m}^{-2} \text{K}^{-1}$ )	$C_0$	albedo type	$\epsilon$
1	bare soil	0.0	0.00	0.00	0.00	0.005	0.5E-4	0.0	0.0	0.00	17	0.94
2	crops, mixed farming	180.0	3.00	1.00	0.00	0.10	0.001	10.0	10.0	0.00	2	0.95
3	short grass	110.0	2.00	1.00	0.00	0.03	0.3E-4	10.0	10.0	0.00	5	0.95
4	evergreen needleleaf trees	500.0	5.00	1.00	0.03	2.00	2.00	20.0	15.0	0.00	6	0.97
5	deciduous needleleaf trees	500.0	5.00	1.00	0.03	2.00	2.00	20.0	15.0	0.00	8	0.97
6	evergreen broadleaf trees	175.0	5.00	1.00	0.03	2.00	2.00	20.0	15.0	0.00	9	0.97
7	deciduous broadleaf trees	240.0	6.00	0.99	0.13	2.00	2.00	20.0	15.0	0.00	7	0.97
8	tall grass	100.0	2.00	0.70	0.00	0.47	0.47E-2	10.0	10.0	0.00	10	0.97
9	desert	250.0	0.05	0.00	0.00	0.013	0.013E-2	15.0	15.0	0.00	11	0.94
10	tundra	80.0	1.00	0.50	0.00	0.034	0.034E-2	10.0	10.0	0.00	13	0.97
11	irrigated crops	180.0	3.00	1.00	0.00	0.5	0.50E-2	10.0	10.0	0.00	2	0.97
12	semidesert	150.0	0.50	0.10	0.00	0.17	0.17E-2	10.0	10.0	0.00	11	0.97
13	ice caps and glaciers	0.0	0.00	0.00	0.00	1.3E-3	1.3E-4	58.0	58.0	0.00	14	0.97
14	bogs and marshes	240.0	4.00	0.60	0.00	0.83	0.83E-2	10.0	10.0	0.00	3	0.97
15	evergreen shrubs	225.0	3.00	0.50	0.00	0.10	0.10E-2	10.0	10.0	0.00	4	0.97
16	deciduous shrubs	225.0	1.50	0.50	0.00	0.25	0.25E-2	10.0	10.0	0.00	5	0.97
17	mixed forest/woodland	250.0	5.00	1.00	0.03	2.00	2.00	20.0	15.0	0.00	10	0.97
18	interrupted forest	175.0	2.50	1.00	0.03	1.10	1.10	20.0	15.0	0.00	7	0.97



**Table 2.** Look-up table for soil parameters.

type	description	$\alpha$	$l$	$n$	$\gamma_{\text{sat}}$ ( $\text{m s}^{-1}$ )	$m_{\text{sat}}$ ( $\text{m}^3 \text{s}^{-3}$ )	$m_{\text{fc}}$ ( $\text{m}^3 \text{s}^{-3}$ )	$m_{\text{wilt}}$ ( $\text{m}^3 \text{s}^{-3}$ )	$m_{\text{res}}$ ( $\text{m}^3 \text{s}^{-3}$ )
1	coarse	3.83	1.150	1.38	6.94E-6	0.403	0.244	0.059	0.025
2	medium	3.14	-2.342	1.28	1.16E-6	0.439	0.347	0.151	0.010
3	medium-fine	0.83	-0.588	1.25	0.26E-6	0.430	0.383	0.133	0.010
4	fine	3.67	-1.977	1.10	2.87E-6	0.520	0.448	0.279	0.010
5	very fine	2.65	2.500	1.10	1.74E-6	0.614	0.541	0.335	0.010
6	organic	1.30	0.400	1.20	1.20E-6	0.766	0.663	0.267	0.010

### 195 2.2.1 Soil heat transport

The Fourier law of diffusion ~~reads is~~

$$(\rho C)_{\text{soil}} \frac{\partial T_{\text{soil}}}{\partial t} = \frac{\partial}{\partial z} \left( \lambda_{\text{T}} \frac{\partial T_{\text{soil}}}{\partial z} \right), \quad (24)$$

with  $(\rho C)_{\text{soil}}$  and  $\lambda_{\text{T}}$  being the volumetric heat capacity and the thermal conductivity of ~~the soil layer in question~~ a soil layer, respectively.  $\lambda_{\text{T}}$  is calculated as

$$\lambda_{\text{T}} = Ke (\lambda_{\text{T,sat}} - \lambda_{\text{T,dry}}) + \lambda_{\text{T,dry}}, \quad (25)$$

with  $\lambda_{\text{T,sat}}$ ,  $\lambda_{\text{T,dry}}$ ,  $\lambda_{\text{T,sat}}$ ,  $\lambda_{\text{T,dry}}$ , and  $Ke$  being the thermal conductivity of saturated soil, the thermal conductivity of dry soil and the Kersten number, respectively.  $\lambda_{\text{T,sat}}$  is given by

$$\lambda_{\text{T,sat}} = \lambda_{\text{T,sm}}^{1-m_{\text{soil,sat}}} \lambda_{\text{m}}. \quad (26)$$

Here,  $\lambda_{\text{T,sm}}$  is the thermal conductivity of the soil matrix and  $\lambda_{\text{m}}$  is the heat conductivity of water. The Kersten number  $(Ke)$  is calculated as

$$Ke = \log_{10} \left[ \max \left( 0.1, \frac{m_{\text{soil}}}{m_{\text{sat}}} \right) \right] + 1. \quad (27)$$

At the bottom boundary a fixed deep soil temperature  $T_{\text{deep}}$  is prescribed (Dirichlet conditions), ~~which is a plausible assumption for short-term simulations covering only a few days.~~ The user must ensure that the soil model reaches deep enough such that atmospheric-driven temperature changes do not propagate down to the boundary condition.

### 210 2.2.2 Soil moisture transport

The vertical transport of water within the soil matrix is calculated using Richards' equation, viz.

$$\frac{\partial m_{\text{soil}}}{\partial t} = \frac{\partial}{\partial z} \left( \lambda_{\text{m}} \frac{\partial m_{\text{soil}}}{\partial z} - \gamma \right) + S_m, \quad (28)$$

**Table 3.** Look-up table for albedo parameters. (\*1) land ice is treated differently than sea ice (\*2) preliminary/dummy values

albedo type	Description	broadband	longwave	shortwave	Notes
1	ocean	0.06	0.06	0.06	
2	mixed farming, tall grassland	0.19	0.28	0.09	
3	tall/medium grassland	0.23	0.33	0.11	
4	evergreen shrubland	0.23	0.33	0.11	
5	short grassland/meadow/shrubland	0.25	0.34	0.14	
6	evergreen needleleaf forest	0.14	0.22	0.06	
7	mixed deciduous forest	0.17	0.27	0.06	
8	deciduous forest	0.19	0.31	0.06	
9	tropical evergreen broadleaved forest	0.14	0.22	0.06	
10	medium/tall grassland/woodland	0.18	0.28	0.06	
11	desert, sandy	0.43	0.51	0.35	
12	desert, rocky	0.32	0.40	0.24	
13	tundra	0.19	0.27	0.10	
14	land ice	0.77	0.65	0.90	*1
15	sea ice	0.77	0.65	0.90	
16	snow	0.82	0.70	0.95	
17	bare soil	0.08	0.08	0.08	
18	asphalt/concrete mix	0.17	0.17	0.17	*2
19	asphalt (asphalt concrete)	0.17	0.17	0.17	*2
20	concrete (Portland concrete)	0.30	0.30	0.30	*2
21	sett	0.17	0.17	0.17	*2
22	paving stones	0.17	0.17	0.17	*2
23	cobblestone	0.17	0.17	0.17	*2
24	metal	0.17	0.17	0.17	*2
25	wood	0.17	0.17	0.17	*2
26	gravel	0.17	0.17	0.17	*2
27	fine gravel	0.17	0.17	0.17	*2
28	pebblestone	0.17	0.17	0.17	*2
29	woodchips	0.17	0.17	0.17	*2
30	tartan (sports)	0.17	0.17	0.17	*2
31	artificial turf (sports)	0.17	0.17	0.17	*2
32	clay (sports)	0.17	0.17	0.17	*2
33	building (dummy)	0.17	0.17	0.17	*2

where  $\lambda_m$ ,  $\lambda_m$ ,  $\gamma$ , and  $S_m$  are the hydraulic diffusion coefficient, hydraulic conductivity, and a sink term due to root extraction, respectively. The hydraulic diffusion coefficient is calculated after Clapp and Hornberger (1978) as

$$215 \quad \lambda_{m,m} = \frac{b\gamma_{\text{sat}}(-\Psi_{\text{sat}})}{m_{\text{sat}}} \left( \frac{m_{\text{soil}}}{m_{\text{sat}}} \right)^{b+2}, \quad (29)$$

with  $b = 6.04$  being a fixed parameter,  $\gamma_{\text{sat}}$  being the hydraulic conductivity at saturation, and  $\Psi_{\text{sat}} = -338$  m being the soil matrix potential at saturation. The hydraulic conductivity ( $\gamma$ ) is calculated after van Genuchten (1980) (as in HTESSEL):

$$\gamma = \gamma_{\text{sat}} \frac{[(1 + (\alpha h)^n)^{1-1/n} - (\alpha h)^{n-1}]^2}{(1 + (\alpha h)^n)^{(1-1/n)(l+2)}}. \quad (30)$$

220 Here,  $\alpha$ ,  $n$ , and  $l$  are van Genuchten coefficients that depend on the soil type (see Table 2).  $h$  is the pressure head, which is calculated via rearrangement of

$$m_{\text{soil}}(h) = m_{\text{res}} + \frac{m_{\text{sat}} - m_{\text{res}}}{(1 + (\alpha h)^n)^{1-1/n}}. \quad (31)$$

The root extraction of water from the respective soil layer  $S_{m,k}$ ,  $S_{m,k}$  is calculated as follows:

$$S_{m,k,m,k} = \frac{LE_{\text{veg}}}{\rho_l l_v} \frac{R_{\text{fr},k}}{\Delta z_{\text{soil},k}} \frac{m_{\text{soil},k}}{m_{\text{total}}}, \quad (32)$$

where  $m_{\text{total}}$  is the total water content of the soil,

$$225 \quad m_{\text{total}} = \sum_{k=1}^N R_{\text{fr},k} m_{\text{soil},k}, \quad (33)$$

with  $R_{\text{fr},k}$  being the root fraction in soil layer  $k$ . Only ~~those layer are summed up which the layers that~~ have a soil moisture above ~~wilting point (the wilting point are used in Eq. (33) (i.e. plants are not able, plants are unable~~ to withdraw water from layers with soil moisture below ~~the~~ wilting point). The root distribution within the soil must ~~chosen in such a way be chosen~~ ~~such~~ that

$$230 \quad \sum_{k=1}^N R_{\text{fr},k} = 1. \quad (34)$$

There are two options ~~available for the bottom boundary conditions for soil moisture for the soil moisture bottom boundary condition~~. The bottom surface can be set to either bedrock, i.e. water ~~can not be drained off and is not drained. Instead it is~~ accumulated in the lowest soil layer (water content conservation), ~~or it~~, ~~Alternatively, the bottom boundary~~ can be set to free drainage, i.e. an open bottom where soil water is continuously lost by drainage ~~(water content is not conserved)~~.

### 235 2.2.3 Treatment of pavements

Pavements are treated ~~as a common identically to a~~ soil (allowing varying number and depths of the pavement layers) but with ~~the~~ physical properties of the pavement material. The pavement layer is impermeable to water ~~and, which~~ prohibits the vertical transport of soil moisture. Soil layers are placed below the pavement layers.

## 2.2.4 Treatment of water bodies

240 For water surfaces, PALM currently only allows ~~for prescribing~~ prescription of a bulk water temperature. The energy balance is then solved as for land surfaces, but without evapotranspiration from vegetation and bare soil (see above). A skin layer is adopted ~~so that  $C_0 = 0$  and  $\Lambda = 1 \times 10^{11}$  in order~~ to calculate the heat flux into the water body, with  $C_0 = 0 \text{ J m}^{-2} \text{ K}^{-1}$  and  $\Lambda = 1 \times 10^{11} \text{ W m}^2 \text{ K}^{-1}$ .

## 2.3 Numerical methods

245 ~~In order to solve for the~~ In case of an ocean surface, a Charnock parameterization can be switched on to account for the effect of waves on the surface friction in terms of a modification of the surface roughness length as described by Beljaars (1994). For details, see Maronga et al. (2020).

## 2.3 Numerical methods

250 To solve the energy balance for the surface temperature ( $T_0$ ), Eq. (1) is first linearized around  $T_0$  at the current time step and then discretized in time using PALM's default Runge-Kutta third-order ~~time stepping scheme. In this way~~ time stepping scheme. With this method, an iterative procedure to solve the energy balance is avoided ~~and the~~. The prognostic equation then reads

$$T_0^{t+1} = \frac{A\Delta t + C_0 T_0^t}{C_0 + B\Delta t}, \quad (35)$$

where  $t$  is the time index, ~~and~~  $\Delta t$  is the current time step.  $A$  and  $B$  are coefficients given by

$$255 \quad A = R_n + 3\sigma T_0^4 + \left(\frac{\rho c_p}{r_a}\right) \theta_{\text{mo}} + \left(\frac{\rho l_v}{r_a + r_s}\right) \left(q_{\text{mo}} - q_{\text{sat}} + \frac{dq_{\text{sat}}}{dT} T_0\right) + \Lambda T_{\text{soil},1} \quad (36)$$

and

$$B = \Lambda + \left(\frac{\rho l_v}{r_a + r_s}\right) \frac{dq_{\text{sat}}}{dT} + \left(\frac{\rho c_p}{r_a \Pi}\right) + 4\sigma T_0^3. \quad (37)$$

Here,  $\sigma = 5.67037 \times 10^{-8}$  is the Stefan-Boltzmann constant. For vegetated surfaces, where  $C_0$  is zero, Eq. (35) reduces to a diagnostic relationship viz.

$$260 \quad T_0^{t+1} = \frac{A}{B}. \quad (38)$$

## 3 CESAR observations

For model evaluation we ~~chose to~~ simulate two consecutive clear-sky days observed on 5<sup>th</sup> and 6<sup>th</sup> of May, 2008, at the CESAR site ~~at Cabauw. The period was chosen near Cabauw. This period is used~~ because the forcing from the surface was dominant and larger-scale advection ~~played a minor role. We used direct measurements of temperature, humidity and wind and derived~~

observation of sensible, latent and ground heat flux as well as net radiation is limited. We decided to simulate a two-day period, to study the behavior of the model over a full diurnal cycle and also look at how the first day affects the following day. With a longer period, model drift must be taken care of by e.g. adding nudging to the forcing and/or data assimilation. This would add additional uncertainty, because height-dependent advection, needed to drive the LES model, is difficult or impossible to obtain from observations, particularly within the boundary layer where turbulence dominates.

As part of the IMPACT-EUCAARI campaign in May 2008 (Intensive Measurement Period at the Cabauw Tower within the European Integrated project on Aerosol Cloud Climate and Air Quality Interactions), radiosondes were launched daily at 05:00, 10:00, and 16:00 UTC, which we use to initialize the simulations. CESAR features a 213 m high measurement mast with instruments in 1.5, 10, 20, 40, 80, 140 and 200 m (Bosveld, 2020b). For temperature, humidity and wind, our model comparison, we use 10-minute averages of temperature, specific humidity and horizontal wind speed. The measurement accuracy for temperature and humidity is 0.1 K and 3.5 %, respectively (Meijer, 2000). The accuracy for horizontal wind speed is the largest of either 1 % or  $0.1 \text{ m s}^{-1}$ . Soil temperature is observed at a depth of 0, 2, 4, 6, 8, 12, 20, 30 and 50 cm. Additionally, as part of the IMPACT-EUCAARI campaign in May 2008 (Intensive Measurement Period at the Cabauw Tower within the European Integrated project on Aerosol Cloud Climate and Air Quality Interactions), radiosondes were launched daily at 05:00, 10:00, and 16:00 UTC.

The surface soil heat flux is derived from soil temperature measurements computed from soil heat flux measurements at 5 and 10 cm depth by means of a Fourier extrapolation (see Bosveld, 2020b, variable FG0 in their Ch. 19). Net radiation is derived from the budget of the four radiation components (see Eq. 2). Turbulent fluxes of sensible and latent heat are derived by means of the eddy-covariance (EC) method. In the process of calculating EC-fluxes from raw turbulence data it is unavoidable that low frequency contributions to the flux are getting lost. For the current Cabauw data, the 10 minute intervals are used with simple subtraction of the mean means are subtracted from the raw turbulent timeseries; afterwards, a low frequency correction is applied based on the spectra by Kaimal et al. (1972), taking into account with dependence on wind speed and stability (Bosveld, 2020a). The method of low-frequency loss correction assumes that all turbulence characteristics follow surface-layer scaling. However, this is not always true, as for example, horizontal advection by organized turbulent structures (Eder et al., 2015; De Roo and Mauder, 2018) may add further low-frequency contributions which is that are not accounted for in the surface-layer scaling. For further information on the instrumentation of the Cabauw site we refer to see Bosveld (2020b).

Cabauw is located in the western part of the Netherlands and surrounded by mostly meadows with ditches passing through as well as villages, orchards and lines of trees. The CESAR tower itself is installed over an area of short grass that is kept at a height of about 8 cm. The immediate surroundings of the measuring tower are free of significant heterogeneities for a few hundreds of meters. During the simulation period from 5<sup>th</sup> to 6<sup>th</sup> of May, 2008 the prevailing wind direction was from the south-east and the 10 m average wind speed ranged from 2 to  $6 \text{ m s}^{-1}$ . The convective boundary layer reached a height of around 2 km. The groundwater level is in 1.3 m depth. The soil temperature,  $T_{\text{soil}}$ , and soil moisture,  $m_{\text{soil}}$ , at 2008-05-05 05:00 UTC is depicted in Fig. 1 (see Sect. 4). The profiles of temperature  $\theta$ , humidity  $q$  and horizontal wind  $v_h$  retrieved from radiosounding are shown in Fig. 2 together

with tower measurements from all available levels (see cf. details in Sect. 5). The soil temperature  $T_{\text{soil}}$  and soil moisture  $m_{\text{soil}}$  at 2008-05-05 05:00 UTC is depicted in Fig. 1 (see Sect. 4).

#### 4 Simulation setup

To evaluate a land surface parameterization scheme the relevant information is required. With regards to this, the CESAR site is well described in literature (e.g. Beljaars and Bosveld, 1997) (e.g. Beljaars and Bosveld, 1997) which provides a good opportunity to evaluate the land surface parameterization proposed in the present study. The reference simulation (hereafter referred to as case REF) was set up as a best guess with the Cabauw land surface parameters according to Beljaars and Bosveld (1997), Ek and Holtslag (2004) and Maronga and Bosveld (2017). The vegetation type of the surface is ‘short grass’ with some modifications. The roughness length for momentum is set to 0.15 m, which is representative for a few kilometers of upstream terrain from the Cabauw tower and the roughness length for temperature is set to  $2.35 \times 10^{-5}$  m (Ek and Holtslag, 2004). The leaf area index is  $1.7 \text{ m}^2 \text{ m}^{-2}$ , the minimum canopy resistance is  $110 \text{ s m}^{-1}$  and we chose a heat conductivity between skin layer and soil of  $4 \text{ W m}^{-2} \text{ K}^{-1}$ . The heat capacity of the skin layer is set to  $0 \text{ J m}^{-2} \text{ K}^{-1}$ , the surface emissivity  $\epsilon = 0.97$  and the surface is covered to 100 % with vegetation. The soil layers are defined at depth depths of 0.005, 0.02, 0.04, 0.065, 0.1, 0.15, 0.24, 0.45, 0.675, 1.125 and 2.25 m. The soil parameters for field capacity and wilting point are 0.491 and 0.314, respectively, which according to the ECMWF-IFS classification would best be described as is consistent with a very fine soil. The residual moisture is set to  $0.01 \text{ m}^3 \text{ m}^{-3}$  and the minimum soil resistance to  $50 \text{ s m}^{-1}$ . The texture. The deep soil temperature is fixed at 283.19 K which is a valid assumption, because the lower two soil levels are not reached by diurnal temperature variations (lower part of Fig. 1 does not change over time). The van Genuchten coefficients, the hydraulic conductivity at saturation and the porosity vary with soil depth. In the uppermost 24 cm, the parameters are set to match medium-fine soil (type 3, in Table 2), the layer between 24-60 cm is the same as the identical to the uppermost layer except for the porosity which had to be, which is increased due to observed large values of soil moisture, and between. Between 60-225 cm, the parameters are set to organic soil according to Table 2. In all simulations, the land surface and soil parameters are homogeneous over the model domain. This means that buildings of the small town of Lopik, west of the CESAR tower, are neglected, as well as are the small ditches which that cross the observation site.

The root fraction and initial soil profiles of temperature and moisture are shown in Fig. 1. In 3 cm depth begins a layer of relatively high root density,  $R_{\text{dens}}$ , extends from 3 cm below the surface down to 18 cm followed by a layer of relatively low root density down to 60 cm depth. According to Jager et al. (1976), no roots are found near the surface ( $< 3$  cm) and in the deep soil layers ( $> 60$  cm). The initial soil temperature and moisture profiles are taken from measurement data at 2008-05-05 05:00 UTC (shown in Fig. 1). A summary of the land surface scheme configuration and its used values are listed in Table 4.

Vertical soil layer setup of root density ( $R_{\text{dens}}$ , left) as well as initial profiles of temperature (middle) and moisture (right). Note the broken vertical axis with a changed linear increment in the deeper layers. The root fraction of each soil layer (cf.  $R_{\text{fr}}$  in Table 4) is the difference of the cumulative root fraction (Eq. (34), shown on the vertical axis) between two layers.

**Table 4.** Overview of the land surface scheme configuration for case REF.

Parameter	Value	Description
Skin layer parameters		
$C_0$	$0 \text{ J m}^{-2} \text{ K}^{-1}$	Heat capacity of the skin layer
$c_{\text{veg}}$	100%	Vegetation coverage of the surface
$LAI$	<del>1.7</del> $1.7 \text{ m}^2 \text{ m}^{-2}$	Leaf area index
$r_{c,\text{min}}$	$110 \text{ s m}^{-1}$	Minimum canopy resistance
$z_0$	0.15 m	Roughness length for momentum
$z_{0,h}$	$2.35 \times 10^{-5} \text{ m}$	Roughness length for temperature
$\Lambda_{\text{skin}}$	<del>4.0</del> $4.0 \text{ W m}^2 \text{ K}^{-1}$	Heat conductivity between skin layer and soil
$\epsilon$	<del>0.97</del> $0.97$	Surface emissivity
Soil parameters		
$m_{\text{res}}$	$0.010 \text{ m}^3 \text{ s}^{-3}$	Residual volumetric soil moisture
$r_{\text{soil},\text{min}}$	$50 \text{ s m}^{-1}$	Minimum soil resistance
$T_{\text{deep}}$	283.19 K	Deep soil temperature
$m_{\text{fc}}$	$0.491 \text{ m}^3 \text{ s}^{-3}$	Volumetric soil moisture at field capacity
$m_{\text{wilt}}$	$0.314 \text{ m}^3 \text{ s}^{-3}$	Volumetric soil moisture at permanent wilting point
Height dependent soil parameters (0-24 cm, 24-60 cm, 60-225 cm)		
$\alpha$	0.83, 0.83, 1.30	van Genuchten coefficient
$l$	-0.588, -0.588, 0.400	van Genuchten coefficient
$n$	1.25, 1.25, 1.20	van Genuchten coefficient
$\gamma_{\text{sat}}$	<del>0.26, 0.26, <math>1.2 \times 10^{-6} \text{ m s}^{-1}</math></del> $0.26, 0.26, 1.2 \times 10^{-6} \text{ m s}^{-1}$	Hydraulic conductivity of the soil at saturation
$m_{\text{sat}}$	<del>0.430, 0.766, <math>0.766 \text{ m}^3 \text{ s}^{-3}</math></del> $0.430, 0.766, 0.766 \text{ m}^3 \text{ s}^{-3}$	Volumetric soil moisture at saturation (porosity)
Initial soil profiles		
$T_{\text{soil},k}$	283.96, 284.00, 284.62, 284.59, 284.70, 284.77, 284.55, 283.50, 283.19, 283.19, 283.19 K	Soil temperature at depth level $k$ ( $k \in 1, 11$ )
$m_{\text{soil},k}$	0.324, 0.324, 0.324, 0.332, 0.352, 0.380, 0.477, 0.605, 0.670, 0.721, <del>0.721</del> $0.721 \text{ m}^3 \text{ m}^{-3}$ $0.721 \text{ m}^3 \text{ m}^{-3}$	Soil moisture at depth level $k$ ( $k \in 1, 11$ )
$R_{\text{fr},k}$	0, 0, 0.2, 0.2, 0.2, 0.2, 0.1, 0.1, 0, 0, 0	Root fraction at depth level $k$ ( $k \in 1, 11$ )

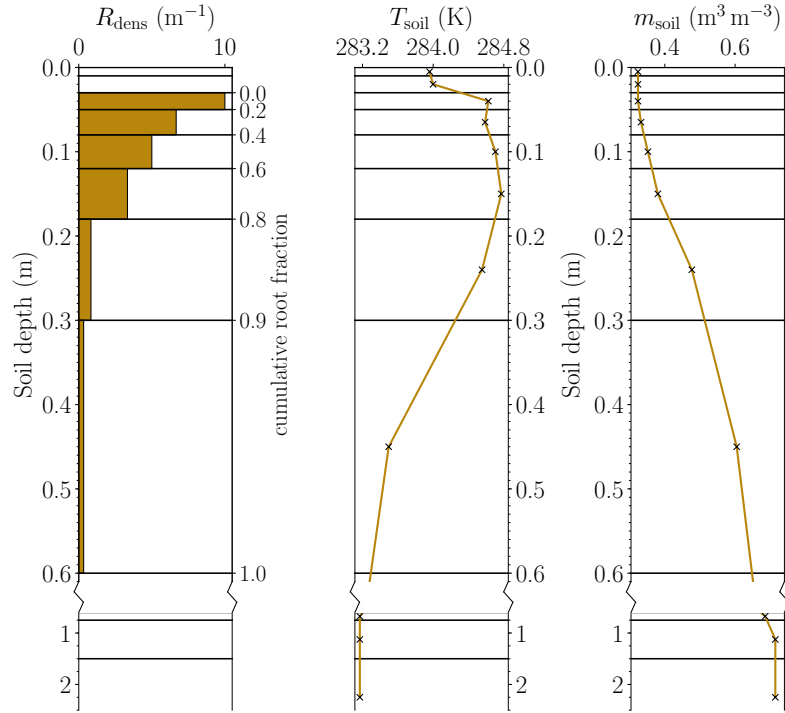
**Table 5.** Overview of the case study and their changes relative to case REF.

Case	changes to case REF
REF	-
ALBE_2404	shortwave albedo (at 60°) of <del>0.24</del> 0.04
ALBE_4424	shortwave albedo (at 60°) of <del>0.44</del> 0.24
ADV_tq	advection of $T$ and $q$ <del>in</del> at all heights according to mean change in radiosounding data between 2.5 km and 4 km
CAP_2e4	$C_0 = 2 \times 10^4 \text{ J m}^{-2} \text{ K}^{-1}$
CLEARSKY	clear-sky radiation model
COND_2	$\Lambda_{\text{skin}} = 2 \text{ W m}^2 \text{ K}^{-1}$
COND_6	$\Lambda_{\text{skin}} = 6 \text{ W m}^2 \text{ K}^{-1}$
Dz_2	$\Delta z = 2 \text{ m}$ and $\Delta x = \Delta y = 5 \text{ m}$
EMIS_95	$\epsilon = 0.95$
EMIS_100	$\epsilon = 1.00$
HUMID_dry	initialization with $q_{v,k} = 0$
HUMID_sat	initialization with $q_{v,k} = q_{v,\text{sat}}$
LAI_05	$LAI = 0.5 \text{ m}^2 \text{ m}^{-2}$
LAI_3	$LAI = 3 \text{ m}^2 \text{ m}^{-2}$
ROUGH_01	$z_0 = 0.01 \text{ m}$ and $z_{0,h} = 1.57 \times 10^{-6} \text{ m}$
ROUGH_001	$z_0 = 0.1 \text{ m}$ and $z_{0,h} = 1.57 \times 10^{-7} \text{ m}$
RRTMG	Rapid Radiation Transfer Model for Global Models (RRTMG)
SOIL_2	$\alpha_{\text{VG}}, l_{\text{VG}}, n_{\text{VG}}, \gamma_{\text{sat}}$ as in soil type 2 (in the uppermost 60 cm)
SOIL_4	$\alpha_{\text{VG}}, l_{\text{VG}}, n_{\text{VG}}, \gamma_{\text{sat}}$ as in soil type 4 (in the uppermost 60 cm)
TEMP_9	initialization with $T_0 = 280.15$ (ca. 9 °C)
TEMP_11	initialization with $T_0 = 286.15$ (ca. 11 °C)

Case REF is driven by external forcing of incoming short- and longwave radiation taken from the Cabauw measurements. ~~In this way, the~~ The effects of high altitude aerosols, moisture and clouds are included in this forcing. Accordingly, the ~~degree of freedom is~~ degrees of freedom are reduced and we can focus on parameters of the LSM, rather than additional uncertainties of a radiation model. Nonetheless, we performed sensitivity tests using the Rapid Radiation Transfer Model for Global Models (RRTMG, Clough et al., 2005) as well as a clear-sky radiation parameterization, which are described in detail in Maronga et al. (2020). The longwave outgoing radiation of the surface is calculated from the skin-layer temperature using the Stefan-Boltzmann law. The long- and shortwave albedos ~~of~~ for diffusive radiation are set to 0.34 and 0.14, respectively, to fit the dominating grassland. Albedos ~~of~~ for direct radiation are calculated according to Briegleb (1992) considering a weak solar zenith-angle dependence as such that their direct values equal the diffusive ones at 60°.

The model domain for case REF is  $(x \times y \times z)$  2000 m  $\times$  2000 m  $\times$  4317 m with a horizontal and vertical grid spacing of 50 m and 10 m, respectively. A grid sensitivity study ~~was~~ is carried out to justify this ~~choice~~ (see grid spacing (cf. Sect. 5). Starting

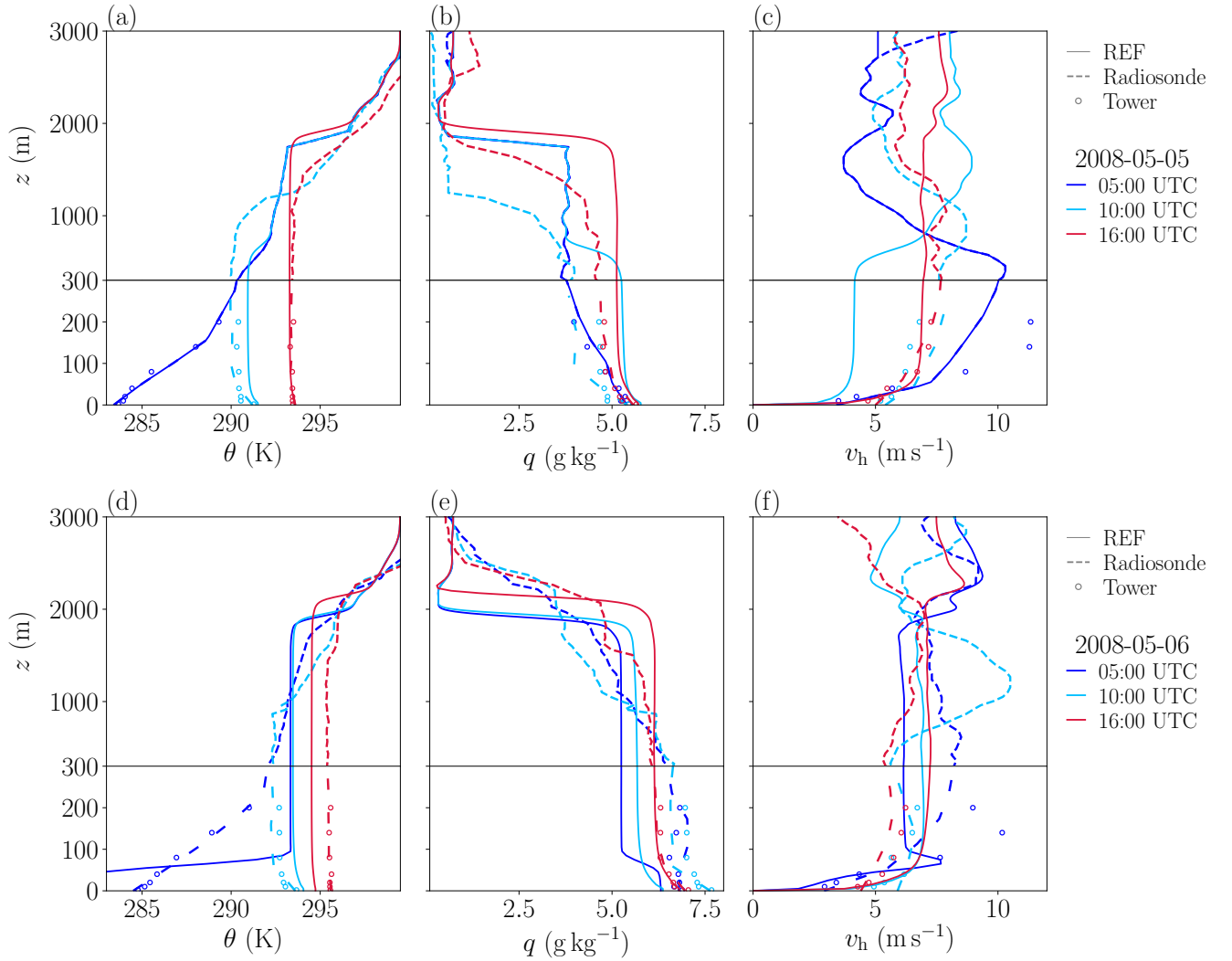




**Figure 1.** Vertical soil layer setup of root density ( $R_{\text{dens}}$ , left) as well as initial profiles of temperature (middle) and moisture (right). Note the broken vertical axis with a changed linear increment in the deeper layers. The root fraction of each soil layer (cf.  $R_{\text{fr}}$  in Table 4) is the difference of the cumulative root fraction (Eq. (34), shown on the vertical axis) between two layers.

at 2000 m, i.e. above the boundary-layer top ( $> 2000 \text{ m}$ ), a vertical grid stretching is applied with a stretching factor of 1.08 and a maximum vertical grid spacing of 100 m. Initial profiles of temperature, humidity and horizontal wind of case REF are taken from radiosounding data and are shown in Fig. 2. The horizontal wind equals the initial  $u$  component of the wind vector, i.e. at the beginning of the simulation there is no wind turning with height due to Coriolis force. The geostrophic wind at the domain top of the domain is set to  $7 \text{ m s}^{-1}$  and  $0 \text{ m s}^{-1}$  for the  $u$  and  $v$  component, respectively. The lateral boundary conditions are cyclic.

In addition to simulation REF sensitivity simulations are performed. All sensitivity simulations are based on the setup of the reference simulation and only differ by the respective parameter to be analyzed. one specific parameter that is varied in a reasonable range. With these sensitivity simulations we do not intend to give a comprehensive parameter study which usually covers a wide range of parameters. Nonetheless, the simulations provide an idea on how sensitive the model reacts on specific parameters or processes included. This is mainly motivated by the fact that in many simulation setups the respective input data for the land-surface parameters are often not or only roughly available, posing a rather high uncertainty in these data. An overview of the sensitivity simulations and their well-defined change compared to case REF is given in Table 5.



**Figure 2.** Vertical profiles of  $\theta$ ,  $q$  and  $v_h$  measured by radiosonde (dashed lines) and tower (point markers) at the CESAR site as well as simulated profiles of case REF (continuous lines) during the two days period. Note that the lower 300m are shown with a higher vertical resolution than the layers above to better visualize individual tower measurement (black horizontal line indicates the break). **In the left panel, the orange line is partially hidden behind the blue one.**

## 5 Results

### 5.1 Boundary-layer profiles

At first, we will look at the vertical profiles of the radiosonde and the tower measurements. Figure 2 shows vertical profiles of  $\theta$ ,  $q$  and  $v_h$  indicating the evolution of the boundary layer during the considered period of time in Cabauw at the times

360 of the radiosonde ascents. Both nights show a stable nocturnal boundary layer before sunrise (at 05:00 UTC). A nighttime low-level jet is observed in the horizontal wind. Profiles of potential temperature and water vapor mixing ratio show a well-mixed convective boundary layer at 10:00 UTC and 16:00 UTC of either day. Over the course of the two days, the depth of the convective boundary layer increases. The profiles of  $v_h$  show a mismatch between the radiosounding measurements and the tower data. This is explained by different analysis timescales. ~~One-the-one-hand, the~~ The radiosonde records instantaneous values with a sampling frequency of 0.1 Hz, which results in five recording heights below 300 m. ~~On-the-other-hand~~ Conversely, data from the cup anemometers at the mast ~~is-are~~ temporally averaged over 10 min, based on a 3 s running mean calculated with an update frequency of 4 Hz (Bosveld, 2020b). Thus, the mean tower profiles can be over- or underestimated by the instantaneous values of the radiosonde. Note that we initialize the simulations with data from the radiosonde ascent, because it reaches high enough, but ~~the-comparison-later-is-~~ comparisons are made against tower data due to ~~its-the~~ higher temporal resolution.

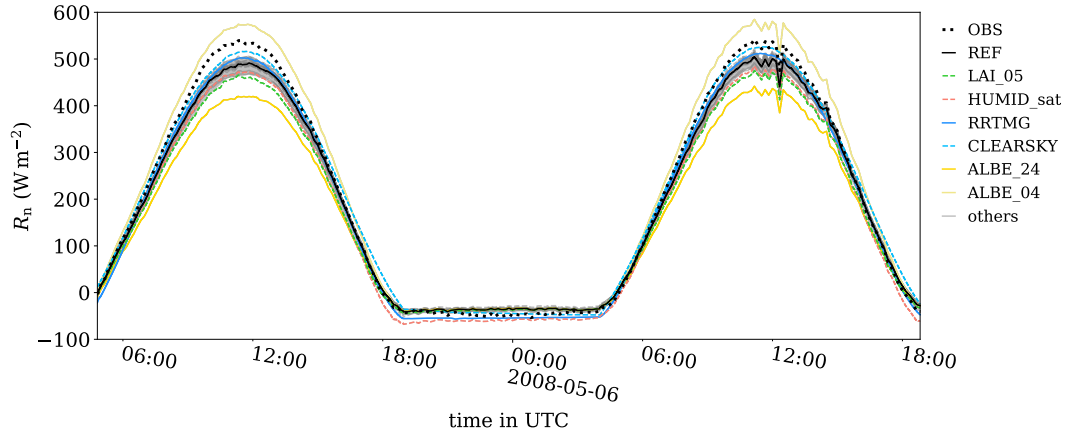
Figure 2 also shows the vertical profiles of the reference simulation, which are temporally averaged over 15 min and horizontally averaged over the whole domain. A comparison of observed profiles of  $\theta$  and  $q$  with those of the LES show that ~~on-the-first-day, the-simulated-boundary-layer-is-too-shallow-~~ the boundary layer depth is underestimated by the simulation at 10:00 UTC on the first day. One hypothesis to explain this is that the turbulence development during model spin-up, i.e. ~~in-on~~ the first morning, ~~is-slower-~~ takes longer than in reality. ~~At-this-time, the-~~ The horizontal wind speed ~~shows-no-good-agreement-between-~~ of the model and observations do not agree at 10 :00 UTC on the first day. At 16:00 UTC of the first day, the simulation slightly overestimates the boundary layer depth despite a fairly good agreement of model and observations regarding mixed-layer temperature and humidity values. The wind profiles, ~~however, agree well-~~ agree well with observations. At night (2008-05-06 05:00 UTC), the near-surface temperature is significantly lower (out of range at ca. 281 K) than measured. ~~Another-difference-is-that-the-nocturnal-boundary-layer-is-too-shallow-in-the-simulation-~~ Similar to the mixing layer, the simulation also underestimates the nocturnal boundary layer depth. At the same time, the near-surface humidity shows small differences between model and reality. The simulated horizontal wind speed also depicts a low-level jet, but compared with observations it occurs closer to the surface, in accordance with the simulated nocturnal boundary-layer depth. Even though the LES cannot reproduce the nocturnal ~~situation-very-well-~~ observations precisely, the mixed-layer quickly develops ~~in-~~ the next morning (2008-05-06 10:00 UTC), which is in agreement with findings of van Stratum and Stevens (2015). At 10:00 UTC of the second day, the simulation indicates a warmer and deeper boundary layer compared to the ~~observation-~~ observations, which could be caused by advection processes in reality modifying the residual layer and thus the boundary-layer evolution during the morning transition. The wind profiles agree fairly well with observations. The temperature profile at 2008-05-06 16:00 UTC shows that the simulation underestimates the mixed-layer temperature ~~but-instead-has-a-higher-boundary-layer-top,-~~ suggesting that the, but overestimates the depth of the boundary layer. This overestimation suggests that the total energy input into the boundary layer is similar to the observations, but distributed over a deeper layer. The humidity profiles agree well with observations in the lowest 1000 m, but deviate above this height in accordance with the differences in boundary-layer depth. The horizontal wind is slightly overestimated at 16 :00 UTC on the second day. In general, the simulated profiles are much more constant with height in the well-mixed layer, because they depict domain averages, as opposed to the local measurements,

395 so that a direct comparison is inherently improper. Above the boundary layer, in the free atmosphere, synoptic-scale processes dominate in reality. ~~Since-Because~~ we did not consider these processes in our simulations, profiles may deviate. Given the ~~fact-that-the~~ surface forcing in this particular case was the dominant forcing for the development of the boundary layer, this deficiency should not compromise the present evaluation study.

## 5.2 Evaluation of energy-balance components

### 400 5.2.1 Net radiation

Figure 3 shows the time series of surface net radiation which follows a diurnal cycle typical for clear-sky conditions. At noon, the reference simulation underestimates  $R_n$ , whereas it overestimates  $R_n$  at night, i.e. it is less negative by approximately  $30 \text{ W m}^{-2}$ . The nocturnal differences are due to much lower surface temperatures of 280 K in case REF vs. 285 K in the observations (not shown). ~~Since-we prescribe the incoming radiation in case REF, the~~ The reason for this is that the long-wave outgoing radiation flux is underestimated, because the incoming radiation in case REF is prescribed. At around 12:00 UTC of the second day, clouds influence the net radiation, which is indicated by fluctuations in the curve. ~~Since-Because~~ case REF is driven with  $SW_{\downarrow}$  taken from CESAR data, this cloud effect on the surface net radiation is included and can also be seen in the other surface fluxes (cf. Fig. 4, 5, 6). The maximum and minimum values of net radiation of case REF indicate little horizontal variation of approximately  $\pm 10 \text{ W m}^{-2}$  at noon. Figure 3 also depicts the effect of different land-surface properties and simulation setups on the surface radiation. ~~Please-note-that-At this point we want to emphasize again that the discussion of the sensitivity simulations is not intended to find the perfect parameter combination but to give an estimate of the sensitivity of the modeled energy balance components on specific land-surface parameters and outline their complex interactions among each other. Note that~~ we will only highlight the cases ~~which-that~~ lead to the largest differences compared to case REF for the respective variable. In Fig. 3, these are mostly the changes to the albedo and the radiation models, as well as cases LAI\_05 and HUMID\_sat which have the largest impact on the surface net radiation. The most obvious differences ~~are seen in case of a changed albedo occur in the albedo sensitivity simulations~~. An increase of the albedo (case ALBE\_44) ~~leads to a decrease in-24) decreases~~  $R_n$  and vice versa (case ALBE\_24)04), with maximum deviation ~~to-from~~ case REF of ~~about~~ approximately  $\pm 100 \text{ W m}^{-2}$  at noon, ~~indicating-~~ This deviation suggests that mismatches in the estimated albedo cause significant errors on the surface net radiation during daytime. However, ~~also~~ the spread of the non-highlighted simulations (gray) is ~~about~~ approximately  $50 \text{ W m}^{-2}$  at noon, ~~meaning-that-also-~~ This spread implies that variations in the surface parameters (e.g., emissivity, roughness length) also affect the surface net radiation significantly. ~~Using-the~~ Besides this, cases EMIS\_95, EMIS\_100, ROUGH\_01, ROUGH\_001, SOIL\_2, and SOIL\_4 are always among the non-highlighted cases in Figs. 4, 5, and 6. With RRTMG, radiative fluxes are calculated for each horizontal grid box of the LES based on timeand-geographical-locationas well-as-, geographical location, air pressure, and-as well as local profiles of temperature and humidity (Clough et al., 2005) instead of being taken from the measurements. ~~One-A~~ A key feature of the RRTMG ~~;-which-is-neither-included-in-observations nor-in-the-clear-sky-model, is the-is the~~ direct cooling of air due to longwave radiative flux divergence, particularly during nighttime. This ~~can-result~~ results in cooling rates ~~in-on~~ the order of  $0.1$  to  $0.3 \text{ K h}^{-1}$  (see, e.g. Maronga and Bosveld, 2017,

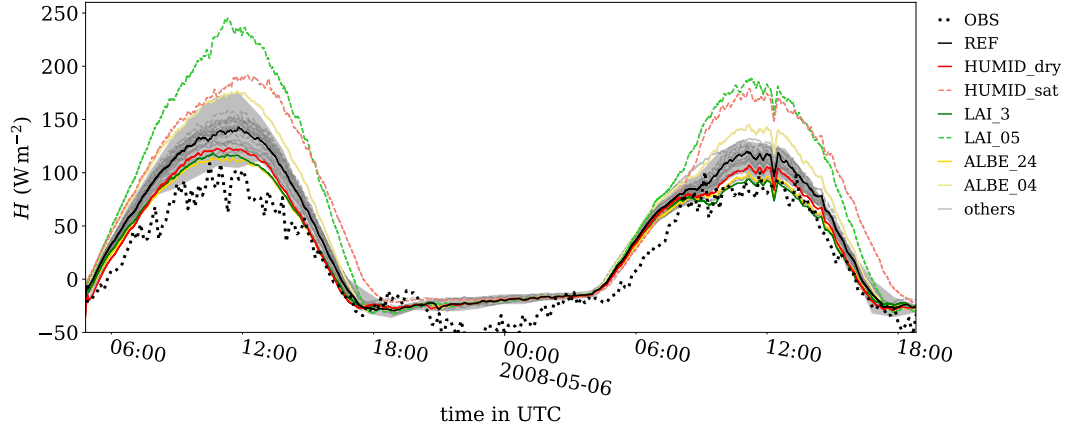


**Figure 3.** Timeseries of surface net radiation  $R_n$  as measured at Cabauw (OBS) and domain-averaged  $R_n$  for all ~~simulation~~ simulated cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray. The gray, transparent area shows minimum and maximum values of the net radiation of case REF based on instantaneous hourly horizontal cross-sections.

their Fig. 1). This effect is not included in case REF, where incoming radiation is prescribed, and in case CLEARSKY. In case RRTMG, the simulated  $R_n$  approaches that of the observations during day and night, though it slightly underestimates the surface net radiation at noon by ~~about~~ approximately  $30 \text{ W m}^{-2}$ . ~~However, we have to note~~ Note that we use default input of the RRTMG with standard profiles of water vapor, other trace gases, and aerosol concentration above model top, which ~~do~~ does not necessarily reflect the real conditions at Cabauw during the simulated time period and may impact the simulated surface net radiation, ~~too.~~ Using the as well. The clear-sky model ~~, the simulated~~ simulates  $R_n$  ~~becomes even better during~~ more accurately during the day, but also marginally overestimates net radiation during the night by 10 to  $20 \text{ W m}^{-2}$ . Case LAI\_05 has a smaller  $R_n$  ~~compared to~~ than case REF, because a smaller  $LAI$  significantly decreases  $LE$  and increases  $H$ , which leads to higher surface temperatures during ~~day~~ resulting the day. Higher temperatures result in higher  $LW_{\uparrow}$ . In the case of a saturated humidity mixing-ratio (case HUMID\_sat)  ~~$LE$ ,  $H$  and  $G$  are significantly altered in a way that higher~~ surface temperature is ~~simulated during~~ higher during both, day and night. ~~This, due to significant alterations to  $LE$ ,  $H$  and  $G$ .~~ This higher surface temperature again leads to higher  $LW_{\uparrow}$  and ~~hence,~~ hence, lower  $R_n$ . Except for ALBE\_24\_04, the simulations ~~tend to underestimate the surface net radiation during daytime,~~ while the simulated surface net radiation tend to overestimate the observed ~~one during~~ surface net radiation during the night.

### 5.2.2 Sensible heat flux

Figure 4 shows that the observed surface sensible heat flux reaches ~~at maximum a~~ maximum of approx.  $100 \text{ W m}^{-2}$  and ~~at minimum a~~ minimum of  $-50 \text{ W m}^{-2}$ . Compared to observations (OBS), case REF significantly overestimates  $H$ , ~~especially~~ at noon by up to approximately  $40 \text{ W m}^{-2}$ , with the largest overestimation at noon. Moreover, the ~~simulated model~~ overestimates  $H$  ~~differs from observations~~ in the morning and afternoon hours, where it is positive earlier and later, respec-



**Figure 4.** Timeseries of surface sensible heat flux  $H$  as measured at Cabauw (OBS) and domain averaged  $H$  for all ~~simulation~~ simulated cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray. The gray, transparent area shows minimum and maximum values of the sensible heat flux of case REF based on instantaneous hourly horizontal cross-sections.

tively. At night, case REF shows a constantly increasing  $H$  from  $-25$  to  $-15 \text{ W m}^{-2}$ , whereas the ~~observation~~ observations show a secondary minimum of ca.  $-50 \text{ W m}^{-2}$  between 20:00 UTC and 23:00 UTC. Here, we note that the depicted sensible heat flux from the simulations are domain-averaged values. Considered locally in the simulation, we can detect similar

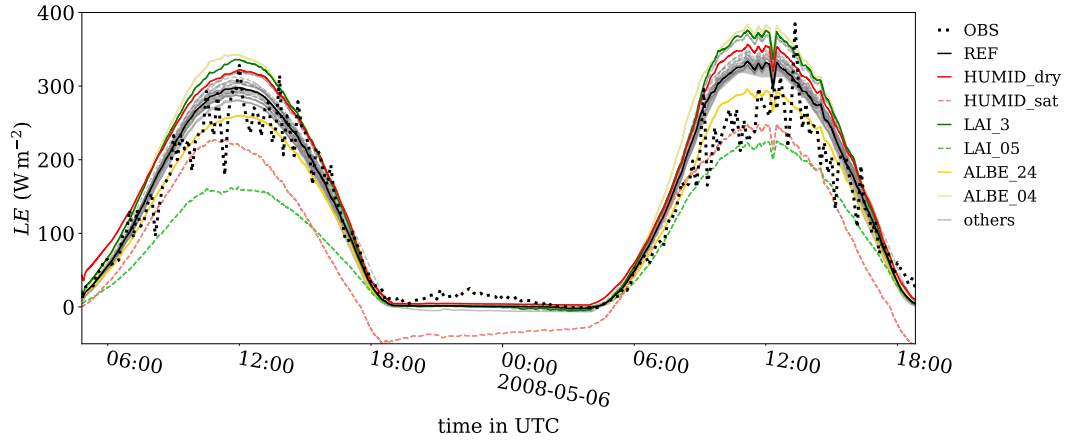
450 temporal fluctuations in  $H$ , though with a smaller amplitude than those observed at the CESAR site. The horizontal variation (maximum and minimum values) of the sensible heat flux of case REF is approximately  $\pm 30 \text{ W m}^{-2}$  at noon. Like case REF, all sensitivity simulations tend to overestimate the observed heat flux ~~;~~ while and the spread among the considered cases is significant. At noon, ~~where~~ when the spread among all simulations is largest, ~~it reaches values up to differences in simulated heat fluxes approach~~ 120  $\text{W m}^{-2}$ . ALBE\_44, LAI\_3, and HUMID\_dry ~~best meet~~ show the best agreement with the observa-

455 tions. With a higher albedo (ALBE\_44) 24, the available energy at the surface ~~becomes lower (cf. is lower (see Fig. 3) leading to~~ -. The lower available surface energy results in smaller fluxes of  $H$  and  $LE$ . Besides changing the albedo, case LAI\_3 and LAI\_05 show the largest impact on  $H$ , ~~as~~ because with lower and higher  $LAI$  the available energy is preferentially partitioned into  $H$  and  $LE$ , respectively. Similarly, the available energy is also preferentially partitioned into  $H$  and  $LE$  for humid and dry air, indicated by HUMID\_sat and HUMID\_dry, respectively.

### 460 5.2.3 Latent heat flux

Figure 5 shows timeseries of the latent heat flux. The observations range from 0 during ~~night to about the night to approximately~~ 300  $\text{W m}^{-2}$  at noon. Case REF matches the observation reasonably well during day- and nighttime, even though it overestimates  $LE$  during the second day. The maximum and minimum values of latent heat flux of case REF indicate little horizontal variation of approximately  $\pm 10 \text{ W m}^{-2}$  at noon. Having a lower  $LE$  than case REF, ALBE\_44 ~~best meets~~ 24 best matches

465 the observed  $LE$  of the second day. ~~Besides this~~ Additionally, case LAI\_05 significantly underestimates  $LE$  by preferentially



**Figure 5.** Timeseries of latent heat flux  $LE$  as measured at Cabauw (OBS) and domain-averaged  $LE$  for all simulated cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray. The gray, transparent area shows minimum and maximum values of the latent heat flux of case REF based on instantaneous hourly horizontal cross-sections.

partitioning the available energy into  $H$  (Fig. 4). Moreover, case HUMID\_sat stands out because it shows a negative  $LE$  during night, ~~which is explained by due to~~ dew formation in the model. ~~Compared with the observations, where negative  $LE$  is not observed, this suggests~~ Observations suggest that dew formation ~~has not been was not~~ observed in Cabauw. ~~The peak, because there is no negative  $LE$ . The maximum~~ spread of all non-highlighted sensitivity studies is ~~around approximately~~ 50 W m<sup>-2</sup> and   
 470 . This spread is up to 170 W m<sup>-2</sup> of the highlighted for the highlighted sensitivity studies. From Fig. 4 and Fig. 5 we ~~can~~ calculate the daytime Bowen ratio  $\beta_0 = \frac{H}{LE}$  (not shown). We find that the difference between the Bowen ratio of the simulations and that observed at Cabauw ( $\beta_{0,\text{sim}} - \beta_{0,\text{obs}}$ ) ranges from  $-0.1$  to  $+0.2$ , except for cases HUMID\_sat and LAI\_05, which ~~show significantly lower~~ have significantly higher Bowen ratios compared to observations. As opposed to other models, which were shown to systematically overestimate the Bowen ratio measured in Cabauw during summer (Schulz et al., 1998; Sheppard and   
 475 Wild, 2002), we cannot identify a systematic bias for Bowen ratio of PALM's LSM by means of the analyzed period.

Taking into account that modeled surface net radiation is underestimated but the observed modeled latent and sensible fluxes both exceed the observed values, while the modeled Bowen ratio matches with the observed values, one can conclude that the remaining energy is included the ground heat flux, which will be discussed in the next paragraph.

#### 5.2.4 Ground heat flux

480 Figure 6 shows the timeseries of the ground heat flux, which reaches  $60 \text{ W m}^{-2}$  during daytime in the observations. With respect to the amplitude of  $G$ , case REF shows good agreement with the observation. However, the shape of the timeseries is discernibly different. ~~At During~~ daytime, OBS ~~shows a more sinusoidal has a pseudo-sinusoidal~~ shape, whereas the simulations ~~show have~~ a more humped-shaped diurnal variation. We attribute this to the method used to derive the observed ground heat flux, which involves the average soil heat flux in at 10 cm and 5 cm depth ~~and the soil temperature difference between 0 and~~



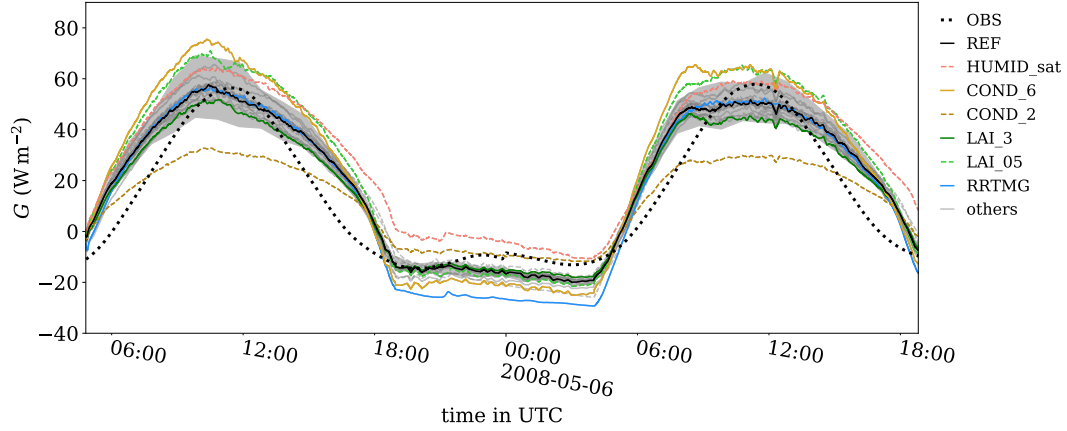
485 ~~2 cm (Bosveld, 2020b)~~ (details in Sec. 3). In the model, ~~on the other hand, the~~ the ground heat flux is parameterized according to Eq. (6) and thus only the temperature gradient between the surface and 0.5 cm are ~~taken into account~~ accounted for. Hence, the simulated  $G$  resembles the diurnal cycle of the surface temperature, while the observed  $G$  correlates with the soil temperature. Observations and case REF agree well at noon with values of  $G$  around  $55 \text{ W m}^{-2}$ , whereas in the afternoon at 15:00 UTC case REF is almost  $10 \text{ W m}^{-2}$  higher than the observations. ~~It strikes that~~ The horizontal variation (maximum and minimum values)  
490 of the ground heat flux of case REF is approximately  $\pm 10 \text{ W m}^{-2}$  at noon. Remarkably, all simulations have a fairly small spread at ~~that point in this~~ time. During the day, changing the heat conductivity between the skin layer and uppermost soil layer has the largest impact on the ground heat flux, with smaller and larger  $G$  observed in case COND\_2 and COND\_6, respectively. This can be attributed to the linear relationship between  $G$  and  $\Lambda$  in Eq. (6). Also, the ~~simulation cases with a different LAI~~ sensitivity cases show a relatively large deviation from case REF at noon, with increased (decreased)  $G$  for LAI\_05 (LAI\_3).  
495 ~~For example, with a smaller leaf area and thus less transpiration the available energy is preferentially partitioned into  $G$  and  $H$  (see Fig. 4) rather than into  $LE$ .~~ At noon, the spread among all simulations is ~~about~~ approximately  $40 \text{ W m}^{-2}$ , whereas the non-highlighted cases show a maximum deviation from REF of no more than  $\pm 10 \text{ W m}^{-2}$ . In the night, cases HUMID\_sat and RRTMG stand out. In both cases the controlling variable is the (near-) surface temperature (see Fig. 9). In case RRTMG, the surface cools ~~down~~, while the soil temperature does not change in the same amount (see Fig. 8), which leads to a strong ground  
500 heat flux directed towards the surface. In case HUMID\_sat, the surface ~~stays~~ remains relatively warm, ~~therefore it which~~ results in a small negative  $G$ . Except for cases HUMID\_sat and COND\_2, the model tends to simulate a stronger upward directed ground heat flux at night compared to the observations. Compared to the observations, case COND\_2 overestimates  $G$  during ~~night the night~~, whereas it underestimates  $G$  at noon, which points to a stability dependence of the conductivity  ~~$\Lambda$~~   $\Lambda$ . Even though this is technically realized in the code, the standard values for short grass do not differ between stable and unstable  
505 conditions due to a lack of knowledge of the correct relationship between  ~~$\Lambda$~~   $\Lambda$  and stability. As cases COND\_2 and COND\_6 do not have a significant impact on any of the other variables, we can infer that uncertainties in modeling  $G$  ~~is~~ are relatively small. Nonetheless, ~~please~~ note that we only analyzed a short period of two days. For longer simulation times covering multiple days, e.g., for heat-wave scenarios where heat storage within the soil becomes important, ~~this it~~ might become relevant ~~again to~~ accurately simulate  $G$ . Furthermore, ~~it shall be mentioned that~~ the ground heat flux is small compared to the surface net  
510 radiation, as well as the surface sensible and latent heat fluxes.

As the modeled and observed ground heat flux are in a similar range but the modeled surface net radiation underestimates the observed values, the available energy in the simulation is smaller compared to the observation. However, at the same time the modeled fluxes, in which the available energy is partitioned into, exceed the observed values. One possible explanation of this discrepancy could be mismatches coming with the observed values of the energy-balance components which will be  
515 discussed in the next paragraph.

### 5.2.5 Uncertainty in model-observation comparison due to energy-balance non-closure

At this point we want to briefly address the well-known problem of energy-balance non-closure and how this affects our model-observation comparison. Similar to many other sites, ~~also~~ eddy-covariance measurements at the Cabauw site suffer





**Figure 6.** Timeseries of ground heat flux  $G$  as measured at Cabauw (OBS) and domain-averaged  $G$  for all simulated cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray. The gray, transparent area shows minimum and maximum values of the ground heat flux of case REF based on instantaneous hourly horizontal cross-sections.

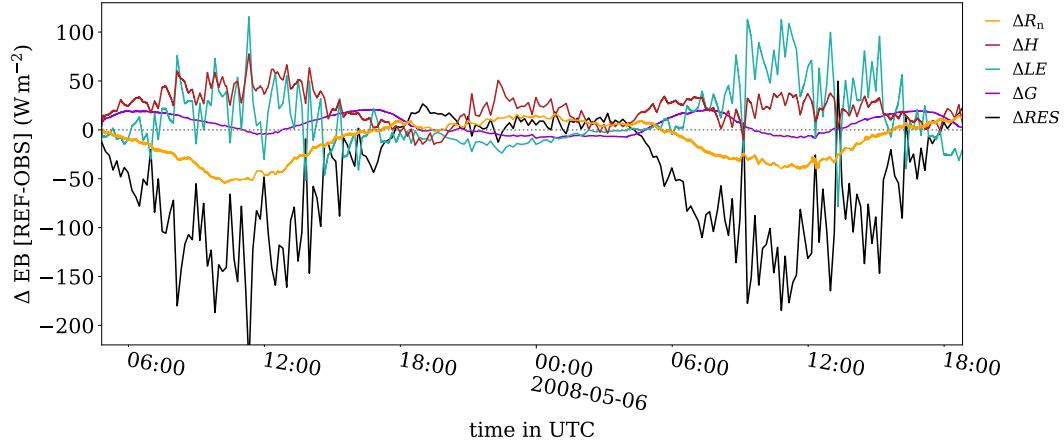
from ~~the well-known problem of~~ energy-balance non-closure (de Roode et al., 2010). To date, the leading hypothesis is that

520 low-frequency contributions are inherently not captured by the eddy-covariance method, leading to the situation that the sum of the surface heat fluxes underestimates the available energy by 10 to 30 % (Foken et al., 2011). Figure 7 shows the ~~timeseries of the residual~~ residual of the energy balance in the observations (RES), as well as the ~~individual particular~~ differences between the simulated (~~case REF~~) and ~~reference case REF~~ and the observed energy balance components. Note that the ~~shown residual is effectively the residual of the observations, because the energy balance in the model is zero at all times our simulations is closed~~ (according to Eq. (1) with  $C_0 = 0 \text{ J m}^{-2} \text{ K}^{-1}$ ), i.e. no storage and thus no residual for the simulations is shown.

The residual in the observation exhibits a diurnal cycle with only small positive values of ~~about approximately~~  $15 \text{ W m}^{-2}$  during nighttime, while during daytime the residual indicates that up to  $150 \text{ W m}^{-2}$  are missing in order to close the surface energy balance. The differences between the simulated and observed  $H$  and  $LE$  correlate negatively with the residual, ~~which indicates that the simulated~~. This suggests that either the model overestimates  $H$  and  $LE$  are overestimated compared to the

530 ~~observations. In reverse, however, this may also suggest that the observed~~, or conversely, the observations underestimate  $H$  and  $LE$  are underestimated, even though we explicitly note that we cannot know from the simulations how the missing energy is partitioned onto the individual energy-balance components in reality as both, observation and simulation may contain a bias. Taking this into account due to low-frequency losses in the EC-method. Figure 7 thus emphasizes the uncertainty inherent in such a model-observation comparison, especially during daytime. In this respect, we summarize that the four components of the

535 energy balance are represented reasonably well by the LES-LSM interface. ~~Also, the daytime Bowen ratio agrees fairly wells with that observed at Cabauw. By contrast, climate models have often shown to overestimate the Bowen ratio in summertime (Schulz et al., 1998; Sheppard and Wild, 2002).~~

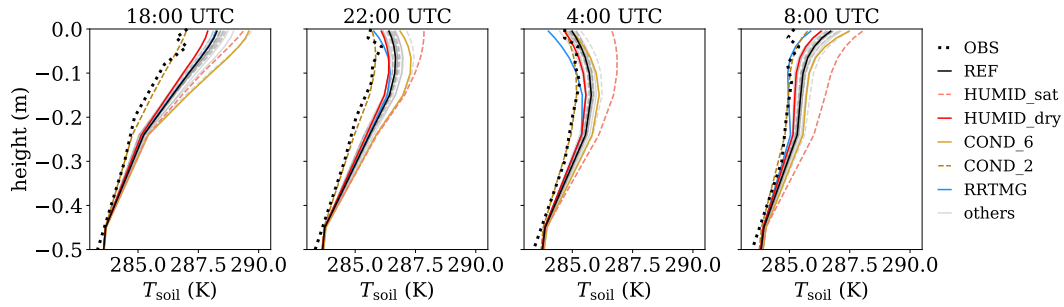


**Figure 7.** Timeseries of the individual differences between case REF and observations for the energy balance (EB) components  $R_n$ ,  $H$ ,  $LE$ , and  $G$ , as well as the residual  $RES = R_n - H - LE - G$  from the Cabauw observations.

### 5.3 Soil temperature

Figure 8 shows mean profiles of soil temperature,  $T_{\text{soil}}$ , averaged over the whole domain and over 15 min. ~~Please note~~ Note that there are doubts about the quality of the CESAR soil temperature observations, because of problems during calibration (Bosveld, 2020b) and the speculation that the sensors had sunken deeper into the soil over time (Bosveld, 2020a). Hence, observations ~~cannot be used as~~ are not a reliable reference to evaluate the model. Nevertheless, we ~~will~~ discuss the variation among the different simulations. In the late afternoon of the first day (18:00 UTC), case REF shows a continuous decrease with depth from 288 K close to the surface to 283.5 K ~~in at~~ at  $-50$  cm. During ~~the~~ the night, the soil begins to cool down starting at the surface, where it reaches 285 K at 04:00 UTC. The maximum soil temperature ~~is now found in~~ occurs at a depth of  $-15$  cm. ~~As soon as the net radiation becomes positive~~ The near-surface soil temperature responds quickly to positive net radiation (at 08:00 UTC) ~~the soil temperature close to the surface follows~~ such that the maximum ~~is again found~~ temperature is again close to the surface. The parameters with the largest spread - up to 4 K during day and 2 K during night - are, like for  $G$ , the conductivity and atmospheric humidity. The non-highlighted cases have a spread of ~~about~~ approximately 1 K ~~indication that~~ This spread indicates that, e.g. changing radiation or vegetation parameters have only minor effect on the soil temperature. A higher conductivity (COND\_6) ~~leads to results in~~ higher soil temperatures, especially in the late afternoon (18:00 UTC), because the heat of the surface net radiation is more easily conducted into the soil (cf. Fig. 6); correspondingly COND\_2 has smaller  $T_{\text{soil}}$ . During ~~the~~ the night, the lower atmosphere in case HUMID\_sat ~~does not cool down as much as~~ cools less rapidly than case REF (cf. Fig. 9), therefore the soil remains relatively warm, too.

### 5.4 Evaluation of atmospheric quantities



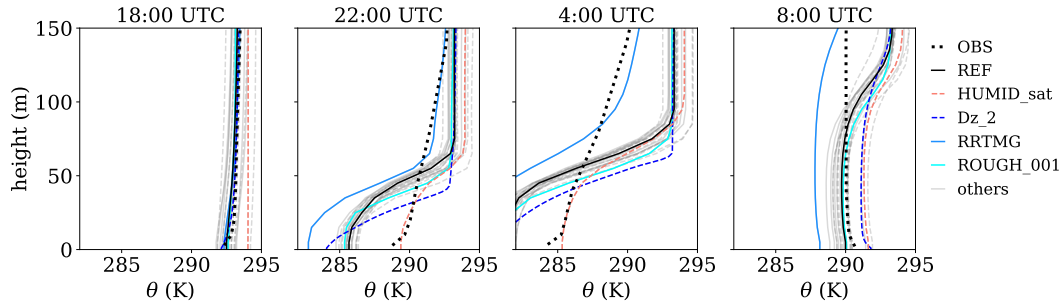
**Figure 8.** Vertical profiles of soil temperature as measured at Cabauw (OBS) and domain-averaged soil temperature for all simulated cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.

~~Next, we will~~ In this section, we evaluate the quantities characterizing the lower atmosphere, that is temperature, humidity, and wind speed. First, we ~~will~~ discuss the nighttime situation followed by the daytime situation. ~~Finally, results from~~ Then, we discuss the results of the sensitivity study ~~will be discussed~~.

#### 5.4.1 Stability of the nocturnal boundary layer

Figure 9 shows the vertical profiles of potential temperature during the night. At 18:00 UTC, the lowest levels start to cool, while the upper levels ~~still indicate~~ remain in a vertically well-mixed state; observation and case REF agree well. At 22:00 UTC, the stable boundary layer of case REF ~~has reached~~ reaches a height of ~~about~~ approximately 75m and has a mean potential temperature gradient of  $0.1 \text{ K m}^{-1}$ . Tower measurements, however, ~~show that it was~~ suggests the atmosphere is less stable ( $0.02 \text{ K m}^{-1}$ ) and the stable layer is at least 200m deep (not shown), ~~above~~. Above the mast we have no data to compare it the simulated temperature profile with. The radiosounding profiles at 05:00 UTC in Fig. 2 show that there is no residual layer, but the stable layer extends up to the capping inversion at ~~about~~ approximately  $z = 1800 \text{ m}$ . The next morning, at 08:00 UTC, observations ~~shows~~ show an already vertically well-mixed lower boundary layer, while in the simulations a stable layer at ~~about~~ approximately 100m is still present and gets eroded due to surface heating and mixing processes. This is consistent with Fig. 11, which indicates a faster temperature increase at  $z = 40 \text{ m}$  compared to the ~~observation~~ observations.

~~One reason for~~ It could be suspected that the misrepresentation of the nocturnal stable layer could be too coarse grid spacing. However, compared to case REF, case Dz\_2 with smaller grid spacing is even more stable during the night hours, ~~having~~. The case with smaller grid spacing has a lower boundary layer depth and cooler near-surface temperature than case REF in agreement with van Stratum and Stevens (2015). ~~This suggests~~ These results suggest that the coarse resolution might not explain the misrepresentation of the stable boundary layer, ~~nonetheless with~~ though. With only one grid sensitivity simulation, a non-linear relationship cannot be fully ruled out. For more information on grid spacing of models in stable conditions, see Sullivan et al. (2016) and Dai et al. (2021). Another parameter affecting the diffusivity is the subgrid-scale scheme of the LES. Yet, the subgrid-scale scheme of ~~Dai et al. (2020)~~ Dai et al. (2021), which is more diffusive in the middle of the stable boundary layer but less diffusive towards the surface, does not change the stability of the simulation (not shown). Alternatively,



**Figure 9.** Vertical profiles of mean potential temperature as measured at Cabauw (OBS) and for all cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.

too low diffusivity could be due to low wind speed, however the simulated wind speed in Fig. 13 agrees well with observations.

580 Note, as  $LW_{\downarrow}$  is prescribed in case REF we can exclude all factors affecting the longwave-incoming radiation, e.g., humidity or clouds, as a possible reason for the misrepresentation of the nighttime stable layer. At 02:00 UTC, case RRMTG, where cooling of the air column by vertical divergence of radiative fluxes is directly considered, reveals an even cooler and more stable (compared to REF) nighttime boundary layer. Too much cooling indicates too small humidity in the atmosphere, which is in agreement to Fig. 12, where humidity is slightly smaller in case RRMTG compared to than the observations between

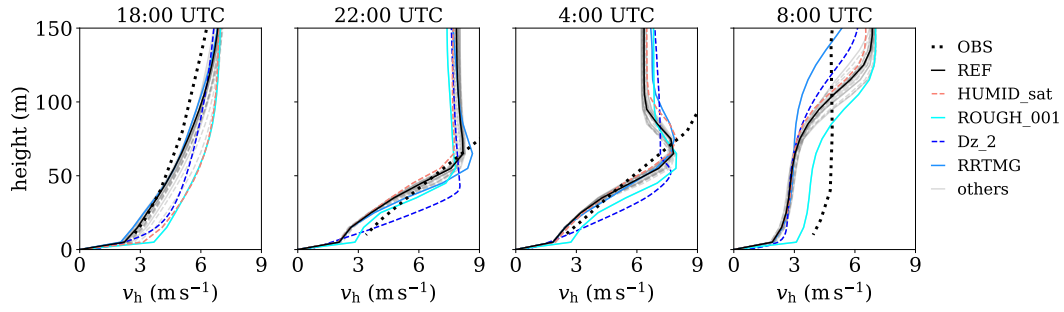
585 00:00 UTC and 06:00 UTC. Furthermore, compared to the ~~observation~~ observations, the simulations indicate a less negative  $H$  (see Fig. 4), though we note that the simulated  $H$  shows is a domain-averaged value, while the observed value is taken from a point measurement with more temporal fluctuation. Nevertheless, according to the less negative  $H$ , the near-surface layer should be less stable in the simulations, ~~which, however, is in contrast to the profiles shown in.~~ This is not the case (see profiles from Fig. 11, ~~which would suggest that the simulated  $H$  should show larger negative values~~). It is unclear, why the nocturnal

590 boundary layer is misrepresented in the LES.

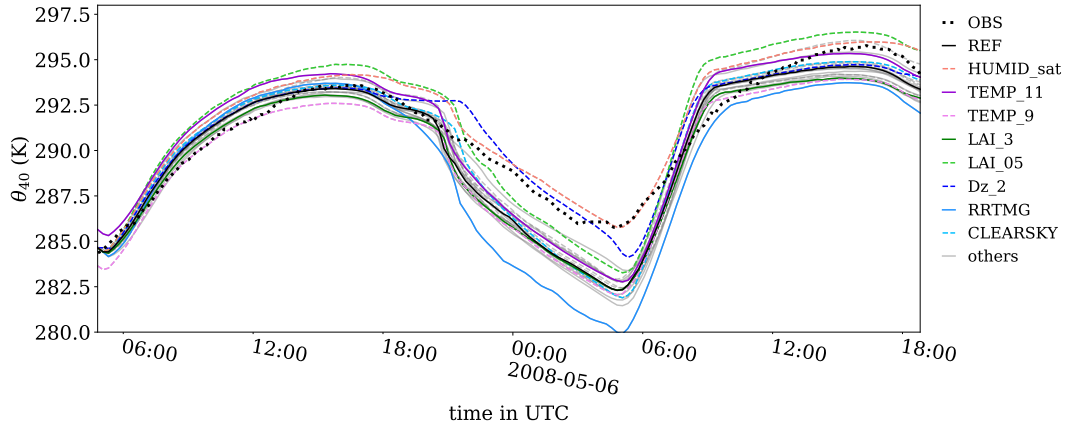
Figure 10 shows vertical profiles of horizontal wind speed  $v_h$ . Again, we find that case REF and all sensitivity studies simulate a shallower nocturnal boundary layer than observed at the CESAR tower (22:00 UTC and 04:00 UTC). A low-level jet develops in reality, as well as in the simulations. ~~From the observational data it is deduced that the maximum wind speed occurs in a height above the tower.~~ According to the radiosounding measurements at 2008-05-06 05:00 UTC (Fig. 2), the

595 maximum wind speed occurs ~~in at~~ approximately 500 m. In the LES (case REF), however, the low-level jet has ~~its maximum between a maximum velocity between a height of 60 m and 70 m height~~ (cf. Fig. 2, 10). In the morning, when convective mixing sets in, the wind profiles become more well-mixed. At 08:00 UTC, we already find a mixed-layer in reality. By contrast, the stable layer has not been fully eroded in the LES, ~~therefore.~~ Therefore, the wind speed in the residual layer (above ca. 100 m) is still higher than that of the lower 100 m. Even though case RRTMG shows a large influence on the whole temperature profile

600 (Fig. 9), ~~its the~~ wind profile only deviates from case REF above the low-level jet and during the morning transition ~~where when~~ the stable layer is already further eroded. Likewise, case HUMID\_sat does not deviate much ~~in from case REF~~ (see Fig. 10).



**Figure 10.** Vertical profiles of mean horizontal wind speed as measured at Cabauw (OBS) and for all cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.



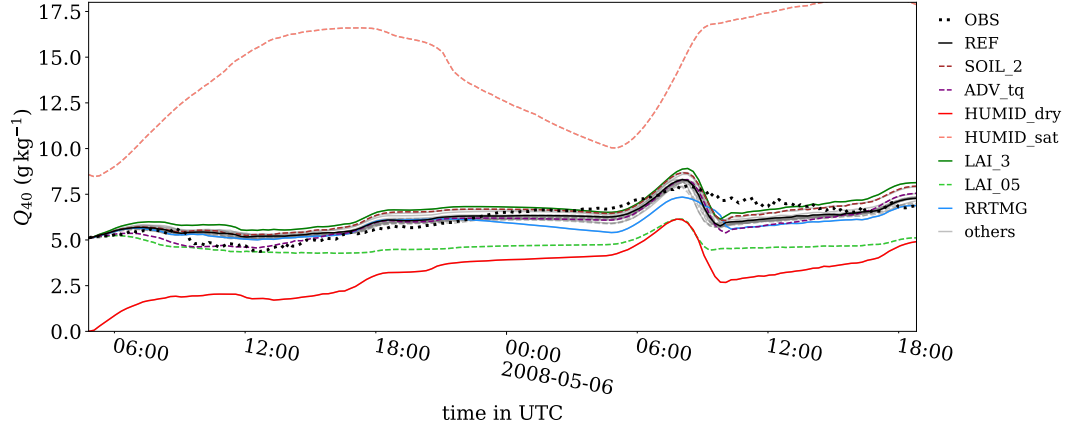
**Figure 11.** Timeseries of mean potential temperature at 40m as measured at Cabauw (OBS) and for all cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.

Conversely, case ROUGH\_001 deviates significantly from case REF in the wind profiles but does not show significant changes in the temperature profiles (Fig. 9). Only case Dz\_2 shows a strong interconnection between temperature and wind profiles.

#### 5.4.2 Potential temperature

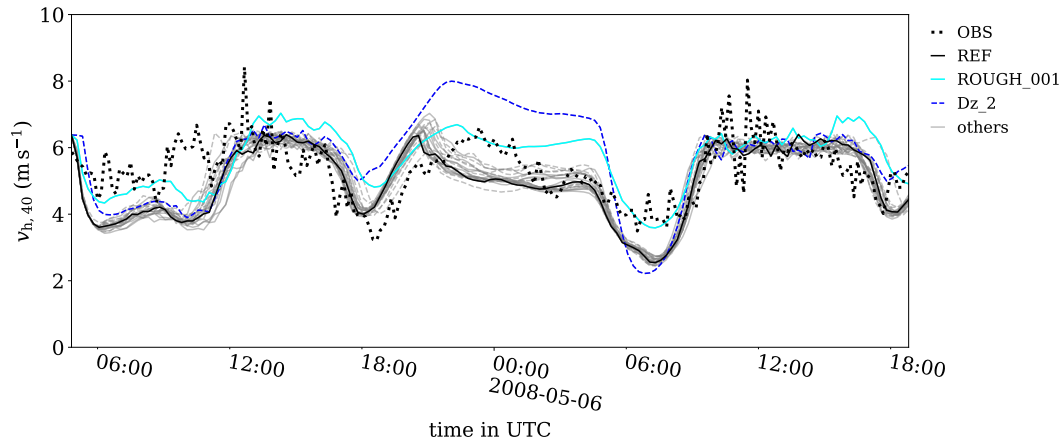
605 The CESAR tower samples temperature, humidity and horizontal wind speed in-at 10, 40, and 80m height-heights (wind speed is not available at 1.5m). We ~~chose to~~ compare timeseries of the 40 m level, because ~~on-the-one-hand~~, 10 m lies between the first and second grid level and ~~on-the-other-hand~~, 80 m ~~might already be~~ is likely above the nocturnal boundary layer (see discussion of Fig. 9). Figure 11 shows the diurnal cycle of temperature at  $z = 40$  m, which, at first sight, agrees fairly well between observations and case REF. The temperature amplitude is around 10 K and peak temperatures reach 295 K and 296 K between

610 16:00 and 17:00 UTC ~~peak-temperatures-of-295-K-and-296-K-are-reached~~ on the first and second day, respectively. Significant differences ~~are, however, found during night~~ exist during the night, as well as during the second day: During ~~night~~, ~~case~~



**Figure 12.** Timeseries of mean total water mixing ratio at 40 m as measured at Cabauw (OBS) and for all cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.

REF becomes more stable compared to the night, the temperature at 40 m height of case REF is lower than the observations. Hence, the simulated boundary-layer is more stable than the observations (cf. Fig. 9), hence, the temperature in 40 m height is lower compared to the observations, as discussed earlier. The following morning, the boundary layer is heated up quickly between roughly warms quickly between approximately 06:00 UTC and 08:00 UTC in the simulation, whereas in reality the temperature increases slower. Again, the reason is that case REF the model develops a much shallower stable boundary layer (with a residual layer above, cf. Fig. 2) than the observations. Once the stable layer is eroded surface heating in the morning has eroded the stable layer up to 40 m the temperature can increase, the temperature increases rapidly at that level as it is heated from the surface as well as from above, where warmer air from the residual layer is mixed into the shallow layer. In reality, where the stable layer is much deeper, the air in at 40 m height is first only heated from the surface and not by entrainment. The sensitivity study shows that during the night case HUMID\_sat is in agreement with the observations at that height (cf. Fig. 9). This agreement is because the humidity reaches saturation and condenses on the vegetation as dew, and heat from condensation is released to the air by different partitioning of the surface fluxes. However, dew formation was not observed at the CESAR site, as shown by timeseries of  $LE$  (Fig. 5), thus. Thus humidity is not suitable to explain the misrepresentation of near-surface temperature. In Fig. 11, also case Dz\_2 also seems to agree with observations. However, as shown in Fig. 9, this is only true for the depicted height of 40 m, e.g. For example, below 20 m, the nocturnal air temperature is lower compared to case REF. Even though case Dz\_2 has a much higher resolution than case REF, we cannot be certain that turbulence in the nocturnal boundary layer is sufficiently resolved and hence the resolution is still a possible explanation for the misrepresentation of the stable boundary layer in the LES. Another parameter to point out in Fig. 11 is case RRTMG, which shows significant differences to case REF during the night, i.e. lower temperature and higher boundary layer. The following day, the mixed-layer temperature of case RRTMG is persistently smaller compared to case REF due to the additional cooling of the air volume at night, but similar boundary-layer heating. This misrepresentation propagates into the next day.



**Figure 13.** Timeseries of mean horizontal wind speed at 40m as measured at Cabauw (OBS) and for all cases listed in Table 5. Only case REF and some relevant cases are highlighted, all others are shown in gray.

### 5.4.3 Specific humidity

Figure 12 shows that the observed humidity ~~in~~ at 40m has only small diurnal variations. The reference simulation mostly agrees with the observations. ~~The Case ADV\_tq reproduces the~~ The Case ADV\_tq reproduces the small decrease in the measurement data around 09:00 UTC of the first day ~~can be reproduced by case ADV\_tq, which~~. This case includes advection tendencies of  $\theta$  and  $q$  according to the mean change well above the boundary layer. However, during the morning transition of the second day (around 08:00 UTC) the humidity in the simulations first rises and then drops, while near-surface air ~~is mixed~~ mixes with dry air from above. By contrast, this is not observed in the measurement data, most likely because no strong stable layer had developed during nighttime. The majority of the simulated cases do not show a large spread (~~about~~ approximately  $0.5 \text{ g kg}^{-1}$ ), except for cases HUMID\_dry and HUMID\_sat. Case HUMID\_dry, which is initialized with zero humidity, but becomes humid through the latent heat flux from the surface, follows a similar curve as case REF but with persistently lower values. ~~On the other hand,~~ case Case HUMID\_sat, initialized with saturation moisture, shows a diurnal cycle similar to that of potential temperature. Apart from the extreme cases of initialized humidity, the choice of the radiation model and  $LAI$  reveal the most significant deviation from case REF. If the diurnal cycles of temperature and humidity are placed in context with those of sensible and latent heat fluxes, ~~it is found that while then~~ the the fluxes are consistently overestimated during ~~day, the day, but~~ corresponding higher values of  $\theta$  or  $q$  are not observed. ~~This can be explained with differences between observation and simulation in the boundary-layer depth. Instead, the~~ The observed differences in the atmosphere ~~occur during that occur during the~~ night or in the morning ~~and~~ are attributed to a misrepresentation of the nocturnal boundary layer.



#### 650 5.4.4 Horizontal wind speed

The timeseries of observed horizontal wind speed, depicted in Fig. 13, show a significant amount of fluctuation during day-and  
some-during-night. Since the day and less significant fluctuations during the night. Because the LES data shown is horizontally  
averaged over the entire horizontal domain, the fluctuations are less prominent in the simulations. Nevertheless, the mean day-  
time and nighttime magnitude agrees reasonably well with the observations, including the decreases and consecutive increases  
655 in the evening (ca. 18:00 UTC) and morning (ca. 06:00 UTC). Relevant deviations from case REF are found if the roughness  
length is reduced, i.e. case ROUGH\_001. As expected, this leads-to-results in higher wind speed close to the surface. Reducing  
the grid size (case DZ\_2) has the same effect mostly-shown-at-night(most discernable at night).

## 6 Summary

In this paper we gave a description of the land-surface model embedded into the PALM model system which is applied to model  
660 the surface energy balance at vegetated or paved land surfaces as well as at water surfaces. We evaluated the land-surface model  
implementation against in-situ observations of the energy-balance components as well as near-surface wind-, temperature-, and  
humidity-profiles taken at Cabauw over a quasi homogeneous flat grass site (Monna and Bosveld, 2013) for two consecutive  
diurnal cycles. A sensitivity study showed the relative importance of the choice of land-surface input parameters and thereby  
gives valuable reference for the user.

665 The diurnal cycles of surface latent and sensible heat flux are well-represented. Even though the model reasonable well  
represented considering all uncertainties included. During daytime, the model tends to underestimate the surface net radiation,  
while during nighttime the surface net radiation shows smaller negative values than the observations, indicating underestimated  
longwave radiative cooling. The model seems to overestimate the fluxes, sensible and latent heat fluxes. However, we assume  
that the discrepancies are due to the differences-can-be-explained-by-the non-closure of observed energy balance, i.e. missing  
670 energy in  $H$  and  $LE$  observed at the CESAR site. During daytime, the simulated Bowen ratio agrees reasonably well with  
the observed one, whereas climate models Bowen ratio. This is in contrast to climate models, which overestimate the summer  
Bowen ratios observed in Cabauw (e.g. Schulz et al., 1998; Sheppard and Wild, 2002). The diurnal cycle of the modeled ground  
heat flux agrees with the observations, though the modeled flux overestimates the observed flux during the morning and evening  
hours. During nighttime, the modeled ground heat flux shows slightly larger negative values compared-to-the-observationthan  
675 the observations. Due to its-the relatively small contribution of this flux compared to the surface net radiation or the surface  
latent and sensible heat fluxes, these mismatches-inaccuracies do not affect boundary-layer development too-much, if only one  
or two days are simulated significantly on a two day timespan. However, for longer simulation periods heat storage in the soil  
may become an important factor, e.g., when heat waves built-up over days (Miralles et al., 2014). During daytime the model  
tends to underestimate the surface net radiation, while during nighttime the surface net radiation shows slightly less negative  
680 values compared to the observations, indicating underestimated longwave radiative cooling.

The near-surface temperature matches well with the observed one-near-surface temperature at the first simulated day. During  
nighttime, however, the model underestimates the near-surface temperature is-underestimated-in-the-model, the-and-the noctur-



nal boundary layer is too stable and too shallow compared to the observations. The diurnal cycle of the near-surface wind is well represented, though the low-level jet in the model occurs much closer to the surface.

685 In addition to the evaluation against observations, we carried out a ~~comprehensive sensitivity study. Land-surface sensitivity study, in which land-surface~~ parameters and the initial state of the atmosphere ~~were~~ are varied within a typical range for the respective quantity. This is motivated mainly by two reasons: First, to test the embedded land-surface model for a wider range of parameters, and second, to estimate the scatter in model results depending on the choice of specific input parameters that often lack appropriate observations. The net radiation is significantly ~~influence~~ influenced by the albedo, the radiation  
690 model ~~as well as~~ and many of the land surface parameters. The ground heat flux, though not as important as the other energy-balance components, is ~~mostly~~ primarily influenced by the soil thermal conductivity. The distribution of the available energy into surface sensible and latent heat fluxes depends mostly on the leaf area index, as well as the initial atmospheric humidity. ~~Within the~~ The sensible and latent heat fluxes deviate from the reference simulation by up to 50 % within the investigated range of *LAI* values for short grass, ~~differences of up to 50 % are possible~~. Moderately less relative deviation is found for the  
695 range of completely dry to saturated initial atmospheric humidity. ~~Potential temperature is influenced by a change~~ A change of  $\pm 2 \text{ m}^2 \text{ m}^{-2}$  in *LAI* resulted in the same potential temperature as a change of  $\pm 2 \text{ m}^2 \text{ m}^{-2}$  as much as it is influenced by an initial temperature deviation in the whole atmosphere of  $\pm 1 \text{ K}$  to the initial, entire atmosphere. While most of the sensitivity studies show the most significant difference during day, the choice of initial humidity, grid size, roughness length and radiation model ~~plays an important role~~ are significant at night. Overall, we could not identify a single parameter as being the most  
700 sensitive in all quantities at the same time. In fact, different parameters become relevant if different quantities are analyzed.

In order to evaluate and possibly improve land-surface schemes ~~also~~ for different types of surfaces, e.g. ~~pavements or different~~ , pavements or various vegetation coverage, reliable measurements of the energy-balance terms ~~at on~~ at different surfaces types are required. However, eddy-covariance measurements often suffer from the well-known problem of energy-balance non-closure, where the sum of surface sensible, latent, and ground heat flux underestimates the surface net radiation by ~~about~~ approximately  
705 10 to 30 % (Foken et al., 2011). Besides measurement uncertainties and footprint biases, one leading hypothesis is that low-frequency contributions from organized turbulent structures or surface heterogeneity-induced circulations are inherently not captured by the eddy-covariance method (e.g. Finnigan et al., 2003; Foken, 2008; Eder et al., 2015). ~~As~~ Because the modeled energy balance is closed by definition, it is ~~thus~~ difficult to draw ~~final conclusions that may point directions for further conclusions that facilitate~~ improvements of the land-surface parameterizations ~~as both, the model as well as the observation may contain a bias.~~

As the description of the land-surface model embedded in PALM only reflects its current state, a short outlook into future development is given below. ~~Until now~~ Currently, the LSM implementation does not incorporate a tile approach ~~(as the embedded building-surface model (Resler et al., 2017) does), such that for water surfaces;~~ a land-surface grid cell in PALM is either classified as water or as pavement-/vegetation-covered. Particularly for coarser grids, ~~however,~~ patchy landscapes such  
715 as e.g. Cabauw (interspersed with small ditches like in Cabauw, ) might be filtered, ~~meaning that.~~ This filtering would make the relative contributions of surface types and thus the area-averaged energy-balance terms ~~become~~ a function of the horizontal grid size. ~~In order to avoid this~~ To avoid this dependence, one of our next steps will be to implement ~~a tile approach~~

into the land-surface model the tile approach also for water surfaces. Furthermore, the current LSM implementation does not include heat storage within water bodies the change of the water temperature due to the heat flux into the body, or penetration of shortwave radiation and absorption within different water layers. However, especially for multi-day simulations, e.g., for heat-wave scenarios, this heat storage in water bodies might become important to accurately represent the cool-air production in urban environments. Hence, we plan to improve the representation of water surfaces by implementing a lake parameterization (e.g. Mironov et al., 2010). Further lines of future development will be the implementation of a snow parameterization as well as a parameterization to consider of phase transitions to also consider simulate frozen soil. In this study we only considered a homogeneously flat surface. Nonetheless, heterogeneous surfaces as well as are supported by the LSM as well, with individual classification of the surface type at each horizontal grid cell. Moreover, the LSM can also be applied in simulations with non-flat surfaces. With PALM's Cartesian topography approach (see Maronga et al., 2015, 2020) where a grid cell belongs either to the atmosphere or an obstacle, this currently results in a step-like orography is implemented in the LSM. In addition, we plan to implement an immersed boundary method (Mason and Sykes, 1978) where representation of the surface. We plan to revise PALM's topography implementation such that elevation changes can also be represented by slanted surfaces sloped surfaces according to a cut-cell approach (e.g. Shaw and Weller, 2016). This will also have consequences for the modeling of the surface-energy balance as the slope of the surface needs to be considered.

*Code availability.* The PALM model system is freely available from <http://palm-model.org> and distributed under the GNU General Public License v3 (<http://www.gnu.org/copyleft/gpl.html>).

*Author contributions.* Implementation of the LSM: BM; Conceptualization of the study: BM; investigation, formal analysis, and visualization: KFG; methodology, software, writing original draft, review and editing of writing: all authors.

*Competing interests.* The authors declare that they have no conflict of interest.

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