# libcloudph++ 2.0: aqueous phase chemistry extension of the particle-based cloud microphysics scheme

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**Abstract.** This paper introduces a new scheme available in the library of algorithms for representing cloud microphysics in numerical models named *libcloudph*++. The scheme extends the particle-based microphysics scheme with a Monte-Carlo coalescence available in *libcloudph*++ to the aqueous phase chemical processes occurring within cloud droplets. The representation of chemical processes focuses on the aqueous phase oxidation of the dissolved SO<sub>2</sub> by O<sub>3</sub> and H<sub>2</sub>O<sub>2</sub>. The particle-based microphysics and chemistry scheme allows tracking the changes in the cloud condensation nuclei (CCN) distribution caused by both collisions between cloud droplets and aqueous phase oxidation.

The scheme is implemented in C++ and equipped with bindings to Python. The scheme can be used on either CPU or GPU, and is distributed under the GPLv3 license. Here, the particle-based microphysics and chemistry scheme is tested in a simple 0-dimensional adiabatic parcel model and then used in a 2-dimensional prescribed flow framework. The results are discussed with the focus on changes to the CCN sizes and compared with other model simulations discussed in the literature.

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#### 1 Introduction

libcloudph++ is an open-source library of schemes for representing cloud microphysics in numerical models. It was first introduced in Arabas et al. (2015) where the authors present the different microphysics schemes available in the library, show its programming interface, and discuss its performance. The flagship component of libcloudph++ is the particle-based (i.e. particle tracking or "Lagrangian-in-droplet-radius-and-space") microphysics scheme. The scheme resolves the evolution of aerosol, cloud droplet, and rain drop<sup>1</sup> size spectrum. It allows representing from the first principles cloud microphysical processes and is especially well suited to track changes in the CCN size distribution that are caused by clouds (i.e. cloud-aerosol processing). The scheme can be used in models of any dimensionality or dynamical core, and can be run on both CPU and GPU. The main software design principle employed while developing libcloudph++ core code is the separation of concerns. The code is open-source and its programming interface is documented in Arabas et al. (2015). All those features facilitate further development

<sup>&</sup>lt;sup>1</sup> For convenience cloud droplets and rain drops will be often labeled together as water drops

Sulfate aerosols cool Earth's climate by scattering sunlight and thus increasing Earth's shortwave albedo (direct radiative forcing) and by changing radiative properties of clouds (cloud albedo effect). According to the chapter 8 of IPCC<sup>2</sup> Assessment Report (Myhre et al., 2013), the range of effective radiative forcings for all aerosol-radiation interactions is -0.95 to 0.05 W/m<sup>2</sup> and for aerosol-cloud interactions is -1.2 to 0.0 W/m<sup>2</sup>. The level of scientific understanding in that report for the cloud albedo effect is still marked as "low". From the air quality perspective, in extreme cases sulfur chemistry may lead to the creation of acid rain or acid fog (Dianwu et al., 1988; Wang et al., 2016). Based on analysis of 20 modeling studies, the review by Faloona (2009) marks wet deposition of aerosol sulfate, dry deposition of  $SO_2$  and heterogeneous (aqueous phase) oxidation of  $SO_2$  in aerosol particles and clouds as the most challenging to quantify in models. For an overview of the representation of sulfur oxidation in regional and global models see Ervens (2015). The aqueous phase oxidation is reported as a dominant mechanism of production of sulfate: A numerical study by Barth et al. (2000) reports that for the in-cloud conditions, aqueous phase reactions accounts for 81% of sulfate production rate. According to their study total of  $\sim 50\% - 60\%$  of sulfate burden in the troposphere is produced by aqueous phase chemistry. The gas phase  $SO_2$  is oxidized in a matter of days by gas phase reactions or within minutes or a few hours within clouds by aqueous phase reactions, see the review by Faloona (2009).

From the cloud microphysics stand point, aqueous phase oxidation of sulfur is interesting because it affects the CCN within water drops. Noteworthy, sulfate is a common component of aerosol particles (10%-67% of the submicron particle mass is made of sulfate, 32% on average; see Zhang et al., 2007). The aqueous phase oxidation of sulfur is irreversible, meaning that the produced sulfate remains within water drops and increases the dissolved CCN mass. Collisions and the subsequent coalescence of water drops as well as collisions between aerosol particles and water drops are another in-cloud irreversible process that affects aerosol particles. As water drops collide and coalesce, the created new water drop carries within the combined CCN mass of all of its colliding predecessors. Efficient collisions between cloud droplets may quickly lead to the onset of precipitation which can in turn effectively cleanse the atmosphere from aerosol particles and water soluble trace gases. In non-precipitating clouds, aerosol particles that served as CCN are altered by cloud microphysical and chemical processes and then returned to the atmosphere after water drops evaporate (the process is referred to as the CCN deactivation, aerosol regeneration, aerosol recycling or aerosol resuspension). The cloud-processed aerosol particles can be observed in measurements (Hoppel et al. 1986, 1994, Werner et al. 2014, Hudson et al. 2015). The case without precipitation reaching the surface is especially interesting as it allows for aerosol-cloud interactions to loop for several cloud life-cycles without removing the altered aerosol particles. The cloud-processed aerosol particles may again serve as CCN and influence microphysical properties of the next generation of clouds. The study of Pruppacher and Jaenicke (1995) estimates that on global average an atmospheric aerosol particle has been cycled 3 times by cloud systems. The Particle-Based Microphysics and Chemistry (PBMC) scheme introduced here offers a chance to represent the effects of such cloud-processing on CCN sizes stemming from both collisions between water drops and aqueous phase oxidation reaction within water drops. The PBMC can be used in multidimensional simulations with a fully coupled dynamics model, which has not been possible before. To the authors knowledge, the presented

<sup>&</sup>lt;sup>2</sup> Intergovernmental Panel on Climate Change, see http://www.ipcc.ch/

scheme is the first to represent the impact of both collisions and aqueous phase chemistry on the aerosol size spectrum in the particle-based microphysics framework.

This paper documents the extension of the particle-based microphysics scheme with a numerical scheme that represents aqueous phase chemical reactions inside cloud droplets and the uptake of the trace gases into cloud droplets. The representation of chemical reactions includes only the aqueous phase processes (i.e., no gas phase chemical reactions) and revolves around oxidation of sulfur dissolved in water drops to sulfate. Two reaction paths are considered – the oxidation by ozone and by hydrogen peroxide. In total, six trace gases are included in the chemistry description: sulfur dioxide ( $SO_2$ ), ozone ( $O_3$ ), hydrogen peroxide ( $H_2O_2$ ), carbon dioxide ( $CO_2$ ), nitric acid ( $H_2O_3$ ), and ammonia ( $NH_3$ ). Their dissolution, and if applicable, dissociation is resolved. The structure of the presented work is as follows: Section 2 presents briefly the particle-based scheme available in *libcloudph+++*. Section 3 discusses the design of the new aqueous chemistry scheme and Sec. A in the Appendix describes the new programming interface. Section 4 compares the results from the new scheme with the results from moving-bin schemes. Section 5 discusses the results from simulations where the PBMC scheme is incorporated into a simple model of a stratocumulus cloud. The effects of both collisions between water drops and aqueous phase oxidation of sulfur on the aerosol particle size distribution are presented.

#### 2 Particle-based microphysics scheme

The particle-based scheme used in this work is described in detail in Arabas et al. (2015) and this section only briefly summarizes its major concepts. In the particle-based approach to modeling cloud microphysics the computational domain is filled with "numerical point particles" representing a specified number (called multiplicity) of real particles (aerosol particles, cloud droplets or rain drops) of the same properties. Following the nomenclature introduced by Shima et al. (2009), the "numerical particles" are labeled here as super-droplets (SD). Each SD has a set of attributes describing the properties of the aerosol particles or water drops it represents. As discussed in Arabas et al. (2015), for microphysical purposes, the required attributes are: the multiplicity (N), the position of SD in the computational domain, the wet radius ( $r_w$ ), the dry radius<sup>3</sup> ( $r_d$ ) and the hygroscopicity parameter<sup>4</sup> ( $\kappa$ ). The aqueous phase chemistry scheme extends the list of required attributes by masses of chemical compounds dissolved in droplets. Eight new attributes are needed, see Sec. 3 for details.

The key attribute of the particle-based microphysics scheme is the SD multiplicity. The multiplicity defines the number of aerosol particles or water drops represented by a given SD. All particles represented by one SD are assumed to be identical. Using multiplicites allows to reduce the complexity of the problem and enables efficient numerical computations.

The particle-based scheme used here requires no division into artificial categories of aerosol particles, cloud and rain water, as it is often done in the bulk schemes, for example Kessler (1995), Seifert and Beheng (2001), Morrison and Grabowski (2007). All the modeled microphysical processes are represented by calculating the changes to the SD attributes. The equation

<sup>&</sup>lt;sup>3</sup> It is a volume equivalent radius for solute in the water drop.

<sup>&</sup>lt;sup>4</sup> Following Ghan et al. (2001) and Petters and Kreidenweis (2007) it is a single parameter representing the hygroscopicity of the solvent. In this work we use the notation from Petters and Kreidenweis (2007).

of condensational growth is solved for each SDs wet radius (see Sec. 5.1.3 in Arabas et al. (2015) for details). The process of condensational growth from deliquescent aerosol particles to cloud droplets is thus resolved and no additional parametrisation of cloud droplet activation is required as it is again often done in the bulk microphysics schemes, see for example Morrison and Grabowski (2007).

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Following Shima et al. (2009) the collisions between SDs are represented using a Monte-Carlo scheme (see Sec. 5.1.4 in Arabas et al. (2015) for details). Information about the SD attributes is retained within the model throughout the whole simulation. This means that the size distribution of both water drops and aerosol particles in each computational grid-cell can be easily obtained by taking into account the SD attributes of  $r_w$ ,  $r_d$  and N. As a result, the particle-based scheme is capable of resolving the changes to both aerosol and water drop size distributions. The same functionality is offered by the 2-dimensional bin schemes, for example Ovchinnikov and Easter (2010) or Lebo and Seinfeld (2011). However, the particle-based approach greatly reduces the numerical diffusion errors. As discussed in Unterstrasser et al. (2017), it does introduce the statistical errors, i.e. fluctuations between different realizations of the same collision/coalescence scenario. These errors are easier to minimize than diffusion numerical errors, for example by increasing the number of SDs in the computational domain or by averaging over an ensemble of simulation runs. Dziekan and Pawlowska (2017) showed that for high SD concentrations the super-droplet method accurately represents collisions between the drops (with regard to the expected value and the standard deviation of the autoconversion time). An interesting comparison between the bin and the particle-based schemes is provided by Li et al. (2017). On a side note, Grabowski and Abade (2017) show that an additional scheme modeling the broadening of droplet spectra due to supersaturation fluctuations might be necessary for the particle-based schemes.

The collision efficiency used in this study is based on Hall (1980) and Pinsky et al. (2008). It is well suited for representing the collisions between water drops. An additional collision efficiency look-up table based for example on Ladino et al. (2011) or Ardon-Dryer et al. (2015) should be used to study the collection of submicron aerosol particles by droplets. Similarly, additional collision efficiency corrections based for example on Chen et al. (2018) should be applied to study the effects of turbulence on the aerosol size distribution.

The Particle-based methods are becoming a well known tool for studying cloud microphysics in both warm-clouds (Shima et al. (2009), Andrejczuk et al. (2010), Riechelmann et al. (2012), Arabas and Shima (2013), Andrejczuk et al. (2014), Lee et al. (2014), Naumann and Seifert (2015), Sato et al. (2017), Hoffmann (2017), Grabowski et al. (2018), Sardina et al. (2018) ); and ice-phase clouds (Sölch and Kärcher (2010), Unterstrasser and Sölch (2014)). None of the above however, included description of the aqueous phase chemical reactions happening within cloud droplets.

# 3 Aqueous phase chemistry scheme

In order to represent the chemical composition of water drops the aqueous-phase chemistry scheme extends the list of SD attributes. The additional attributes are defined as the total mass of each of the chemical compounds in a given SD (including both the dissolved and, if applicable, dissociated fraction). An additional variable – the mass of the H<sup>+</sup> ions – is also added,

in order to keep track of SD's acidity. This results in eight new SD attributes needed for simulations with aqueous phase chemistry:

- the total mass of dissolved O<sub>3</sub>,
- the total mass of dissolved H<sub>2</sub>O<sub>2</sub>,
- the total mass of dissolved  $SO_2$  (including:  $SO_2 * H_2O$ ,  $HSO_3^-$  and  $SO_3^{2-}$ ),
  - the total mass of dissolved  $CO_2$  (including:  $CO_2 * H_2O$ ,  $HCO_3$  and  $CO_3^{2-}$ ),
  - the total mass of dissolved NH<sub>3</sub> (including: NH<sub>3</sub> \* H<sub>2</sub>O and NH<sub>4</sub>),
  - the total mass of dissolved HNO<sub>3</sub> (including: HNO<sub>3</sub>(aq) and NO<sub>2</sub>),
  - the total mass of created  $H_2SO_4$  (including:  $HSO_4^-$  and  $SO_4^{2-}$ )
- the total mass of H<sup>+</sup> ions.

The scheme needs to be coupled to a driver model (i.e. a dynamical core) that provides information about the environment in which SD are immersed (i.e. temperature, humidity, trace gas mixing ratio, and air wind field). The representation of aqueous phase chemistry more than doubles the number of required SD attributes and significantly increases the computational time. On the other hand, thanks to the added attributes, the mass of any ion for any SD can be easily diagnosed using just a dissociation constant. This, in turn, allows for a very straightforward representation of the aqueous chemical processes and does not call for any additional parametrisation.

All aqueous-phase chemistry included in the scheme is formulated under the assumption that solution droplets are diluted. Therefore, in the PBMC scheme, chemical processes are only performed for the SDs with ionic strength smaller than 0.02 moles/liter (the same criterion is used for example in Ovchinnikov and Easter, 2010). In practice, this condition results in excluding from aqueous chemistry calculations the SDs with small wet radii (i.e. the SDs representing haze particles and very small cloud droplets). Exclusion of the SDs with small wet radii also prevents numerical issues during condensation procedure when changes in dry radius caused by oxidation could prevent convergence of the condensation scheme during the initial rapid growth of cloud droplets during activation.

Combining the particle-based microphysics scheme with aqueous phase chemistry is straightforward. Condensation/evaporation does not affect the chemical attributes of SDs. During collisions the mass of chemical compounds is summed when recalculating SD attributes (it is an extensive parameter). In principle the  $\kappa$  attribute should be recalculated in every time-step based on the new chemical composition of each SD. However, the  $\kappa$  values relevant for this study are very similar - the  $\kappa$  value of ammonium bisulfate is 0.61 (Petters and Kreidenweis, 2007) and of sulfuric acid is 0.64 (Kim et al., 2016). Therefore, the hygroscopicity parameter is assumed to be constant.

#### 3.1 Dissociation

Dissociation is a reversible process of splitting of molecules dissolved in water drops into ions. It is treated as an equilibrium process and is described using the dissociation constants. The dissociation constant of chemical compound A is denoted here by  $\mathbb{K}_A$ . The dissociation constants are corrected for temperature using the formula of van 't Hoff (1885):

$$5 \quad \mathbb{K}_A(T) = \mathbb{K}_A(T_0) \exp\left(\frac{-\Delta H_D}{R} \left(\frac{1}{T} - \frac{1}{T_0}\right)\right),\tag{1}$$

where  $\Delta H_D$  denotes the reaction enthalpy of dissociation at constant temperature and pressure, T is the temperature of air and R is the gas constant. The list of considered dissociation constants as well as their temperature dependence coefficients is available in Tab. C2. The dissociation of water, although very small, is also taken into account<sup>5</sup>. It is assumed that the water dissociation constant does not vary with temperature.

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It is assumed that there is no electric charge of water drops and therefore the concentrations of positive and negative ions created during dissociation should balance each other. Using the dissociation constants (see Tab. C2), all ion concentrations can be expressed as a function of the total concentration of the dissolved chemical compounds and the concentration of H<sup>+</sup> ions. For example derivation see Chap. 7.6.2 in Seinfeld and Pandis (2016). The neutral charge condition can be expressed as

$$\underbrace{[H^{+}] + \frac{[N^{-III}][H^{+}]\mathbb{K}_{NH3}}{\mathbb{K}_{H2O} + \mathbb{K}_{NH3}[H^{+}]}}_{\text{log}} = \underbrace{\frac{[N^{V}]\mathbb{K}_{HNO3}}{[H^{+}] + \mathbb{K}_{HNO3}}}_{\text{log}} + \underbrace{\frac{[S^{IV}]\mathbb{K}_{SO2}([H^{+}] + 2\mathbb{K}_{HSO3})}{([H^{+}]^{2} + [H^{+}]\mathbb{K}_{SO2} + \mathbb{K}_{SO2}\mathbb{K}_{HSO3})}}_{\text{log}} + \underbrace{\frac{\mathbb{K}_{H2O}}{[H^{+}]}}_{\text{log}} + \underbrace{\frac{[S^{VI}]([H^{+}] + 2\mathbb{K}_{H2SO4})}{([H^{+}] + \mathbb{K}_{H2SO4})}}_{\text{negative ions}} + \underbrace{\frac{[C^{IV}]\mathbb{K}_{CO2}([H^{+}] + 2\mathbb{K}_{HCO3})}{([H^{+}]^{2} + [H^{+}]\mathbb{K}_{CO2} + \mathbb{K}_{CO2}\mathbb{K}_{HCO3})}}_{\text{negative ions}}.$$
(2)

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The [ ] brackets denote the concentration of each of the chemical compounds, (traditionally defined in units of moles per liter), capital letters denote the chemical compound and roman numbers mark its oxidation state. In Eq. (2) the dissociation constants of  $SO_3^{2-}$ ,  $CO_3^{2-}$ , and  $SO_4^{2-}$  ions (i.e.  $\mathbb{K}_{HSO3}$ ,  $\mathbb{K}_{HCO3}$ , and  $\mathbb{K}_{H2SO4}$ ) are multiplied by a factor of two, to take into account the respective electric charge number of those ions.

Equation (2) has only one unknown variable – the new equilibrium concentration of the H<sup>+</sup> ions. The new concentration is obtained iteratively using a numerical root-finding algorithm<sup>6</sup>. The algorithm searches for solution between pH = -1 and pH = 9. The lower bound for the pH scale is unrealistically low and is only necessary at the start of the simulation when the initial SDs have very small volume and are highly acidic. The upper bound is set arbitrarily, but is sufficient for the expected pH of the modeled droplets. At the end of the dissociation procedure the mass of H<sup>+</sup> ions is updated based on the new equilibrium concentration.

When the SD wet radius is quickly changing, for example during the initial condensational growth of cloud droplet or

<sup>&</sup>lt;sup>5</sup> The concentration of undissociated water molecules is so big that it is usually assumed constant and it traditionally multiplies the dissociation constant of water. This leads to a different definition of the dissociation constant for water:  $\mathbb{K}_{H2O} = [H^+][OH^-]$ 

<sup>&</sup>lt;sup>6</sup> TOMS 748 algorithm from *Boost* library. See www.boost.org/doc for documentation and Alefeld et al. (1995) for derivation

rain drop evaporation, the dissociation procedure requires small time-steps to reach convergence. The time-step used in the dissociation procedure can be divided into user-specified number of sub-steps in order to prevent limiting the overall simulation time-step by dissociation.

# 3.2 Dissolution

The amount of the chemical compound that can dissolve into water drop from the gas phase is proportional to its partial pressure above the surface of the drop. Due to the longer timescale of the process, in contrast to dissociation, the transfer between the gas and liquid phase is not treated as an instantaneous process. Assuming that the water drop is internally mixed, the gas-liquid transfer is limited by the diffusion of gas phase particles to the drop surface (gas-phase limitation) and the probability that the molecule will enter the drop after collision (interfacial limitation). Following chapter 8.4.2 in Warneck (1999), for a chemical compound "A" the rate of transfer from the gas phase to the aqueous phase equals

$$\frac{d[\mathbf{A}]}{dt} = \left(\frac{4r_w}{3 < v > \alpha_{M_A}} + \frac{r_w^2}{3D_A}\right)^{-1} \left(c_{A\infty} - \frac{[\mathbf{A}]}{\mathbb{H}_A^{eff}RT}\right) \tag{3}$$

where  $D_A$  and  $\alpha_{M_A}$  are the diffusion and mass accommodation coefficients for the chemical compound "A",  $\langle v \rangle = \sqrt{\frac{8RT}{\pi M_A}}$  is the average velocity of the molecules calculated from the Maxwell-Boltzmann distribution function,  $M_A$  is the molar mass of the chemical compound "A",  $c_{A\infty}$  is the ambient concentration of the trace gas "A" and  $H_A^{eff}$  is the effective Henry's law constant of the chemical compound "A" (i.e. the equilibrium dissolution constant). The Henry's law constants depend on the temperature following similar relation as for dissociation Eq. (1). Table C3 shows the Henry's law constants and their temperature dependencies and Tab. C4 presents the diffusion and mass accommodation coefficients. The term "effective" marks that the dissolution constants take into account the increase of the efficiency due to dissociation (see Seinfeld and Pandis (2016) chap. 7.3 for the exact equations). Equation (3) is solved for each SD and for each of the considered trace gases. It is solved implicitly with respect to the aqueous-phase concentration and explicitly with respect to the gas-phase concentration. The input ambient trace gas concentration is calculated from the trace gas mixing ratio provided by the driver model to which the PBMC scheme is coupled. Obtained aqueous phase concentration is recalculated to the mass of dissolved chemical compounds and the corresponding SD attribute is updated. The changes in the ambient trace gas mixing ratios are calculated by PBMC scheme by summing the changes in chemical composition in all SDs in a given grid-cell and then subtracting them from the trace gas mixing ratio of the driver model. To ensure that the sum of sinks from each SD does not exceed the available ambient trace gas mixing ratio, a relatively short time-steps should be applied. If necessary the user can divide the model time-step into sub-steps.

### 3.3 Oxidation

The reaction rates of oxidation by ozone and hydrogen peroxide can be described as (Hoffmann and Calvert, 1985):

$$\mathbb{R}_{O3} = \frac{d[S^{VI}]}{dt} \Big|_{O3} = \left(k_0 + \frac{k_1 \mathbb{K}_{SO2}}{[H^+]} + \frac{k_2 \mathbb{K}_{SO2} \mathbb{K}_{HSO3}}{[H^+]^2}\right) [O_3][SO_2 * H_2O]$$
(4)

$$\mathbb{R}_{\text{H2O2}} = \frac{d[S^{\text{VI}}]}{dt} \bigg|_{\text{H2O2}} = \frac{k_3 \mathbb{K}_{\text{SO2}}}{1 + k_4 [\text{H}^+]} [\text{H}_2 \text{O}_2] [\text{SO}_2 * \text{H}_2 \text{O}]$$
 (5)

where  $\mathbb{R}_A$  is the reaction rate of the chemical compound "A" and  $k_{0,\dots,4}$  are the reaction rate coefficients.  $k_{0,\dots,4}$  depend on the temperature following similar relation as for dissociation Eq. (1). Table C5 shows the values of the reaction rate coefficients and their temperature dependence coefficients.

Equations (4) and (5) return the new concentration of  $S^{VI}$  created in each SD in each time-step. Based on the new concentration, the new mass of  $S^{VI}$  and the new dry radius are calculated and the corresponding SD attributes are updated. The dry particle density of 1.8 g/cm<sup>3</sup> is assumed while evaluating the dry radius from the  $S^{VI}$  mass.

For the typical atmospheric conditions, say pH between 3 and 6 (i.e.  $[H^+]$  between  $10^{-3}$ M and  $10^{-6}$ M), it can be said that the rate of oxidation by  $H_2O_2$  depends very weakly on pH. In contrast, oxidation by ozone depends strongly on pH of the solution and can become very fast if pH is high. For example, increasing pH by 1 point results in approximately 100 increase in  $O_3$  reaction rate.

#### 3.4 Initialization

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The initial aerosol is assumed to be ammonium bisulfate (NH<sub>4</sub>HSO<sub>4</sub>), with a dry particle density of 1.8 g cm<sup>-3</sup>. Using the dry particle density and the dry radius of each SD, the initial mass of H<sup>+</sup>, NH<sub>4</sub><sup>+</sup> and SO<sub>4</sub><sup>2-</sup> ions is calculated. The initial mass of other molecules and ions is set to zero and is therefore not in equilibrium with the initial ambient trace gas conditions. For the initial conditions where supersaturation is present in the environment it is advisable to allow for a spin-up period with only condensation/evaporation and the equilibrium chemical processes enabled, to allow the model to reach equilibrium. Such initial conditions are mostly relevant for the kinematic models.

# 4 Comparison with moving-bin schemes

The PBMC scheme is set to reproduce results from the model intercomparison study by Kreidenweis et al. (2003), where several bulk and moving-bin schemes representing cloud microphysics and aqueous-phase chemistry were tested in an adiabatic parcel model setup. The parcel model used here is a 0-dimensional model that represents an idealized scenario of a finite volume of air rising adiabatically with a constant vertical velocity. As the parcel of air raises, its temperature decreases leading to supersaturation. This results in activation and further condensational growth of cloud droplets. For the studied oxidation reaction, the presence of liquid water enables aqueous-phase chemical reactions and leads to creation of sulfuric acid within cloud droplets. The collisions between cloud droplets are not included in the parcel simulations to allow an easy comparison with Kreidenweis et al. (2003).

The initial conditions are the same as in Kreidenweis et al. (2003) and are provided for convenience in Tab. 1. The simulation starts below cloud base (i.e. with subsaturation). The initial aerosol is ammonium bisulfate and the initial aerosol particle size

**Table 1.** Initial conditions for the adiabatic parcel test.

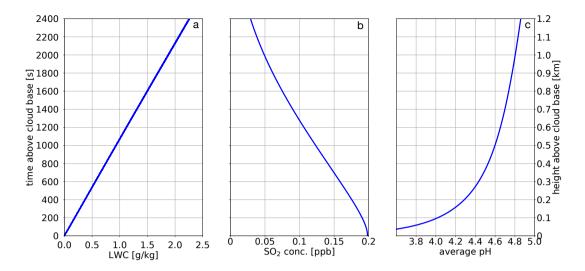
factor	value	units
number of super-droplets	1024	-
time-step	0.1	s
temperature at $t = 0$	285.2	K
pressure at $t = 0$	950	hPa
relative humidity at $t = 0$	95	%
updraft velocity	0.5	m/s
median radius	0.4	$\mu$ m
geometric standard deviation	2	-
total aerosol number concentration	566	$\mathrm{cm}^{-3}$
dry particle density	1.8	g/cm <sup>3</sup>
hygroscopicity	0.61	-
concentration of $SO_2$ at $t = 0$	0.2	ppb-v
concentration of $O_3$ at $t = 0$	50	ppb-v
concentration of $H_2O_2$ at $t = 0$	0.5	ppb-v
concentration of $CO_2$ at $t = 0$	360	ppm-v
concentration of $HNO_3$ at $t = 0$	0.1	ppb-v
concentration of $NH_3$ at $t = 0$	0.1	ppb-v

distribution is assumed to be lognormal with one mode

$$n(r_d) = \frac{n_{tot}}{r_d \sqrt{2\pi} ln(\sigma_g)} exp\left(-\frac{(ln(r_d) - ln(\overline{r_d}))^2}{2ln^2(\sigma_g)}\right)$$
 (6)

where  $n(r_d)$  is the spectral density function of aerosol particle sizes,  $n_{tot}$  is the total aerosol concentration,  $\overline{r_d}$  is median radius and  $\sigma_g$  is the geometric standard deviation. See Sec. 5.1.6 in Arabas et al. (2015) for the details on how the super-droplet dry and wet radii are initialized.

The parcel model employed in this study uses dry air density  $\rho_d$ , dry air potential temperature  $\theta$ , water vapor mixing ratio  $r_v$ , and mixing ratios of ambient trace gases as model variables. In order to calculate  $\rho_d$  at each time-level (or each height-level of the parcel ascent) the model needs to assume a vertical profile of pressure. In the presented simulations the pressure profile is obtained by integrating the hydrostatic equation and assuming that the density of air is constant and equal to 1.15 kg m<sup>-3</sup>. The assumed density is based on the density provided in Tab. 3 in Kreidenweis et al. (2003). Then, at each level,  $\rho_d$  is calculated



**Figure 1.** Physical and chemical conditions in the adiabatic parcel model. Panel a shows the liquid water mixing ratio (LWC), panel b shows the SO<sub>2</sub> concentration (both gas phase and dissolved) panel c shows the water volume weighted average pH of the simulated population of water drops.

from the ideal gas law taking into account the current  $r_v$  and  $\theta$ :

$$\rho_d(p,\theta,r_v) = \frac{p - p_v(p,r_v)}{R_d \theta(\frac{p}{p_{1000}})^{\frac{R_d}{c_{pd}}}},\tag{7}$$

where:  $p_v$  represents the partial pressure of water vapor,  $p_{1000}$  stands for the pressure equal 1000 hPa that comes from the definition of potential temperature,  $R_d$  is the gas constant for dry air and  $c_{pd}$  is the specific heat at constant pressure for dry air. Because the simulated air parcel is assumed to be adiabatic, only the processes resolved by the particle-based scheme can change  $\theta$ ,  $r_v$  and other trace gas mixing ratios. In each model time-step, the particle-based microphysics scheme changes  $\theta$  and  $r_v$  according to Eq. 25 and 26 from Arabas et al. (2015). The changes in the trace gas mixing ratios are resolved following the procedure discussed in Sec. 3.2. It is assumed that the initial mass of dry air within the parcel is 1 kg.

Figure 1 shows the general physical and chemical conditions from the cloud base up to the end of the test run 1.2 km above the cloud base. Two vertical axes are used, representing either the time or the height above the cloud base. Figure 1a shows the liquid water mixing ratio (LWC). The increase in LWC is linear and the LWC reaches above 2 g/kg at height 1.2km above the cloud base. Figure 1b shows the total SO<sub>2</sub> concentration (both in gas phase and dissolved in water) in ppb units. The concentration of SO<sub>2</sub> is decreasing due to oxidation taking place in the cloud droplets. Figure 1c shows the water volume weighted average pH of the cloud droplets. The pH near the cloud base is very low due to the acidic nature of the assumed initial aerosol and the small size of the activated cloud droplets. As the drops grow bigger and become more diluted, the average pH increases. Figure 1 compares well with Fig. 1 from Kreidenweis et al. (2003).

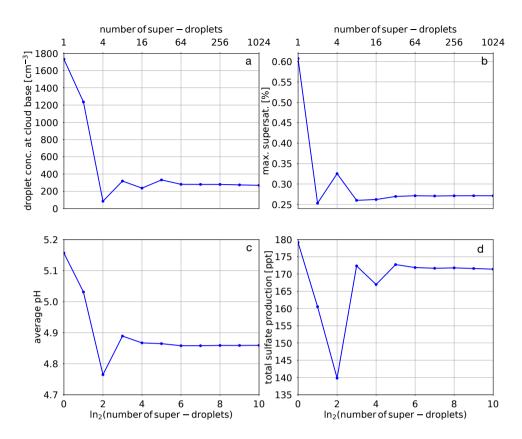
At the end of the test simulation, 85% of SO<sub>2</sub> is converted into S<sup>VI</sup> and the final water volume weighted average pH is equal to 4.86. The total sulfate production is 171 ppt with 99 ppt produced by the  $H_2O_2$  reaction path and 72 ppt produced by the  $O_3$ 

reaction path. Based on Fig. 2 in Kreidenweis et al. (2003), the range of average pH values reported by different size resolving (moving-bin) schemes was between 4.82 and 4.85, and the range of total sulfate production values was between 170 - 180 ppt. Based on Fig. 3 in Kreidenweis et al. (2003) the production by  $H_2O_2$  ranged between 85 and 105 ppt, and by  $O_3$  between 70 and 85 ppt for the size resolving schemes. In short, the results from the particle-based scheme are close to the range of values reported by the moving-bin schemes.

The microphysics schemes taking part in the Kreidenweis et al. (2003) intercomparison study reported significant differences between the number of activated cloud droplets. Based on Tab. 2 in Kreidenweis et al. (2003) the droplet number concentration at the cloud base varied between 275 and 358 cm<sup>-3</sup>. One of the differences between the moving-bin schemes responsible for causing this discrepancy is the different water vapor mass accommodation coefficient  $\alpha_M$  leading to different predicted maximum supersaturation. Figure 6 in Kreidenweis et al. (2003) shows that the observed maximum supersaturations were lower (between 0.23% – 0.26%) for schemes using high values of  $\alpha_M$  (either 0.5 or 1). In contrast, a scheme using  $\alpha_M = 0.042$  predicted maximum supersaturation equal to 0.37%. The particle-based scheme used in this study reports concentration 269 cm<sup>-3</sup> at the level of maximum supersaturation. The maximum supersaturation is equal to 0.27%. The particle-based scheme assumes  $\alpha_M$  equal to unity and therefore fits with the trend of high  $\alpha_M$  causing lower supersaturation presented in Fig. 6 in Kreidenweis et al. (2003).

Another cause for the discrepancy between the bin schemes listed in the intercomparison study are the different sizes and locations of bins in different models, see also the discussion in Arabas and Pawlowska (2011). Along those lines, here it is tested how sensitive the particle-based scheme is to the number of SDs. The results of this test are summarized in Fig. 2 showing the cloud droplet concentration at the cloud base (a), the maximum supersaturation (b), the average pH (c) and the total sulfate production (d). The results are plotted against the logarithm of base two of the number of SDs in the computational domain (meaning that "0" represents one SD and "10" represents 1024 SDs). All values seem to converge for SD numbers greater than 64. The average pH, maximum supersaturation and total sulfate production do not change for those four testruns. The concentration of droplets at the cloud base varies little (between 269 cm<sup>-3</sup> and 281 cm<sup>-3</sup>). The concentrations from simulations with SD number between 512 and 1024 vary between 274 cm<sup>-3</sup> and 269 cm<sup>-3</sup>. For SD numbers between 32 and 64 there are insignificant changes in the maximum supersaturation. The values of pH vary by 0.01 and the total sulfate production increases by 1 ppt. There are, however, large differences between the number of droplets at the cloud base (between 281 cm<sup>-3</sup> and 332 cm<sup>-3</sup>). This confirms the observations from Kreidenweis et al. (2003) that the predicted cloud droplet number concentration strongly depends on the representation of the size distribution of modeled aerosol particles and cloud droplets and that this may become a major source of uncertainties in the microphysics representation. Decrease in the SD number below 32 leads to a big variance in the cloud droplet concentration as well as other parameters.

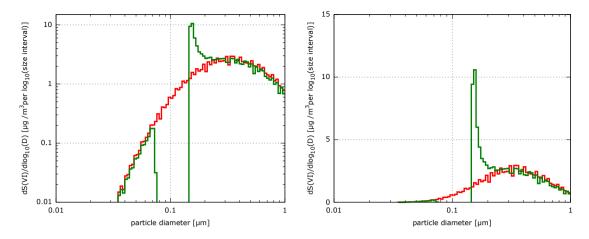
Figure 3 shows the simulated modification of the aerosol size distribution. The red line depicts the initial distribution and the green line shows model state at the end of the adiabatic parcel test. For convenience, Fig. 3 uses both logarithmic (left panel) and linear (right panel) scale on ordinate. The change in the aerosol size distribution is caused by oxidation. The aerosol particles that are too small to become cloud droplets are not affected by aqueous phase oxidation and do not grow in size.



**Figure 2.** Results of the convergence test for the adiabatic parcel simulations. All figures show how a given parameter depends on the number of SDs (shown on the abscissa as the logarithm of base two of the number of SDs). Panel a the cloud droplet concentration at the cloud base, panel b the maximum supersaturation, panel c the water volume weighted average pH at the end of simulation and panel d the total sulfate production.

The large aerosol particles grow in size due to  $S^{VI}$  production during oxidation, but the increase in size is small compared to their initial size. The smallest activated aerosol particles are affected most by oxidation. The increase in their size due to the produced  $S^{VI}$  is the largest compared to their initial size. In short, oxidation produces a "gap", often labeled the "Hoppel minimum", between the CCN processed by the cloud and the smaller unactivated aerosol particles.

The effect of in-cloud sulfate production on the aerosol particle size distribution presented in Fig. 3, combined with other tests presented in this chapter, documents the correctness of the implementation of the aqueous chemistry in the particle-based scheme. The formation of the "Hoppel minimum" was reported by many observational studies, see Hoppel et al. (1994), Bower et al. (1997) and Hudson et al. (2015). Figure 3 compares well with the aerosol size distribution plots from the intercomparison study shown in Fig. 9 in Kreidenweis et al. (2003). Other numerical schemes also reported the formation of the "Hoppel minimum", see for example Flossmann (1994), Feingold and Kreidenweis (2000, 2002) and Ovchinnikov and Easter (2010).



**Figure 3.** Modification of the dry aerosol sulfate mass. The red line shows the initial condition and the green line shows the final model state. The left panel uses logarithmic ordinate and the right panel uses linear ordinate.

The work by Cantrell et al. (1999) shows that the maximum supersaturation and the cloud droplet concentration in the clouds which processed the aerosol particles can be inferred based on the location of the aerosol size distribution minimum.

#### 5 Example simulations

#### 5.1 2D kinematic model

- The kinematic model mimics a single 2D eddy spanning a stratocumulus cloud deck and a boundary layer below it. The model is based on a test scenario from the 8<sup>th</sup> International Cloud Modeling Workshop (ICMW; Muhlbauer et al., 2013, case 1). The velocity field is prescribed as in Szumowski et al. (1998); Morrison and Grabowski (2007); Rasinski et al. (2011). The same model was used when presenting the initial release of *libcloudph++*, see Sec. 2 in Arabas et al. (2015) for the details of the model formulation. The kinematic model is based on the open-source library of parallel MPDATA-based solvers for systems of generalised transport equations, see Jaruga et al. (2015). The temperature, moisture, and trace gas fields are discretized on the Eulerian grid and are advected using the prescribed velocity field. Then, the model variables are passed to the PBMC scheme, where the microphysical and chemical processes are resolved. Finally, the source and sink terms due to microphysics and chemistry are calculated and applied in each model grid-cell as described in Sec. 2 and 3.
- The collisions between water drops are represented using the geometric kernel with collision efficiency for big drops (i.e. radius greater than 20  $\mu$ m) from Hall (1980) and for small droplets (i.e. radius smaller than 20  $\mu$ m) from Pinsky et al. (2008). For big drops, the collision efficiencies were obtained from the fit to measurements, see Hall (1980). For small droplets, the collision efficiencies were based on numerical simulations taking into account turbulence typical for stratocumulus clouds, see Pinsky et al. (2008). The collision efficiencies are provided via a look-up tables for different drop sizes.

The initial conditions are summarized in Tab. 2. The computational domain size is 1.5 km in both directions and the computational grid is composed of  $75\times75$  cells of equal size (the grid lengths are  $20\,\mathrm{m}$ ) and is periodic in the horizontal direction. The initial air density profile corresponds to the hydrostatic equilibrium with the pressure of  $1015\,\mathrm{hPa}$  at the bottom of the domain. At the beginning of the simulation it is assumed that there is no condensed water, and the initial profiles of  $\theta$  and  $r_{\nu}$  are constant with altitude. To keep the simulation setup simple and due to a relatively low vertical extent of the computational domain, the initial trace gas volume fractions are also assumed to be constant with altitude. This unrealistic initial condition results in very high initial supersaturation in the upper part of the domain. As a consequence a  $10^5$  second ( $\sim$  2h 45min) spin-up period is necessary to allow for the simulated water drops to reach equilibrium with their environment. During the spin-up only the reversible processes (condensation and evaporation, dissolving of trace gases and dissociation into ions) are allowed and the supersaturation is limited to 5% (RH=1.05). After spin-up the simulations are run for 30 minutes. The chosen simulation time is enough to deplete the SO<sub>2</sub> available in the cloudy part of the domain as well as to create precipitation.

Similarly to the adiabatic parcel test, the initial aerosol is ammonium bisulfate and the aerosol particle size distribution is lognormal with one mode. The initial condition for trace gases is defined in terms of volume fractions and then translated to mixing ratios that serve as the model variables. The initial  $SO_2$ ,  $O_3$  and  $H_2O_2$  volume fractions are taken from the simulation setup used in Ovchinnikov and Easter (2010). The values for  $SO_2$  and  $O_3$  are based on the measurements from MASE campaign (Wang et al., 2008) and the value for  $H_2O_2$  is based on the representative values for the Eastern Pacific Ocean (Genfa et al., 1999). The  $NH_3$ ,  $HNO_3$  and  $CO_2$  volume fractions are the same as in the parcel test from Sec. 4.

The setup detailed in Tab. 2 corresponds to "very clean conditions" (i.e. low aerosol particle concentrations). The initial aerosol particle sizes are also relatively small. Three additional simulation cases are studied to check the sensitivity of the model to different conditions. In case1 the reversible chemical processes are allowed, but the oxidation reaction is prohibited. In case2 the initial volume fraction of NH<sub>3</sub> is increased and in case3 the initial aerosol size distribution is changed. The conditions for all the sensitivity simulation cases are summarized in Tab. 3.

As discussed in Flossmann (1994), the initial chemical scenario is idealized. For instance, although the initial conditions represent clean maritime environment, the setup lacks sea salt aerosol particles. As discussed by Twohy et al. (1989), sea salt aerosol particles are alkaline, which may in turn increase the pH of water drops and thus affect the oxidation rate. On the other hand, a study by von Glasow and Sander (2001) indicates that alkaline sea salt particles are quickly converted to acidic due to the uptake of HCl vapor. More importantly, including sea salt would result in aerosol particles with very different hygroscopicity values ( $\kappa$  of ammonium bisulfate is 0.61, whereas  $\kappa$  of NaCl is 1.28; Petters and Kreidenweis, 2007). Including sea salt would also result in the initial bi-modal size distribution with one mode representing smaller ammonium bisulfate aerosol particles and the second mode representing larger sea salt particles. In general, including sea salt should result in a very different condensational growth of aerosol particles. The setup used in this study also lacks other particles containing sulfate, such as ammonium sulfate or sulfuric acid aerosol particles.

The initial aerosol size distribution parameters are based on the test cases studied in Feingold and Kreidenweis (2002). The discussion presented in their study introduced two regimes for oxidation with regard to the mean aerosol size  $\overline{r_d}$  and precipitation: (i) for relatively small initial  $\overline{r_d}$  production of sulfate enhances precipitation, (ii) for relatively big initial  $\overline{r_d}$ 

**Table 2.** Initial conditions for base case of 2-dimensional kinematic model.

factor	value	units
number of super-droplets	256	#/grid-cell
model time-step	1	s
particle-based scheme time-step	0.1	s
dry air potential temperature at $t = 0$	289	K
water vapor mixing ratio at $t = 0$	7.5	g/kg
pressure at $z = 0$	1015	hPa
median radius	0.05	μm
geometric standard deviation	1.8	-
total aerosol number concentration	50	$cm^{-3}$
dry particle density	1.8	g/cm <sup>3</sup>
hygroscopicity	0.61	-
concentration of SO <sub>2</sub> at t=0	0.2	ppb-v
concentration of O <sub>3</sub> at t=0	25	ppb-v
concentration of H <sub>2</sub> O <sub>2</sub> at t=0	0.4	ppb-v
concentration of CO <sub>2</sub> at t=0	360	ppm-v
concentration of HNO <sub>3</sub> at t=0	0.1	ppb-v
concentration of NH <sub>3</sub> at t=0	0.1	ppb-v

**Table 3.** Initial conditions for sensitivity test cases of 2-dimensional kinematic model. Specified are: aqueous phase chemistry choice, initial volume fraction of NH<sub>3</sub>, mean radius of the assumed lognormal aerosol particle size distribution  $\overline{r_d}$ , total aerosol concentration  $n_{tot}$  and geometric standard deviation  $\sigma_g$ . Other parameters for each case are the same as in base case (Tab. 2) The parameters that distinguish each sensitivity test case are marked in bold.

case	oxidation reaction	NH <sub>3</sub> [ppb-v]	$\overline{r_d}$ [ $\mu$ m]	$n_{tot}$ [cm <sup>-3</sup> ]	$\sigma_g$
case1	off	0.1	0.05	50	1.8
case2	on	0.4	0.05	50	1.8
case3	on	0.1	0.05	150	1.8

production of sulfate suppresses precipitation. The overall impact depends strongly on the initial concentration of aerosol particles, see Feingold and Kreidenweis (2002) for the discussion. The short simulation time used in this study hinders analysis of the impact of oxidation on the overall precipitation. The work presented here focuses on the evolution of aerosol particle sizes and pH values of cloud and drizzle droplets. Future LES simulations should focus on the impacts of aqueous chemistry on precipitation, cloud lifetime, and cloud dynamics.

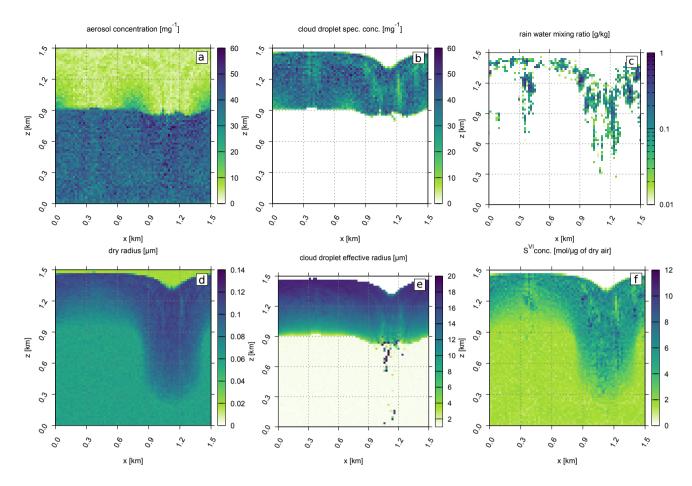


Figure 4. Base case setup (see Tab. 2). All panels depict the model state after 30 minutes simulation time (excluding the spin-up) and show: unactivated aerosol concentration (a), cloud droplet specific concentration (b), rain water mixing ratio (c), mean dry radius (d), cloud droplet effective radius (e) and concentration of  $S^{VI}$  molecules (f). The thresholds for particle radii are: unactivated aerosol < 1  $\mu$ m; 1  $\mu$ m < cloud < 25  $\mu$ m; rain > 25  $\mu$ m. Note the logarithmic scale for the rain water mixing ratio plot.

The kinematic setup precludes any links between cloud microphysical processes and dynamics of the air motion. The setup limits the study to the smooth velocity and therefore smooth saturation fields and prevents mixing between air parcels with different trajectories and properties. On the other hand, the kinematic setup has low computational cost and allows easy testing and sensitivity analysis. Prescribing the velocity ensures that all changes to the aerosol particle and water drop size distributions are caused by the cloud microphysics and aqueous-phase chemistry alone. Moreover, the kinematic setup allows for a straightforward selection of the updraft and downdraft regions, further simplifying the analysis of the microphysical processes.

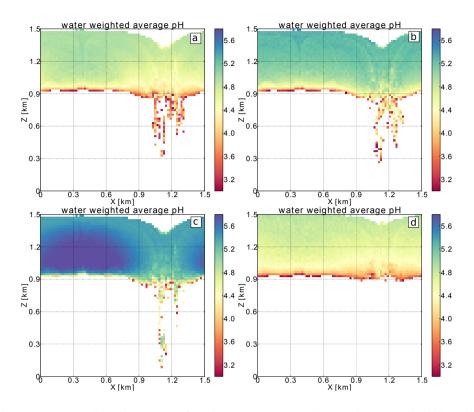
#### 5.2 Results

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Figure 4 shows the model state after 30 minutes of simulation from base case (see Tab. 2). Figure 4a shows the concentration of the unactivated aerosol particles (defined as the SDs with wet radius smaller than 1 µm). The lower part of the plot (below 900 m) shows cloud-free conditions and corresponds to the initial concentration of aerosol particles. The upper part of the plot shows the interstitial aerosol particles, i.e. those in-cloud aerosol particles that did not activate. The difference between the upper and lower parts of Fig. 4a shows the impact of nucleation scavenging on aerosol population. The regions with slightly higher concentration of the in-cloud aerosol particles near the cloud base correspond to regions with low vertical velocities, lower supersaturations and thus lower concentrations of the cloud droplets. Figure 4b shows the concentration of the cloud droplets (defined as the SDs with wet radii between 1 and 25  $\mu$ m). It is nearly constant with height, that agrees with the observations in stratocumulus clouds (e.g. Pawlowska et al., 2000). The regions with lower cloud droplet concentrations correspond to the regions with drizzle (see Fig. 4c). Figure 4c shows the rain water mixing ratio (water drops with wet radius greater than 25  $\mu$ m) using a logarithmic color scale. Rain forms quickly in the simulation due to the relatively high values of cloud droplet radii after the spin-up caused by the low initial aerosol particle concentration. The footprint of precipitation can be seen in Fig. 4a and f where the cloud droplet concentration is depleted in regions of drizzle. Figure 4d shows the mean dry radius of all particles (both the aerosol particles and water drops). The mean dry radius is increasing due to oxidation. In the updraft (left-hand side of panel d) the environmental aerosol particles that have not been affected by cloud are advected into the cloudy region. Once the cloud droplets are formed, the aqueous phase oxidation starts to produce sulfate and changes the CCN size distribution. In the downdraft (right-hand side of panel d) cloud droplets are advected out of the cloud and they evaporate. The cloud-processed CCN are returned to the environment and change the ambient air aerosol particle size distribution. Figure 4e depicts the cloud droplet effective radius. As expected, the effective radius increases with height. At the top of the cloud the effective radius reaches 20 µm, which is linked to the small cloud droplet concentration. High effective radii imply efficient drizzle production after the spin-up (usually water drop radius  $\sim 12 \mu m$  is reported as the threshold value for efficient collisions between water drops and the production of precipitation, for example Rosenfeld and Gutman (1994); Pawlowska and Brenguier (2003)). Figure 4f shows the concentration of SVI molecules (all molecules containing sulfur at +6 oxidation state) and represents molecules from the initial ammonium bisulfate (NH<sub>4</sub>HSO<sub>4</sub>) aerosol and the molecules created during oxidation. It corresponds to the mean dry radius plotted in Fig. 4d. Additionally, some effects of collisions and precipitation can be seen when comparing the irregular features from Fig. 4f with rain water mixing ratio in Fig. 4c. Precipitation displaces the largest water drops, which causes the irregular distribution of SVI molecules in cloudy grid-cells. Figure 4f also shows that the particle-based scheme can track the dissolved chemical compounds in the evaporating rain drops below the cloud base.

Figure 4b, d, e and f show a layer of very clean air above the cloud which is caused by sedimentation of cloud droplets. In the downdraft region, the prescribed velocity field advects the clean layer into the domain. This feature is not present in the aerosol concentration plot (Fig. 4a) because the clean layer contains small aerosol particles with small sedimentation velocity. The depicted clean layer is an artifact caused by the prescribed velocity field and the absence of aerosol sources in the compu-

tational domain. The relatively short simulation time is chosen to minimize the impact of the clean layer on the simulation.



**Figure 5.** The liquid water volume weighted average pH from base case (a), case1 (b), case2 (c), case3 (d). See Tab. 2 and Tab. 3 for a definition of simulation setups.

Figure 5 shows the liquid water volume weighted average pH in each computational grid-cell from base case (a) and sensitivity test cases (b-d). In order to better adjust the color scale to the in-cloud pH variability, pH values below 3 that correspond to very acidic aerosol particles below the cloud base have been clipped. Figure 5 captures the pH of cloud droplets as well as the pH of some evaporating rain drops below the cloud base. The droplets in the downdraft of Fig.5a are more acidic due to  $S^{VI}$  created during aqueous phase oxidation. For the base case, Fig. 5a, pH increases with height above the cloud base. Initially the water drops are very acidic, but as they grow in size they become more diluted. Even though  $S^{VI}$  is created during oxidation, the average pH still increases with height due to dilution. The same behavior is shown in the adiabatic parcel tests discussed in Sec. 4 and shown in Fig. 1. The increase of pH with height is also observed in a 1-dimensional model representing processing of sulfur in small cumuli in marine environment used by (Alfonso and Raga, 2002). Due to the pH variability shown in Fig. 5a oxidation by  $O_3$  happens mostly near the cloud top in the base case. As discussed in Sec. 3.3, the rate of oxidation by  $O_3$  increases significantly with increasing pH (see Eq. 4) whereas oxidation by  $H_2O_2$  depends very weakly on the acidity (Eq. 5). The study by Walcek and Taylor (1986) also reported that the pH of droplets increased with height due to dilution despite the

production of sulfuric acid. In turn, increased pH promotes oxidation by  $O_3$  in the upper parts of the cloud, whereas oxidation by  $H_2O_2$  dominates in lower parts of the cloud, according to their study.

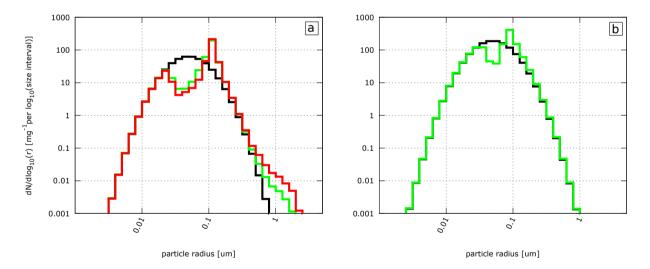
Case1 shown in Fig. 5b represents a hypothetical "no oxidation" scenario where all physical and chemical conditions are the same as in base case, the reversible chemical processes are allowed, and the oxidation reaction is prohibited. The scenario without oxidation is overall less acidic than base case (Fig. 5a). Additionally, without oxidation there is no difference between the pH values in the updraft and downdraft in Fig. 5b. Without oxidation, all the chemical processes are reversible and the dissolved chemical compounds are outgassed to the atmosphere as the cloud droplets evaporate in the downdraft.

Case2 differs from the base case by increasing the initial NH<sub>3</sub> volume fraction from 0.1 ppb-v to 0.4 ppb-v (see Tab. 2 and Tab. 3). Because the initial aerosol particle size distribution is the same as in the base case, the mean aerosol and droplet sizes and concentrations at the end of the simulation are not different from base case (not shown). Figure 5c shows the liquid water volume weighted average pH for case2. The average pH in case2 (Fig. 5c) is higher than in the base case (Fig. 5a), that is, both cloud droplets and rain drops are less acidic in case2 than in the base case. In contrast to the base case, in the updraft (left-hand side of the plots), the pH in case2 actually decreases with height above the cloud base. This is because the higher initial NH<sub>3</sub> volume fraction increases its uptake and counters the low pH values caused by the initial acidic aerosol particles, see Eq. (2). Then, as the water drops are advected upwards, oxidation produces sulfuric acid and the average pH decreases. Near the cloud top, the NH<sub>3</sub> is degassed back to the environment. The case2 results are in agreement with the trajectory ensemble model simulations by Zhang et al. (1999). In their study, the initial aerosol size distribution is the same as in the base case and case2. However, their initial trace gas volume fractions are much higher and aim to represent a "moderately polluted marine environment" (their base case NH<sub>3</sub> volume fraction is ten times larger than base case value assumed here). As in case2 presented here, the high initial NH<sub>3</sub> volume fractions in Zhang et al. (1999) increase the pH near the cloud base and promote oxidation by O<sub>3</sub> during the first minutes after the simulated parcels entered the cloud. Because the sulfuric acid was produced, the pH dropped and oxidation by H<sub>2</sub>O<sub>2</sub> becomes dominant in the higher regions of the cloud, as reported in their study.

Case3 increases the initial aerosol concentration to 150 cm<sup>-3</sup>, while keeping all other initial conditions the same as in the base case (see Tab. 2 and Tab. 3). In general, higher initial aerosol particle concentration results in higher cloud droplet concentrations. This in turn creates smaller cloud droplet effective radii that virtually prohibits the onset of precipitation during the 30 minutes simulation time (not shown). Figure 5d shows the liquid water volume weighted average pH for case3. Similar to the base case (Fig. 5a), the pH increases with height due to the dilution and the downdraft droplets are more acidic due to the ongoing oxidation. However, case3 is more acidic than the base case because the overall droplet sizes are smaller and they are therefore less diluted.

At the end of the base case simulation, 18% of the total available  $S^{IV}$  is oxidized. As a result, 0.14  $\mu g/m^3$  of the dry particulate matter are created during oxidation (an average value for the whole computational domain reported in relation to the dry air volume). In total, 40% of the final dry particulate matter is created due to oxidation and 60% originates from the initial aerosol mass. The oxidation is a significant source of the dry particulate matter because the initial aerosol mass is very low (only 0.21  $\mu g/m^3$  dry air). Oxidation by  $H_2O_2$  is the dominant path: 92% of  $S^{VI}$  originates from  $S^{IV}$  oxidation by  $H_2O_2$ . More

alkaline conditions of case2 enhance the efficiency of oxidation. At the end of case2, 21% of available  $S^{IV}$  is oxidized. As a result, oxidation produces 0.16  $\mu$ g of dry particulate matter per m³ of dry air (average over the whole computational domain). For case2, 44% of the final dry particulate matter is created due to oxidation and 56% originates from the initial ammonium bisulfate aerosol. Similarly to base case, the significance of oxidation as a source of dry particulate matter is caused by a very low initial aerosol mass. Due to more alkaline conditions, oxidation by  $O_3$  becomes more important than in base case. At the end of case3 simulation, 39% of the  $S^{VI}$  originates from  $S^{VI}$  oxidation by  $O_3$  and 61% by  $H_2O_2$ . In contrast, more acidic conditions of case3 hinder the  $O_3$  reaction path. Virtually all molecules of sulfate that are created during oxidation are oxidized by  $H_2O_2$ . As a result, the conversion of sulfur to sulfate is slightly less effective in case3. At the end of case3 simulation, 17% of available  $S^{IV}$  is oxidized. As a result 0.13  $\mu$ g of dry particulate matter are created per m³ of dry air. At the end of case3 simulation, 17% of the dry particulate matter is created by oxidation and 83% originates from the initial aerosol. The initial aerosol mass is larger in case3 than in base case due to the higher initial aerosol concentration (case3 contains initially 0.61  $\mu$ g/m³ of dry particulate matter). Due to the simple kinematic setup chosen in this study the values reported here cannot be treated as representative for the atmospheric conditions. They are shown to allow comparison between the base case and the sensitivity test cases.



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**Figure 6.** The size distributions of dry radii for the base case (a) and case 3 (b). The initial dry radius size distribution is marked in black, the final dry radius size distribution from grid-cells with  $r_c > 0.01$  g/kg in green and from grid-cells with  $r_r > 0.01$  g/kg in red. See Tab. 2 and Tab. 3 for a definition of simulation setups.

Finally, the impact of collisions and aqueous phase oxidation of sulfur on the aerosol and water drop size distributions is examined. For this purpose, the aerosol particle size distributions from the base case (Fig. 6a) and case3 (Fig. 6b) are compared. The black line represents the initial aerosol size distribution and the green and red lines represent the final aerosol size distribution for the in-cloud ( $r_c > 0.01$  g/kg) and precipitating ( $r_r > 0.01$  g/kg) grid-cells. The two cases are chosen

because they have different initial aerosol size distributions. In both cases the cloud-processed aerosol size distributions (green and red lines) have a bi-modal shape. This is a footprint of oxidation that creates the Hoppel minimum in the dry radius size distribution. The same effect is obtained in the adiabatic parcel tests discussed in Sec. 4 and shown in Fig. 3. Moreover, the efficient collisions between water drops in base case create a tail of bigger aerosol sizes in Fig. 6a. The effect is stronger for the precipitating grid-cells (red line). In case3 fewer collisions between water drops occur than in base case and therefore no precipitation and no tail of big aerosol particles is created. Also, in case3, the change in size distribution of aerosol particles caused by oxidation is smaller because the produced sulfate is divided among a larger number of aerosol particles.

# 6 Summary and outlook

The work presented here describes an extension of the *libcloudph++* library that allows to represent the aqueous phase chemical reactions within water drops in the particle-based microphysics scheme. The extension covers the aqueous phase oxidation of sulfur to sulfate. The modular way in which the library is implemented along with the provided documentation should allow, if needed, further development to cover more chemical compounds and reactions. The particle-based microphysics and chemistry scheme is used in 0-dimensional and 2-dimensional modelling setups. The former setup tests the new scheme against the previous numerical studies that used moving-bin microphysics and aqueous phase chemistry schemes. The latter setup focuses on the cloud effects on the aerosol particle size distribution (cloud-aerosol processing). Additionally, the changes in the programming interface due to the aqueous chemistry extension are described in Sec. A in the Appendix. Section C in the Appendix completes the description with a list of chemical constants used in the library and chemical reactions included.

The models used in this study to test the chemistry scheme provide a simplified view of the macrophysical cloud properties. They enable testing of the particle-based scheme but do not provide a good balance between the representation of cloud microphysics and dynamics. As a next step, the particle-based scheme needs to be coupled to an eddy-resolving model. This would allow quantifying how microphysical and chemical processes affect precipitation in the model and how they affect the cloud lifetimes simulated by the model.

Code availability. The libcloudph++ library along with the aqueous-phase chemistry extension, the parcel model and the 2D kinematic model are released under GNU General Public License v3.0. The version of libcloudph++ accompanying this publication is tagged as "2.0.0" at the project repository and is also available as an electronic supplement to this paper. libcloudph++ and the 2D slice model are available at: https://github.com/igfuw/libcloudphxx and the parcel model is available at: https://github.com/igfuw/parcel.

The supported platforms are: Linux with GNU g++, Linux with LLVM clang++ and Apple OSX with the Apple clang++. The code requires c++14 support. The compilation is tested using Travis continuous integration framework.

```
template<typename real_t>
struct opts_t
{
   // process toggling for chemistry
bool chem_dsl, chem_dsc, chem_rct;
   // ...
```

Figure A1. lgrngn::opts\_t definition

#### **Appendix A: Programming Interface**

The programming interface of the particle-based microphysics scheme of libcloudph++ is presented in Sec 5.2. in Arabas et al. (2015). Here, additional information related to the new aqueous phase chemistry scheme is provided. The libcloudph++ is implemented in C++ and therefore some nomenclature related to this programming language is used. For a thorough introduction to C++ programming language see Stroustrup (2013).

The aqueous chemistry module is implemented as an optional extension to the particle-based microphysics scheme in *libcloudph*++. It uses the same **libcloudph**xx::**lgrngn** namespace as the original scheme. Again the template parameter **real\_t** selects between floating point formats of simulations. The particle-based microphysics scheme options are grouped into the structure named **lgrngn::opts\_t**. Chemistry module adds three Boolean fields to this structure: **chem\_dsl chem\_dsc** and **chem\_rct**, see code listing in Fig. A1. When set to true by the user, they switch on dissolving of trace gases into water drops, dissociation of chemical compounds in water drops and oxidation reaction, respectively. The parameters in **lgrngn::opts\_t** can be changed during simulation. For example during the 2-dimensional kinematic simulations from Sec. 5, oxidation is enabled by setting the **chem\_rct** parameter to true at the end of spin-up. Other parameters that cannot be changed during simulation are encapsulated in the **lgrngn::opts\_init\_t** structure. Chemistry module adds three fields to this structure: (i) The Boolean **chem\_switch** field that enables memory allocation for additional variables needed for chemistry representation. (ii) The integer **sstp\_chem** field that defines the number of sub-steps to be carried out in aqueous chemistry calculations. (iii) The **real\_t chem\_rho** field that defines the dry aerosol density, see code listing in Fig. A2.

The names of chemical compounds available in the aqueous phase chemistry module are stored in the **chem\_species\_t** enum, see code listing in Fig. A3. The state of all variables used by the particle-based scheme is stored in an instance of the **lgrngn::particles\_t** structure shown in code listing in Fig. A4. The second template parameter of that structure selects between CPU and GPU calculations (see the discussion in Sec. 5.2 in Arabas et al., 2015, for details). The initialization, time-stepping and output from the particle-based scheme are done using the methods of the **lgrngn::particles\_t** structure. Their signatures are provided in code listing in Fig. A4.

```
template<typename real_t>
struct opts_init_t
{
    // if false no chemical reactions are allowed
    // (no memory allocation)
    bool chem_switch;
    // substeps for chemistry
    int sstp_chem;
    // assumed dry aerosol density
    real_t chem_rho;
```

Figure A2. lgrngn::opts\_init\_t definition

```
enum chem_species_t
{
    // both gas and total dissolved chem species
    HN03, NH3, C02, S02, H202, 03,
    // created sulfate
    S_VI,
    // additional H+ for pH
    H
}
```

Figure A3. lgrngn::chem\_species\_t definition

The **init()** method performs initialization and should be called first. As discussed in Arabas et al. (2015) the first three arguments are obligatory and should point to the dry air potential temperature, water vapor mixing ratio and dry air density fields of the driver model that uses the *libcloudph++*. The next three arguments should point to the Courant number field components. They are optional and depend on the dimensionality of the solved problem. For example, for the parcel model tests from Sec. 4 none are necessary, whereas for the 2-dimensional kinematic model from Sec. 5 two arguments are specified in order to describe the velocity field. The last argument of **init()** is a map with keys from the **chem\_species\_t** enum and values pointing to the corresponding trace gas mixing ratio fields from the driver model. This is an optional argument for simulations with aqueous phase chemistry.

During time-stepping the particle-based scheme computations are performed by the **step\_sync()** and **step\_async()** methods. The first one gathers all the processes that affect the driver model fields (such as condensation/evaporation or aqueous phase chemistry) and the second one gathers all the processes that can be calculated asynchronously (for example collisions or sedimentation). The list of arguments of the **step\_sync()** method is extended by the chemistry module. Similar to the **init()** method, a map linking the **chem\_species\_t** enum items with the driver model mixing ratio fields needs to be provided as the last optional argument. The particle-based scheme overwrites the driver model fields during simulation. The signature of the

```
template <typename real_t, backend_t backend>
struct particles_t: particles_proto_t<real_t>
  // initialisation
  void init(
    const arrinfo_t<real_t> th.
    const arrinfo_t<real_t> rv,
    const arrinfo_t<real_t> rhod,
    const arrinfo_t<real_t> courant_x,
    const arrinfo_t<real_t> courant_y,
    const arrinfo_t<real_t> courant_z,
    const std::map<</pre>
      enum chem_species_t,
      const arrinfo_t<real_t>
   > ambient_chem
  );
  // time-stepping methods
  void step_sync(
    const opts_t<real_t> &,
    arrinfo_t<real_t> th,
    arrinfo_t<real_t> rv,
    const arrinfo_t<real_t> rhod,
    const arrinfo_t<real_t> courant_x,
    const arrinfo_t<real_t> courant_y,
    const arrinfo_t<real_t> courant_z,
    std::map<
      enum chem_species_t,
      arrinfo_t<real_t>
    > ambient_chem,
  );
  void step_async(
    const opts_t<real_t> &
  // diagnostic methods
  // ...
  void diag_chem(const enum chem_species_t&);
```

Figure A4. lgrngn::particles\_t definition

**step\_async** method is not changed by the new chemistry module.

As discussed in Arabas et al. (2015), the **lgrngn::particles\_t** structure provides many methods for obtaining statistical information on the SD parameters (prefixed with **diag**). The chemistry model adds to them the **diag\_chem** method that outputs the total mass of a chemical compound dissolved into droplets. The chemical compound is selected using the **chem\_species\_t** enum items. See the discussion in Sec. 5.2 in Arabas et al. (2015) for the details on how to select the size ranges of droplets specified for output or how to output other statistical parameters.

# Appendix B: Glossary

$\alpha_{M}$	mass accommodation coefficient of water vapor
$\alpha_{M_A}$	mass accommodation coefficient of the chemical compound "A"
$c_{A\infty}$	ambient concentration of the trace gas "A"
$c_{pd}$	specific heat at constant pressure for dry air
$D_A$	diffusion coefficient of the chemical compound "A"
E	reaction activation energy
$\mathbb{H}_{\mathrm{A}}^{eff}$	effective Henry's law constant of the chemical compound "A"
$\Delta H_D$	reaction enthalpy of dissociation at constant temperature and pressure
$\Delta H_H$	reaction enthalpy of dissolution at constant temperature and pressure
$\mathbb{K}_A$	dissociation constant
k <sub>0,,4</sub>	reaction rate coefficients
К	hygroscopicity parameter
$M_A$	molar mass of the chemical compound "A"
$n(r_d)$	spectral density function of aerosol particle sizes
$n_{tot}$	total aerosol concentration
$\mathcal{N}$	super-droplet multiplicity
$\theta$	dry air potential temperature
$p_{v}$	partial pressure of water vapor
$p_{1000}$	pressure equal 1000 hPa
R	gas constant
$R_d$	gas constant for dry air
$\mathbb{R}_{A}$	reaction rate of the chemical compound "A"
$ ho_d$	dry air density
$r_d$	dry radius
$\overline{r_d}$	mean radius of the assumed lognormal aerosol particle size distribution
$r_w$	drop radius
$r_c$	cloud water mixing ratio
$r_r$	rain water mixing ratio
$r_v$	water vapor mixing ratio
$\sigma_g$	geometric standard deviation
T	temperature of air
< <i>v</i> >	average velocity of the molecules

# Appendix C: List of chemical compounds and constants

**Table C1.** Chemical compounds considered in this work.

chemical compound	formula	molar mass (g moles <sup>-1</sup> )	source
ammonia	NH <sub>3</sub>	17	trace gas
carbon dioxide	CO <sub>2</sub>	44	trace gas
hydrogen peroxide	H <sub>2</sub> O <sub>2</sub>	34	trace gas
nitric acid	$HNO_3$	63	trace gas
ozone	$O_3$	48	trace gas
sulfur dioxide	$SO_2$	64	trace gas
sulfuric acid	H <sub>2</sub> SO <sub>4</sub>	98	oxidation reaction product
ammonium bisulfate	NH <sub>4</sub> HSO <sub>4</sub>	115	initial aerosol

**Table C2.** Dissociation constants and their temperature dependence coefficients (taken from Kreidenweis et al., 2003). Dissociation of H<sub>2</sub>SO<sub>4</sub> is taken from Tab. 7A.1 in Seinfeld and Pandis (2016).

	equilibrium reaction	dissociation constant at 298K (moles liter <sup>-1</sup> )	temp. dep. $\frac{-\Delta H_D}{R}$ (K)
$\mathbb{K}_{HNO3}$	$HNO_3(aq) <> H^+ + NO_3^-$	15.4	8700
$\mathbb{K}_{SO2}$	$SO_2 * H_2O <> H^+ + HSO_3^-$	$1.3 \times 10^{-2}$	1960
$\mathbb{K}_{NH3}$	$NH_3 * H_2O <> NH_4^+ + OH^-$	$1.7 \times 10^{-5}$	-450
$\mathbb{K}_{\text{CO2}}$	$CO_2 * H_2O <> H^+ + HCO_3^-$	$4.3 \times 10^{-7}$	-1000
$\mathbb{K}_{HSO3}$	$HSO_3^- <> H^+ + SO_3^{2-}$	$6.6 \times 10^{-8}$	1500
$\mathbb{K}_{HCO3}$	$HCO_3^- <> H^+ + CO_3^{2-}$	$4.68 \times 10^{-11}$	-1760
	$H_2SO_4(aq) <> H^+ + HSO_4^-$	00	-
$\mathbb{K}_{\text{H2SO4}}$	$H_2SO_4(aq) <> H^+ + HSO_4^-$ $HSO_4^- <> H^+ + SO_4^{2-}$	$1.2 \times 10^{-2}$	2720

Table C3. Dissolution constants and their temperature dependence coefficients (taken from Kreidenweis et al., 2003).

	equilibrium reaction	Henry's law constant at 298K (moles liter <sup>-1</sup> atm <sup>-1</sup> )	temp. dep. $\frac{-\Delta H_H}{R}$ (K)
$\mathbb{H}_{HNO3}$	$HNO_3(g) <> HNO_3(aq)$	$2.10 \times 10^5$	-
$\mathbb{H}_{\mathrm{H2O2}}$	$H_2O_2(g) <> H_2O_2(aq)$	$7.45 \times 10^4$	7300
$\mathbb{H}_{\text{NH3}}$	$NH_3(g) <> NH_3 * H_2O$	62	4110
$\mathbb{H}_{SO2}$	$SO_2(g) < > SO_2 * H_2O$	1.23	3150
$\mathbb{H}_{\text{CO2}}$	$CO_2(g) <> CO_2 * H_2O$	$3.40 \times 10^{-2}$	2440
$\mathbb{H}_{\mathrm{O3}}$	$O_3(g) <> O_3(aq)$	$1.13 \times 10^{-2}$	2540

**Table C4.** Diffusion constants (Massman, 1998; Tang et al., 2014) and accommodation coefficients (Kreidenweis et al., 2003) for relevant chemical compounds.

	diffusion coeff. $D_A$ (m <sup>2</sup> /s)	mass accommodation coeff. $\alpha_{M_A}$
$HNO_3$	$65.25 \times 10^{-6}$	0.05
$H_2O_2$	$87.00 \times 10^{-6}$	0.018
$NH_3$	$19.78 \times 10^{-6}$	0.05
$SO_2$	$10.89 \times 10^{-6}$	0.035
$CO_2$	$13.81 \times 10^{-6}$	0.05
$O_3$	$14.44 \times 10^{-6}$	0.00053

Table C5. Reaction rate coefficients and their temperature dependence coefficients (taken from Kreidenweis et al., 2003).

oxidation reaction path	reaction rate coefficient (liter moles <sup>-1</sup> s <sup>-1</sup> ) at 298K	temperature dependence $\frac{-E}{R}$ (K)
a ( ) a vi a aVI		
$O_3(aq) + SO_2 * H_2O - > S^{VI}$	$k_0 = 2.4 \times 10^4$	0
$O_3(aq) + HSO_3^ > S^{VI}$	$k_1 = 3.5 \times 10^5$	-5530
$O_3(aq) + SO_3^{2-} - > S^{VI}$	$k_2 = 1.5 \times 10^9$	-5280
$\mathrm{H}_2\mathrm{O}_2(\mathrm{aq}) + \mathrm{HSO}_3^ > \mathrm{S}^{\mathrm{VI}}$	$k_3 = 7.45 \times 10^7$	-4430

Competing interests. The authors declare that they have no competing interests

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#### References

5

15

- Alefeld, G. E., Potra, F. A., and Shi, Y.: Algorithm 748: Enclosing Zeros of Continuous Functions, ACM Trans. Math. Softw., 21, 327–344, https://doi.org/10.1145/210089.210111, 1995.
- Alfonso, L. and Raga, G.: Estimating the impact of natural and anthropogenic emissions on cloud chemistry: Part I. Sulfur cycle, Atmos. Res., 62, 33–55, https://doi.org/10.1016/S0169-8095(02)00022-4, 2002.
- Andrejczuk, M., Grabowski, W. W., Reisner, J., and Gadian, A.: Cloud-aerosol interactions for boundary layer stratocumulus in the Lagrangian Cloud Model, J. Geophys. Res: Atmospheres, 115, https://doi.org/10.1029/2010JD014248, 2010.
- Andrejczuk, M., Gadian, A., and Blyth, A.: Numerical simulations of stratocumulus cloud response to aerosol perturbation, Atmos. Res., 140–141, 76–84, https://doi.org/10.1016/j.atmosres.2014.01.006, 2014.
- Arabas, S. and Pawlowska, H.: Adaptive method of lines for multi-component aerosol condensational growth and CCN activation, Geosci. Model Dev., 4, 15–31, https://doi.org/10.5194/gmd-4-15-2011, 2011.
  - Arabas, S. and Shima, S.-I.: Large-Eddy Simulations of Trade Wind Cumuli Using Particle-Based Microphysics with Monte Carlo Coalescence, J. Atmos. Sci., 70, 2768–2777, https://doi.org/10.1175/JAS-D-12-0295.1, 2013.
  - Arabas, S., Jaruga, A., Pawlowska, H., and Grabowski, W. W.: libcloudph++ 1.0: a single-moment bulk, double-moment bulk, and particle-based warm-rain microphysics library in C++, Geosci. Model Dev., 8, 1677–1707, https://doi.org/10.5194/gmd-8-1677-2015, 2015.
  - Ardon-Dryer, K., Huang, Y.-W., and Cziczo, D. J.: Laboratory studies of collection efficiency of sub-micrometer aerosol particles by cloud droplets on a single-droplet basis, Atmos. Chem. Phys., 15, 9159–9171, https://doi.org/10.5194/acp-15-9159-2015, 2015.
  - Barth, M. C., Rasch, P. J., Kiehl, J. T., Benkovitz, C. M., and Schwartz, S. E.: Sulfur chemistry in the National Center for Atmospheric Research Community Climate Model: Description, evaluation, features, and sensitivity to aqueous chemistry, J. Geophys. Res: Atmospheres, 105, 1387–1415, https://doi.org/10.1029/1999JD900773, 2000.
  - Bower, K., Choularton, T., Gallagher, M., Colvile, R., Wells, M., Beswick, K., Wiedensohler, A., Hansson, H.-C., Svenningsson, B., Swietlicki, E., Wendisch, M., Berner, A., Kruisz, C., Laj, P., Facchini, M., Fuzzi, S., Bizjak, M., Dollard, G., Jones, B., Acker, K., Wieprecht, W., Preiss, M., Sutton, M., Hargreaves, K., Storeton-West, R., Cape, J., and Arends, B.: Observations and modelling of the processing of aerosol by a hill cap cloud, Atmos. Environ., 31, 2527 2543, https://doi.org/10.1016/S1352-2310(96)00317-2, 1997.
- 25 Cantrell, W., Shaw, G., and Benner, R.: Cloud properties inferred from bimodal aerosol number distributions, J. Geophys. Res, 104, 27615–27624, https://doi.org/10.1029/1999JD900252, 1999.
  - Chen, S., Yau, M. K., and Bartello, P.: Turbulence Effects of Collision Efficiency and Broadening of Droplet Size Distribution in Cumulus Clouds, J. Atmos. Sci., 75, 203–217, https://doi.org/10.1175/JAS-D-17-0123.1, 2018.
- Dianwu, Z., Jiling, X., Yu, X., and Chan, W. H.: Acid rain in southwestern China, Atmos. Environ., 22, 349–358, 30 https://doi.org/10.1016/0004-6981(88)90040-6, 1988.
  - Dziekan, P. and Pawlowska, H.: Stochastic coalescence in Lagrangian cloud microphysics, Atmospheric Chemistry and Physics, 17, 13509–13520, https://doi.org/10.5194/acp-17-13509-2017, 2017.
  - Ervens, B.: Modeling the Processing of Aerosol and Trace Gases in Clouds and Fogs, Chem. Rev., 115, 4157–4198, https://doi.org/10.1021/cr5005887, 2015.
- Faloona, I.: Sulfur processing in the marine atmospheric boundary layer: A review and critical assessment of modeling uncertainties, Atmos. Environ., 43, 2841–2854, https://doi.org/10.1016/j.atmosenv.2009.02.043, 2009.

- Feingold, G. and Kreidenweis, S. M.: Does cloud processing of aerosol enhance droplet concentrations?, J. Geophys. Res.: Atmospheres, 105, 24351–24361, https://doi.org/10.1029/2000JD900369, 2000.
- Feingold, G. and Kreidenweis, S. M.: Cloud processing of aerosol as modeled by a large eddy simulation with coupled microphysics and aqueous chemistry, J. Geophys. Res: Atmospheres, 107, AAC 6–1–AAC 6–15, https://doi.org/10.1029/2002JD002054, 2002.
- Flossmann, A.: A 2-D spectral model simulation of the scavenging of gaseous and particulate sulfate by a warm marine cloud, Atmos. Res., 32, 233–248, https://doi.org/10.1016/0169-8095(94)90063-9, 1994.
  - Genfa, Z., Dasgupta, P. K., Frick, G. M., and Hoppel, W. A.: Airship Measurements of Hydrogen Peroxide and Related Parameters in the Marine Atmosphere Along the Western U.S. Coast, Microchem. J., 62, 99–113, https://doi.org/10.1006/mchj.1999.1715, 1999.
  - Ghan, S., Easter, R., Hudson, J., and Bréon, F.-M.: Evaluation of aerosol indirect radiative forcing in MIRAGE, J. Geophys. Res., 106, 5317–5334, https://doi.org/10.1029/2000JD900501, 2001.

20

25

- Grabowski, W. W. and Abade, G. C.: Broadening of Cloud Droplet Spectra through Eddy Hopping: Turbulent Adiabatic Parcel Simulations, Journal of the Atmospheric Sciences, 74, 1485–1493, https://doi.org/10.1175/JAS-D-17-0043.1, 2017.
- Grabowski, W. W., Dziekan, P., and Pawlowska, H.: Lagrangian condensation microphysics with Twomey CCN activation, Geosci. Model Dev., 11, 103–120, https://doi.org/10.5194/gmd-11-103-2018, 2018.
- Hall, W.: A Detailed Microphysical Model Within a Two-Dimensional Dynamic Framework: Model Description and Preliminary Results, J. Atmos. Sci., 37, 2486–2507, https://doi.org/10.1175/1520-0469(1980)037<2486:ADMMWA>2.0.CO;2, 1980.
  - Hoffmann, F.: On the limits of Köhler activation theory: how do collision and coalescence affect the activation of aerosols?, Atmos. Chem. Phys., 17, 8343–8356, https://doi.org/10.5194/acp-17-8343-2017, 2017.
  - Hoffmann, M. and Calvert, J.: Chemical Transformation Modules for Eulerian Acid Deposition Models: Volume II, the Aqueous-phase Chemistry, U.S. Environmental Protection Agency, Research Triangle Park, NC, 1985.
  - Hoppel, W. A., Frick, G. M., and Larson, R. E.: Effect of nonprecipitating clouds on the aerosol size distribution in the marine boundary layer, Geophys. Res. Lett., 13, 125–128, https://doi.org/10.1029/GL013i002p00125, 1986.
  - Hoppel, W. A., Frick, G. M., Fitzgerald, J. W., and Larson, R. E.: Marine boundary layer measurements of new particle formation and the effects nonprecipitating clouds have on aerosol size distribution, J. Geophys. Res: Atmospheres, 99, 14443–14459, https://doi.org/10.1029/94JD00797, 1994.
  - Hudson, J. G., Noble, S., and Tabor, S.: Cloud supersaturations from CCN spectra Hoppel minima, J. of Geophys. Res.: Atmospheres, 120, 3436 3452, https://doi.org/10.1002/2014JD022669, 2015.
  - Jaruga, A., Arabas, S., Jarecka, D., Pawlowska, H., Smolarkiewicz, P. K., and Waruszewski, M.: libmpdata++ 1.0: a library of parallel MPDATA solvers for systems of generalised transport equations, Geosci. Model Dev., 8, 1005–1032, https://doi.org/10.5194/gmd-8-1005-2015, 2015.
  - Kessler, E.: On the continuity and distribution of water substance in atmospheric circulations, Atmos. Res., 38, 109–145, https://doi.org/10.1016/0169-8095(94)00090-Z, 1995.
  - Kim, J., Ahlm, L., Yli-Juuti, T., Lawler, M., Keskinen, H., Tröstl, J., Schobesberger, S., Duplissy, J., Amorim, A., Bianchi, F., Donahue, N. M., Flagan, R. C., Hakala, J., Heinritzi, M., Jokinen, T., Kürten, A., Laaksonen, A., Lehtipalo, K., Miettinen, P., Petäjä, T., Rissanen,
- M. P., Rondo, L., Sengupta, K., Simon, M., Tomé, A., Williamson, C., Wimmer, D., Winkler, P. M., Ehrhart, S., Ye, P., Kirkby, J., Curtius, J., Baltensperger, U., Kulmala, M., Lehtinen, K. E. J., Smith, J. N., Riipinen, I., and Virtanen, A.: Hygroscopicity of nanoparticles produced from homogeneous nucleation in the CLOUD experiments, Atmos. Chem. Phys., 16, 293–304, https://doi.org/10.5194/acp-16-293-2016, 2016.

- Kreidenweis, S., Walcek, C., Feingold, G., Gong, W., Jacobson, M., Kim, C., Liu, X., Penner, J., Nenes, A., and Seinfeld, J.: Modification of aerosol mass and size distribution due to aqueous-phase SO2 oxidation in clouds: Comparisons of several models, J. Geophys. Res: Atmospheres, 108, https://doi.org/10.1029/2002JD002697, 2003.
- Ladino, L., Stetzer, O., Hattendorf, B., Günther, D., Croft, B., and Lohmann, U.: Experimental Study of Collection Efficiencies between Submicron Aerosols and Cloud Droplets, J. Atmos. Sci., 68, 1853–1864, https://doi.org/10.1175/JAS-D-11-012.1, 2011.

30

- Lebo, Z. J. and Seinfeld, J. H.: A continuous spectral aerosol-droplet microphysics model, Atmos. Chem. Phys., 11, 12297–12316, https://doi.org/10.5194/acp-11-12297-2011, 2011.
- Lee, J., Noh, Y., Raasch, S., Riechelmann, T., and Wang, L.: Investigation of droplet dynamics in a convective cloud using a Lagrangian cloud model, Meteor. Atmos. Phys., 124, 1–21, https://doi.org/10.1007/s00703-014-0311-y, 2014.
- Li, X.-Y., Brandenburg, A., Haugen, N. E. L., and Svensson, G.: Eulerian and Lagrangian approaches to multidimensional condensation and collection, J. Adv. Model. Earth Syst., 9, 1116–1137, https://doi.org/10.1002/2017MS000930, 2017.
  - Massman, W. J.: A review of the molecular diffusivities of H2O, CO2, CH4, CO, O3, SO2, NH3, N2O, NO, and NO2 in air, O2 and N2 near STP, Atmos. Environ., 32, 1111–1127, https://doi.org/10.1016/S1352-2310(97)00391-9, 1998.
- Morrison, H. and Grabowski, W.: Comparison of Bulk and Bin Warm-Rain Microphysics Models Using a Kinematic Framework, J. Atmos. Sci., 64, 2839–2861, https://doi.org/10.1175/JAS3980, 2007.
  - Muhlbauer, A., Grabowski, W. W., Malinowski, S. P., Ackerman, T. P., Bryan, G. H., Lebo, Z. J., Milbrandt, J. A., Morrison, H., Ovchinnikov, M., Tessendorf, S., Thériault, J. M., and Thompson, G.: Reexamination of the State-of-the-art of Cloud Modeling Shows Real Improvements, Bull. Amer. Meteor. Soc., 94, ES45 ES48, https://doi.org/10.1175/BAMS-D-12-00188.1, 2013.
- Myhre, G., Shindell, D., Bréon, F.-M., Collins, W., Fuglestvedt, J., Huang, J., Koch, D., Lamarque, J.-F., Lee, D., Mendoza, B., Nakajima,
   T., Robock, A., Stephens, G., Takemura, T., and Zhang, H.: Anthropogenic and Natural Radiative Forcing, book section 8, pp. 659–740,
   Cambridge University Press, Cambridge, United Kingdom and New York, NY, USA, https://doi.org/10.1017/CBO9781107415324.018,
   2013.
  - Naumann, A. K. and Seifert, A.: A Lagrangian drop model to study warm rain microphysical processes in shallow cumulus, J. Adv. Model. Earth Syst., 7, 1136–1154, https://doi.org/10.1002/2015MS000456, 2015.
- Ovchinnikov, M. and Easter, R.: Modeling aerosol growth by aqueous chemistry in a nonprecipitating stratiform cloud, J. Geophys. Res: Atmospheres, 115, https://doi.org/10.1029/2009JD012816, 2010.
  - Pawlowska, H. and Brenguier, J.-L.: An observational study of drizzle formation in stratocumulus clouds for general circulation model (GCM) parameterizations, J. Geophys. Res.: Atmospheres, 108, https://doi.org/10.1029/2002JD002679, 2003.
  - Pawlowska, H., Brenguier, J., and Burnet, F.: Microphysical properties of stratocumulus clouds, Atmos. Res., 55, 15 33, https://doi.org/10.1016/S0169-8095(00)00054-5, 2000.
  - Petters, M. and Kreidenweis, S.: A single parameter representation of hygroscopic growth and cloud condensation nucleus activity, Atmos. Chem. Phys., 7, 1961–1971, https://doi.org/10.5194/acp-8-6273-2008, 2007.
  - Pinsky, M., Khain, A., and Krugliak, H.: Collisions of Cloud Droplets in a Turbulent Flow. Part V: Application of Detailed Tables of Turbulent Collision Rate Enhancement to Simulation of Droplet Spectra Evolution, J. Atmos. Sci., 65, 357–374, https://doi.org/10.1175/2007JAS2358.1, 2008.
  - Pruppacher, H. and Jaenicke, R.: The processing of water vapor and aerosols by atmospheric clouds, a global estimate, Atmos. Res., 38, 283–295, https://doi.org/10.1016/0169-8095(94)00098-X, 1995.

- Rasinski, P., Pawlowska, H., and Grabowski, W.: Observations and kinematic modeling of drizzling marine stratocumulus, Atmos. Res., 102, 120 135, https://doi.org/10.1016/j.atmosres.2011.06.020, 2011.
- Riechelmann, T., Noh, Y., and Raasch, S.: A new method for large-eddy simulations of clouds with Lagrangian droplets including the effects of turbulent collision, New J. Phys., 14, 065 008, https://doi.org/10.1088/1367-2630/14/6/065008, 2012.
- Rosenfeld, D. and Gutman, G.: Retrieving microphysical properties near the tops of potential rain clouds by multispectral analysis of AVHRR data, Atmos. Res., 34, 259–283, https://doi.org/10.1016/0169-8095(94)90096-5, 1994.
  - Sardina, G., Poulain, S., Brandt, L., and Caballero, R.: Broadening of Cloud Droplet Size Spectra by Stochastic Condensation: Effects of Mean Updraft Velocity and CCN Activation, J. Atmos. Sci., 75, 451–467, https://doi.org/10.1175/JAS-D-17-0241.1, 2018.
  - Sato, Y., Shima, S., and Tomita, H.: A grid refinement study of trade wind cumuli simulated by a Lagrangian cloud microphysical model: the super-droplet method, Atmos. Sci. Lett., 18, 359–365, https://doi.org/10.1002/asl.764, 2017.

- Seifert, A. and Beheng, K. D.: A double-moment parameterization for simulating autoconversion, accretion and selfcollection, Atmos. Res., 59–60, 265–281, https://doi.org/http://dx.doi.org/10.1016/S0169-8095(01)00126-0, 2001.
- Seinfeld, J. and Pandis, S. N.: Atmospheric Chemistry and Physics: From Air Pollution to Climate Change, J. Wiley, New York, 3rd edn., 2016.
- 15 Shima, S., Kusano, K., Kawano, A., Sugiyama, T., and Kawahara, S.: The super-droplet method for the numerical simulation of clouds and precipitation: a particle-based and probabilistic microphysics model coupled with a non-hydrostatic model, Quart. J. Roy. Meteor. Soc., 135, 1307–1320, https://doi.org/10.1002/qj.441, 2009.
  - Sölch, I. and Kärcher, B.: A large-eddy model for cirrus clouds with explicit aerosol and ice microphysics and Lagrangian ice particle tracking, Quart. J. Roy. Meteor. Soc., 136, 2074–2093, https://doi.org/10.1002/qj.689, 2010.
- 20 Stroustrup, B.: C++ Programming Language, Addison-Wesley Professional, fourth edn., http://www.stroustrup.com/4th.html, 2013.
  - Szumowski, M., Grabowski, W., and Ochs III, H.: Simple two-dimensional kinematic framework designed to test warm rain microphysical models, Atmos. Res., 45, 299 326, https://doi.org/10.1016/S0169-8095(97)00082-3, 1998.
  - Tang, M. J., Cox, R. A., and Kalberer, M.: Compilation and evaluation of gas phase diffusion coefficients of reactive trace gases in the atmosphere: volume 1. Inorganic compounds, Atmos. Chem. Phys., 14, 9233–9247, https://doi.org/10.5194/acp-14-9233-2014, 2014.
- Twohy, C. H., Austin, P. H., and Charlson, R. J.: Chemical consequences of the initial diffusional growth of cloud droplets: a clean marine case, Tellus B, 41B, 51–60, https://doi.org/10.1111/j.1600-0889.1989.tb00124.x, 1989.
  - Unterstrasser, S. and Sölch, I.: Optimisation of the simulation particle number in a Lagrangian ice microphysical model, Geosci. Model Dev., 7, 695–709, https://doi.org/10.5194/gmd-7-695-2014, 2014.
  - Unterstrasser, S., Hoffmann, F., and Lerch, M.: Collection/aggregation algorithms in Lagrangian cloud microphysical models: rigorous evaluation in box model simulations, Geosci. Model Dev., 10, 1521–1548, https://doi.org/10.5194/gmd-10-1521-2017, 2017.
  - van 't Hoff, J. H.: L'équilibre chimique dans les systèmes gazeux on dissous à l'état dilué, Recueil des Travaux Chimiques des Pays-Bas, 4, 424–427, https://doi.org/10.1002/recl.18850041207, 1885.
  - von Glasow, R. and Sander, R.: Variation of sea salt aerosol pH with relative humidity, Geophys. Res. Lett., 28, 247–250, https://doi.org/10.1029/2000GL012387, 2001.
- Walcek, C. J. and Taylor, G. R.: A Theoretical Method for Computing Vertical Distributions of Acidity and Sulfate Production within Cumulus Clouds, J. Atmos. Sci., 43, 339–355, https://doi.org/10.1175/1520-0469(1986)043<0339:ATMFCV>2.0.CO;2, 1986.
  - Wang, G., Zhang, R., Gomez, M. E., Yang, L., Levy Zamora, M., Hu, M., Lin, Y., Peng, J., Guo, S., Meng, J., Li, J., Cheng, C., Hu, T., Ren, Y., Wang, Y., Gao, J., Cao, J., An, Z., Zhou, W., Li, G., Wang, J., Tian, P., Marrero-Ortiz, W., Secrest, J., Du, Z., Zheng, J., Shang, D.,

- Zeng, L., Shao, M., Wang, W., Huang, Y., Wang, Y., Zhu, Y., Li, Y., Hu, J., Pan, B., Cai, L., Cheng, Y., Ji, Y., Zhang, F., Rosenfeld, D., Liss, P. S., Duce, R. A., Kolb, C. E., and Molina, M. J.: Persistent sulfate formation from London Fog to Chinese haze, Proc. Natl. Acad. Sci. USA, https://doi.org/10.1073/pnas.1616540113, 2016.
- Wang, J., Lee, Y.-N., Daum, P. H., Jayne, J., and Alexander, M. L.: Effects of aerosol organics on cloud condensation nucleus (CCN) concentration and first indirect aerosol effect, Atmos. Chem. Phys., 8, 6325–6339, https://doi.org/10.5194/acp-8-6325-2008, 2008.
- Warneck, P.: Chemistry of the Natural Atmosphere, vol. 71 of International Geophysics Series, Academic Press, second edn., 1999.

- Werner, F., Ditas, F., Siebert, H., Simmel, M., Wehner, B., Pilewskie, P., Schmeissner, T., Shaw, R. A., Hartmann, S., Wex, H., Roberts, G. C., and Wendisch, M.: Twomey effect observed from collocated microphysical and remote sensing measurements over shallow cumulus, J. Geophys. Res: Atmospheres, 119, 1534–1545, https://doi.org/10.1002/2013JD020131, 2014.
- Zhang, Q., Jimenez, J. L., Canagaratna, M. R., Allan, J. D., Coe, H., Ulbrich, I., Alfarra, M. R., Takami, A., Middlebrook, A. M., Sun, Y. L., Dzepina, K., Dunlea, E., Docherty, K., DeCarlo, P. F., Salcedo, D., Onasch, T., Jayne, J. T., Miyoshi, T., Shimono, A., Hatakeyama, S., Takegawa, N., Kondo, Y., Schneider, J., Drewnick, F., Borrmann, S., Weimer, S., Demerjian, K., Williams, P., Bower, K., Bahreini, R., Cottrell, L., Griffin, R. J., Rautiainen, J., Sun, J. Y., Zhang, Y. M., and Worsnop, D. R.: Ubiquity and dominance of oxygenated species in organic aerosols in anthropogenically-influenced Northern Hemisphere midlatitudes, Geophys. Res. Lett., 34, https://doi.org/10.1029/2007GL029979, 113801, 2007.
  - Zhang, Y., Kreidenweis, S. M., and Feingold, G.: Stratocumulus processing of gases and cloud condensation nuclei: 2. Chemistry sensitivity analysis, J. Geophys. Res: Atmospheres, 104, 16 061–16 080, https://doi.org/11.1029/1999JD900206, 1999.