



Global Simulation of Semivolatile Organic Compounds – Development and Evaluation of the MESSy Submodel SVOC (v1.0)

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Abstract

The new submodel SVOC for the Modular Earth Submodel System (MESSy) was developed and applied within the ECHAM5/MESSy Atmospheric Chemistry (EMAC) model to simulate the atmospheric cycling and air-surface exchange processes of semivolatile organic pollutants. Our focus is on four polycyclic aromatic hydrocarbons (PAHs) of largely varying properties. Some new features in input and physics parameterizations of tracers were tested: emission seasonality, the size discretization of particulate-phase tracers, the application of poly-parameter linear free energy relationships in gas-particle partitioning, and re-volatilization from land and sea surfaces. The results indicate that the predicted global distribution of the 3-ring PAH phenanthrene is sensitive to the seasonality of its emissions, followed by the effects from considering re-volatilization from surfaces. The predicted distributions of the 4-ring PAHs fluoranthene and pyrene, and the 5-ring PAH benzo(a)pyrene are found sensitive to the combinations of factors with their synergistic effects being stronger than the direct effects of the individual factors. The model was validated against observations of PAH concentrations and aerosol particulate mass fraction. The annual mean concentrations are simulated to the right order of magnitude for most cases and the model well captures the species and regional variations. However, large underestimation is found in the ocean. It is found that the particulate mass fraction of the benzo(a)pyrene is well simulated whereas those of other species are lower than observed.

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1 Introduction

The atmospheric cycling of semivolatile organic compounds (SOCs) is particularly complex, because of partitioning across phases and air-surface exchange processes, including multihopping (or 'grasshopper effect'; Semeena and Lammel, 2005) and accumulation in ground compartments such as seawater, soil, vegetation, and ice/snow. Many SOC do resist degradation in environmental compartments, hence, are persistent. In regulation of chemical substances and in international chemicals legislation (e.g. UNEP, 2017), model-based quantifications of the overall environmental residence time (persistence) and the long-range transport potential are requested or encouraged to be applied.

Global and regional distribution and transport of SOC has been studied using multimedia fate (box) models and chemistry transport models (CTMs) (Scheringer and Wania, 2003). The multimedia models describe the whole or part of the globe as a few zones of homogeneous environmental characteristics (Wania and Mackay, 1999; Mackay, 2010). These models are used as tools to assess the influences of environmental parameters and change on pollutant levels in multiple compartments (Dalla Valle et al., 2007; MacLeod et al., 2005; Lamon et al., 2009). On the other hand, CTMs generally imply the application of three-dimensional Eulerian models coupled with surface and chemistry modules (e.g. Ma et al., 2003; Hansen et al., 2004; Malanichev et al., 2004; Gusev et al., 2005; Semeena et al., 2006; Gong et al., 2007;



Friedman and Selin, 2012; Galarneau et al., 2014; Shrivastava et al., 2017). The addition of a
40 surface module aims to describe air–surface exchange processes and biogeochemical cycles of
contaminants whereas a chemistry module describes the changes in air concentrations due
to phase partitioning and chemical transformations. Compared to the multimedia models,
CTMs have better spatial and temporal resolution but require more computational effort.
They are suitable for use to investigate the variability and episodic character of environmental
45 fate and transport. To date, pollutants addressed in model studies were persistent organic
pollutants, such as dichlorodiphenyltrichloroethane (DDT), polychlorinated biphenyls (PCBs),
hexachlorocyclohexanes (HCHs), polycyclic aromatic hydrocarbons (PAHs), and more recent
so-called emerging pollutants (e.g. MacLeod et al., 2011).

The sensitivity of distributions to specific processes of SOC cycling and related input
50 parameters has been the focus of CTM-based studies (Semeena et al., 2006; Sehili and Lammel,
2007; Friedman and Selin, 2012; Galarneau et al., 2014; Thackray et al., 2015). Sehili and
Lammel (2007), for instance, suggest that the gas–particle partitioning and particulate-phase
oxidation scenarios have significant influences on the long-range atmospheric transport of
PAHs. This finding is supported by Friedman and Selin (2012), who, furthermore, concluded
55 that the effects are higher than those of irreversible partitioning and of increased aerosol
concentrations.

This study presents the new multicompartment module (submodel) SVOC for the Modular
Earth Submodel System (MESSy; Jöckel et al., 2006, 2010). MESSy provides a modular
framework for simulations accounting for various degree of complexity and to facilitate a
60 continuous future submodel improvements. The submodel has been applied using the general
circulation model ECHAM5 (Roeckner et al., 2003, 2006) as a base model. In connection
with the ECHAM5/MESSy (EMAC) model, SVOC encompasses a 3D atmosphere and
2D surface compartments (soil, vegetation, snow, and ocean mixed layer), and considers
multicompartment fate and exchange processes, such as emission, phase partitioning, wet and
65 dry deposition of gases and particles, degradation, and air–surface gas exchange, including
re-volatilization. The focus of this submodel development is the global distribution of four
PAH species of largely varying properties. PAHs enter the atmospheric environment as
by-products of all technological combustion processes (Shen et al., 2013) and of open fires
(Gullett et al., 2008). They are ubiquitous pollutants of particular environmental and health
70 concern (WHO, 2003; Laender et al., 2011; Lammel, 2015) and due to their continuous global
emissions. Here we describe submodel development, compare the results to observations, and
assess the significance of four model features to PAH distributions and fate. These features
are the temporal resolution of emissions, the size discretization of particulate-phase tracers
(bulk or modal), the choice of the gas–particle partitioning scheme, and re-volatilization from
75 surfaces.



2 Methods

2.1 Model descriptions

The global model applied in this study is the ECHAM/MESSy Atmospheric Chemistry–Climate model (EMAC), a three-dimensional Eulerian model for the simulations
80 of meteorological variables, gaseous, aerosols, clouds, and other climate-related parameters. EMAC combines the general circulation model ECHAM5 (here, version 5.3.02) (Roeckner et al., 2003, 2006) with the Modular Earth Submodel System (MESSy version 2.50; Jöckel et al., 2006, 2010). The atmospheric component ECHAM5 derives four prognostic variables, namely, vorticity, divergence, temperature, and the logarithm of surface pressure in truncated
85 series of spherical harmonics, whereas specific humidity, cloud water, and cloud ice are represented in grid point space. MESSy provides a modular framework to define atmospheric dynamics, chemistry, transport, and radiative transfer processes. For a more detailed description of the EMAC model, evaluation and relevant studies, refer to Jöckel et al. (2006, 2010) and <http://www.messy-interface.org>. List of MESSy process-based modules, hereinafter
90 submodels, applied in the study are summarized in Table 1.

The new MESSy submodel SVOC for simulating the fate and cycling of SOCs in the global environment is presented. Processes involved in the submodel include gas–particle partitioning, volatilization from the surface, dry and wet depositions, chemical and biotic degradations. These processes are connected to other MESSy submodels. For example,
95 deposition of gas-phase SOCs are calculated by the submodels SCAV and DDEP, aerosol microphysics by GMXe, gas-phase chemistry mechanisms by MECCA, and ocean–air flux exchange by AIRSEA. Figure 1 illustrates the SVOC structure within EMAC system and its interactions with other MESSy submodels. More details on some process parameterizations are given in the following section. A user manual can be found in the Supplement with the
100 list of submodel input and output variables.

[Insert Table 1](#)

[Insert Figure 1](#)

2.2 Parameterizations of cycling processes in multiple compartments

2.2.1 Representation of SOC in particulate phase

105 The parameterizations of aerosol microphysical processes for SOCs such as gas-to-particle partitioning and dry and wet deposition depend on the way the particulate phase is represented in the model. Here, there are two approaches employed in the submodel to represent the particulate-phase SOC: 1) it is assumed as a bulk species, or 2) the particle sizes are resolved into n continuous (modal) distributions. The former will be hereinafter referred to as the
110 *bulk* scheme, and the latter referred to as the *modal* scheme. In the *modal* scheme, n is equal to the 7 log-normal modes of the GMXe submodel, four with hydrophilic coating (ns :



nucleation soluble, *ks*: Aitken soluble, *as*: accumulation soluble, *cs*: coarse soluble) and three hydrophobic (*ki*: Aitken insoluble, *ai*: accumulation insoluble, *ci*: coarse insoluble) (Pringle et al., 2010). Each mode is treated as an individual tracer.

115 2.2.2 Partitioning between gas phase and particulate phase

Gas–particle partitioning is assumed to take place when SOC is at equilibrium between the gas and particulate phases. The concentration of the species that is bound to particles (C_{particle}) is calculated with

$$C_{\text{particle}} = \theta \times (C_{\text{particle}} + C_{\text{gas}}), \quad (1)$$

and the particulate mass fraction (θ) is defined as

$$\theta = \frac{C_{\text{particle}}}{C_{\text{particle}} + C_{\text{gas}}} = \frac{K_p \times C_{\text{PM}}}{1 + K_p \times C_{\text{PM}}} = \frac{K'_p}{1 + K'_p}, \quad (2)$$

120 where C_{PM} is the concentration of particulate matter or PM ($\mu\text{g m}^{-3}$), K_p is the temperature-dependent particle–air partition coefficient ($\text{m}^3 \mu\text{g}^{-1}$), and K'_p is the dimensionless K_p .

In a model configuration using size-resolved particles (viz. the *modal* scheme), each SOC tracer is introduced in the model as eight different species, seven aerosol particles in *ns*, *ks*, *as*, *cs*, *ki*, *ai*, *ci* modes and one in the gas phase. The particulate fraction of the species in mode *i* (θ_i) is calculated using Equation 2 with C_{PM_i} and K_{p_i} being the PM mass concentration and aerosol–air partition coefficient in the corresponding mode, respectively. The gaseous concentration C_{gas} is calculated using the sum of K_p values across modes, as well as total C_{particle} and C_{PM} :

$$C_{\text{gas}} = \frac{\sum_{i=1}^7 C_{\text{particle}_i} / \sum_{i=1}^7 C_{\text{PM}_i}}{\sum_{i=1}^7 K_{p_i}}. \quad (3)$$

130 It is noted that this approach may not hold the constraint of mass consistency, and is thus subject to further corrections. For the current study, the effects from this problem are expected to be minimal, given the fact that PAHs in the particulate phase are mainly distributed in the accumulation mode (Lammel et al., 2010, and references therein).

For K_p calculation four options of gas–particle partitioning schemes are available in SVOC, they are: (1) a parameterization that is based on adsorption onto aerosol surface (Junge, 1977; Pankow, 1987), (2) absorption into organic matter (Finizio et al., 1997), (3) a combination of two ways of organic matter absorption and black carbon adsorption (Lohmann and Lammel, 2004), and (4) multiple phase of the two-ways sorption system (Goss and Schwarzenbach, 2001; Endo and Goss, 2014; Shahpoury et al., 2016). Two schemes used in this study are described below.



140 Lohmann–Lammel scheme

The Lohmann–Lammel scheme takes into account an adsorption onto black carbon (BC) surface in addition to absorption into OM (Lohmann and Lammel, 2004). This dual sorption theory empirically calculates K_p according to the following relation

$$K_p = 10^{-12} \left(f_{OM} \frac{\gamma_{oct} MW_{oct}}{\gamma_{OM} MW_{OM} \rho_{oct}} K_{oa} + f_{BC} \frac{a_{atm-BC}}{a_{soot} \rho_{BC}} K_{sa} \right), \quad (4)$$

where ρ_{BC} is the density of BC (assumed as 1 kg L^{-1}), ρ_{oct} is the density of octanol (0.82 kg L^{-1} at 20°), K_{sa} is the partition coefficient between diesel soot and air, a_{atm-BC} is the available surface of atmospheric BC ($\text{m}^2 \text{ g}^{-1}$), and a_{soot} is the specific surface area of diesel soot ($\text{m}^2 \text{ g}^{-1}$). The adsorptive properties of diesel soot are selected to represent the atmospheric BC because this material is considered the most significant type of BC in polluted air.

The K_{sa} value is calculated as a function of sub-cooled liquid vapor pressure p_L^0 using an estimate suggested by van Noort (2003),

$$\log K_{sa} = -0.85 \log p_L^0 + 8.94 - \log \left(\frac{998}{a_{soot}} \right), \quad (5)$$

where a_{soot} in the model is set as $18.21 \text{ m}^2 \text{ g}^{-1}$.

Poly-Parameter Linear Free Energy Relationships (ppLFER) scheme

The concept of poly-parameter linear free energy relationships (ppLFER) for the prediction of equilibrium partition coefficients is introduced by Goss and Schwarzenbach (2001), and its application in environmental chemistry has been reviewed by Endo and Goss (2014). This approach can describe a composite of different types of interactions between gas-phase species and aerosols. In contrast, single-parameter LFERs only correlates the partition coefficient to the sub-cooled liquid vapor pressure or the octanol–air partition coefficient of the species, hence only valid within the group of compounds for which they were developed.

In the study, ppLFER scheme is incorporated into SVOC in which it defines K_p as the sum of individual partition coefficients representing surface adsorption and bulk-phase absorption processes to inorganic and organic aerosols. The formulation of K_p is adopted from Shahpoury et al. (2016) and is described as follows



$$\begin{aligned}
 K_p &= \frac{K'_p}{C_{PM}} \\
 K'_p &= K_{EC} \times a_{EC} \times C_{EC} \times 10^{-6} + \\
 &K_{(NH_4)_2SO_4} \times a_{(NH_4)_2SO_4} \times C_{(NH_4)_2SO_4} \times 10^{-6} + \\
 &K_{NaCl} \times a_{NaCl} \times C_{NaCl} \times 10^{-6} + \\
 &K_{DMSO} \times \frac{C_{WSOM}}{\rho_{DMSO}} \times 10^{-6} + \\
 &K_{PU} \times 0.2 \times C_{WDOM} \times 10^{-12} + \\
 &K_{hexadecane} \times 0.8 \times \frac{C_{WDOM}}{\rho_{hexadecane}} \times 10^{-6}
 \end{aligned} \tag{6}$$

where K_{EC} , $K_{(NH_4)_2SO_4}$, and K_{NaCl} are the substance partition (adsorption) coefficients
 165 ($m^3_{air} m^{-2}_{surface}$) for elemental-carbon/diesel soot, ammonium sulfate, and sodium chloride
 aerosol surface-air systems, respectively. K_{DMSO} is the substance partition (absorption)
 coefficient for dimethyl sulfoxide-air system ($L_{air} L^{-1}_{DMSO}$). K_{PU} is the substance partition
 (absorption) coefficient for polyurethane-air system ($m^3_{air} kg^{-1}_{PU}$). $K_{hexadecane}$ is the substance
 partition (absorption) coefficient for hexadecane-air system ($L_{air} L^{-1}_{hexadecane}$). a_{EC} , $a_{(NH_4)_2SO_4}$,
 170 and a_{NaCl} are the adsorbent specific surface areas: 18.21, 0.1, and 0.1 $m^2_{surface} g^{-1}_{adsorbent}$,
 respectively. ρ_{DMSO} and $\rho_{hexadecane}$ are the dimethyl sulfoxide and hexadecane densities:
 1.1×10^6 and 0.77×10^6 $g m^{-3}$, respectively. C_{EC} , $C_{(NH_4)_2SO_4}$, C_{NaCl} , C_{WSOM} , C_{WDOM} are the
 concentration ($\mu g_{substance} m^{-3}_{air}$) of elemental carbon (here, black carbon), ammonium sulfate,
 sodium chloride, water-soluble organic matter, and water-insoluble organic matter, respectively.

175 The ppLFER scheme calculates the sorptive partition coefficient for every aerosol system,
 as summarized in Table S1. Each of the coefficients requires information on system parameters
 (e , s , a , b , v , l), and the constant c , as shown in Table S2. The Abraham solute descriptors
 (E , S , A , B , V , and L) are substance specific, and for the species selected in this study, refer
 to Table S6. All the predicted partition constants are adjusted to environmental temperature
 180 using the van't Hoff equation

$$\ln K_{(T)} = \ln K_{(T_0)} - \frac{\Delta H}{R} \left(\frac{1}{T} - \frac{1}{T_0} \right), \tag{7}$$

where ΔH is the enthalpy of solvent-air phase transfer in $J mol^{-1}$. This variable is system
 specific and calculated by applying the ppLFER equations given in Table S3 and input
 parameters given in Table S4. The sequence of K_p calculation from ppLFER analysis in
 SVOC is illustrated in Figure S1.

185 2.2.5 Volatilization

Soil

For soil volatilization, two parameterization schemes are implemented in the SVOC submodel,
 that is, the Jury scheme (Jury et al., 1983, 1990), and the Smit scheme (Smit et al., 1997).
 The latter was applied in the study which is based on the volatilization of pesticides from the



190 surface of fallow soils (Smit et al., 1997). The volatilization occurs upon partitioning over
 three soil phases (solid, gas, and liquid). The concentration of the chemical in the soil system
 (kg m⁻³) is formulated as

$$C_{\text{soil}} = Q \times C_{\text{vapor}} , \quad (8)$$

and the capacity factor Q is given by

$$Q = \psi + \varphi K_{\text{wa}} + \rho_{\text{soil}} K_{\text{wa}} K_{\text{sl}} , \quad (9)$$

where ψ and φ are the volume fractions of air and moisture, respectively, ρ_{soil} is the soil density,
 195 K_{wa} is the water–air partition coefficient where $K_{\text{wa}} = 1/K_{\text{aw}}$, and K_{sl} is the solid–liquid
 partition coefficient. K_{aw} is calculated based on the Henry’s Law constant:

$$K_{\text{aw}} = 1/(H R T) , \quad \text{and} \quad (10)$$

$$H = k_{\text{H}}^{\ominus} \times \exp \left[\frac{-\Delta H_{\text{soln}}}{R} \left(\frac{1}{T} - \frac{1}{T_0} \right) \right] , \quad (11)$$

where H is the temperature-adjusted Henry coefficient (M atm⁻¹), R is the dry air gas
 constant (= 8.314 J mol K⁻¹), T is the environment temperature (K), ΔH_{soln} is the enthalpy
 of dissolution (J mol⁻¹).

200 K_{sl} can be set equal to the sorption coefficient to soil organic matter K_{om} times the
 fraction of organic carbon in soil f_{OMs} . Since K_{om} data is not available, the coefficient for
 sorption to soil organic carbon K_{oc} was used to estimate K_{sl} :

$$K_{\text{sl}} = 0.56 K_{\text{oc}} f_{\text{OMs}} \quad (12)$$

Mackay and Boethling (2000) has suggested a reasonably good regression relationship
 between K_{oc} and octanol–water partition coefficient K_{ow} for PAHs:

$$\log(K_{\text{oc}}/1000) = 0.823 \log(K_{\text{ow}}/1000) - 0.727 . \quad (13)$$

205 where the factor of 1000 is needed because K_{oc} and K_{ow} are expressed in m³ kg⁻¹ whereas,
 in the original regression, they used mL g⁻¹.

Once Q is computed, the dimensionless fraction of the chemical in the gas phase F_{gas} is
 then calculated as

$$F_{\text{gas}} = \frac{\psi}{Q} . \quad (14)$$



In the Smit scheme, an empirical relation was established between F_{gas} and cumulative
 210 volatilization (CV in % of substance deposit). CV was determined based on field and
 greenhouse experiments with numerous pesticides at 21 days after application. For normal to
 moist field conditions, CV is expressed as

$$CV = 71.9 + 11.6 \log(100 F_{\text{gas}}) ; 6.33 \times 10^{-9} < F_{\text{gas}} \leq 1 , \quad (15)$$

and for dry field conditions,

$$CV = 42.3 + 9.0 \log(100 F_{\text{gas}}) ; 0.2 \times 10^{-6} < F_{\text{gas}} \leq 1 . \quad (16)$$

215 Vegetation

Smit et al. (1998) derived an equation for the cumulative volatilization CV from plants against
 vapor pressure P_v (mPa) at seven days after application based on field and climate chamber
 experiments of pesticide volatilization (Equation 17).

$$CV = 10^{1.528+0.466 \log P_v} ; P_v \leq 10.3 . \quad (17)$$

For compounds with P_v above 10.3 mPa, CV is set at 100 % of deposit. Temperature
 220 adjustments were made for P_v using the Clausius–Clapeyron equation:

$$\frac{d(\ln P_v)}{dT} = - \frac{\Delta H_{\text{vap}}}{R T^2} . \quad (18)$$

Snow and glaciers

The parameterization of substance loss by volatilization from snow pack follows Wania (1997)
 whereby the process is calculated using a consecutive cycle of an equilibrium partitioning
 225 among four phases followed by a contaminant loss. The four phases considered are liquid
 water, organic matter contained in the snowpack, snow pores (air), and an ice–air interface.
 Fugacity capacity factors for these phases are expressed with the following relations:

$$\text{air (mol m}^{-3} \text{ Pa}^{-1}) \quad Z_a = 1/RT \quad (19a)$$

$$\text{water (mol m}^{-3} \text{ Pa}^{-1}) \quad Z_l = K_{\text{wa}}/RT = K_{\text{wa}} Z_a \quad (19b)$$

$$\text{organic carbon (mol m}^{-3} \text{ Pa}^{-1}) \quad Z_o = Z_l 0.41 K'_{\text{ow}} \quad (19c)$$

$$\text{ice–air interface (mol m}^{-2} \text{ Pa}^{-1}) \quad z_i = K_{\text{ia}}/RT = K_{\text{ia}} Z_a \quad (19d)$$

where R is the dry air gas constant ($8.312 \text{ J mol}^{-1} \text{ K}^{-1}$), T is the air temperature (K),
 K_{wa} is the water–air partition coefficient (unitless), K'_{ow} is the dimensionless octanol–water
 230 partition coefficient, and K_{ia} is the ice surface–air partition coefficient (m). K_{ia} at 20°C is
 estimated using K_{wa} and water solubility C_w^s (mol m^{-3}),



$$\log K_{ia}(20^\circ\text{C}) = -0.769 \log C_w^s - 5.97 + \log K_{wa}, \quad (20)$$

and further extrapolated to other temperatures using enthalpy of condensation of solid (ΔH_{subl} in J mol^{-1}),

$$\log K_{ia}(T) = \log K_{ia}(20^\circ\text{C}) + \frac{0.878 \Delta H_{\text{subl}}}{2.303 R} \left(\frac{1}{T} - \frac{1}{293} \right). \quad (21)$$

An equilibrium fugacity f_s is thereby determined by

$$f_s = \frac{M_{\text{sp}}}{Z_a v_a + Z_l v_l + z_i A_{\text{snow}} \rho_{\text{mw}} + Z_o v_o}, \quad (22)$$

235 where M_{sp} is the amount of chemical contained in snowpack (the model here applies snow
 burden of the chemical in kg m^{-2}), v_a , v_l , and v_o is the volume fraction of air, liquid water,
 and organic matter in snowpack ($\text{m}^3 \text{m}^{-3}$). For this study, v_a and v_l values are set to 0.3 and
 0.1, respectively, whereas v_o is zero assuming no polluted snow. A_{snow} is the specific snow
 surface area ($\text{m}^2 \text{g}^{-1}$). In Daly and Wania (2004), a value of $0.1 \text{ m}^2 \text{g}^{-1}$ for A_{snow} was used
 240 for snow accumulation period and a linear decrease from 0.1 to $0.01 \text{ m}^2 \text{g}^{-1}$ was used during
 the snowmelt period. In SVOC submodel, a value of $0.025 \text{ m}^2 \text{g}^{-1}$ is adopted for A_{snow} to
 represent fairly aged snowpack. ρ_{mw} is the density of snowmelt water and here is taken as
 $7 \times 10^5 \text{ g m}^{-3}$.

Volatilization rate ($\text{kg m}^{-2} \text{s}^{-1}$) is calculated by applying

$$\frac{dM_{\text{sp}}}{dt} = \frac{1}{\frac{1}{U_7 Z_a} + \frac{1}{U_5 Z_l} + \frac{1}{U_6 Z_a}} \times \frac{f_s}{h_s}, \quad (23)$$

245 where h_s is the snow depth (m), U_5 is the snow–water phase diffusion mass transfer coefficient
 (m h^{-1}), U_6 is the snow–air phase diffusion mass transfer coefficient (m h^{-1}), and U_7 is the
 snow–air boundary layer mass transfer coefficient (m h^{-1}). U_5 and U_6 are calculated from
 molecular diffusivities in air and water (Equations 24 and 25), whereas a typical value of 5 m
 h^{-1} is adopted for U_7 .

$$U_5 = B_w \frac{v_l^{10/3} / (v_a + v_l)^2}{\ln 2h_s}, \quad (24)$$

$$U_6 = B_a \frac{v_a^{10/3} / (v_a + v_l)^2}{\ln 2h_s}, \quad (25)$$

250 where B_w and B_a are the molecular diffusivities ($\text{m}^2 \text{h}^{-1}$) in water and air respectively. In
 the model, B_a is derived from the molecular weight (MW) as $B_a = \frac{1.55}{\text{MW}^{0.65}} \text{ cm}^2 \text{s}^{-1}$, whereas
 B_w is set as 1×10^4 less than B_a , following Schwarzenbach et al. (2005).

Ocean



In the study, the ocean is treated as a two-dimensional compartment (no vertical layer). The
 255 flux of chemicals from the ocean to the atmosphere is parameterized based on the two-film
 model of Liss and Slater (1974) and is calculated within the AIRSEA submodel (Pozzer et al.,
 2006). Note that no ocean and sea-ice dynamics were included in the simulations.

2.2.10 Dry deposition

Dry deposition is simulated using deposition velocities. For gas-phase SOCs, the velocities are
 260 calculated by the DDEP submodel (Kerkweg et al., 2006a), whereas particulate-bound SOCs
 are assumed to deposit at similar rates to other aerosols whose velocities are also computed
 by DDEP. If the *modal* scheme is selected (see Section 2.2.1), the particle deposition velocity
 v_d^{SOC} at mode i is equal to the aerosol deposition velocity v_d^{aer} at the respective mode. On the
 other hand, for the *bulk* scheme, v_d^{SOC} is computed as a weighted average of v_d^{aer} from the four
 265 BC modes (*ki*, *ks*, *as*, and *cs*) where the weight is the surface area of BC. This approach is
 most relevant for PAHs as they are assumed to be predominantly transported by sorption to
 BC. The above relations are formulated as follows

$$\text{modal scheme:} \quad v_{d,i}^{\text{SOC}} = v_{d,i}^{\text{aer}} \quad (26a)$$

$$\text{bulk scheme:} \quad v_{d,\text{bulk}}^{\text{SOC}} = \frac{\sum_{i=1}^4 S_{\text{BC}_i} \times v_{d,i}^{\text{aer}}}{\sum_{i=1}^4 S_{\text{BC}_i}} \quad (26b)$$

and the BC surface area per unit volume S_{BC} ($\text{cm}^2 \text{ cm}^{-3}$) is given by

$$S_{\text{BC}_i} = 4\pi [r_i \exp(\ln^2 \sigma_{g_i})]^2 N_i \times \frac{C_{\text{BC}_i}}{C_{\text{aer}_i}}, \quad (27)$$

where N_i is the number concentration for mode i (cm^{-3}), r_i is the number radius (cm), σ_{g_i}
 270 is the geometric standard deviation, C_{BC_i} is the BC concentration ($\mu\text{g m}^{-3}$) in mode i , and
 C_{aer_i} is the sum of aerosol concentrations in the same mode ($\mu\text{g m}^{-3}$).

2.2.11 Wet deposition

Wet deposition is applied to both gas and particulate SOCs. The gaseous fraction is scavenged
 into cloud and rain droplets according to diffusion limitation, Henry's law equilibrium, and
 275 accommodation coefficient, and this process is parameterized and solved empirically in the
 SCAV submodel (Tost et al., 2006a). Particulate-phase SOCs are scavenged in convective
 updrafts, rainout and washout, and cloud evaporation, with the rate being proportional to
 BC wet scavenging; hence the change in SOC concentration is described as



$$\text{modal scheme:} \quad \frac{\Delta C_{\text{SOC}_i}}{\Delta t} = \frac{\mu_{\text{SOC}_i}}{\mu_{\text{BC}_i}} \times \frac{\Delta C_{\text{BC}_i}}{\Delta t} \quad (28a)$$

$$\text{bulk scheme:} \quad \frac{\Delta C_{\text{SOC}}}{\Delta t} = \frac{\mu_{\text{SOC}}}{\sum_{i=1}^4 \mu_{\text{BC}_i}} \times \sum_{i=1}^4 \frac{\Delta C_{\text{BC}_i}}{\Delta t} \quad (28b)$$

where μ is the particle volume mixing ratio ($\text{mol}_{\text{SOC/BC}} \text{mol}_{\text{air}}^{-1}$) and Δt is the model time
 280 step (s). Note that Equation 28a imposes a restrictive prerequisite, namely, BC and the
 particle-bound SOC have similar size distributions. When this condition is not met, there is
 a high possibility of an artificial mass being produced, usually to the largest aerosol mode.
 To solve this problem, a correction factor is applied and defined as a function of the ratio of
 positive fluxes to negative fluxes integrated across levels and modes.

285 2.2.12 Atmospheric degradation

The atmospheric degradation of SOCs in the gas phase as well as within aerosol particles
 are explicitly treated in SVOC. The gas-phase chemical mechanism is calculated within the
 MECCA submodel (Sander et al., 2011). SOC gaseous degradation is from photochemical
 reactions with OH, NO₃, and O₃ radicals which follow a 2nd-order transformation, with the
 290 rate constants $k^{(2)}$ obtained from laboratory studies. $k_{\text{OH}}^{(2)}$ value is typically higher, suggesting
 that oxidation with OH radical is the dominant loss pathway.

Most models do not consider oxidation rate of particulate-phase SOCs as experimental
 aerosols studied in laboratory cover only a small part of atmospheric relevant aerosols. For
 PAHs, such as benzo(a)pyrene which stays mostly in the particulate phase, the degradation
 295 is more efficient by surface reactions with O₃ (Shiraiwa et al., 2009, and references therein)
 with the rate depending on the substrate. The SVOC submodel includes degradation pro-
 cess of PAHs on aerosol particles from the O₃ reaction with one assumption, that is, the
 heterogeneous reaction does not lead to a change in the oxidant concentration. Due to a
 limited number of kinetic studies of heterogeneous reactions of PAHs, only two species are
 300 considered (phenanthrene and benzo(a)pyrene). Nevertheless, the submodel structure provides
 a relatively straightforward approach to allow more species in the future.

The reaction rate coefficient for particulate-phase PHE with O₃ at aerosol surfaces was
 derived from laboratory experiments using chemically unspecific model aerosol (silica) with
 PAH surface coverage of less than a monolayer (Perraudin et al., 2007). To this end, the
 305 second order rate coefficient, $k^{(2)}$, in $\text{cm}^4 \text{molec}^{-1} \text{s}^{-1}$ was derived from the reported PHE
 decay kinetics, $k_{\text{O}_3, \text{het}}^{(2)}$ ($\text{cm}^3 \text{molec}^{-1} \text{s}^{-1}$), as $k^{(2)} = k_{\text{O}_3, \text{het}}^{(2)} / \left(\frac{S}{V}\right) = (6.2 \pm 4.8) \times 10^{-17} \text{cm}^4$
 $\text{molec}^{-1} \text{s}^{-1}$, with $\frac{S}{V}$ (cm^{-1}) being the experimental aerosol surface concentration (0.56 ± 0.43
 cm^{-1} in Perraudin et al. (2007)). In the submodel, $k_{\text{O}_3, \text{het}}^{(2)}$ ($\text{cm}^3 \text{molec}^{-1} \text{s}^{-1}$) is calculated
 using the ambient aerosol surface concentration. As for BaP, the pseudo-first-order rate
 310 coefficient, $k_{\text{O}_3, \text{het}}^{(1)}$ in s^{-1} , was derived from surface-adsorbed BaP reaction with O₃ on solid
 organic and salt aerosols following the Langmuir–Hinshelwood mechanism (Kwamena et al.,
 2004).



2.2.13 Biotic and abiotic degradations

Biotic and abiotic processes in surface compartments contribute to the degradation of chemicals
315 and are strongly dependent on local environmental conditions, for example, nutrient contents,
water, temperature, PH, and light. In SVOC, these factors are not explicitly quantified. The
degradation is alternatively described as following a first-order decay law (Equation 29), and
the rate is assumed to double for every 10 K warming (Equation 30).

$$\frac{\partial C_{\text{SOC}_s}}{\partial t} = -k_{\text{sfc}} \times C_{\text{SOC}_s}, \quad (29)$$

$$k_{\text{sfc}}(T) = k_{\text{sfc}}(T_{\text{ref}}) \times 2^{\frac{T-T_{\text{ref}}}{10}}, \quad (30)$$

where C_{SOC_s} is the substance concentration (kg m^{-3}) in surface compartments (that is,
320 soil, vegetation, or ocean) and k_{sfc} is the first-order decay rate (s^{-1}). T_{ref} is the reference
temperature, that is, 298 K for soil and 273 K for ocean. Note that the degradation in
vegetation is calculated assuming the same k_{sfc} for the soil compartment.

2.3 Input data

2.3.1 Kinetic and physicochemical properties

325 The model simulations were performed for four PAH species: phenanthrene (PHE), pyrene
(PYR), fluoranthene (FLT), and benzo(a)pyrene (BaP). To simulate the fate and environmental
distribution of these species, the model requires some physicochemical properties as summarized
in Table S5 of the supplement. These include equilibrium partition coefficients and their
related energies of phase transfer. The characteristics from PHE to BaP are indicated by
330 decreasing volatility (as molar mass increases), increasing K_{oa} and K_{ow} , and decreasing
water solubility (as C_w^s and Henry's coefficients decrease). The properties also include the
second-rate coefficients for homogeneous oxidation with OH, O₃, and NO₃ except for BaP
where the gaseous reaction is switched off. Heterogeneous oxidation by O₃ is simulated only
for PHE and BaP. Furthermore, the model also requires compound solute parameters for
335 simulations using the ppLFR gas–particle partitioning scheme (Table S6).

2.3.2 Emissions and other model input

As model input, several emission datasets were employed in the study. Emission estimates for
PAHs were obtained from the annual mean inventory of Shen et al. (2013) for the year 2008.
They applied regression and technology split methods to construct country-level emissions for
340 six categories (coal, petroleum, natural gas, solid wastes, biomass, and an industrial process
category) or six sectors (energy production, industry, transportation, commercial/residential
sources, agriculture, and deforestation/wildfire) before further regridding the emissions to a
 $0.1^\circ \times 0.1^\circ$ grid.



Emissions of aerosol species such as organic carbon (OC), black carbon (BC), mineral
345 dust (DU), and sea salt (SS) were included. For BC and OC, the Representative Con-
centration Pathway (RCP) 6.0 emission scenario of the IPCC (Intergovernmental Panel
on Climate Change) (van Vuuren et al., 2011) was used and accessible via [ftp://ftp-ipcc.fz-
juelich.de/pub/emissions/gridded_netcdf](ftp://ftp-ipcc.fz-juelich.de/pub/emissions/gridded_netcdf). Emissions are calculated for anthropogenic, biomass
burning, ship, and aircraft. The RCP database provides a seasonality only for the biomass
350 burning and ship emissions. In this study, seasonal scale factors were applied to the an-
thropogenic emission whereby the seasonality was based upon the monthly variation of the
Hemispheric Transport of Air Pollutants (HTAP) v2.2 anthropogenic emission inventory
(Janssens-Maenhout et al., 2015). BC emissions from all sectors were assumed to be hydropho-
bic. For OC, it was assumed to be 65% hydrophilic and 35% hydrophobic upon biomass
355 burning emissions and to be 100% hydrophobic upon anthropogenic and ship emissions. Both
OC and BC were emitted at Aitken mode which spans the size range from about 5 to 50
nanometer in diameter. A factor of 1.724 was used to scale the OC emissions to primary
organic matter (OM).

DU and SS emissions were computed online by the ONEMIS submodel (Kerkweg et al.,
360 2006b). DU emission flux is calculated based on wind speed at 10 m altitude and soil
parameters (Schulz et al., 1998). The emission of SS particles by bubble bursting is described
as wind-speed dependent particle mass and number fluxes at accumulation (50–500 nm)
and coarse (>500 nm) modes. In ONEMIS, the fluxes are determined from pre-calculated
lookup tables following the Guelle et al. (2001) parameterization. SVOC submodel accounts
365 for OM fraction in the SS mass fluxes (J_{SS}), and the fraction is estimated using a 10 m wind
(v_{10})-dependent empirical relationship derived from Figure 2a in Gantt et al. (2011). Equation
31 below is used to calculate the OM mass fluxes, J_{OM} , in $\text{kg m}^{-2} \text{s}^{-1}$.

$$J_{OM} = J_{SS} \times \frac{1}{2} \left(\frac{0.78}{1 + 0.03 \exp(0.48v_{10})} + \frac{0.24}{1 + 0.05 \exp(0.38v_{10})} \right). \quad (31)$$

Emissions of other gases including volatile organic species (SO_2 , CO , NH_3 , NO , CH_4 , and
NMHC) were prescribed using the IPCC RCP6.0 dataset (van Vuuren et al., 2011). The
370 model also requires soil properties data (i.e., dry bulk density and organic matter content)
and ocean mixed layer depths (de Boyer Montégut et al., 2004).

2.4 Observational data

The observation data used for model performance evaluation were collected from several
surface monitoring networks: the European Monitoring and Evaluation Programme (EMEP)
375 (Tørseth et al., 2012), the Arctic Monitoring and Assessment Program (AMAP) (Hung et al.,
2005), the Great Lakes Integrated Atmospheric Deposition Network (IADN) (IADN, 2014),
the Department for Environment, Food & Rural Affairs (DEFRA) UK-AIR (Air Information
Resource) program (DEFRA, 2010), and the MONitoring NETwork in the African continent
(MONET-Africa) (Klánová et al., 2008). These data were screened and quality controlled
380 according to the description in Supplement SIII. Final stations with reliable monthly data are



depicted in Figure 2 wherein shows 3 stations in the Arctic, 19 in the northern mid-latitudes, and 6 in the tropics. The availability of data differs by station, species, and variable of interest; see the site-specific information in Table S10 for total concentration and Table S11 for particulate mass fraction (θ).

385 Insert Figure 2

The study also compared simulated concentrations in the marine atmosphere to two ship cruise measurement campaigns: 1) on a West to East transect across the tropical Atlantic Ocean (Lohmann et al., 2013), and 2) along the Asian marginal seas, the Indian, and the
390 Pacific Oceans (Liu et al., 2014). The monthly mean modeled values were compared to daily measurements at each sampling points.

2.5 Experiment designs

2.5.1 Model configuration

The model was run on a spectral T42 grid in the horizontal (approximately 2.8° in a lat–lon
395 grid) and 19 unevenly distributed layers in the vertical with the top level at 10 hPa. The vertical layers are discretized using a hybrid coordinate (the lowest level follows the terrain and becomes surfaces of constant pressure in the stratosphere). All simulations were run for a three-year period (i.e., 2007–2009), with a one-year spin-up (i.e., 2006), and nudged toward the European Centre for Medium-range Weather Forecasts (ECMWF) reanalysis data (Dee
400 et al., 2011). Note that the simulation period was selected based on the representative year of PAH emissions (i.e., 2008) and the availability of reliable observation data (see Supplement SIII).

2.5.2 Sensitivity to the temporal resolution of emissions and process parameterizations

405 Factor separation analysis (Stein and Alpert, 1993) was used to quantitatively evaluate the contributions to changes in a particular output variable that result from changing components of model input and physics parameterizations. The model sensitivity to four model components (hereinafter, "factors") was tested. The four factors were:

1. Temporal resolution of emissions (hereinafter, *fac1*)
410 The PAH emission inventory of Shen et al. (2013) was based on 2008 annual emission totals from all sectors (see Section 2.3.2). The emissions were divided over the year using monthly factors derived from BC anthropogenic emissions. Two sets of simulations were carried out to test the sensitivity of model output to the seasonal profile of emission.



415 The first set used constant emissions throughout the simulations whereas the second set used monthly emission interval.

2. The size-discretization of particulate-phase PAHs (hereinafter, *fac2*)

420 The two options for this factor was tested: *bulk* versus *modal* (Section 2.2.1). Note that with regards to BaP, 95% of the emissions were assumed to be in particulate phase and for the *modal*-scheme scenario, all of the emitted particles are treated as the hydrophobic Aitken (*ki*) tracers.

3. The choice of gas–particle partitioning scheme (hereinafter, *fac3*)

The present study focuses on the comparison between the Lohmann–Lammel and ppLFFER schemes for gas–particle partitioning.

4. The influence of re-volatilization (hereinafter, *fac4*)

425 Model runs with volatilization process switched off are compared to those runs which have volatilization switched on.

The factor separation technique is described in Supplement SIV including the equations used to compute the model sensitivity to four factors. A total of 16 (or 2^n) experiments summarized in 3 are necessary to supply the complete solution for the factor analysis. The 430 ABLN experiment is designed to be the *base* simulation (f_0), in which annual emission (A), the *bulk* scheme (B), the Lohmann–Lammel scheme (L), and no re-volatilization (N) were applied. SMPW is referred to as the *target* simulation (f_{1234}) in which the more sophisticated choice of the four features (factors) were tested, i.e., seasonal emissions + *modal* scheme + ppLFFER scheme + with re-volatilization. The total (gas + particle) concentration at the 435 lowest model level was selected as the selected output for its higher relevance with all the factors (compared to, e.g., atmospheric burden) and to facilitate direct comparison with observations.

3 Results and discussion

3.1 Sensitivity tests

440 The analysis of the factor separation results is given below. For each factor, the analysis includes the assessment of direct effects (\hat{f}_i) and total interaction effects ($\Sigma \hat{f}_{ij} + \Sigma \hat{f}_{ijk} + \hat{f}_{1234}$) on near-surface PAH concentrations in two seasons, i.e., December–January–February (DJF) and June–July–August (JJA). Figure 4 shows the respective effects for all factors as relative to the seasonal means of the *base* experiment (f_0). A positive value indicates a concentration 445 increase with respect to f_0 , whereas a negative indicates a decrease. The spatial distributions of f_0 seasonal mean concentrations for the four species are shown in Supplement SVI, Figure S4.

We studied the relative effects in five climate zones (Arctic, northern mid-latitudes, Tropics, southern mid-latitudes, and Antarctica) The global distributions of the relative effects are 450 presented in Figures S5–S12 whereas Figures S13–S20 present the relative interaction effects



from the individual combination of factors. In the following, we do not look to interpret concentration responses to each interaction term. Reasons for this are that (1) accounting for all such interactions is complicated given the number of factors and (2) higher-order interactions (combinations of more than two factors) is hard to physically justify.

455 We further investigate the factor effects on model performance by comparing the predicted seasonal mean near-surface concentrations from 16 experiments against observation data in the Arctic and northern mid-latitudes (Supplement SVII).

3.1.1 Effects of seasonality of emissions

Figure 4a shows that using monthly emissions increases PAH concentrations in DJF and
460 decreases the concentrations in JJA over the areas from the middle to high latitudes of NH. This result is expected and is attributed to emissions during the northern winter (summer) are higher (lower) than annual means and photochemistry being less (more) active. Over the Arctic, the relative changes (\hat{f}_1/f_0) in DJF show a median increase of 30% for PHE, PYR and FLT, and 7% for BaP, whereas \hat{f}_1/f_0 in JJA are weaker in magnitude for PHE and PYR
465 (−16%) but comparable for FLT (−28%) and BaP (−5%). Accounting for seasonality leads to generally lower bias for PHE, and this effect is more pronounced in middle than in high latitudes (Supplement SVII, Figure S21).

In general, \hat{f}_1/f_0 becomes smaller over the northern mid-latitudes by around half. The upper (lower) quartile of \hat{f}_1/f_0 in DJF (JJA) indicates about one-quarter areas of the temperate
470 and polar regions experience at least 40% of an increase (decrease), most were located in northeastern Eurasia (see the left panels of Figures S5 and S6). Note \hat{f}_1/f_0 over the tropics are small-to-negligible ($\pm 1\%$) mainly due to little variation in emissions from anthropogenic sectors. PAH concentrations may be higher in dry season due to increased amounts of biomass burning, but they are poorly represented in the current inventory. In southern mid- and
475 high-latitudes, the direct effects of emission change are substantially opposite in sign to the effects seen in the northern latitudes, being negative in DJF (median ranges from −4% to −32%) and positive in JJA (7%–25%).

Insert Figure 4

The total interactions between *fac1* and other factors generally produce opposite signals
480 to \hat{f}_1 over middle and high latitudes in the two seasons. This result indicates that the changes in other factors tend to buffer the influence of monthly emission on increasing or decreasing f_0 concentrations. Some exceptions are seen over parts of East Asia in DJF for all species (Figure S5, right panels) and over the Southern Ocean in JJA for BaP (Figure S6, right panels) where the interactions work to reinforce the direct effects. In DJF, the degree of interactions is
485 smaller or comparable to the size of \hat{f}_1 for the Arctic and northern mid-latitudes but becomes stronger by at least double for Antarctica. The opposite tendency is seen in JJA but only applies to PYR and FLT. In agreement to \hat{f}_1 , the interaction effects are less apparent over the



tropics. Note that the positive effects in \hat{f}_{14} during local summer tend to be more dominant than the effects in other combinations for PHE, PYR, and FLT (Figures S17-S19). In the
490 simulation, the presence of re-volatilization in summer tends to suppress \hat{f}_1 by promoting more gases available for long-range transport, thus implies a negative feedback.

3.1.2 Effects of size-discretization of particulate-phase tracer

The direct effects of the *modal* scheme (\hat{f}_2) vary among species (Figure 4c). \hat{f}_2 is almost absent for PHE as the species resides almost completely in the gas phase. For PYR and
495 FLT, \hat{f}_2 is negative during DJF over northern mid-latitudes (\hat{f}_2/f_0 quartiles range from -5% to -30%) and the Arctic (-50% to -75%) whereas it is hardly visible in JJA or over other regions. Further analysis reveals stronger particle deposition results when the aerosol phase is discretized into different modes (not shown). In long-range transport under modal aerosol representation, the aerosols are more associated with larger particles hence particle
500 deposition becomes more effective. The choice of size discretization has only minor effects for atmospheric levels, except for BaP, especially during DJF, for which overestimates are significantly compensated for (Figure S24). Actually, for BaP, the *modal* scheme generally decreases the concentrations in the Arctic (as median, -35% in DJF and -15% during JJA) and increases (approx. 5%) those over mid- and low-latitude landmass (Figures S7d and S8d,
505 left panels).

As is the case for the direct effects, the interaction contributions are peculiar to individual species (Figure 4d). For PHE, the interaction effects in DJF are reflected in negative concentration responses over the Arctic (18%) and positive over Antarctica (7%), in contrast to relatively mild influences over other regions or in JJA. For PYR and FLT, the effects are
510 negative over the Arctic both in DJF (quartiles vary from -20% to -75%) and in JJA (-6% to -120%). It is interesting to note that the interaction between the *modal* and ppLFER schemes has a major influence on the negative signal (Figures S14, S15, S18, S19), suggesting that the decrease in simulated concentration associated with the change from *bulk* to *modal* could be intensified when the ppLFER scheme is used. In the remaining areas, the interaction effects
515 vary in sign spatially as illustrated in the right panels of Figures S7b,c–S8b,c. Nevertheless, it shows for both species that maximum influences occur over the Southern Ocean in DJF (where the effects may reach two orders of magnitude) and mid-latitude landmass in JJA (more than a factor of five). As for BaP, the median effects are negative (-7% to -30%) in both seasons, although some positive signals are apparent in parts of high latitudes while
520 the tropical oceans bear small synergistic effects. Similar to other species, the degree of interactions are stronger than \hat{f}_2 by more than a factor of three for the majority of grid cells (Figures S7d–S8d, right panels). The large fractions of the effects are dominated by two-factor and three-factor combinations related to the interaction with the ppLFER scheme and/or re-volatilization (Figures S16, S20).



525 3.1.3 Effects of the choice of gas–particle partitioning scheme

Figure 4e shows that the direct effects of the ppLFFER scheme (\hat{f}_3) show little spatial heterogeneities in both seasons and for all species. The effects are barely important for PHE due to low gas–aerosol partition constant (K_p). \hat{f}_3 is positive for PYR and FLT over polar regions and northern mid-latitudes especially in winter when low temperature favors partitioning to aerosols (higher K_p). The median of \hat{f}_3/f_0 varies from 1% to 25% with some parts of Antarctica showing an increase larger than 50%. For BaP, the effects are overall negative (by at least -5%) with \hat{f}_3/f_0 reflecting a positive north–south gradient (increasing from the Arctic to Antarctica), associated in part with stronger signals over oceans (Figures S9d and S10d, left panels). In particular under the modal size discretization, the choice of gas–particle partitioning scheme has only minor effects for atmospheric levels, except for BaP for which model overestimates are compensated by the choice of the ppLFFER scheme (Figure S24). Under the bulk size discretization, the ppLFFER scheme tends to enhance some of the overestimate in the Arctic summer (FLT, PYR; Figures S22–S23). The application of ppLFFER increases K_p as this module is calculated from not only interaction with BC and OM (as in Lohmann–Lammel scheme) but also with some other aerosol matrices. Higher K_p indicates higher particle mass fraction. For PYR and FLT, this leads to an increase in total atmospheric lifetime as the aerosol phase is not degraded, and can, therefore, be transported over a larger distance. For BaP, the additional particles are subject to depositions and heterogeneous oxidation by ozone, particularly in regions away from sources. The factor influence is notably too small for PHE as oxidations occur in both phases.

The effects from *fac3* interactions vary by region and are relatively stronger than \hat{f}_3 (Figure 4f). This finding is common to all species and seasons. The degree of effects is weaker for PHE compared to that for other species. However, the interactions increase polar concentrations in local summer, by 20% to a factor of five, mainly associated with the coupled effect of ppLFFER and volatilization (\hat{f}_{34} , Figures S13 and S17). For PYR and FLT, there is a high spatial variability over extratropical regions in local summer, as indicated by the interquartile range (distance between the 3rd and 1st quartiles). With regard to synergistic terms, ppLFFER interactions with the *modal* scheme and re-volatilization, in two- or three-factor combinations, are more important than other contributions (Figures S14–S15 and S18–S19). For BaP, the interaction effects show negative signals similar to \hat{f}_3 , suggesting a positive feedback. The interactions exert a stronger influence on the concentrations of the oceans than on that of land, except in the tropics (Figures S9d and S10d, right panels). The median of relative effects ranges from -1% to a factor of -10 , minimum (maximum) in the northern (southern) extratropics. Two second-order interactions likely make major contributions, that is, \hat{f}_{34} which dominates the response over oceans, and \hat{f}_{23} which dominates over land (Figure S16 and S20).

3.1.4 Effects of re-volatilization

The direct effects of re-volatilization (\hat{f}_4) are illustrated in Figure 4g. \hat{f}_4 is positive in the tropics in both seasons, with the median \hat{f}_4/f_0 ranging from 5% to 50%. Intensive re-



volatilization in this region would increase net surface fluxes, thereby increasing concentrations.
565 For PHE, positive \hat{f}_4 values are more localized over the tropical landmass, whereas negative
 \hat{f}_4 values are predicted over the tropical ocean (Figures S11a and S12a; left panels). The
positive (negative) effects over land (ocean) areas are also apparent at higher latitudes during
most of the year. This reflects the fact that the negative effects on concentrations over
ocean act contrary to the positive effects on net surface fluxes, mainly caused by the non-
570 linear relationships of air–sea gas exchange (deposition and volatilization), air and surface
burden, atmospheric oxidation, and emissions. Accounting for re-volatilization compensates
for a significant part of underestimates of PHE in the Arctic during summer, but adds to
overestimates in mid-latitudes (Figure S21).

For the studied species of mid semivolatility, PYR and FLT, a positive signal is apparent
575 over the high and middle latitudes during local summer in contrast to a negative signal
during local winter (Figure 4g). Similar to PHE, the negative signal is confined over oceans
(Figures S11b–c and S12b–c; left panels). The summer increases are stronger (20% to a
factor of ten) than the winter decreases (–10% to –60%) and the magnitudes are higher
in FLT than in PYR. The near-ground concentrations of PYR and FLT are estimated by
580 ≈ 30 –80% in mid-latitudes of which ≈ 30 % are attributable to re-volatilization. In the Arctic,
re-volatilization compensates for ≈ 60 % of PYR underestimation (Figure S22) and explains
most of ≈ 60 –80% of FLT overestimation (Figure S23). For BaP, \hat{f}_4 is positive consistently
across regions and seasons (\hat{f}_4/f_0 ranges from 20% to a factor of ten), with substantial effects
occurring over oceans (Figures S11d–S12d; left panels). Accounting for re-volatilization
585 creates some overestimates in the Arctic during summer (Figure S24). It should be noted
that the parameterization adopted here to describe volatilization from soils (the Smit scheme)
is derived from an experimental study on mid-polar to polar pesticides and there is a need to
validate and eventually sophisticate the parameterization to apolar substances.

The interactions generally point toward positive effects for the high-to-medium volatility
590 species (Figure 4h), despite some negative effects present over parts of the southern
(northern) oceans in DJF (JJA) (Figures S11–S12; right panels). As for BaP, the effects
are uniformly negative, inferring the interactions work in opposition to \hat{f}_4 . The negative
response is almost entirely caused by the negative \hat{f}_{34} , that is, the two-factor interaction
between re-volatilization and the ppLFER scheme (Figures S16,S20). Compared to \hat{f}_4 , the
595 degree of interactions are weaker for PHE, except in polar regions during local summer
where the interactions could amplify \hat{f}_4 . The above implies that \hat{f}_4 may point in the right
direction regardless of the influences from other factor changes. On the contrary, the degree
of interactions is overall comparable to \hat{f}_4 for the other species.

600 3.2 Model evaluation

Model performance using the sophisticated realization of the four features (factors), i.e.,
Seasonal emission + *Modal* scheme + ppLFER scheme + With re-volatilization (SMPW)
is presented below. Two predicted variables are evaluated, that is, total (gas+particle)



concentrations and aerosol particulate mass fraction at the lowest model level. The metrics
605 applied is listed in Supplement SV.

3.2.1 Near-surface air concentration

Comparison to land monitoring stations

Central tendency. Table 2 shows statistical indices for near-surface concentrations of atmospheric PAHs from observations and simulations and their comparisons, averaged across
610 stations in the Arctic, northern mid-latitudes, and the tropics. We can see that mean observed concentrations are higher for PHE and smaller for BaP over all regions. Furthermore, the Arctic concentrations are lower than those in the northern mid-latitudes by a factor of around 20 and those in the tropics by approx. two orders of magnitude. The model captures well these species and regional variations, but the magnitudes are both under- and over-estimated. In
615 the Arctic, it underestimates PHE ($MB = -0.060 \text{ ng m}^{-3}$) and BaP ($MB = -0.006 \text{ ng m}^{-3}$) concentrations but slightly overestimates PYR ($MB = 0.001 \text{ ng m}^{-3}$) and FLT ($MB = 0.04 \text{ ng m}^{-3}$). In the NH mid-latitudes, the model overestimates the three species predominantly occur in the gas phase ($MB = 0.077\text{--}0.867 \text{ ng m}^{-3}$) but underestimates BaP ($MB = -0.58 \text{ ng m}^{-3}$). Negative bias is seen in the tropics for three PAHs ($MB = -3.443$ to -6.851 ng m^{-3}).
620 Nevertheless, the comparison of model and observations at individual monitoring station can be different from the regional mean statistics, as described in Supplement SVIII. Comparing all four PAHs, a larger degree of bias is found for BaP which increases from the northern mid-latitudes ($NMB = -0.58$, $NMBF = -1.40$, $FAC2 = 0.31$, $FAC10 = 0.79$) to the Arctic ($NMB = -0.92$, $NMBF = -12.17$, $FAC2 = 0.17$, $FAC10 = 0.33$).

625 Insert Table 2

Dispersion of monthly concentrations. In the following, the coefficient of variation (CoV) is used to compare the dispersion of concentrations among species of different ranges. CoV was calculated by dividing standard deviation (SD) of all data points by its mean value (\bar{x}). The observations show high variability ($CoV > 1$) with CoV ranging between 1.12 and
630 2.14. The simulated concentrations appear to be less dispersed than the observations ($CoV = 0.78\text{--}1.93$) except for the Arctic PHE and PYR concentrations. The degree of underestimation is larger in the tropics with CoV being 30%–50% smaller than the observations. Furthermore, correlations between predicted and observed concentrations are weaker than those in other regions where r varies between 0.29–0.63 (the model reproduces 8%–40% of the variance in observed concentrations). Comparing the four species, the simulated BaP shows greater
635 underpredictions of the variability where CoV values are less than half of those observed and correlations are less than 0.2 (accounting not over than 4% of observed variance). Higher variability in BaP measurements (than in model results) can be influenced by strongly varying emissions in source regions that are not reflected in emission inventory (Matthias et al., 2009).



640 *Seasonal variation.* Figure 5 compares simulated and observed seasonal cycle of average
concentrations for different species and regions. The observed mean concentrations are largest
in winter and lowest during summer because of less emission and the strong presence of
OH for oxidation. The winter maximum to summer minimum ratio (amplitude) is more
pronounced (by more than a factor of two) in the Arctic than that in the NH mid-latitudes.
645 The seasonality between model and observations is in a qualitative agreement, particularly
over the Arctic (except in summer) and mid-latitudes. In the Arctic, the model overestimates
the seasonal amplitude of PHE and BaP and underestimates their mean concentrations. The
contrast is seen for PYR and FLT. FLT concentration is overestimated by up to a factor
of three in summer while PYR is quite well predicted. In the NH mid-latitudes, the model
650 underestimates the amplitude but overestimates the concentrations of PHE, PYR, and FLT
(by typically a factor of two), whereas a systematic negative bias is found for BaP. In the
tropics, both the amplitude and magnitude are too low in the model (for magnitude, by a
factor of 2–5).

Insert Figure 5

655 *Intermodel comparison.* A comparison of simulated PAH concentrations from the new EMAC
model (with SVOC submodel) for the year 2007–2009 and from GEOS-Chem model for
2005–2009 (Friedman and Selin, 2012) revealed that both models overall overestimate observed
PHE and PYR over mid-latitudes whereas the Arctic concentrations for both species are
underestimated by EMAC but overestimated by GEOS-Chem in winter. The discrepancies
660 are more likely to be caused by different emission inventory and by high spatial variability
in the emissions which results in changes in PAH concentrations on scales not captured
by both GCMs with a grid spacing of >100 km. It is also noteworthy that the simulation
carried out with GEOS-Chem differs in complexity to EMAC, that is, the model neglected
volatilization and seasonal variation in emissions and applied a dual BC adsorption and OM
665 absorption gas–particle partitioning. Nevertheless, both models systematically underestimate
the BaP concentrations over all regions, reflecting a difficulty in modeling the behavior of this
species. Model bias is usually larger for the Arctic sites than for the mid-latitude sites and the
agreement with observations is more satisfactory in summer. BaP underestimation in remote
regions has also been reported in other global modeling studies (Sehili and Lammel, 2007;
670 Shrivastava et al., 2017). Despite uncertainties in the emissions, which seasonality has been
taken into account in our simulations, this finding suggests that current models underestimate
the role of temperature for the long-range transport of BaP. A GEOS-Chem configuration
which neglected heterogeneous oxidations for BaP yielded predicted concentrations higher
than measurements over NH mid-latitudes but rather in a good agreement over the Arctic
675 (Friedman and Selin, 2012), suggesting a strong dependence of BaP long-range transport on the
temperature sensitivity of particulate-phase oxidation. Similarly, agreement/overprediction
of BaP was found using ECHAM5/HAM model (Lammel et al., 2009) when setting $k_{\text{O}_3, \text{het}}^{(2)}$
to zero (BAPOB scenario; agreement for the Zeppelin station, overestimation for the Alert



station). In addition, the latest global BaP modeling study carried out with CAM5 model
680 (Shrivastava et al., 2017) also found that temperature and humidity are important factors
to affect heterogeneous oxidation kinetics of BaP and influence the long-range transport.
Their default simulation underestimates observed BaP concentrations by approximately 77%
whereas the predictions improved when the model assumed that BaP underwent a complete
shielding by highly viscous organic aerosols under cool/dry conditions. It is important to
685 understand further the effect of the variation of temperature and humidity on reactive rate
constants of the heterogeneous oxidations during the long-range transport. This topic is
discussed in details in Mu et al. (2018).

3.2.3 Comparison to ship cruise measurements

Measurements of PHE, PYR, and FLT concentrations over the Atlantic Ocean were taken
690 during a cruise in July 2009 (Lohmann et al., 2013). Figure 6 shows the ship sample
concentrations overlaying the simulated PAH concentrations. Sample arithmetic (geometric)
means during the whole cruise transect are 322 (209), 95 (88), and 128 (111) pg m^{-3} for PHE,
PYR, and FLT respectively. The model poorly reproduces the remote marine environments
and overall underestimates the observations, except at 3 locations along the North American
695 coast. The simulated means across sampling positions are 23 (7), 20 (3), and 39 (2) pg m^{-3}
respectively and the underestimation ranges from a factor of 2 to 1000. The degree of bias is
most apparent over the tropical South Atlantic at latitude bands 5°S–15°S.

As reported in Liu et al. (2014), the measured concentrations of BaP over the Asian
marginal seas, the Indian Ocean, the South and North Pacific Oceans are 131 (45), 14 (3),
700 9 (2), and 8 (3) pg m^{-3} , respectively, for the arithmetic (geometric) means of all samples.
Similar to other species, the model also underestimates the BaP concentrations with mean
values being 75 (15), 4 (0.05), 0.09 (0.03), and 0.2 (0.06) pg m^{-3} , respectively. The discrepancy
appears relatively smaller over the Asian marginal seas as compared to other locations (Figure
7). A substantial degree of bias is seen over the Indian Ocean covering approximately the area
705 bounded by 70°E–90°E and 10°S–30°S, with simulated values being more than two orders of
magnitude smaller than the observed.

Insert Figure 6

Insert Figure 7

The model tendency to underestimate the marine air concentrations may likely be due to
710 several factors: (a) The grid resolution is not sufficient to reproduce fine-scale processes at the
grid points close to shipping tracks; (b) Great uncertainties associated with the air–sea gas
exchange parameterizations still exist, most notably in the estimation of gas transfer velocity;
(c) The global inventory (Shen et al., 2013) may significantly underestimate emissions from
ocean shipping and do not well characterize the spatial and temporal variability of biomass
715 burning plumes as another potential origins of pollutants in the marine air (Nizzetto et al.,



2008); (d) PAH concentration over remote oceans is controlled by atmospheric components (e.g., temperature, wind speed, boundary layer height, photochemical degradation) and the dynamical and biogeochemical components of the ocean. However, the ocean components have not been covered in the simulation; (e) The particulate-bound PAHs may undergo too fast heterogeneous oxidation (most relevant for BaP), leading to short atmospheric lifetimes and weaker long-range transport.

3.2.4 Particulate mass fraction

Measurements of particulate mass fraction (θ) were available only from E3 station in Europe and IADN stations (I1-I7) in North America (see Table S11). Table 3 presents summary statistics on monthly mean θ from observations and simulations including some performance metrics. The observed mean θ is smaller for PHE (0.051 ± 0.035) and higher for BaP (0.949 ± 0.067). This result is expected as volatility decreases (hence θ increases) from (lighter) PHE to (heavier) BaP. The θ values for PYR and FLT are larger by over five times than those for PHE and lower by around one-third than those for BaP. The model reproduces well the distinct differences among species but underestimates the observed θ for PHE, PYR and FLT. The degree of negative bias is relatively large in PHE (NMB = -0.910 and NMBF = -10.145), whereas for the isomer pair of PYR and FLT, the model exhibits a similar performance with a slight improvement in PYR (NMB = -0.410 and NMBF = -0.694). With regard to BaP, there is a satisfactorily small bias (MB = 0.015 , RMSE = 0.074 , NMB = 0.016 , and NMBF = 0.016) although the observed and simulated values have a very weak correlation ($r = 0.03$).

Figure 8 shows the seasonal mean θ averaged over three years for all PAHs. Observations show that θ for BaP varies less than those for 3–4 ring PAHs. Although the model adequately reproduces this feature as well as seasonal variation of individual species, the simulated θ of PHE, PYR, and FLT are generally lower than the observations (except for PYR in winter). For BaP, differences between model and observations are less than 10% in all months. The SVOC submodel describes the gas–particle partitioning of atmospheric SOCs as a function of temperature and aerosol phase composition. The underestimation might be related to the fact that the submodel assumes the particle to be fully in equilibrium with the gas phase at all times. It neglects kinetic limitations of molecular diffusivity that could lead to the trapping of particles inside viscous (or semisolid) organic aerosol coating. This shielding effect increases equilibration times of the particles, thereby reducing part of θ from the mass available for gas–particle partitioning. Deviations from measurements can also be partly attributed to the locations of some stations that are within, or close to, residential and industrial area (namely, I4, I6, and I7) where the scale and gradient in anthropogenic emissions are not resolved by the model grid resolution nor represented by the emission inventory.

[Insert Table 3](#)

[Insert Figure 8](#)



4 Summary and conclusions

The submodel SVOC has been developed and operated within the EMAC model for the application to global distribution and environmental fate of SOCs. In this first development, the focus was set on the predictions of four PAH species: phenanthrene (PHE), pyrene (PYR), fluoranthene (FLT), and benzo(a)pyrene (BaP). Multicompartmental fate and air–surface exchange processes were included in SVOC. Some novel features in PAH modeling were tested, including seasonality in emissions, the *modal* scheme for particulate-phase tracer representation, the ppLFER scheme for gas–particle partitioning, and re-volatilization from surfaces. The results indicate that using seasonal emission compensates for model biases in the predictions of more volatile species (PHE) whereas the effects of the *modal* and ppLFER schemes are of less significance. Re-volatilization increases the near-ground concentrations in air, which is found most significant for species of mid semivolatility (PYR and FLT). Hereby attribution of model response to individual features (factors) is blurred by the non-linear interactions between two and more factors. The effects of these interactions are found to both reinforce (positive feedback) and suppress (hence, negative feedback) the effects of the individual factors.

For near-surface concentrations, model bias varies by region and/or species, being negative (positive) in the Arctic within typically a factor of 2–13 (6% to a factor of two) for PHE and BaP (PYR and FLT), positive in the northern mid-latitudes for PHE, PYR, and FLT by up to a factor of three, negative in the Tropics (by a factor of 2–3) and largely over ocean up to a factor of three orders of magnitude. The model adequately reproduces the seasonal variation of the particulate mass fraction (θ), but underestimates θ for high-to-medium volatility PAHs. This might be related to a systematic underestimation of OC by the model, which neglects secondary organic aerosols (SOA). The latter may cause significant underestimation of the overall atmospheric aerosol burden and θ of SOCs, in particular over ocean. Since recently a MESSy submodel, ORACLE, dedicated to the simulations of SOA (Tsimpidi et al., 2014) based on lumping organic species in volatility bins is available. It should be included in future SOC simulations using EMAC.

Moreover, SVOC implicit assumption of instantaneous gas–particle equilibrium may cause both over- and underestimates of θ , as inter-phase mass transfer may be kinetically limited to gaseous sources (hence, overestimate of θ) or within the particle bulk (hence, underestimate θ), as the PAHs may become trapped within particles during transport (Friedman et al., 2014; Zelenyuk et al., 2012; Mu et al., 2018). For multidecadal studies, the coupling of a 3D ocean model (coupled with a marine biogeochemistry module) would be needed since the present model application does not allow for horizontal and vertical transports in the deep ocean. For the same reason, contaminant remobilization within deep soil layers should also be introduced. To this end, a multi-layer (3D) soil compartment would be needed to replace the 2D soil compartment used here.



Code availability

SVOC submodel presented here has been based on the Modular Earth Submodel System (MESSy) version 2.50 and the global atmospheric model ECHAM version 5.3.02. MESSy is continuously further developed and applied by a consortium of institutions. The usage
795 of MESSy and access to the source code is licensed to all affiliates of institutions which are members of the MESSy Consortium. Institutions can be a member of the MESSy Consortium by signing the MESSy Memorandum of Understanding. More information can be found on the MESSy Consortium website (<http://www.messy-interface.org>). The SVOC submodel will be incorporated into the next released version of the ECHAM/MESSy (EMAC) model (v2.55)
800 and will therefore be made publicly available (with respect to the EMAC license regulations).

Author contributions

M.O. and G.L. conceived the study and designed the experiments. M.O. developed the SVOC submodel with input from all co-authors. M.O. performed model simulations and data analyses. M.O and G.L discussed the results. M.O. wrote the manuscript with contributions
805 from all co-authors.

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Figures and Tables

Table 1. Summary of MESSy process submodels used in the study

Submodel	Purpose	Reference
AEROPT	Aerosol optical properties	Lauer et al. (2007)
AIRSEA	Air–sea exchange	Pozzer et al. (2006)
CLOUD	ECHAM5 cloud and precipitation scheme as MESSy submodel	Roeckner et al. (2006) and references therein
CONVECT	Convection parameterizations	Tost et al. (2006b, 2010)
CVTRANS	Convective tracer transport	Tost (2006)
DDEP	Dry deposition of gases and aerosols	Kerkweg et al. (2006a)
GMXe	Aerosol dynamics and thermodynamics	Pringle et al. (2010)
JVAL	Rate of photolysis	based on Landgraf and Crutzen (1998)
LNOX	NO _x production from lightning	Tost et al. (2007)
MECCA	Tropospheric and stratospheric chemistry	Sander et al. (2011)
OFFEMIS	Offline emissions	Kerkweg et al. (2006b)
ONEMIS	Online emissions	Kerkweg et al. (2006b)
RAD	ECHAM5 radiation scheme as MESSy submodel	Roeckner et al. (2006); Jöckel et al. (2006)
SCAV	Scavenging of gases and aerosols	Tost et al. (2006a)
SEDI	Aerosol sedimentation	Kerkweg et al. (2006a)

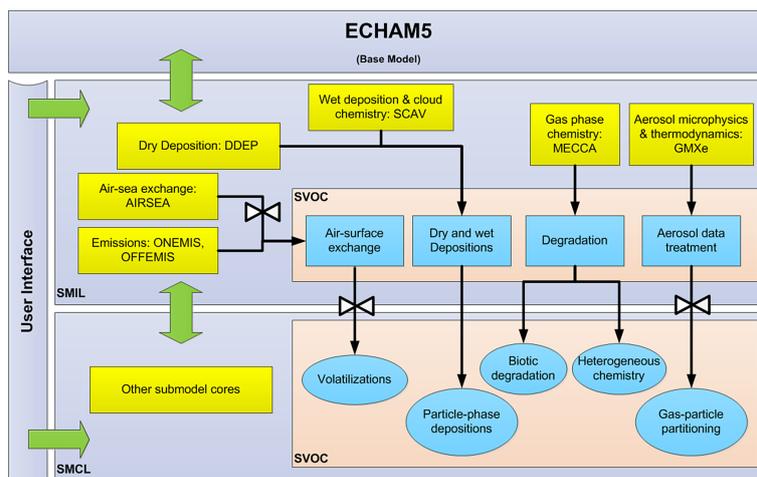


Figure 1. Overview of EMAC-SVOC model structure, the cycling processes in SVOC submodel and its interaction with other MESSy submodels. SMIL (submodel interface layer) and SMCL (submodel core layer) are components of MESSy coding standard, see (Jöckel et al., 2006) for further details.

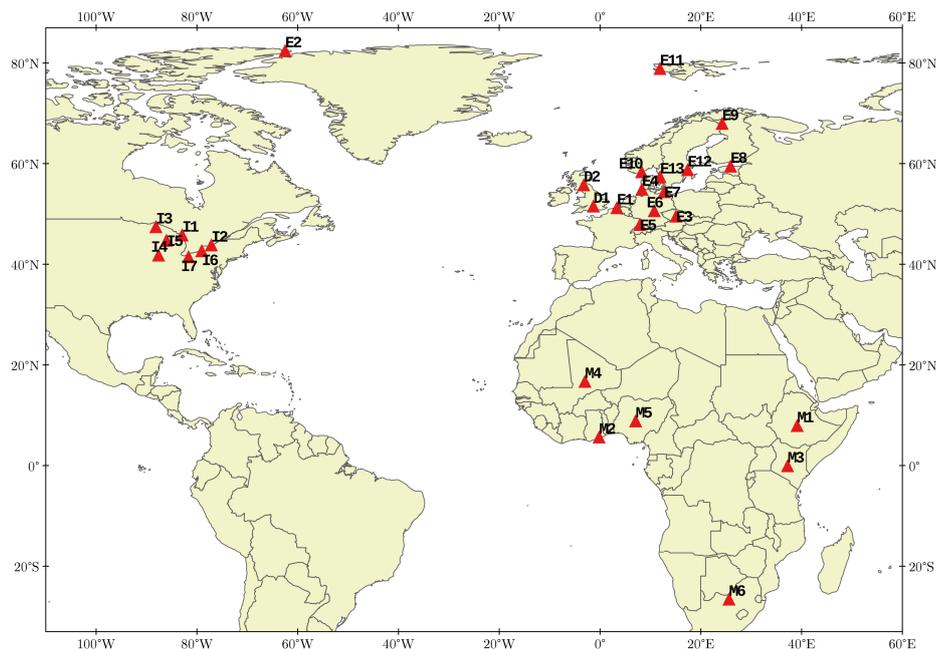


Figure 2. Locations of monitoring stations used in the study. The initial letter of each station ID refers to the individual monitoring network (E: EMEP & AMAP; D: DEFRA, I: IADN, M: MONET-Africa)



	Emission	Particle Representation	Gas-Particle Partitioning	Volatilization
f_0	Annual	Bulk	L+L	No
f_1	Seasonal	Bulk	L+L	No
f_2	Annual	Mode	L+L	No
f_3	Annual	Bulk	PpLFER	No
f_4	Annual	Bulk	L+L	With
f_{12}	Seasonal	Mode	L+L	No
f_{23}	Annual	Mode	PpLFER	No
f_{34}	Annual	Bulk	PpLFER	With
f_{13}	Seasonal	Bulk	PpLFER	No
f_{14}	Seasonal	Bulk	L+L	With
f_{24}	Annual	Mode	L+L	With
f_{123}	Seasonal	Mode	PpLFER	No
f_{124}	Seasonal	Mode	L+L	With
f_{134}	Seasonal	Bulk	PpLFER	With
f_{234}	Annual	Mode	PpLFER	With
f_{1234}	Seasonal	Mode	PpLFER	With

Figure 3. List of experiments performed for the factor separation analysis to study sensitivity to temporal variation in emission and process parameterizations (particulate-phase representation, gas–particle partitioning scheme, and volatilization). L+L: Lohmann–Lammel, ppLFER: Poly Parameter Linear Free Energy Relationships.

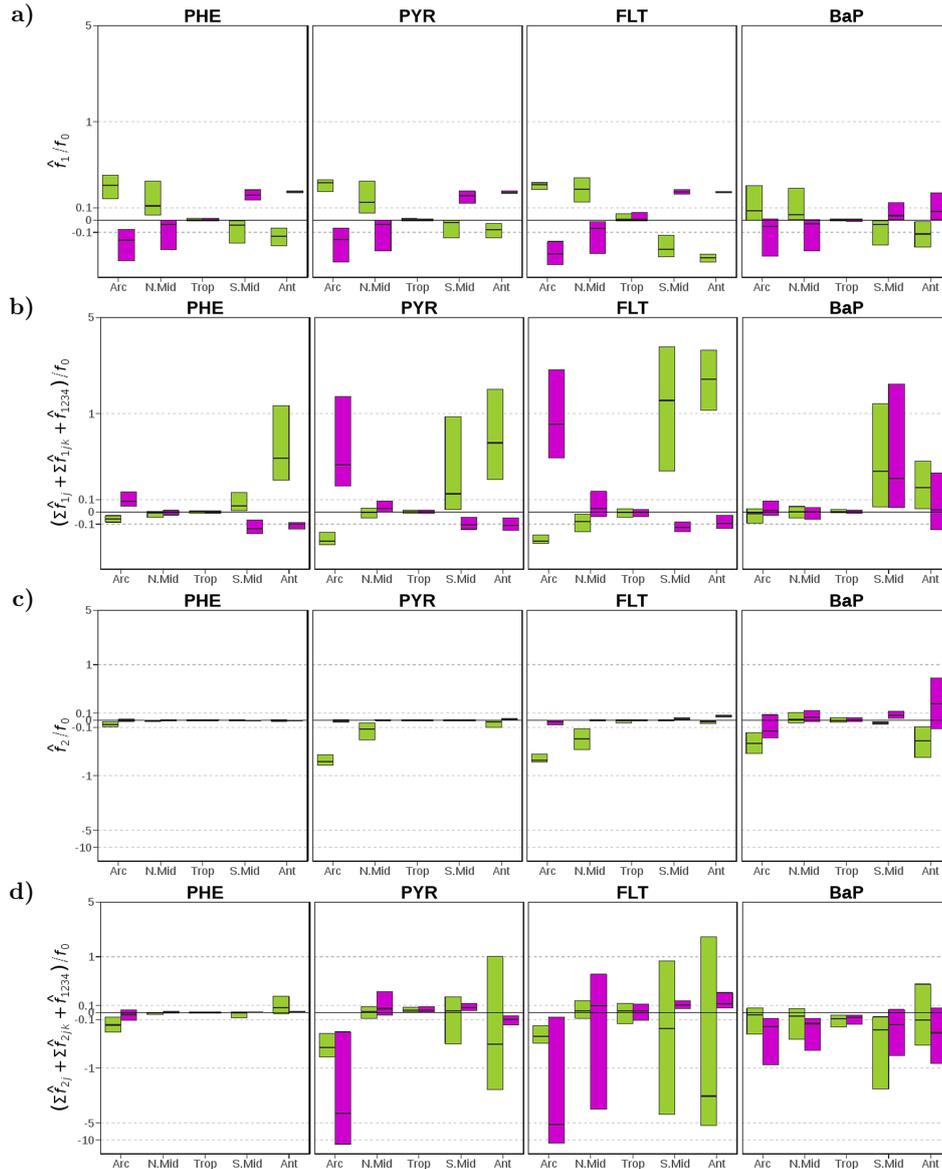


Figure 4. Direct and interaction effects on seasonal-mean near-surface PAH concentrations of (a,b) monthly emissions ($i = 1$), (c,d) the *modal* scheme ($i = 2$), (e,f) the ppLFFER scheme ($i = 3$), and (g,h) volatilization ($i = 4$). The direct effects (a,c,e,g) are expressed as the difference between two distributions ($\hat{f}_i = f_i - f_0$) whereas the interaction effects (b,d,f,h) are expressed as the sum of two ($\Sigma \hat{f}_{ij}$, $i \neq j$), three ($\Sigma \hat{f}_{ijk}$, $i \neq j \neq k$), and all (\hat{f}_{1234}) factor interactions. They are presented as relative to concentrations from the *base* (f_0) simulation. The figures display the median, 25th and 75th percentiles of the relative effects over each of five main climatic regions. Note the inverse hyperbolic sine function has been used in scaling the *y* axes.

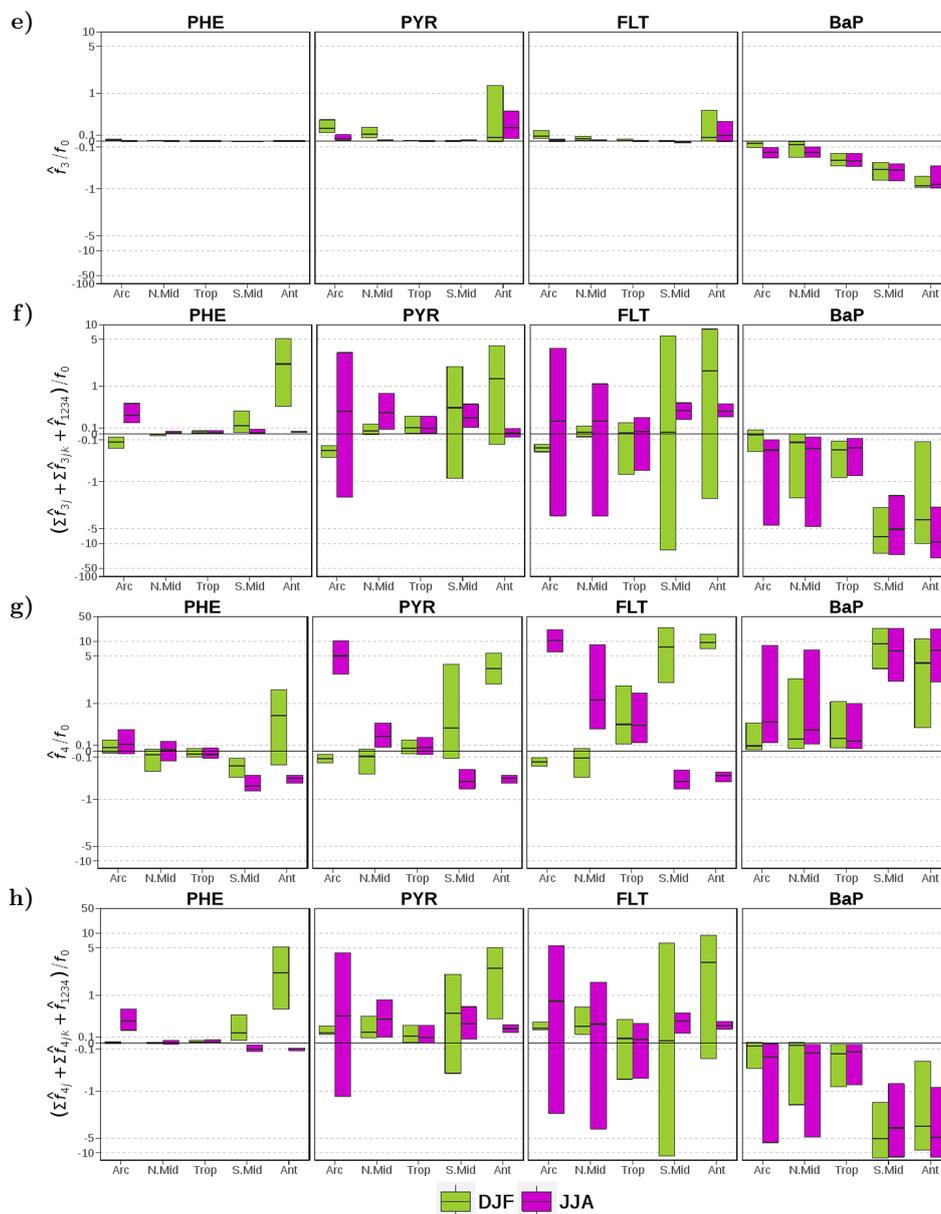


Figure 4. continued



Table 2. Statistics comparison of model simulation and observations of total (gas+particle) concentrations of PAHs from stations in the Arctic, northern mid-latitudes and tropics. N : Number of observed-simulated monthly data pairs; \bar{x} : Mean; $Q2_x$: Median; SD_x : Standard deviation; GM_x : Geometric mean; x : Simulated (M) or Observed (O) data; MB: Mean bias; RMSE: Root mean square error; NMB: Normalized mean bias; NMBF: Normalized mean bias factor; FAC2: Factor of 2; FAC10: Factor of 10; r : Correlation coefficient.

Metrics	Unit	Arctic						NH mid-latitudes						Tropics			
		PHE	PYR	FLT	BaP	PHE	FLT	PHE	PYR	FLT	BaP	PHE	FLT	PYR	PHE	PYR	FLT
N	months	89	89	89	46	361	328	372	405	34	34	34	34	34	34	34	34
$N_{<LOQ}$	months	0	0	0	30	0	0	0	0	0	0	0	0	0	0	0	0
\bar{O}	ng m ⁻³	0.107	0.024	0.039	0.007	2.193	0.408	0.803	0.141	11.818	6.431	6.843	6.843	6.843	6.431	6.843	6.843
$Q2_O$	ng m ⁻³	0.034	0.014	0.012	0.002	1.301	0.194	0.360	0.037	3.608	2.106	2.181	2.181	2.106	2.106	2.181	2.181
SD_O	ng m ⁻³	0.162	0.027	0.054	0.015	2.956	0.582	1.135	0.253	16.598	10.141	10.217	10.217	10.141	10.141	10.217	10.217
GM_O	ng m ⁻³	0.051	0.014	0.018	0.003	0.968	0.221	0.383	0.046	3.733	1.369	1.726	1.726	1.369	1.369	1.726	1.726
\bar{M}	ng m ⁻³	0.046	0.025	0.079	5.2E-4	2.270	1.086	1.670	0.059	4.966	2.005	3.400	3.400	2.005	2.005	3.400	3.400
$Q2_M$	ng m ⁻³	0.010	0.007	0.034	1.9E-5	0.840	0.500	0.736	0.022	4.274	1.236	2.012	2.012	1.236	1.236	2.012	2.012
SD_M	ng m ⁻³	0.089	0.040	0.099	6.5E-4	2.955	1.225	2.007	0.085	3.897	2.019	3.462	3.462	2.019	2.019	3.462	3.462
GM_M	ng m ⁻³	0.012	0.008	0.041	3.4E-5	1.144	0.635	0.913	0.028	2.816	1.038	1.788	1.788	1.038	1.038	1.788	1.788
MB	ng m ⁻³	-0.060	0.001	0.040	-0.006	0.077	0.679	0.867	-0.083	-6.851	-4.426	-3.443	-3.443	-4.426	-4.426	-3.443	-3.443
RMSE	ng m ⁻³	0.118	0.038	0.099	0.016	3.564	1.404	2.383	0.279	16.005	10.631	10.392	10.392	10.631	10.631	10.392	10.392
NMB	-	-0.56	0.06	1.04	-0.92	0.04	1.66	1.08	-0.58	-0.58	-0.69	-0.50	-0.50	-0.58	-0.69	-0.50	-0.50
NMBF	-	-1.30	0.06	1.04	-12.17	0.04	1.66	1.08	-1.40	-1.38	-2.21	-1.01	-1.01	-1.38	-2.21	-1.01	-1.01
FAC2	-	0.20	0.28	0.30	0.17	0.40	0.39	0.30	0.31	0.26	0.24	0.29	0.29	0.26	0.24	0.29	0.29
FAC10	-	0.82	0.90	0.94	0.33	0.90	0.84	0.83	0.79	0.97	0.79	0.85	0.85	0.97	0.79	0.85	0.85
r	-	0.83	0.42	0.42	0.16	0.27	0.23	0.09	0.01	0.63	0.33	0.29	0.29	0.63	0.33	0.29	0.29

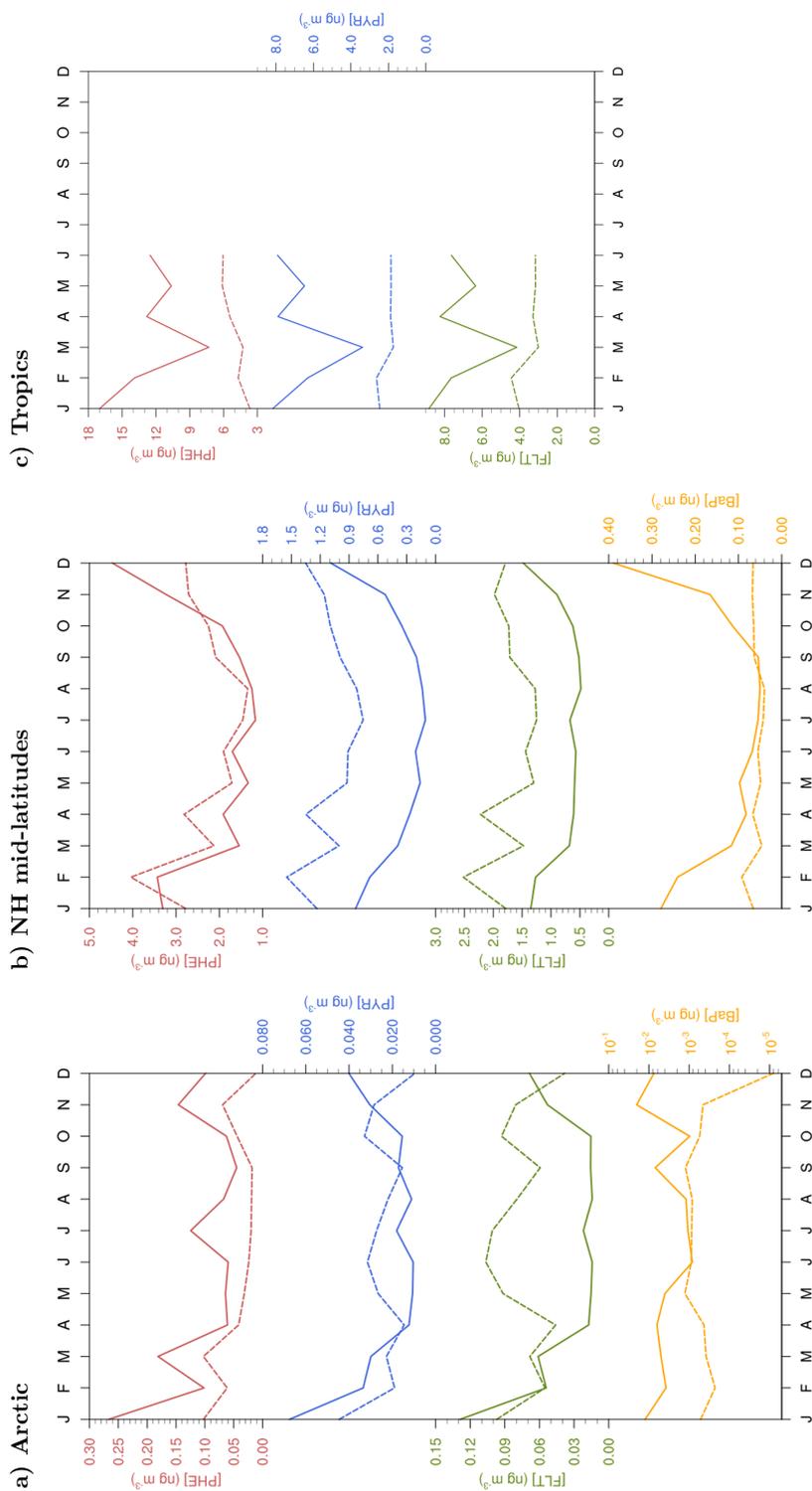


Figure 5. Seasonal mean total (gas+particle) concentrations of PAHs (ng m^{-3}) from observations (dashed lines) and simulations (solid lines) averaged over all stations in the (a) Arctic, (b) northern mid-latitudes, and (c) tropics. Note that logarithmic scale has been used for BaP concentrations in the Arctic.

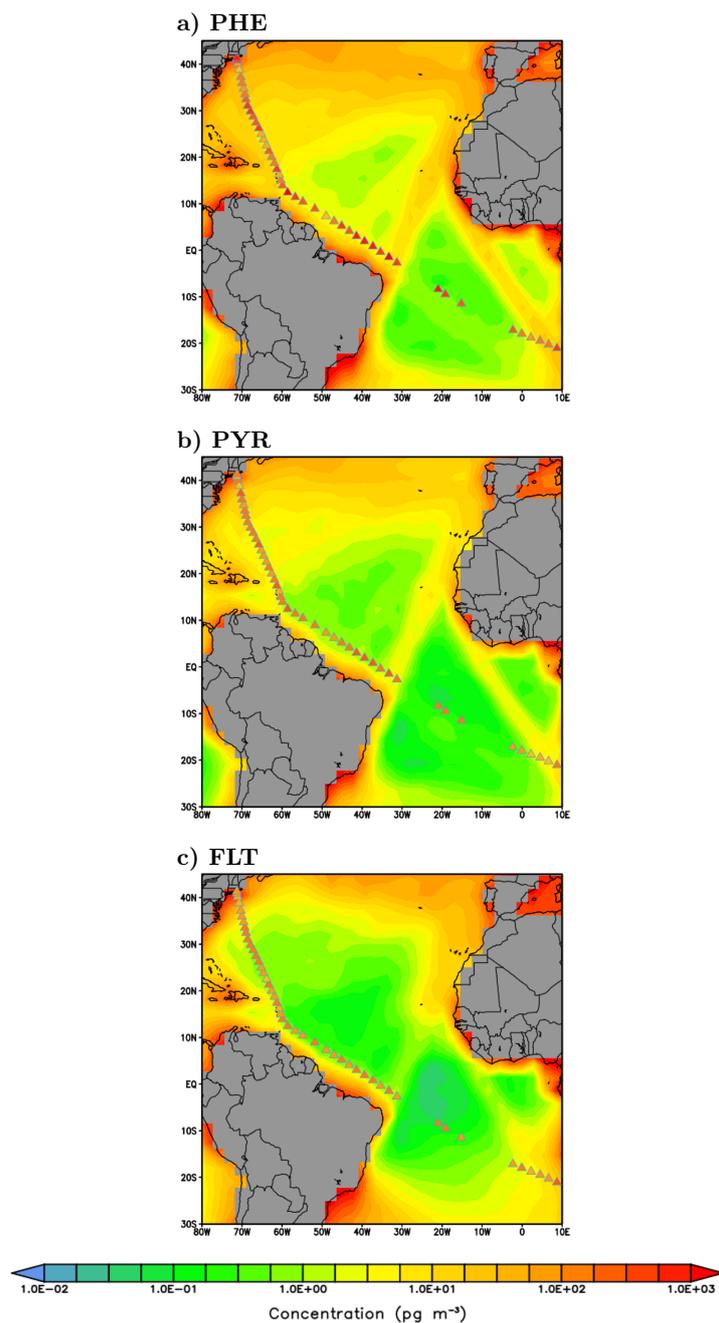


Figure 6. Simulated concentrations of PHE, PYR, and FLT (pg m^{-3}) over the Atlantic ocean overlaid with concentrations from a ship cruise measurement campaign during July 2009 (triangles). Land grid cells are depicted in gray shades.

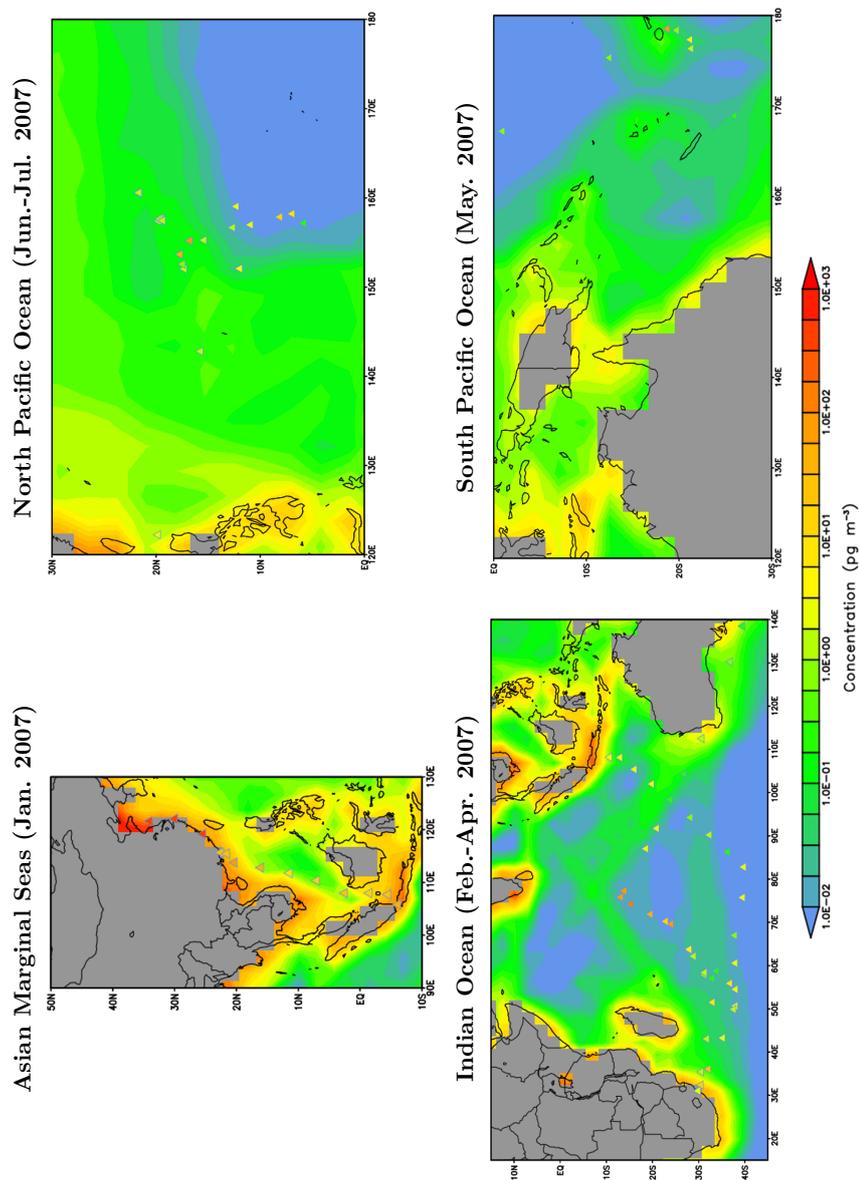


Figure 7. Simulated BaP concentrations (pg m^{-3}) over the four ocean margins overlaid with concentrations from a ship cruise measurement campaign (triangles). Land grid cells are depicted in gray shades.



Table 3. Statistics comparison of model simulation and observations of particulate mass fraction (θ) from a subset of surface stations, as listed in Table S11. N : Number of observed-simulated monthly data pairs; \bar{x} : Mean; SD_x : Standard deviation; x : Simulated (M) or Observed (O) data; MB: Mean bias; RMSE: Root mean square error; NMB: Normalized mean bias; NMBF: Normalized mean bias factor; FAC2: Factor of 2; FAC10: Factor of 10; r : Correlation coefficient.

Metrics	PHE	PYR	FLT	BaP
N	63	63	99	93
\bar{O}	0.051	0.359	0.268	0.949
SD_O	0.035	0.150	0.162	0.067
\bar{M}	0.005	0.212	0.106	0.964
SD_M	0.005	0.138	0.086	0.027
MB	-0.046	-0.147	-0.162	0.015
RMSE	0.057	0.214	0.225	0.074
NMB	-0.910	-0.410	-0.604	0.016
NMBF	-10.145	-0.694	-1.523	0.016
FAC2	0.00	0.56	0.30	1.00
FAC10	0.38	1.00	0.94	1.00
r	0.42	0.42	0.33	0.03

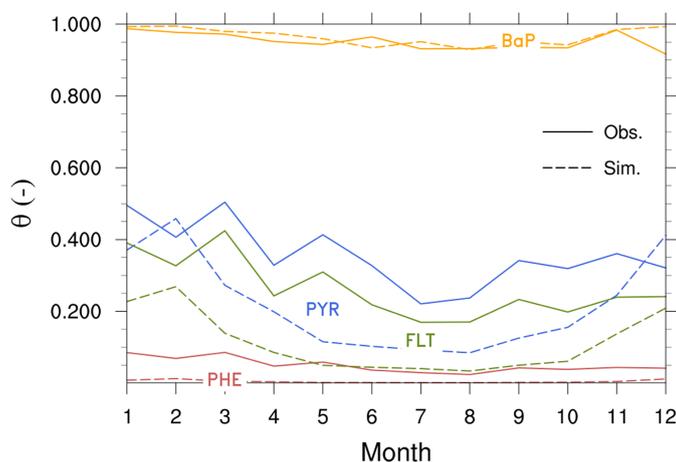


Figure 8. Seasonal mean particulate mass fraction (θ ; unitless) from observations (solid lines) and simulations (dashed lines)