



1 **Developing a monthly albedo change radiative forcing kernel from satellite climatologies**
2 **of Earth's shortwave radiation budget: CACK v1.0**

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13 **Abstract**

14 Due to the potential for land use / land cover change (LULCC) to alter surface albedo, there is
15 need within the LULCC science community for simple and transparent tools for predicting
16 radiative forcings (ΔF) from surface albedo changes ($\Delta\alpha_s$). To that end, the radiative kernel
17 technique – developed by the climate modeling community to diagnose internal feedbacks
18 within general circulation models (GCMs) – has been adopted by the LULCC science
19 community as a tool to perform offline ΔF calculations for $\Delta\alpha_s$. However, the GCM codes
20 are not readily transparent and the atmospheric state variables used as model input are limited
21 to single years, thus being sensitive to anomalous weather conditions that may have occurred
22 in those simulated years. Observation-based kernels founded on longer-term climatologies of
23 Earth's atmospheric state offer an attractive alternative to GCM-based kernels and could be
24 updated annually at relatively low costs. Here, we evaluate simplified models of shortwave
25 radiative transfer as candidates for an albedo change kernel founded on the Clouds and the



26 Earth's Radiant Energy System (CERES) Energy Balance and Filled (EBAF) products. We
27 find that a new, simple model supported by statistical analyses gives remarkable agreement
28 when benchmarked to the mean of four GCM kernels and to two GCM kernels following
29 emulation with their own boundary fluxes as input. Our findings lend support to its candidacy
30 as a satellite-based alternative to GCM kernels and to its application in land-climate studies.

31

32 **1. Introduction**

33 Diagnosing changes to the shortwave radiation balance at the top-of-the-atmosphere (TOA)
34 resulting from changes to albedo at the surface ($\Delta\alpha_s$) is an important step in predicting
35 climate change. However, outside the climate science community, many researchers do not
36 have the tools to convert $\Delta\alpha$ to the climate-relevant ΔF measure (Bright, 2015; Jones et al.,
37 2015), which requires a detailed representation of the atmospheric constituents that absorb or
38 scatter solar radiation (e.g. cloud, aerosols, and gases) and a sophisticated radiative transfer
39 code. For single points in space or for small regions, these calculations are typically
40 performed offline – meaning without feedbacks to the atmosphere (e.g., (Randerson et al.
41 2006)). Large-scale investigations (e.g. Amazonian or pan-boreal LULCC (Dickinson and
42 Henderson-Sellers, 1988; Bonan et al., 1992)) typically prescribe the land surface layer in a
43 GCM with initial and perturbed states, allowing the radiative transfer code to interact with the
44 rest of the model. While this has the benefit of allowing interaction and feedbacks between
45 surface albedo and scattering or absorbing components of the model, such an approach is
46 computationally expensive and thereby restricts the number of LULCC scenarios that can be
47 investigated (Atwood et al., 2016). Consequently, this method does not meet the needs of
48 some modern LULCC studies which may require millions of individual land cover transitions
49 to be evaluated cost effectively (Lutz and Howarth, 2015; Ghimire et al., 2014).

50



51 Within the LULCC science community, two methods have primarily met the need for
52 efficient ΔF calculations from $\Delta\alpha_s$: simplified parameterizations of atmospheric transfer of
53 shortwave radiation (Bright and Kvalevåg, 2013;Cherubini et al., 2012;Bozzi et al.,
54 2015;Muñoz et al., 2010;Caiazzo et al., 2014;Carrer et al., 2018), and radiative kernels
55 (Ghimire et al., 2014;O'Halloran et al., 2012;Vanderhoof et al., 2013) derived from
56 sophisticated radiative transfer schemes embedded in GCMs (Soden et al., 2008;Shell et al.,
57 2008;Pendergrass et al., 2018;Block and Mauritsen, 2014). Simplified parameterizations of
58 the LULCC science community have not been evaluated comprehensively in space and time.
59 Bright & Kvalevåg (2013) evaluated the shortwave ΔF parameterization of Cherubini *et al.*
60 (2012) when applied at several sites distributed globally on land, finding inconsistencies in
61 performance at individual sites despite good overall cross-site performance. Radiative kernels
62 (Soden et al., 2008;Shell et al., 2008;Pendergrass et al., 2018;Block and Mauritsen,
63 2014;Smith et al.) – while being based on state-of-the-art models of radiative transfer – have
64 the downside of being model-dependent and not readily transparent. While the radiative
65 transfer codes behind them are well-documented, the scattering components (i.e. aerosols,
66 gases, and clouds) affecting transmission have many simplifying parameterizations, vary
67 widely across models, and may contain significant biases (Dolinar et al., 2015;Wang and Su,
68 2013). An additional downside is the that atmospheric state variables used as model input are
69 limited to single years, thus being sensitive to anomalous weather conditions that may have
70 occurred in those years. Further, the application of a state-dependent GCM kernel may be
71 undesirable in regions undergoing rapid changes in cloud cover or aerosol optical depth, such
72 as in the northwest United States (Free and Sun, 2014) and in southern and eastern Asia (Zhao
73 et al., 2018;Srivastava, 2017), respectively. A kernel based on remotely-sensed observations
74 could be updated annually to capture changes in atmospheric state at relatively low costs.

75



76 Within the atmospheric science community, simplified radiative transfer frameworks have
77 been developed, either to diagnose effective surface and atmospheric optical properties from
78 climate model outputs, or to study the relative contributions of changes to these properties on
79 shortwave flux changes at the top and bottom of the atmosphere (Rasool and Schneider,
80 1971;Winton, 2005;Winton, 2006;Taylor et al., 2007;Donohoe and Battisti, 2011;Atwood et
81 al., 2016;Kashimura et al., 2017;Qu and Hall, 2006). These frameworks differ by whether or
82 not the reflection and transmission properties of the atmospheric layer are assumed to have a
83 directional dependency (Stephens et al., 2015) and by the number of variables required as
84 input (Qu and Hall, 2006). Winton (2005) presented a four-parameter optical model to
85 account for the directional dependency of up- and downwelling shortwave fluxes through a
86 one-layer atmosphere and found good agreement (RMSE < 2% globally) when benchmarked
87 to online radiative transfer calculations. Also considering a directional dependency of the
88 atmospheric optical properties, Taylor et al. (2007) presented a two-parameter model where
89 atmospheric absorption was assumed to occur at a level above atmospheric reflection.
90 Donohoe and Battisti (2011) subsequently relaxed the directional dependency assumption and
91 found the atmospheric attenuation of the surface albedo contribution to planetary albedo to be
92 8% higher than the model of Taylor et al. (2007). Elsewhere, Qu & Hall (2006) developed a
93 framework making use of additional known atmospheric properties such as cloud cover
94 fraction, cloud optical thickness, and the clear-sky planetary albedo which proved highly
95 accurate when model estimates of planetary albedo were evaluated against climate models
96 and satellite-based datasets.

97 Here, our primary research objective is to thoroughly evaluate a variety of shortwave kernels
98 derived both analytically and statistically from satellite-based climatologies of Earth's
99 shortwave radiation budget. To this end, we employ a 16-yr. time series of Earth's monthly
100 mean radiation budget at both TOA (Loeb et al., 2017) and at the surface (Kato et al., 2012)



101 as input to simplified models linking $\Delta\alpha_s$ to changes in the outgoing shortwave radiation flux
102 at TOA. An initial performance screening is implemented where the six observation-driven
103 kernels are first assessed both qualitatively and quantitatively against the mean of four GCM
104 kernels (Shell et al., 2008; Soden et al., 2008; Pendergrass et al., 2018; Block and Mauritsen,
105 2014). Top performers are then subjected to a more rigorous evaluation where they are
106 applied to emulate the GCM kernels using the GCM's own boundary fluxes as input, which
107 eliminates any bias related to differences in the GCM representation of clouds or other
108 atmosphere state variables. Our results elucidate the merits and uncertainties of empirical
109 alternatives to those based on GCMs.

110

111 We start in Section 2 by introducing the satellite-based energy balance product and the
112 variables derived from them utilized in this study. We then provide a brief overview of the
113 GCM-based kernels and of the methods currently being applied within the LULCC science
114 community to estimate instantaneous radiative forcings from surface albedo change. Section
115 3 details the methods applied to derive candidate GCM kernel alternatives from the radiative
116 fluxes at Earth's upper and lower boundaries. We then present results of a comparative
117 analysis in Section 4 and conclude with a brief discussion surrounding the merits and
118 uncertainties of albedo change kernels based on satellite remote sensing.

119

120 **2 Review of existing approaches**

121 The NASA Clouds and the Earth's Radiant Energy System (CERES) Energy Balance and
122 Filled (EBAF) products provide the monthly mean boundary fluxes and atmospheric state
123 information necessary to derive our GCM kernel alternatives (CERES Science Team, 2018a,
124 b). The latest EBAF-TOA Ed4.0 (version 4.0) products have many improvements with
125 respect to the previous version (version 2.8, Loeb et al. 2009), including the use of advanced



126 and more consistent input data, retrieval of cloud properties, and instrument calibration (Loeb
 127 et al. 2018). The temporal extent of the EBAF dataset employed in our analysis spans the
 128 sixteen full calendar years from January 1, 2001 to December 31, 2016 (retrieved April,
 129 2018). An overview of all CERES inputs used in our analysis is presented in Table 1.

130

131 < Table 1 >

132

133 *a. Shortwave ΔF from $\Delta\alpha_s$*

134 Earth's energy balance (at TOA) in an equilibrium state can be written:

$$135 \quad 0 = F = LW_{\uparrow}^{TOA} - (SW_{\downarrow}^{TOA} - SW_{\uparrow}^{TOA}) \quad (1)$$

136 where the equilibrium flux F is a balance between the net solar energy inputs (
 137 $SW_{\downarrow}^{TOA} - SW_{\uparrow}^{TOA}$) and thermal energy output (LW_{\uparrow}^{TOA}). Perturbing this balance results in a
 138 radiative forcing ΔF , while perturbing the shortwave component is referred to as a shortwave
 139 radiative forcing and may be written as:

$$140 \quad \Delta F = \Delta(SW_{\downarrow}^{TOA} - SW_{\uparrow}^{TOA}) = \Delta SW_{\downarrow}^{TOA} \left(1 - \frac{SW_{\uparrow}^{TOA}}{SW_{\downarrow}^{TOA}} \right) - SW_{\downarrow}^{TOA} \left(\Delta \frac{SW_{\uparrow}^{TOA}}{SW_{\downarrow}^{TOA}} \right) \quad (2)$$

141 where the shortwave radiative forcing results either from changes to solar energy inputs (
 142 $\Delta SW_{\downarrow}^{TOA}$) or from internal perturbations within the Earth system ($\Delta \frac{SW_{\uparrow}^{TOA}}{SW_{\downarrow}^{TOA}}$). The latter can

143 be brought about by changes to the reflective properties of Earth's surface and/or atmosphere
 144 which is the focus in this paper.

145 *b. GCM-based radiative kernels*



146 The radiative kernel technique was developed as a way to assess various climate feedbacks
147 from climate change simulations across multiple climate models in a computationally efficient
148 manner (Shell et al., 2008; Soden et al., 2008). A radiative kernel is defined as the differential
149 response of an outgoing radiation flux at TOA to an incremental change in some climate
150 feedback variable -- such as water vapor, air temperature, or surface albedo (Soden et al.,
151 2008). To generate a radiative kernel for a change in surface albedo $\Delta\alpha$ with a GCM, the
152 prescribed surface albedo is perturbed incrementally by 1% and the response by SW_{\uparrow}^{TOA} is
153 recorded, which can be expressed as:

$$154 \quad \Delta SW_{\uparrow}^{TOA} = SW_{\uparrow}^{TOA}(\alpha_s + \Delta\alpha_s) - SW_{\uparrow}^{TOA}(\alpha_s) = \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta\alpha \equiv K_{\alpha} \Delta\alpha_s \quad (3)$$

155 where K_{α} is the radiative kernel (in Wm^{-2}). The albedo change kernel can then be used with
156 Eq. (1) to estimate an instantaneous shortwave radiative forcing (ΔF) at TOA:

$$157 \quad \begin{aligned} F + \Delta F &= LW_{\uparrow}^{TOA} - (SW_{\downarrow}^{TOA} - SW_{\uparrow}^{TOA} + K_{\alpha_s} \Delta\alpha) \\ \Delta F &= -K_{\alpha_s} \Delta\alpha \end{aligned} \quad (4)$$

158 *c. Simple kernel parameterizations of the LULCC science community*

159 Two simplified parameterizations of shortwave radiative transfer have been applied within the
160 LULCC science community for estimating ΔF from $\Delta\alpha_s$ (Muñoz et al., 2010; Lutz et al.,
161 2015; Bozzi et al., 2015; Caiazzo et al., 2014; Cherubini et al., 2012; Carrer et al., 2018). At the
162 core of these parameterizations is the fundamental assumption that radiative transfer is wholly
163 independent of (or unaffected by) $\Delta\alpha_s$. In other words, they neglect the change in the
164 attenuating effect of multiple reflections between the surface and the atmosphere that
165 accompanies a surface albedo change. Although not referred to as “kernels” in the literature,
166 we present them as such to ensure consistency in notation and terminology henceforth. These
167 are subsequently included in the kernel evaluation exercise presented in Section 4.



168

169 The first simplified kernel presented in Muñoz et al. (2010) makes use of a local two-way
170 transmittance factor based on the local clearness index (defined in Table 1):

$$171 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s \equiv K_{\alpha}^{M10} \Delta \alpha_s = SW_{\downarrow}^{TOA} T^2 \Delta \alpha_s \quad (5)$$

172 where SW_{\downarrow}^{TOA} is the local incoming solar flux at TOA, T is the local clearness index, and
173 $\partial SW_{\uparrow}^{TOA} / \partial \alpha$ is the approximated change in the upwelling shortwave flux at TOA due to a
174 change in albedo at the surface.

175

176 The second simplified kernel proposed in Cherubini et al. (2012) makes direct use of the solar
177 flux incident at the surface SW_{\downarrow}^{SFC} combined with a one-way transmission constant k :

$$178 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s \equiv K_{\alpha}^{C12} \Delta \alpha_s = SW_{\downarrow}^{SFC} k \Delta \alpha_s \quad (6)$$

179 where k is based on the global annual mean share of surface reflected shortwave radiation
180 exiting a clear-sky (Lacis and Hansen, 1974; Lenton and Vaughan, 2009) and is hence
181 temporally and spatially invariant. This value – or 0.85 -- is similar to the global mean ratio
182 of forward-to-total shortwave scattering reported in Iqbal (1983). Bright & Kvalevåg (2013)
183 evaluated Eq. (6) at several locations and found large biases for some regions and months,
184 despite good overall performance globally (normalized RMSE = 7%; $n = 120$ months).

185

186 3. Methods

187 Simple analytical models developed by the climate science community treat the atmosphere
188 as a single layer having various optical properties. These models vary by the number and type
189 of optical properties included, whether these have a directional dependency (i.e., isotropic or



190 anisotropic), or whether inputs other than those derived from the boundary fluxes are required
191 (i.e., cloud properties). These models are adapted here to derive kernels analytically for $\Delta\alpha_s$.

192 *a. CERES isotropic kernel*

193 The surface contribution to the outgoing shortwave flux at TOA $SW_{\uparrow,SFC}^{TOA}$ is given (Stephens et
194 al., 2015; Donohoe and Battisti, 2011; Winton, 2005) as:

$$195 \quad SW_{\uparrow,SFC}^{TOA} = SW_{\downarrow}^{TOA} \alpha_s \frac{(1-r-a)^2}{(1-r\alpha_s)} \quad (7)$$

196 where r is a single pass atmospheric reflection coefficient, a is a single pass atmospheric
197 absorption coefficient, SW_{\downarrow}^{TOA} is the extraterrestrial (downwelling) shortwave flux at TOA,
198 and α_s is the surface albedo (defined in Table 1). The expression in the denominator of the
199 righthand term represents a fraction attenuated by multiple reflections between the surface
200 and the atmosphere. This model assumes that the atmospheric optical properties r and a are
201 insensitive to the origin and direction of shortwave fluxes – or in other words – that they are
202 isotropic.

203 The single-pass reflectance coefficient is calculated from the system boundary fluxes (Table
204 1) following Winton (2005) and Kashimura *et al.* (2017):

$$205 \quad r = \frac{SW_{\downarrow}^{TOA} SW_{\uparrow}^{TOA} - SW_{\downarrow}^{SFC} SW_{\uparrow}^{SFC}}{SW_{\downarrow}^{TOA}{}^2 - SW_{\uparrow}^{SFC}{}^2} \quad (8)$$

206 while the single-pass absorption coefficient a is given as:

$$207 \quad a = 1 - r - T(1 - \alpha_s r) \quad (9)$$



208 where T is the clearness index defined in Table 1. Our interest is in quantifying the $SW_{\uparrow,SFC}^{TOA}$
 209 response to an albedo perturbation at the surface – or the partial derivative of $SW_{\uparrow,SFC}^{TOA}$ with
 210 respect to α in Eq. (7):

$$211 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s = K_{\alpha_s}^{ISO} \Delta \alpha_s = \frac{SW_{\downarrow}^{TOA} (1-r-a)^2}{(1-r\alpha_s)^2} \Delta \alpha_s \quad (10)$$

212 where $K_{\alpha_s}^{ISO}$ is referred to henceforth as the CERES isotropic kernel.

213

214 *b. CERES anisotropic kernel*

215 The second kernel makes use of three directionally-dependent (anisotropic) bulk optical
 216 properties r_{\uparrow} , t_{\uparrow} , and t_{\downarrow} , where the first is the atmospheric reflectivity to upwelling
 217 shortwave radiation and the latter two are the atmospheric transmission coefficients for
 218 upwelling and downwelling shortwave radiation, respectively (Winton, 2005). It is not
 219 possible to derive r_{\uparrow} analytically from the CERES all-sky boundary fluxes; however, Winton
 220 (2005) provides an empirical formula relating upwelling reflectivity r_{\uparrow} to the ratio of all-sky
 221 to clear-sky fluxes incident at surface:

$$222 \quad r_{\uparrow} = 0.05 + 0.85 \left(1 - \frac{SW_{\downarrow}^{SFC}}{SW_{\downarrow,CLR}^{SFC}} \right) \quad (11)$$

223 where $SW_{\downarrow,CLR}^{SFC}$ is the clear-sky shortwave flux incident at the surface.

224 Knowing r_{\uparrow} , we can then solve for the two remaining optical parameters needed to derive our
 225 kernel:

$$226 \quad t_{\downarrow} = \frac{SW_{\downarrow}^{SFC} - r_{\uparrow} SW_{\uparrow}^{SFC}}{SW_{\downarrow}^{TOA}} \quad (11)$$



$$227 \quad t_{\uparrow} = T_a - [t_{\downarrow} - t_{\downarrow}(1 - r_{\uparrow}\alpha_s)] \quad (12)$$

228 where T_a is the effective atmospheric transmittance (Table 1) of the earth system.

229 The anisotropic kernel $K_{\alpha_s}^{ANISO}$ can now be derived as:

$$230 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s = K_{\alpha_s}^{ANISO} \Delta \alpha_s = \frac{SW_{\downarrow}^{TOA} t_{\downarrow} t_{\uparrow}}{(1 - r_{\uparrow}\alpha_s)^2} \Delta \alpha_s \quad (13)$$

231

232 *c. CERES auxiliary input kernel*

233 Qu and Hall (2006) developed an alternative analytical kernel to the two described above.

234 The model makes use of auxiliary cloud property information commonly provided in satellite-

235 based products of Earth's radiation budget – including CERES EBAF – such as cloud cover

236 area fraction, cloud visible optical depth, and clear-sky planetary albedo. The model links all-

237 sky and clear-sky effective atmospheric transmissivities of the earth system through a linear

238 coefficient k relating the logarithm of cloud visible optical depth to the effective all-sky

239 atmospheric transmissivity:

$$240 \quad k = \frac{(T_{a,CLR}) - (T_a)}{\ln(\tau + 1)} \quad (14)$$

241 where $T_{a,CLR}$ is the clear-sky effective system transmissivity, T_a is the all-sky effective system

242 transmissivity, and τ is the cloud visible optical depth. This linear coefficient can then be

243 used together with the cloud cover area fraction to derive a shortwave kernel based on the

244 model of Qu and Hall (2006) – or $K_{\alpha_s}^{QH06}$:

$$245 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s = K_{\alpha_s}^{QH06} \Delta \alpha_s = SW_{\downarrow}^{SFC} [(T_a) - kc \ln(\tau + 1)] \Delta \alpha_s \quad (15)$$

246 where c is the cloud cover area fraction.



247 *d. CERES statistical kernel*

248 To determine whether the GCM-based kernels could be approximated with sufficient fidelity
249 using even simpler model formulations based on the CERES boundary data, we applied
250 machine learning to identify potential model forms using the CERES EBAF all-sky boundary
251 fluxes (or system parameters derived from these fluxes) that minimized the sum of squared
252 residuals between monthly means of four GCM-based kernels (described below) and model
253 estimates. The reference dataset consisted of a random global sample of 50,000 (~50%) 2.8°
254 x 2.8° grid cells, from the multi-GCM mean, of which 50% were used for training and 50%
255 for validation. Models were identified using a form of genetic programming known as
256 symbolic regression (Eureqa®; Nutonian Inc.; (Schmidt and Lipson, 2009, 2010)) which
257 searches for both optimal model structure and coefficients. A parsimonious solution was
258 chosen by minimizing the error metric and model complexity using the Pareto front (Smits
259 and Kotanchek, 2005). Based on the mean squared deviation (MSD) and Akaike's
260 information criterion (AIC), the best model form of the statistical kernel – subsequently
261 referred to as $K_{\alpha_s}^{BO18}$ -- is given as:

$$262 \quad \frac{\partial SW_{\uparrow}^{TOA}}{\partial \alpha_s} \Delta \alpha_s = K_{\alpha_s}^{BO18} \Delta \alpha_s = SW_{\downarrow}^{SFC} \sqrt{T} \Delta \alpha_s \quad (16)$$

263 *d. Initial screening of candidate models for a CERES-based kernel*

264 Four GCM kernels are employed as benchmarks to initially screen the six CERES-based
265 kernel model candidates: the Community Atmosphere Model version 3, or CAM3 (Shell et
266 al., 2008), the Community Atmosphere Model version 5, or CAM5 (Pendergrass et al., 2018),
267 the European Center and Hamburg model version 6, or ECHAM6 (Block and Mauritsen,
268 2014), and the Geophysical Fluid Dynamics Laboratory model version AM2p12b, or GFDL
269 (Soden et al., 2008). The four GCM kernels vary in vertical and horizontal resolution,



270 parameterization of shortwave radiative transfer, and year of atmospheric state (input
 271 variables).

272

273 We compute a skill metric analogous to the “relative error” metric used to evaluate GCMs by
 274 Anav et al. (2013) that takes into account error in the spatial pattern between a model and an
 275 observation. Because we have no true observational reference, our evaluation instead focuses
 276 on the disagreement or deviation between CERES and GCM kernels at the monthly time step.
 277 Given interannual climate variability in the earth system, the challenge of comparing the
 278 multi-year CERES kernel to a single-year GCM kernel can be partially overcome by
 279 averaging the four GCM kernels.

280

281 Using the multi-GCM mean as the reference, we first compute the absolute deviation $AD_{m,p}^X$

282 as:

$$283 \quad AD_{m,p}^X = \left| CERES_{m,p}^X - \overline{GCM}_{m,p} \right| \quad (17)$$

284 where $CERES_{m,p}^X$ is the kernel for CERES model X in month m and pixel p and $\overline{GCM}_{m,p}$ is
 285 the multi-GCM mean of the same pixel and month. $AD_{m,p}^X$ is then normalized to the
 286 maximum absolute deviation of all six CERES kernels for the same pixel and month to obtain
 287 a normalized absolute deviation, $NAD_{m,p}^X$, which is analogous to the “relative error” metric of
 288 Anav et al. (2013) with values ranging between 0 and 1:

$$289 \quad NAD_{m,p}^X = 1 - \frac{AD_{m,p}^X}{\max(AD_{m,p})} \quad (18)$$

290 where $\max(AD_{m,p})$ is the maximum absolute deviation of all six CERES kernels at pixel p
 291 and month m .



292

293 CERES kernel ranking is based on the mean relative absolute deviation in both space and time

294 – or NAD^x :

$$295 \quad NAD^x = \frac{1}{M} \sum_{m=1}^M \frac{1}{P} \sum_{p=1}^P NAD_{m,p}^x \quad (19)$$

296 where M is the total number of months (i.e., 12) and P is the total number of grid cells.

297

298 *e. GCM kernel emulation*

299 In order to eliminate any bias related to differences in the atmospheric state embedded in the

300 GCM and CERES-derived kernels, we re-compute our simple kernels using the same

301 shortwave boundary fluxes used to compute the two most recent albedo change kernels based

302 on ECHAM6 (Block and Mauritsen, 2014) and CAM5 (Pendergrass et al., 2018). This

303 enables a more critical evaluation of the functional form of the simple models in relation to

304 the more sophisticated radiative transfer schemes employed by ECHAM6 (Stevens et al.,

305 2013) and CAM5 (Hurrell et al., 2013).

306

307 **4. Results**

308 *a. Initial kernel performance screening*

309 Seasonally, differences in latitude band means between the CERES and multi-GCM mean

310 kernels are shown in Figure 1.

311

312 < Figure 1 >

313



314 Qualitatively, starting with December-January-February (*DJF*), K_{α}^{BO18} gives the best
 315 agreement with $\overline{K_{\alpha}^{GCM}}$ with the exception of the zone around 55 – 65°S (-55 – -65°), where
 316 K_{α}^{QH06} gives slightly better agreement (Fig. 1A). In March-April-May (*MAM*), K_{α}^{BO18} appears
 317 to give the best overall agreement with the exception of the high Arctic, where K_{α}^{ANISO} and
 318 K_{α}^{C12} give better agreement, and with the exception of the zone around 60 – 65°S (-60 – -65°)
 319 where K_{α}^{QH06} , K_{α}^{ANISO} , and K_{α}^{C12} agree best with $\overline{K_{\alpha}^{GCM}}$ (Fig. 1B). The largest spread in
 320 disagreement across all six CERES kernels is found in June-July-August (*JJA*; Fig. 1C) at
 321 northern high latitudes. K_{α}^{BO18} appears to agree best both here and elsewhere with the
 322 exception of the zone between ~20 – 35°N, where K_{α}^{QH06} gives slightly better agreement.
 323 In September-October-November (*SON*), K_{α}^{BO18} agrees best with $\overline{K_{\alpha}^{GCM}}$ at all latitudes except
 324 the zone between 10 – 25°N and 55 – 65°S where K_{α}^{QH06} agrees slightly better.

325

326 Quantitatively, the proportion of the total variance explained by linear regressions of monthly
 327 $\overline{K_{\alpha}^{GCM}}$ on monthly K_{α}^{CERES} (i.e., R^2) is highest and equal for the CERES kernels based on the
 328 ANISO, QH06, and BO18 models (Fig. 2 B, C, & D). Of these three, K_{α}^{QH06} has a y-intercept
 329 (B_0) closest to 0 and a slope (m) of 1, although the root mean squared deviation (*RMSD*) – an
 330 accuracy measure – is slightly better (lower) for K_{α}^{BO18} . The two CERES kernels with the
 331 lowest R^2 , highest slopes (negative deviations), highest *RMSDs*, and y-intercepts with the
 332 largest absolute difference from zero are those based on the ISO and M10 models (Fig. 2
 333 A&E).

334

335 < Figure 2 >

336



337 Although the y-intercept deviation from 0 for K_{α}^{C12} is relatively low, its *RMSD* is ~50%
338 higher than that of K_{α}^{QH06} , K_{α}^{BO18} , and K_{α}^{ANISO} and leads to notable positive deviation from the
339 multi-GCM mean ($\overline{K_{\alpha}^{GCM}}$) judging by its slope of 0.92.

340

341 *c. Normalized absolute deviation*

342 Globally, NAD^X for the QH06, ANISO, and BO18 kernels are far superior to the ISO, M10,
343 and C12 kernels (Table 2).

344

345 < Table 2 >

346

347 After filtering to remove grid cells for oceans and other water bodies, NAD^X scores for these
348 three kernels decreased; the decrease was smallest for K_{α}^{BO18} (-0.03) and largest for K_{α}^{QH06} (-
349 0.06). Despite constraining the analysis to land surfaces only, the rank order remained
350 unchanged (Table 2).

351

352 *d. GCM kernel emulation and additional performance screening*

353 Because the simple kernel based on the QH06 model (K_{α}^{QH06}) required auxiliary inputs for
354 cloud cover area fraction and cloud optical depth – two atmospheric state variables not
355 provided with the ECHAM6 and CAM5 kernel datasets – it was not possible to emulate these
356 two GCM kernels using the QH06 model. Additional performance evaluation through GCM
357 kernel emulation is therefore restricted to the ANISO and BO18 models.

358 < Figure 3 >



359 Globally, the kernel based on the ANISO model displays larger annual mean biases relative to
360 BO18 when compared to both ECHAM6 and CAM5 kernels (Figure 3). Notable positive
361 biases over land with respect to both ECHAM6 and CAM5 kernels are evident in the northern
362 Andes region of South America, the Tibetan plateau, and the tropical island region comprising
363 Indonesia, Malaysia, and Papua New Guinea (Fig. 3 A & C). Notable negative biases over
364 land with respect to both ECHAM6 and CAM5 kernels are evident over Greenland,
365 Antarctica, northeastern Africa, and the Arabian Peninsula (Fig. 3 A & C).

366 < Figure 4 >

367 Globally, annual biases for BO18 are generally found to be lower than for ANISO and are
368 mostly non-existent in extra-tropical ocean regions (Fig. 3 B & D). Patterns in biases over
369 land are mostly negative with the exception of Saharan Africa where the annual mean bias
370 with respect to both GCMs is positive. For BO18, systematic positive biases – or biases
371 evident with respect to both GCM kernels – appear over eastern tropical and subtropical
372 marine coastal upwelling zones where marine stratocumulus cloud dynamics are difficult for
373 GCMs to resolve (Bretherton et al., 2004; Richter, 2015).

374 < Table 3 >

375 Performance metrics based on regressing monthly kernels from the two GCMs on kernels
376 emulated with both ANISO and BO18 models indicate a greater overall accuracy (or
377 agreement) for BO18 (Figure 4). RMSDs for monthly kernels emulated with BO18 are 9.0
378 and 8.2 W m⁻² with respect to CAM5 and ECHAM6, respectively – which is ~50-60% of the
379 RMSDs emulated with the ANISO model. Focusing henceforth only on the kernel emulated
380 with BO18 model, negative biases are evident in all months (Table 3), with the largest biases
381 (in magnitude) appearing in May (-4.4 W m⁻²) and November (-2.5 W m⁻²) for CAM5 and
382 ECHAM6, respectively. In absolute terms, largest biases of 8.6 W m⁻² and 6.8 W m⁻² appear



383 in June for CAM5 and ECHAM6, respectively. Annually, the mean absolute bias for CAM5
384 and ECHAM6 is 6.8 and 6.1 W m⁻², respectively – a magnitude which seems remarkably low
385 if one compares this to the annual mean disagreement (standard deviation) of 33 W m⁻² across
386 all four GCM kernels (not shown).

387

388 **5. Discussion and conclusions**

389 Motivated by an increasing abundance of climate impact research focusing on land processes
390 in recent years, we comprehensively evaluated six simplified models linking shortwave
391 radiative flux perturbations at TOA with surface albedo changes at the surface. Relative to
392 albedo change kernels based on sophisticated radiative transfer schemes embedded in GCMs,
393 the simplified models evaluated here can be updated frequently at relatively low cost using
394 boundary fluxes obtained from remote sensing-based products of Earth's shortwave energy
395 budget. This allows greater flexibility to meet the needs of research that focuses on longer-
396 term albedo trends or regions currently undergoing rapid change in atmospheric composition.
397 Although some modeling groups have provided recent updates to radiative kernels using the
398 latest GCM versions, the atmospheric state of the boundary conditions used to derive them
399 may be considered outdated or not in sync with that required for some applications.

400

401 Based on both qualitative and quantitative benchmarking against the mean of four GCM
402 kernels, the simple model derived from machine learning, BO18, together with the two
403 analytically derived models, QH06 and ANISO, proved far superior to the M10, C12, and the
404 ISO kernel models. When subjected to additional performance evaluation, however, we
405 found that the BO18 model was able to more robustly emulate the ECHAM6 and CAM5
406 kernels with exceptionally high accuracy, suggesting that this model can serve as a suitable
407 candidate for an albedo change kernel based on CERES boundary fluxes. The RMSD of this



408 kernel – henceforth referred to as the CERES Albedo Change Kernel (CACK v1.0) – was
409 found to be 6.8 and 6.1 W m⁻² when benchmarked to the CAM5 and ECHAM6 kernel,
410 respectively – a magnitude which is only ~20% of the standard deviation found across four
411 GCM kernels (annual mean). CACK’s remarkable simplicity lends support to the idea of
412 using machine learning to explore and detect emergent properties of shortwave radiative
413 transfer in future research.

414

415 Despite the stronger empirical foundation of CACK over a GCM-based kernel, it is important
416 to recognize its limitations. Firstly, the monthly CERES EBAF-Surface product used to
417 define lower atmospheric boundary conditions is not strictly an observation. The space-borne
418 observation platform is not able to directly observe Earth’s surface fluxes under overcast
419 conditions and hence requires model augmentation. However, the energy-balancing step
420 ensures that fluxes are adjusted to match the observed rate of heat accumulation in the climate
421 system (i.e., the oceans) (Hansen et al., 2005). These processes, as well as extensive ground
422 validation and testing, are documented elsewhere (Kato et al., 2013;Loeb et al., 2009).
423 Further, while CACK has a finer spatial resolution than most GCM kernels, it still represents
424 a spatially averaged response rather than a truly local response; in other words, the state
425 variables used to define the response are tied to the coarse spatial (i.e., 1° x 1°) resolution of
426 the CERES EBAF product grids. Lastly, it is important to emphasize that CACK is based on
427 the climate conditions of the present day (2001-2016); hence, caution should be exercised
428 when applying it to estimate ΔF associated with albedo changes occurring outside this range.

429

430 To conclude, we evaluated six simplified albedo change kernels based on CERES shortwave
431 boundary fluxes as candidate alternatives to GCM-based albedo change kernels. Albedo
432 change kernels are useful tools for estimating instantaneous shortwave radiative forcings



433 connected to anthropogenic land use activities. Our results showed that the BO18 model
434 developed and presented in this study is the best candidate for a CERES albedo change kernel
435 -- or CACK. CACK provides a higher spatial resolution, higher transparency alternative to
436 existing kernels based on GCMs. CACK could be easily applied as part of Monitoring,
437 Reporting, and Verification (MRV) frameworks for biogeophysical impacts on land,
438 analogous to those which currently exist for land sector greenhouse gas emissions. Given the
439 extensive time span of the CERES EBAF products, CACK based on a multi-year climatology
440 of Earth's shortwave radiation budget would better-account for internal climate variability in
441 the earth system. However, CACK's flexibility regarding input year should make it broadly
442 appealing across a range of disciplines. One example is the land-based solar radiation
443 management (SRM) research community who frequently calculate ΔF from $\Delta\alpha$ to evaluate
444 climate mitigation strategies (Ridgwell et al., 2009; Seneviratne et al., 2018; Akbari et al.,
445 2009).

446

447 **Code Availability**

448 An open source Octave (Eaton et al., 2018) script file for generating monthly CACK from
449 CERES EBAF data and example input files are included as a Supplement. The script also
450 demonstrates calculating a TOA RF from CACK and monthly surface albedo perturbation at a
451 user-specified location.

452

453 **Data Availability**

454 CERES EBAF data are available for download at:
455 <https://ceres.larc.nasa.gov/products.php?product=EBAF-TOA> . The CAM3 kernel is
456 available at: <http://people.oregonstate.edu/~shellk/kernel.html> . The CAM5 kernel is
457 available at: <https://www.earthsystemgrid.org/ac/guest/secure/sso.html> . The ECHAM5



458 kernel is available at: [https://swiftbrowser.dkrz.de/public/dkrz_0c07783a-0bdc-4d5e-9f3b-](https://swiftbrowser.dkrz.de/public/dkrz_0c07783a-0bdc-4d5e-9f3b-c1b86fac060d/Radiative_kernels/)
459 [c1b86fac060d/Radiative_kernels/](https://swiftbrowser.dkrz.de/public/dkrz_0c07783a-0bdc-4d5e-9f3b-c1b86fac060d/Radiative_kernels/) .

460

461 **Author contributions**

462 T.L.O. conceptualized the study. R.M.B. and T.L.O. developed the methodology, curated the
463 data, designed the computer programs, and carried out the formal analysis. R.M.B. wrote the
464 original draft, and R.M.B. and T.O. reviewed and edited and final manuscript. T.L.O.
465 produced the figures.

466

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686

687



688 **Table 1.** Definition of CERES input variables and other system optical properties derived
689 from CERES inputs. All variables are 2001-2016 monthly means at $1^\circ \times 1^\circ$ spatial resolution.

CERES EBAF v.4 Shortwave Boundary Fluxes

SW_{\downarrow}^{TOA}	Downwelling solar flux at top-of-atmosphere	Wm^{-2}
SW_{\downarrow}^{SFC}	Downwelling solar flux at surface	Wm^{-2}
$SW_{\downarrow,CLR}^{SFC}$	Clear-sky downwelling solar flux at surface	Wm^{-2}
SW_{\uparrow}^{TOA}	Upwelling solar flux at top-of-atmosphere	Wm^{-2}
SW_{\uparrow}^{SFC}	Upwelling solar flux at surface	Wm^{-2}

System Optical Properties

$T = SW_{\downarrow}^{SFC} / SW_{\downarrow}^{TOA}$	Clearness index	unitless
$\alpha_p = SW_{\uparrow}^{TOA} / SW_{\downarrow}^{TOA}$	Planetary albedo	unitless
$\alpha_s = SW_{\uparrow}^{SFC} / SW_{\downarrow}^{SFC}$	Surface albedo	unitless
$A_p = 1 - \alpha_p$	Effective planetary absorption	unitless
$A_s = [SW_{\downarrow}^{SFC} - SW_{\uparrow}^{SFC}] / SW_{\downarrow}^{TOA}$	Effective surface absorption	unitless
$A_a = A_p - A_s$	Effective atmospheric absorption	unitless
$T_a = 1 - A_a$	Effective atmospheric transmission	unitless
$T_{a,CLR} = 1 - A_{a,CLR}$	Clear-sky effective atmospheric transmission	unitless
τ	Cloud visible optical depth	unitless
c	Cloud area fraction	fraction

690

691



692 **Table 2.** Normalized absolute deviation and CERES kernel ranking.

	Global		Land only		Mean Rank
	<i>NAD</i>	Rank	<i>NAD</i>	Rank	
ISO	0.05	6	0.05	6	6
ANISO	0.64	3	0.59	3	3
C12	0.45	4	0.47	4	4
M10	0.26	5	0.34	5	5
QH06	0.66	2	0.60	2	2
BO18	0.67	1	0.64	1	1

693

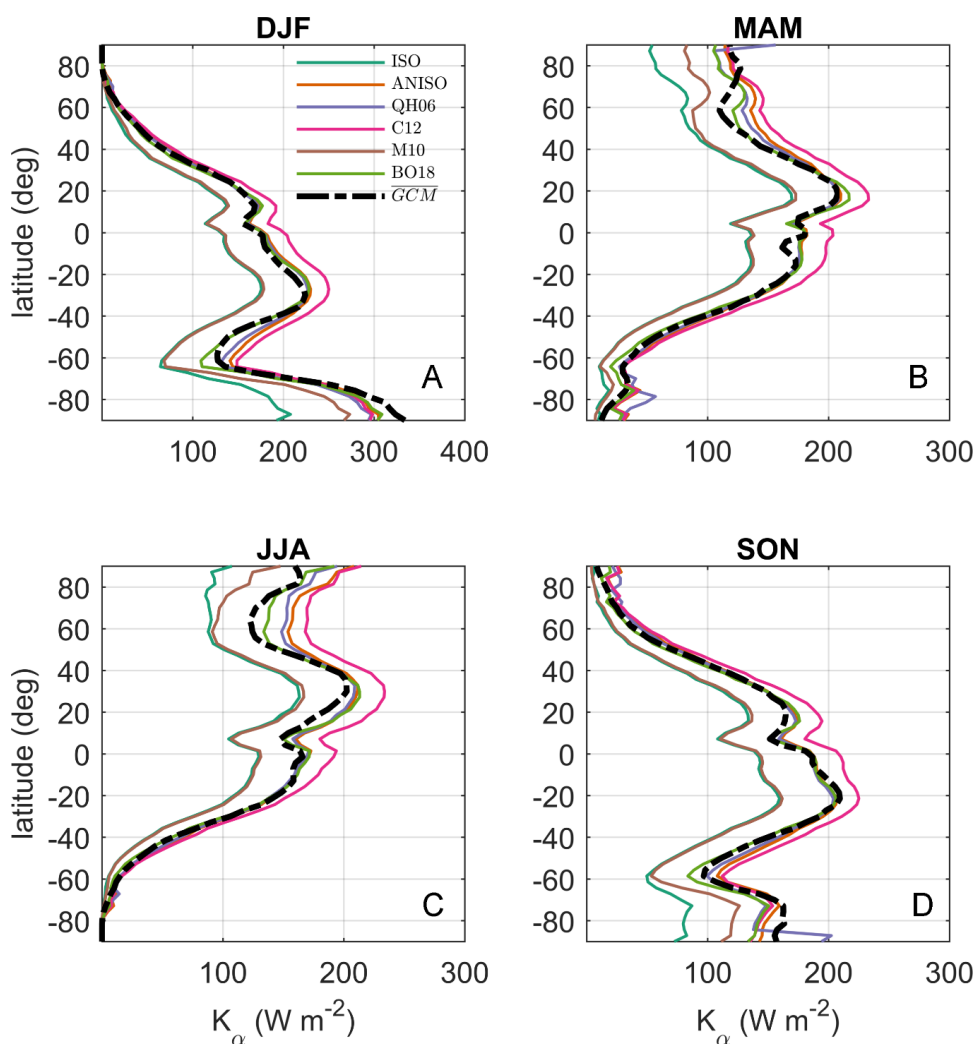
694



Table 3. Global monthly mean bias (*MB*) and mean absolute bias (*MAB*) for K_{α}^{BO18} emulated with T and SW_{\downarrow}^{SFC} from ECHAM6 and CAM5. For reference, the global mean value of K_{α}^{BO18} is 133 W m^{-2} .

	MB (W m^{-2})												
	Jan.	Feb.	Mar.	Apr.	May	Jun.	Jul.	Aug.	Sep.	Oct.	Nov.	Dec.	Ann.
$K_{\alpha}^{BO18} - K_{\alpha}^{CAM5}$	-2.9	-3.4	-3.3	-3.9	-4.4	-3.8	-3.8	-3.7	-3.4	-3.8	-3.7	-3.3	-3.6
$K_{\alpha}^{BO18} - K_{\alpha}^{ECHAM6}$	-1.9	-2.2	-1.8	-1.9	-2.2	-1.5	-1.1	-1.6	-1.7	-2.5	-2.5	-1.8	-1.9
	MAB (W m^{-2})												
	Jan.	Feb.	Mar.	Apr.	May	Jun.	Jul.	Aug.	Sep.	Oct.	Nov.	Dec.	Ann.
$ K_{\alpha}^{BO18} - K_{\alpha}^{CAM5} $	6.9	5.7	5.2	6.8	7.7	8.6	7.9	6.7	5.6	6.1	6.9	6.9	6.8
$ K_{\alpha}^{BO18} - K_{\alpha}^{ECHAM6} $	6.3	5.7	5.0	5.9	6.7	6.8	6.4	5.8	5.3	5.6	6.4	6.7	6.1

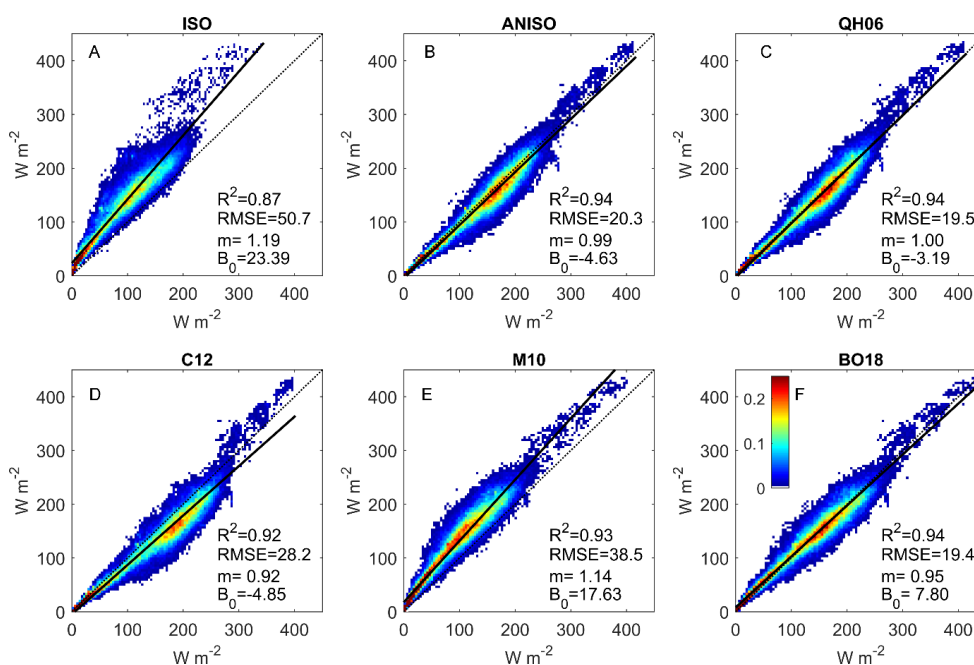
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697 **Figure 1.** Latitudinal (1°) and seasonal means of K_α^{GCM} and K_α^{CERES} for: A) December-
 698 January-February (DJF); B) March-April-May (MAM); C) June-July-August (JJA); D)
 699 September-October-November (SON).

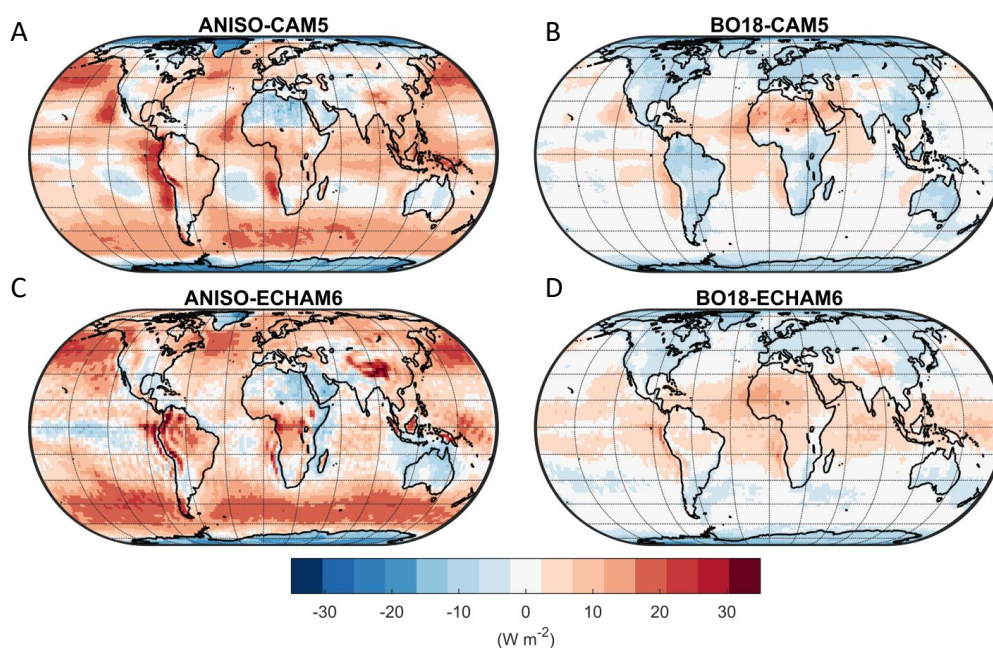
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702 **Figure 2.** A)-F): Scatter-density regressions of global monthly mean K_{α}^{GCM} (y-axis) and
 703 K_{α}^{CERES} (x-axis), with the CERES kernel identifier shown at the top of each sub-panel. “ m ” =
 704 slope; “ B_0 ” = y-intercept. The color scale indicates the percentage of regression points that
 705 fall within a 100×100 sample grid centered on the plotted point.

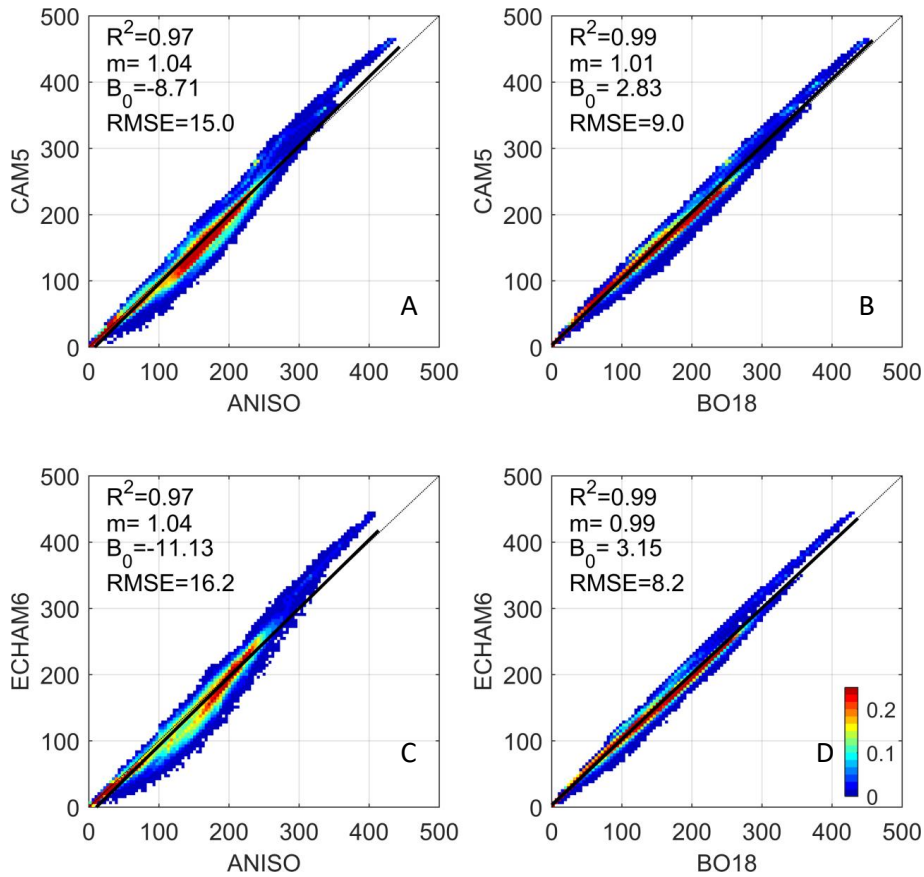
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708 **Figure 3.** A) Mean annual bias of the CAM5 albedo change kernel emulated with the ANISO
709 analytical model; B) Mean annual bias of the CAM5 albedo change kernel emulated with the
710 BO18 parameterization; C) Mean annual bias of the ECHAM6 albedo change kernel emulated
711 with the ANISO analytical model; D) Mean annual bias of the ECHAM6 albedo change
712 kernel emulated with the BO18 parameterization

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715 **Figure 4.** A)-D): Scatter-density regressions of K_α^{GCM} (y-axis) and K_α^{GCM} emulated with the
 716 ANISO model and BO18 parameterization (x-axis); “ m ” = slope; “ B_0 ” = y-intercept. The
 717 color scale indicates the percentage of regression points that fall within a 100×100 sample
 718 grid centered on the plotted point.

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