

Responses to Editor:

Thank you for your constructive suggestions. Our responses to each comment (*italics*) are given below .

- 1) *Please follow the reviewer's advice to motivate your study better in the introduction. I do not think the research question is essential, as it is a model description paper, but nonetheless the purpose of your improved model has to be really clear.*

Response: We have added additional text. The rationale for this study is to extend the applicability of OHM in modelling storage heat flux ΔQ_s for areas (land uses and land covers) where measurements are not available and to cover a wider range of seasonal/meteorological conditions. We have elaborated on this in the manuscript:

- a. The importance of ΔQ_s : page 2 lines 5–17.
- b. The advantages and limitations of OHM: page 3 lines 1–16.
- c. The lack in the physical interpretations of OHM coefficients: page 3 lines 16–18.

- 2) *I would like you to take into account the suggestions of the reviewer concerning the equations. Please make clear where equations are taken from textbooks and where they are your own work.*

Response: We have added additional references to make this clear. Notably, references (Gao et al., 2003; 2010) for the solution (eqns 4 and 5) have been added in the revised manuscript.

- 3) *The reviewer's comments on the outgoing longwave radiation require some additional work, please provide physical justification for ignoring the term.*

Response: First of all, we acknowledge that inclusion of the “re-emitted downwelling longwave radiation” improves the physical rigour of the parameterisation of outgoing longwave radiation L_{\uparrow} . We now include this as part 2 in eqn 10 of the revised manuscript. However, as the omission of this term is well accepted in modelling the outgoing longwave radiation (Bateni and Entekhabi, 2012; Lee et al., 2011; Stensrud, 2007), and its omission greatly enhances the simplicity of the AnOHM formulations, the simplified form of L_{\uparrow} (i.e., $\varepsilon_s \sigma T_s^4$) is still used in this study. The rationale for using the simplified form is presented in Appendix A.

We do not, though, agree with the following comment by the reviewer that “the longwave radiation is not reflected (solely possible for shortwave radiation), but re-emitted.”

We recognise that the concept of reflectivity is only valid for the case of a single wavelength; however, in practise it is referred to in this way for fairly wide wavebands (e.g., such integral reflectivity is referred to as *albedo* for shortwave radiation). However, longwave radiation can be reflected given we are not concerned with ideal blackbody surfaces. The related physics is discussed in section 3 of chapter 1 by Oke (1987).

- 4) *It would also be great if you can explain in your introduction in a single sentence what exactly an OHM is and on which physical principles it is based.*

Response: This explanation has been provided on page 2 lines 29–30 of the introduction: OHM is a model to estimate the storage heat flux ΔQ_S based on a hysteresis relation between ΔQ_S and net radiation Q^* .

5) *It would be useful to have again a native speaker check your paper.*

Response: This has been done.

References:

Bateni, S. M. and Entekhabi, D.: Relative efficiency of land surface energy balance components, *Water Resour. Res.*, 48(4), W04510, doi:10.1029/2011WR011357, 2012.

Gao, Z., Fan, X. G. and Bian, L. G.: An analytical solution to one-dimensional thermal conduction-convection in soil, *Soil Science*, 168(2), 99–107, doi:10.1097/01.ss.0000055305.23789.be, 2003.

Gao, Z., Horton, R. and Liu, H. P.: Impact of wave phase difference between soil surface heat flux and soil surface temperature on soil surface energy balance closure, *J. Geophys. Res.*, 115(D16), D16112, doi:10.1029/2009JD013278, 2010.

Lee, X., Goulden, M. L., Hollinger, D. Y., Barr, A., Black, T. A., Bohrer, G., Bracho, R., Drake, B., Goldstein, A., Gu, L., Katul, G. G., Kolb, T., Law, B. E., Margolis, H., Meyers, T., Monson, R., Munger, W., Oren, R., U, K. T. P., Richardson, A. D., Schmid, H. P., Staebler, R., Wofsy, S. and Zhao, L.: Observed increase in local cooling effect of deforestation at higher latitudes, *Nature*, 479(7373), 384–387, doi:10.1038/nature10588, 2011.

Oke, T. R.: *Boundary Layer Climates*, Taylor & Francis, Abingdon, UK. 1987.

Stensrud, D. J.: *Parameterization schemes: keys to understanding numerical weather prediction models*, Cambridge University Press, Cambridge. 2007.

Responses to Reviewer 2:

We appreciate the comments and constructive suggestions from the reviewer. Our detailed responses are given after each comment (*italics*) below.

Major comment:

- 1) *The authors have revised the manuscript. However, I have the feeling the revised version still contains the main deficiencies that I have mentioned in my previous review. The justification for ignoring the re-emitted downwelling longwave radiation in the upwelling longwave radiation has been addressed, but still ignoring this term is physically wrong. In addition, the readability of the paper has not improved in my view; it is still a long list of equations that are sometimes poorly connected. For example section 3.1 could have literally been copied from a mathematics book, but the link with the AnOHM parameter estimation is missing. I suggest the authors improve the description of all steps that lead to the analytical solutions provided. Moreover the paper lacks a clear research question and some justification for the need to do this research, i.e. what can now be done with the AnOHM that was not possible before with the original OHM, and why an analytical approach is the most feasible to answer the research question.*

Response: We would like to thank the reviewer for his/her constructive comments and critique, which provides us another opportunity to elaborate our thinking.

The reviewer's comments can be summarised as:

1. Ignoring the re-emitted downwelling longwave radiation is physically wrong.
2. The readability of this manuscript needs to be improved.
3. The motivation of this study is unclear.

Our responses to these points are:

1. We recognise that inclusion of the “re-emitted downwelling longwave radiation” improves the physical rigour of the parameterisation of outgoing longwave radiation L_{\uparrow} . This now is included as part 2 in eqn 10 of the revised manuscript.
However, as the omission of this term is very well accepted in modelling the outgoing longwave radiation (Bateni and Entekhabi, 2012; Lee et al., 2011; Stensrud, 2007), and its omission does greatly enhance the simplicity of AnOHM formulations, the simplified form of L_{\uparrow} (i.e., $\varepsilon_s \sigma T_s^4$) is still used in this study. The rationale for using the simplified form is presented in Appendix A.
2. We have checked the steps thoroughly and added explanations where we believe they are necessary.
3. The rationale for this study is to extend the applicability of OHM in modelling storage heat flux ΔQ_S for land covers and time periods where measurements are not available. We have elaborated on this in the manuscript:
 - a. The importance of ΔQ_S : page 2 lines 5–17.
 - b. The advantages and limitations of OHM: page 3 lines 1–16.
 - c. The lack in the physical interpretations of OHM coefficients: page 3 lines 16–18.

Other comments:

- 2) *Eq. 4 and 5 should be correctly referenced, since this is not a solution of your own course.*

Response: References (Gao et al., 2003; 2010) for the solution (eqns 4 and 5) have been added in the revised manuscript.

3) *P13, ln 3: the longwave radiation is not reflected (solely possible for shortwave radiation), but re-emitted.*

Response: we recognise that the concept of reflectivity is only valid for the case of a single wavelength; however, in practise it is often referred to in this way for fairly wide wavebands (e.g., such integral reflectivity is referred to as *albedo*). Longwave radiation can be reflected given we are not concerned with ideal blackbody surfaces. The related physics is discussed (for example) in section 3 of chapter 1 by Oke (1987).

References:

Bateni, S. M. and Entekhabi, D.: Relative efficiency of land surface energy balance components, *Water Resour. Res.*, 48(4), W04510, doi:10.1029/2011WR011357, 2012.

Gao, Z., Fan, X. G. and Bian, L. G.: An analytical solution to one-dimensional thermal conduction-convection in soil, *Soil Science*, 168(2), 99–107, doi:10.1097/01.ss.0000055305.23789.be, 2003.

Gao, Z., Horton, R. and Liu, H. P.: Impact of wave phase difference between soil surface heat flux and soil surface temperature on soil surface energy balance closure, *J. Geophys. Res.*, 115(D16), D16112, doi:10.1029/2009JD013278, 2010.

Lee, X., Goulden, M. L., Hollinger, D. Y., Barr, A., Black, T. A., Bohrer, G., Bracho, R., Drake, B., Goldstein, A., Gu, L., Katul, G. G., Kolb, T., Law, B. E., Margolis, H., Meyers, T., Monson, R., Munger, W., Oren, R., U, K. T. P., Richardson, A. D., Schmid, H. P., Staebler, R., Wofsy, S. and Zhao, L.: Observed increase in local cooling effect of deforestation at higher latitudes, *Nature*, 479(7373), 384–387, doi:10.1038/nature10588, 2011.

Oke, T. R.: *Boundary Layer Climates*, Taylor & Francis, Abingdon, UK. 1987.

Stensrud, D. J.: *Parameterization schemes: keys to understanding numerical weather prediction models*, Cambridge University Press, Cambridge. 2007.

The Analytical Objective Hysteresis Model (AnOHM v1.0):

Methodology to Determine Bulk Storage Heat Flux Coefficients

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Abstract. The net storage heat flux (ΔQ_S) is important in the urban surface energy balance (SEB) but its determination remains a significant challenge. The hysteresis pattern of the diurnal relation between the ΔQ_S and net all-wave radiation (Q^*) has been captured in the Objective Hysteresis Model (OHM) parametrization of ΔQ_S . Although, successfully used in urban areas, the limited availability of coefficients for OHM hampers its application. To facilitate use, and enhance physical interpretations of the OHM coefficients, an analytical solution of the one-dimensional advection-diffusion equation of coupled heat and liquid water transport in conjunction with the SEB is conducted, allowing development of AnOHM (Analytical Objective Hysteresis Model). A sensitivity test of AnOHM to surface properties and hydrometeorological forcing is presented using a stochastic approach (the subset simulation). The sensitivity test suggests that the albedo, Bowen ratio and bulk transfer coefficient, solar radiation and wind speed are most critical. AnOHM, driven by local meteorological conditions at five sites with different land use, is shown to simulate the ΔQ_S flux well (RMSE values of $\sim 30 \text{ W m}^{-2}$). The intra-annual dynamics of OHM coefficients are explored. AnOHM offers significant potential to enhance modelling of the surface energy balance over a wider range of conditions and land covers.

1 Introduction

25 The essential role of an integrated land surface model is to physically predict the land-atmosphere interactions by resolving the transfer of energy, water, and trace gases (Katul et al., 2012; Liang et al., 1994; Sellers et al., 1997). Such land-atmospheric interactions are strongly modulated by the partitioning of solar energy at the land surface (Chen and Dudhia, 2001; McCumber and Pielke, 1981; Yang and Wang, 2014) which can be considered through surface energy balance (SEB) equation (Oke, 1988):

$$Q^* - \Delta Q_S = Q_H + Q_E \quad (1)$$

where Q^* , ΔQ_S , Q_H and Q_E are the net all-wave radiation, net storage, turbulent sensible and latent heat fluxes, respectively. Eqn (1) distinguishes the available energy at the land surface (left hand side) from the heat transfer through turbulent transport (right hand side).

- 5 The turbulent and radiative fluxes (Q^* , Q_H and Q_E) are more readily measured using standard techniques (e.g., eddy-covariance instruments, radiometry, etc.) than ΔQ_S (Offerle et al., 2005; Pauwels and Daly, 2016; Roberts et al., 2006; Wang, 2012). For ΔQ_S the net energy stored or released by changes in sensible heat within the canopy air layer, roughness elements (RE, e.g. vegetation, buildings in an urban environment) and the ground all have to be considered. The volume of interest extends from the top of the roughness sub-layer to the depth in the ground where the daily averaged vertical net heat conduction is zero (see
10 Figure 2 in Masson et al., 2002); this presents very significant challenges of spatial sampling.

Knowledge of ΔQ_S is crucial though to a wide range of processes and applications: from modelling turbulent heat transfer and boundary layer development to predicting soil thermal fields, *etc.* In rural or simple bare soil sites, the flux may be a small fraction of the net all wave radiation (Oliphant et al., 2004). However, in areas where there is more mass, such as cities, the
15 term becomes much more significant (Kotthaus and Grimmond, 2014a) and a key element of the surface energy balance and well-known effects such as the urban heat island.

In urban systems a wide range of techniques have been used to estimate ΔQ_S (Grimmond et al., 1991; Roberts et al., 2006). These include:

- 20 a) Heat conduction approach: the weighted average of heat flows through all urban materials and surfaces by solving heat conduction equations (e.g., buildings, streets, vegetated lands, etc.) (Offerle et al., 2005; Wang et al., 2012; Yang et al., 2014).
- b) Thermal mass scheme: the storage heat is inferred from the changes in thermal mass of all components of the urban system (Kerschgens and Kraus, 1990).
- 25 c) Heat flux plates: combined measurements from grass and paved surfaces (Kerschgens and Hacker 1985; Kerschgens and Draushke 1986)
- d) Parameterisation as a function of Q^* : either as a linear function (Oke et al., 1981); hyperbolic (*cotangent*, *secant*) function (Doll et al., 1985); or a hysteresis relation (Camuffo and Bernardi, 1982). The latter is used in the Objective Hysteresis Model (OHM) (Grimmond et al., 1991).
- 30 e) Residual. Practical difficulties of direct measurement of ΔQ_S in urban areas, result in the SEB residual (i.e., $Q^* + Q_F - (Q_H + Q_E)$) frequently being the “preferred” observations (Ao et al., 2016; Ching et al., 1983; Doll et al., 1985; Li et al., 2015; Oke and Cleugh, 1987) (where Q_F is the anthropogenic heat flux).

The focus here is on the OHM approach, which is forced by Q^* and accounts for the diversity of the surface materials (sub-facets i) in the measurement source area of interest with weightings (f) for their 2 or 3- dimensional extent (Grimmond et al., 1991):

$$\Delta Q_S = \sum_i f_i \left(a_{1,i} Q_i^* + a_{2,i} \frac{\partial Q_i^*}{\partial t} + a_{3,i} \right) \quad (2)$$

where the a_1 , a_2 and a_3 coefficients are for individual facets determined by least square regression between ΔQ_S and Q^* using results from observations (e.g. asphalt road (Anandakumar, 1999), wetlands (Souch et al., 1998), forests (Oliphant et al., 2004)) or numerical modelling (e.g. urban canyons (Arnfield and Grimmond, 1998), roofs (Meyn and Oke, 2009)). These coefficients capture the net behaviour of a facet type in a typical setting, rather than being required to identify the component materials within a facet (e.g. multiple materials making up a roof, wall, with varying thermal connectivity and individual properties). As such, OHM is one of the less demanding parameterisations, yet does capture a more realistic understanding of the relation between ΔQ_S by Q^* compared with other approaches. Despite the shortage of OHM coefficients for the wide range facet types found in cities, OHM captures the urban ΔQ_S overall generally well (Grimmond and Oke, 1999; Järvi et al., 2011; 2014; Karsisto et al., 2015; Roth and Oke, 1995).

OHM is a cornerstone in the urban land surface models, Surface Urban Energy and Water Balance Scheme (SUEWS) (Järvi et al., 2011; 2014; Ward et al., 2016) and Local-scale Urban Meteorological Parameterisation Scheme (LUMPS) (Grimmond and Oke, 2002), and plays an essential role in determining the initial energy partitioning at each time step of the models' simulations. Previous modelling studies (Arnfield and Grimmond, 1998; Meyn and Oke, 2009) have led to better understanding of the OHM coefficients. Solution of the one-dimensional advection-diffusion equation of coupled heat and liquid water transport by Gao et al. (2003; 2008) was used to explore the physical relation of OHM coefficients a_1 and a_2 to the phase lag between ΔQ_S and Q^* . However, insight into a_3 remain unclear (Sun et al., 2013).

In this paper, the solutions of the one-dimensional advection-diffusion equation of coupled heat and liquid water transport (Gao et al., 2003; 2008) are employed with the SEB (eqn 1) to investigate more fully the three OHM coefficients, the outcomes of which lead to development of the Analytical Objective Hysteresis Model (AnOHM) (section 2). The Monte Carlo-based Subset Simulation (Au and Beck, 2001) approach is then used to undertake a sensitivity analysis of AnOHM to surface properties and hydrometeorological conditions (section 3). An offline evaluation of AnOHM's performance for five sites with different land covers (section 4) provides evidence that this is an alternative approach to obtain OHM coefficients, which. Given this allows applications across a much wider range of environments and meteorological conditions, we conclude that AnOHM has important implications for land surface modelling (urban and non-urban).

2 Model Development

2.1 Parameterization of Storage Heat Flux ΔQ_S for a Land Surface

For a given land surface (e.g. bare soil), the governing heat conduction-advection equation can be written (Gao et al., 2003; 2010):

$$\frac{\partial T}{\partial t} = \lambda \frac{\partial^2 T}{\partial z^2} + W \frac{\partial T}{\partial z} \quad (3)$$

- 5 where T is the temperature at a reference depth z (positive downward), t is time, λ is the thermal diffusivity and $W = \partial \lambda / \partial z - (C_w/C_g)w\phi$ is the soil water flux density (Ren et al., 2000), with C_w the volumetric heat capacity of water, C_g the volumetric heat capacity of soil, w the pore water velocity, and ϕ the volumetric soil water content.

The steady-periodic solution of equation (3) corresponding to the principle Earth's rotation frequency ($\omega = 2\pi/24$, in
10 rad h⁻¹), with boundary condition:

$$T_S = A_{T_S} \sin(\omega t - \gamma) + \overline{T_S}, \quad (4)$$

is given by (Gao et al., 2003; 2010):

$$T(z, t) = A_{T_S} \exp(-z/M) \sin(\omega t - z/N - \gamma) + \overline{T_S} \quad (5)$$

where $M = \frac{2\lambda}{\Delta + W}$, $N = \frac{\Delta}{\omega}$ and $\Delta = \sqrt{\frac{W^2 + \sqrt{W^4 + 16\lambda^2 \omega^2}}{2}}$; and $\overline{T_S}$, A_{T_S} and γ denote the daily mean value, amplitude and initial phase of surface temperature, respectively, which need to be determined by the boundary conditions imposed by the SEB.

- 15 From Fourier's law, the soil heat flux is then given by:

$$G(z, t) \equiv -\frac{k \partial T}{\partial z} = k A_{T_S} \frac{\sqrt{M^2 + N^2}}{MN} \exp\left(-\frac{z}{M}\right) \sin\left(\omega t - \frac{z}{N} - \gamma + \delta\right) \quad (6)$$

where $\delta = \arctan\left(\frac{M}{N}\right) = \arctan\left[\frac{2\lambda\omega}{(\Delta + W)\Delta}\right]$ and k is the thermal conductivity. In particular, at the surface $z = 0$, the ground heat flux G_0 is given by:

$$G_0(t) = k A_{T_S} \frac{\sqrt{M^2 + N^2}}{MN} \sin(\omega t - \gamma + \delta), \quad (7)$$

And a simple written form of ΔQ_S (if only one surface) can be given as:

$$\Delta Q_S = G_0 = c_\eta \sin(\omega t + \eta) \quad (8)$$

where $\eta = \delta - \gamma$ and $c_\eta = k A_{T_S} \frac{\sqrt{M^2 + N^2}}{MN}$.

Although the above derivation only considers the land surface made of a single material type, the derived ΔQ_S (eqn 8) can be adapted for surfaces made of composite materials or volumes given appropriate bulk/ensemble properties.

2.2 Parameterization of Net All-wave Radiation Q^* for a Land Surface

- 5 Given the parameterizations of incoming longwave radiation L_\downarrow , outgoing longwave radiation L_\uparrow , sensible heat flux Q_H , latent heat flux Q_E , and storage heat flux ΔQ_S as follows:

$$L_\downarrow = \varepsilon_a \sigma T_a^4, \quad (9)$$

$$L_\uparrow = \underbrace{\varepsilon_s \sigma T_s^4}_1 + \underbrace{(1 - \varepsilon_s) L_\downarrow}_2, \quad (10)$$

$$Q_H = C_h U (T_s - T_a), \quad (11)$$

$$Q_E = Q_H / \beta, \quad (12)$$

$$\Delta Q_S = G_0, \quad (13)$$

the boundary condition imposed by the SEB relation can be rewritten as:

$$(1 - \alpha) K_\downarrow + \varepsilon_a \sigma T_a^4 - \varepsilon_s \sigma T_s^4 = C_h U (1 + \beta^{-1}) (T_s - T_a) + G_0 \quad (14)$$

where the turbulent fluxes Q_H and Q_E are parameterized as functions of temperature gradient $T_s - T_a$ with albedo α , bulk transfer coefficient C_h , wind speed U and Bowen ratio ($\beta = Q_H / Q_E$). Theoretically, the second part of eqn 10 (i.e. $(1 - \varepsilon_s) L_\downarrow$)

- 10 should be accounted for in the estimation of L_\uparrow (Oke, 1987), however, given it is usually less than ~5% of the first part of the equation (see full discussion in Appendix A) for most land covers (Oke, 1987), here it is omitted from consideration and in the development of AnOHM.

- By assuming the incoming solar radiation K_\downarrow and air temperature T_a follow sinusoidal forms through a day as function of the mean value for the day (e.g. $\overline{K_\downarrow}$) (Sun et al., 2013):

$$K_\downarrow = A_K \sin(\omega t) + \overline{K_\downarrow} \quad (15)$$

$$T_a = A_T \sin(\omega t - \tau) + \overline{T_a} \quad (16)$$

and introducing the solar radiation scale:

$$A_K^* = (1 - \alpha) A_K \quad (17)$$

and longwave radiation scale (assuming $\varepsilon_a \approx \varepsilon_s \approx \varepsilon$ as a first order estimate (as AnOHM is insensitive to this parameter see section 3.2); cf. clear sky of ~0.85 (Staley and Jurica, 1972) and urban surfaces of ~0.95 (Kotthaus et al., 2014)):

$$A_T^* = \left(4\varepsilon\sigma\overline{T}_a^3 + (1 + \beta^{-1})C_h U \right) A_T = (f_L + f_T)A_T = fA_T \quad (18)$$

where τ denotes phase differences between T_a and K_\downarrow , the $f = f_L + f_T$ consists of the longwave energy redistribution factor: $f_L = 4\varepsilon\sigma\overline{T}_a^3$ and a turbulent energy redistribution factor: $f_T = (1 + \beta^{-1})C_h U$. Linearizing the fourth-order longwave expressions of temperature at mean daily air temperature \overline{T}_a (Sun et al., 2013), the values of \overline{T}_S and A_{T_S} are obtained:

$$\overline{T}_S = \frac{1 - \alpha}{f} \overline{K}_\downarrow + \overline{T}_a \quad (19)$$

$$\begin{aligned} A_{T_S} &= \frac{fMN \sin(\tau)}{N(fM + k) \sin(\gamma) - kM \cos(\gamma)} A_T \\ &= \frac{1}{\sqrt{M_*^2 + N_*^2}} \frac{\sin(\tau)}{\sin(\gamma - \zeta)} A_T \\ &= \chi_\gamma A_T \end{aligned} \quad (20)$$

where $\zeta = \arctan(N_*/M_*)$, $\gamma = \zeta + \arctan\left(\frac{\sin(\tau)}{\cos(\tau) + A_K^*/A_T^*}\right)$, $M_* = 1 + k/(fM)$, $N_* = k/(fN)$ and $\chi_\gamma = \frac{1}{\sqrt{M_*^2 + N_*^2}} \frac{\sin(\tau)}{\sin(\gamma - \zeta)}$.

5 The net all-wave radiation Q^* is parameterized as:

$$\begin{aligned} Q^* &= (1 - \alpha)K_\downarrow + \varepsilon\sigma T_a^4 - \varepsilon\sigma T_S^4 \\ &= (1 - \alpha)(A_K \sin(\omega t) + \overline{K}_\downarrow) + f_L(T_a - T_S) \\ &= c_\varphi \sin(\omega t + \varphi) + \frac{f_L}{f} (1 - \alpha) \overline{K}_\downarrow \end{aligned} \quad (21)$$

where $\varphi = \arctan\left[\frac{(\chi_\gamma \sin(\gamma) - \sin(\tau))}{[(fA_K^*)/(fLA_T^*) - (\chi_\gamma \cos(\gamma) - \cos(\tau))]} \right]$ and $c_\varphi = \sqrt{\left[\frac{(fA_K^*)^2}{(fLA_T^*)^2} - (\chi_\gamma \cos(\gamma) - \cos(\tau))\right]^2 + [\beta_\gamma \sin(\gamma) - \sin(\tau)]^2}$.

2.3 Derivation of AnOHM coefficients

Based on the above parameterizations of Q^* (eqn 21) and ΔQ_S (eqn 8), together with OHM for a specific surface:

$$\Delta Q_S = a_1 Q^* + a_2 \frac{\partial Q^*}{\partial t} + a_3, \quad (22)$$

the coefficients can be readily derived from the parameterization in section 2.2, as:

$$a_1 = \frac{c_\eta}{c_\varphi} \cos(\eta - \varphi) \quad (23)$$

$$a_2 = \frac{c_\eta}{\omega c_\varphi} \sin(\eta - \varphi) \quad (24)$$

$$\begin{aligned}
a_3 &= -\frac{c_\eta}{c_\varphi} \cos(\eta - \varphi) \cdot \frac{f_T}{f} (1 - \alpha) \overline{K_\downarrow} \\
&= -a_1 \cdot \frac{f_T}{f} (1 - \alpha) \overline{K_\downarrow}
\end{aligned} \tag{25}$$

In the densest parts of cities, the anthropogenic heat (Q_F) often has a large influence on the SEB and it needs to be accounted for (Allen et al., 2011; Chow et al., 2014; Nie et al., 2014; Sailor, 2011). This requires the governing SEB relation (eqn 14) to be rewritten:

$$(1 - \alpha) K_\downarrow + \varepsilon \sigma T_a^4 - \varepsilon \sigma T_s^4 + Q_F = C_h U (1 + \beta^{-1}) (T_s - T_a) + G_0 \tag{26}$$

- 5 Assuming Q_F is diurnally invariant (as a first order estimate e.g. Best and Grimmond (2016)), the derivation (section 2.2) can be extended to include a first order estimate of Q_F to obtain:

$$\begin{aligned}
a_{3F} &= -\frac{c_\eta}{c_\varphi} \cos(\eta - \varphi) \cdot \frac{f_T}{f} (1 - \alpha) \overline{K_\downarrow} - Q_F \\
&= -a_1 \cdot \frac{f_T}{f} (1 - \alpha) \overline{K_\downarrow} - Q_F
\end{aligned} \tag{27}$$

where a_{3F} (subscript ‘F’ indicates the inclusion of Q_F). The other two coefficients remain unchanged.

2.4 Physical Interpretations of AnOHM Coefficients

Based on the parameterizations of AnOHM coefficients (eqns 23, 24, 25/27), physical interpretations can be more fully
10 described compared with OHM:

- a) a_1 characterizes the ratio of ΔQ_S and Q^* and depends on the energy scales (i.e. c_η and c_φ) and their phase difference (i.e. $\eta - \varphi$). The energy scales, representing daily amplitudes of ΔQ_S and Q^* , determine the overall magnitude, while the phase difference moderates the ratio value.
- 15 b) a_2 accounts for the temporal changes in ΔQ_S and Q^* by including the principle Earth’s rotation frequency ω , in addition to the same determinants of a_1 (i.e. c_η , c_φ and $\eta - \varphi$). The complementary sinusoidal functions, with phase difference (i.e., $\sin(\eta - \varphi)$ and $\cos(\eta - \varphi)$), in the formulations of a_1 and a_2 are inversely related with a stronger lag effect from a_2 , and less contribution to ΔQ_S by Q^* (i.e., smaller a_1);
- c) a_3 (or a_{3F}) indicates the baseline ΔQ_S determined by energy redistribution factors (i.e. f_T and f) and energy inputs (i.e. $\overline{K_\downarrow}$, and Q_F if anthropogenic heat is considered) as well as a_1 . It can be inferred from eqn 2 that the nocturnal ΔQ_S is

largely determined by a_3 when the absolute values and variability of Q^* are small at night. A larger daytime energy input (i.e., $\overline{K_d}$, and Q_F if anthropogenic heat is considered) suggests more heat released at night.

3 Sensitivity Analysis

Given to the complex dependence of AnOHM coefficients on surface properties and meteorological forcing (section 2.3), the impacts of these coefficients are assessed further by a sensitivity analysis.

3.1 Subset Simulation

To improve the computational efficiency of undertaking Monte Carlo sensitivity analyses, subset simulation is used (Au and Beck, 2001). This is an adaptive stochastic simulation procedure with particular efficiency in analysing the short-tail of a distribution probability (while also adaptable to long-tail scenarios) (Wang et al., 2011).

If the probability that a critical response Y exceeds a threshold y , $P(Y > y)$, a range of exceedance regions can be specified and sampled using Markov chains. Initially a direct Monte Carlo method is used to choose possible values for the parameter of interest in the anticipated range with a specified distribution (or probability distribution function, PDF) of the uncertainty. From this (level 0), the first exceedance level probability is determined, F_1 at which $P(Y > y_1)$. Then a Markov Chain Monte Carlo (MCMC) procedure is used to generate samples of a given conditional probability p_0 , leading to the exceedance of y_1 in the earlier simulations. This procedure is repeated, for exceedance events F_i at which $P(Y > y_i) = p_0^i$, $i = 1, 2, 3, \dots$, until simulations reach a target exceedance probability, e.g. associated with rare events or risk analysis. Further details of this subset simulation process are provided in Wang et al. (2011).

Subset simulation efficiently generates conditional samples with Metropolis algorithms (Hastings, 1970; Metropolis et al., 1953). This is the basis of MCMC. To generate samples that successively approach a certain conditional probability, a specific Markov chain is designed with the target PDF as its limiting stationary distribution trend as its length increases. The selection of proposal distributions is the key as this controls the next sample generated based on the current one. Ideally, the distribution selection would be automatic but this has an efficiency cost relative to the robustness benefit. For the surface parameters (Table 1a) and hydrometeorological forcing (Table 1b) analyses a normal distribution PDF is used (Au and Beck, 2003; Au et al., 2007), with three conditional levels ($N_{level}=3$) and a conditional probability of $p_0 = 0.1$: i.e., at each level the highest 10% of the outputs are considered to exceed the intermediate threshold. As such, the three-level simulation can effectively capture a rare event with the target exceedance probability of 10^{-4} (i.e., the probability of occurrence is less than 1 in 10000) and generate appropriate samples of different conditional probabilities.

The metric S (in %), used to indicate the sensitivity of the model output Y to a specific uncertainty parameter X (Wang et al., 2011), is:

$$S = \left[\frac{1}{N_{\text{level}}} \sum_{i=1}^{N_{\text{level}}} \frac{E[X|Y > y_i] - E[X]}{E[X]} \right] \times 100 \quad (28)$$

where $i = 1, 2, \dots, N_{\text{level}}$ is the index of conditional sampling level, $E[X]$ is the expectation that the unconditional distribution of a specific uncertainty parameter X , while $E[X|Y > y_i]$ is the expectation of X at conditional level i . A positive (negative) S indicates an increase will lead to increase (decrease) in simulated value. Hence the sign of S indicates the impact of a change in parameter uncertainty. The absolute magnitude of S indicates the sensitivity.

This assessment does not consider if the simulated values have low probability. Later analyses (section 4) consider the simulation results relative to observed fluxes.

10 3.2 Impacts of Surface Properties

Following the sensitivity analysis of AnOHM coefficients to the surface properties, the distributions of conditional samples for thermal conductivity k , bulk heat capacity C_p and emissivity ε are similar to the original proposal distributions (Figure 1), implying weak dependence of a_1 , a_2 and a_3 on these properties. However, for albedo (α) both a_2 and a_3 are sensitive, but a_1 is not; changes in inverse Bowen ratio (β^{-1}) impact all three coefficients; and the bulk transfer coefficient C_h impacts a_1 and a_2 , but has little effect for a_3 .

Using S (eqn 28) to quantify this, it is found that the surface properties (k , C_p and ε) have less sensitivity, with less skewed conditional samples between levels, so S values close to 0 (Figure 2). The S of k is the largest of the three. From the S results for the α sensitivity analysis (Figure 2), it is apparent that an increase in α will increase a_1 while decreasing a_2 and a_3 , whereas the reverse occurs for β^{-1} and C_h (i.e. their decreases leads to larger a_2 and a_3 values but smaller a_1).

From this, the links between the key surface parameters and the storage heat flux can be considered. With an increase in α , there is reduced solar energy in the SEB. This reduces the temporal change in ΔQ_S (smaller a_2) and decreases the baseline value of ΔQ_S (smaller a_3); larger β^{-1} indicates more available energy is dissipated by Q_E than by Q_H , leading to decreased T_s and ΔQ_S (smaller a_1); a smaller portion of Q^* will be dissipated by ΔQ_S (smaller a_1) as the increased C_h can facilitate the turbulent convection and thus increase the total turbulent fluxes.

3.3 Impacts of Hydrometeorological Conditions

Similarly, the sensitivity of AnOHM to hydrometeorological variables is explored (Figure 3). The air temperature (range, mean) and water flux density related variables (i.e. A_T , $\overline{T_a}$ and W) have minimal influence on the skewness of the conditional samples.

In contrast, the incoming shortwave (solar) radiation (range, mean) and wind related variables (i.e. A_K , $\overline{K_l}$ and U) and the phase lag τ between K_l and T_a have large impacts. In terms of the greatest impact on the coefficients (a_1 , a_2 and a_3): A_K and U influences a_1 , τ impacts a_2 , and a_3 responds more to A_K and $\overline{K_l}$ than the other variables.

- 5 Variables that strongly modulate the interactions between ΔQ_S and Q^* can be informed by the S results (Figure 4). For instance, a greater range in K_l (i.e. larger A_K) will occur with larger energy input from solar radiation, leading to stronger heating of the near-surface atmosphere and a smaller portion to ΔQ_S (smaller a_1) but higher baseline ΔQ_S (larger a_3). This is consistent with a reduction in $\overline{K_l}$ having a decrease in a_3 . The temporal change in ΔQ_S is highly correlated with the change in τ , an increase in which implies a slower response of the surface to solar radiation and an overall decrease in ΔQ_S (smaller a_1 , a_2 and a_3). The
- 10 greater sensitivity to τ of a_2 is a key part of the original hysteresis nature of the heating/cooling of a surface. The sensitivity responses of a_1 , a_2 and a_3 to U are very consistent with those to C_h , suggesting the similar pathway that turbulent fluxes (i.e. Q_H and Q_E) modulate ΔQ_S . As W mostly influences the heat conduction-diffusion in the underlying surface as thermal properties (i.e. C_p and k), less dependence is observed on it. This is similar with C_p and k .

4 Model Evaluation

- 15 In this section, the actual ability of AnOHM to determine the storage heat flux relative to observations is evaluated using 30 min observations from five sites of different land use/covers (Table 3). The measurements include turbulent sensible and latent fluxes, along with incoming and outgoing shortwave and longwave radiation and basic meteorological variables (refer to Kotthaus and Grimmond (2014b; 2014a), Klazura et al. (2006), Coulter et al. (2006), Goulden et al. (2006), Scott et al. (2009), and Luo et al. (2007) for details). Anthropogenic heat flux Q_F at the urban site (i.e., UK-Ldn) is estimated using the GreaterQF
- 20 model (Iamarino et al., 2011); the heat storage flux ΔQ_S is thus estimated as the modified residual of urban energy balance as $\Delta Q_S = Q^* + 0.75Q_F - 1.2(Q_H + Q_E)$ (Kotthaus and Grimmond, 2014a; 2014b), which is then used in this evaluation. A similar approach for estimating ΔQ_S (i.e., residual of surface energy balance, $\Delta Q_S = Q^* + Q_F - (Q_H + Q_E)$) is applied at the other (non-urban) sites but with $Q_F = 0$.
- 25 AnOHM is first calibrated with observations under sunny conditions, when the assumptions of AnOHM are best satisfied (i.e., diurnal cycles of K_l and T_a follow sinusoidal forms), to obtain surface properties required by AnOHM (Table 4). As the Bowen ratio β varies daily and monthly (Kotthaus and Grimmond, 2014a; 2014b), β is either determined as the daily value if available, or based on the observation-based monthly climatology (Table 4). The seasonality in albedo α is accounted for also by using its monthly climatology (Table 4). AnOHM is driven by atmospheric forcing (i.e., K_l , T_a , and U) and/or their derived scales
- 30 (A_K , $\overline{K_l}$, A_T , $\overline{T_a}$ and τ) to generate the OHM coefficients (i.e., a_1 , a_2 and a_3 , cf. Figure 5), from which the net heat storage flux ΔQ_S can be predicted (Figure 6) using the observed Q^* with equation 2.

To examine the seasonality of the OHM coefficients, rather than the daily variations in hydrometeorological forcing, LOESS (LOcally wEighted Scatter-plot Smoother, Cleveland and Devlin (1988)) curves are obtained to filter out day-to-day variations in the OHM coefficients (cf. Appendix B for a direct comparison of these coefficients by different modelling and observational regression approaches). Intra-annual variations are found in all the three OHM coefficients (Figure 5), indicating the strong impact of seasonality of meteorological conditions. These controls, as indicated by eqns 23–25(27), are complex and will vary with local conditions. For instance, comparison of OHM coefficients between the AnOHM predictions (LOESS fitted solid lines in Figure 5) and observations at an asphalt road site in Alland, Austria reported in Anandakumar (1999) (empty squares in Figure 5) demonstrates differences in a_1 (Figure 5a) and a_2 (Figure 5b) but general similarity in a_3 (Figure 5c). Compared to a_1 and a_2 , it is noteworthy that, in addition to the S results (cf. Figure 4) given the more explicit mechanism by which the atmospheric conditions moderate a_3 (cf. equation 25/27), such seasonality in a_3 is predicted by AnOHM, and evident in the observations (Figure 5c, also Ward et al. (2013)). Larger $\overline{K_d}$ in warm seasons (May–September) will lead to smaller a_3 (cf. equation 25/27) and *vice versa*.

The AnOHM simulated and observed ΔQ_S agree well at the five different land cover sites, with RMSE values of $\sim 30 \text{ W m}^{-2}$. For comparison purposes, it is noted that the urban land surface model comparison (Best and Grimmond, 2015; Grimmond et al., 2011), found ΔQ_S to be the most poorly represented among all the SEB components with the best RMSE values of 53 W m^{-2} (Lipson et al., 2016). Although the much smaller ΔQ_S RMSE obtained by AnOHM uses a prescribed Bowen ratio in the offline evaluation, such improvement indicates the ability of AnOHM to simulate a more consistent ΔQ_S with observations.

Compared with OHM predictions (orange lines in Figure 6), AnOHM (blue lines in Figure 6) better reproduces the seasonality in ΔQ_S but gives larger bias at two sites with natural land covers (i.e., US-SRM and US-SO4). This can be attributed to the overestimates of nocturnal ΔQ_S by AnOHM. Overall, the evaluation demonstrates good performance of AnOHM in predicting the long-term ΔQ_S with clear seasonality reproduced across a wide range of surface types.

5 Discussion and Concluding Remarks

In this study, an Analytical Objective Hysteresis Model (AnOHM) is developed to obtain OHM coefficients across a wide range of surface and meteorological conditions and to improve physical understanding of the interactions between ΔQ_S and Q^* . The sensitivity of AnOHM to surface properties and hydrometeorological conditions are analysed through Monte Carlo based subset simulations (Au and Beck, 2001). The results highlight the importance of the albedo, the Bowen ratio and the bulk transfer coefficient, and the importance of solar radiation and wind speed in regulating the heat storage. The importance of albedo in modulating the heat storage was also found by Wang et al. (2011), who also used the same subset simulation approach with the Single Layer Urban Canopy Model (SLUCM, details refer to Kusaka et al. (2001)). This demonstrates the

consistency in heat storage modelling between AnOHM and UCM. From the sensitivity results, variations in OHM coefficients of a similar size may arise from either surface property parameters or hydrometeorological forcing that are associated with the same physical processes (cf. bulk transfer coefficient C_h in Figure 2 and wind speed U in Figure 4). This suggests the ability of AnOHM in representing physical processes. An offline evaluation of AnOHM using flux observations from five sites with different land covers demonstrates its ability to predict the intra-annual dynamics of OHM coefficients and shows good agreement between simulated and observed storage heat fluxes. In particular, the seasonality in the OHM coefficient a_3 observed in a previous study (Anandakumar, 1999) is well predicted by AnOHM.

The limitations of AnOHM are important to consider. First, given the assumption that the incoming solar radiation K_{\downarrow} and air temperature T_a diurnal cycles are sinusoidal, optimal performance of AnOHM occurs under clear sky conditions. The current parameterisations of K_{\downarrow} and T_a within AnOHM only consider the harmonics of principal frequencies for formulation simplicity. More frequencies may potentially resolve more realistic diurnal variations in K_{\downarrow} and T_a . As the reflected part of L_{\downarrow} (i.e., $(1 - \varepsilon_s)L_{\downarrow}$) is assumed negligible, and similar emissivity values are assumed for sky and land surface (i.e., $\varepsilon_s \approx \varepsilon_a \approx \varepsilon$), the outgoing longwave radiation is underestimated. These simplifications greatly facilitate the AnOHM formulation without qualitatively changing the final results as the sensitivity analyses (cf. minimal S values for ε in Figure 2) demonstrate. The inclusion of water flux density W equips AnOHM with an ability to investigate the hydrological impacts of the underlying surface on land-atmosphere interactions. However, estimation of W remains challenging (Wang, 2014) and the resulting uncertainty in the final results warrants caution in conducting simulations over land covers with strong soil moisture dynamics (e.g., grassland with high soil moisture under clear sky condition).

Despite these limitations, AnOHM does permit improved modelling of the surface energy balance through its physically-based parameterization scheme for storage heat flux ΔQ_s . Compared to OHM, AnOHM has the benefit of allowing ΔQ_s to be simulated for land covers for which coefficients are not available and to allow for seasonal variability to be accounted for. As AnOHM shares similar hydrometeorological forcing inputs (i.e., K_{\downarrow} , T_a and U) to other land surface models (LSMs), it can potentially be used within in LSMs to estimate ΔQ_s , or if turbulent fluxes are included to be a complete LSM. The overall improvements from adopting AnOHM in modelling land surface processes will be presented in forthcoming work in the SUEWS-AnOHM framework.

Code availability

The Fortran source code for AnOHM can be obtained from the corresponding authors upon request.

Appendix A: Rationale for A Simplified Formulation of Outgoing Longwave Radiation

In the formulation of outgoing longwave radiation L_{\uparrow} , a simplified form (i.e., $\varepsilon_s \sigma T_s^4$) is used for AnOHM by ignoring part 2 of eqn 10 (i.e., $(1 - \varepsilon_s)L_{\downarrow}$). The rationale for such simplification is that given ε_s is usually larger than 0.9, $(1 - \varepsilon_s)L_{\downarrow}$ contributes a relatively small portion to the total longwave component (Oke, 1987) and omission of this part is well accepted in the parameterization of outgoing longwave radiation for land surface modeling across various land covers (Bateni and Entekhabi, 2012; Lee et al., 2011){Stensrud:2007dw}.

Using the parameterisation of incoming longwave radiation in the AnOHM framework (i.e., $L_{\downarrow} = \varepsilon_a \sigma T_a^4 \approx \varepsilon_s \sigma T_a^4$), we conduct a sensitivity analysis of the ratio between the ignored part (i.e., $(1 - \varepsilon_s)L_{\downarrow}$) and total outgoing longwave radiation (i.e., $\varepsilon_s \sigma T_s^4 + (1 - \varepsilon_s)L_{\downarrow}$) at a constant air temperature of 20 °C and find this ratio is generally less than 5% given ε_s ranges between 0.90 and 0.99 (Figure A1).

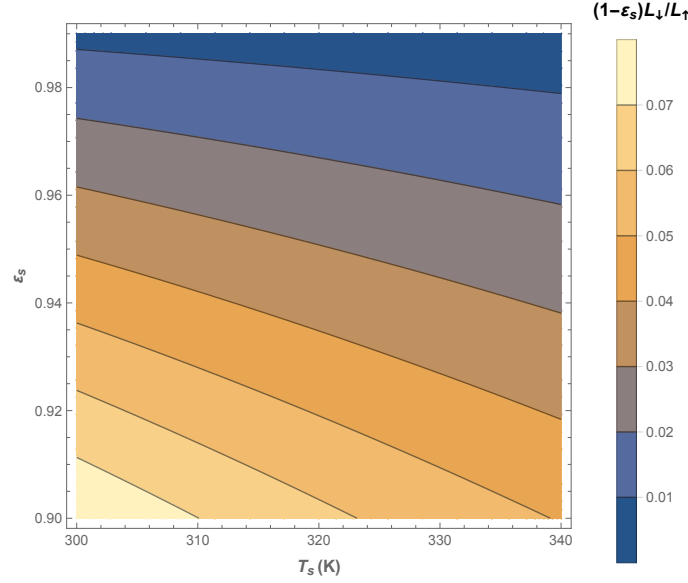


Figure A1 Ratio between the second part of eqn 10 (i.e., $(1 - \varepsilon_s)L_{\downarrow}$) and total outgoing longwave radiation (i.e., $\varepsilon_s \sigma T_s^4 + (1 - \varepsilon_s)L_{\downarrow}$) at a constant air temperature of 20 °C.

Moreover, if $(1 - \varepsilon_s)L_{\downarrow}$ is included in the net longwave radiation, the induced effect can be incorporated into a modified sky emissivity $\varepsilon'_a = \varepsilon_s \varepsilon_a$ as follows:

$$\begin{aligned}
 L_{net} &= L_{\downarrow} - L_{\uparrow} \\
 &= L_{\downarrow} - (\varepsilon_s \sigma T_s^4 + (1 - \varepsilon_s)L_{\downarrow}) \\
 &= \varepsilon_s L_{\downarrow} - \varepsilon_s \sigma T_s^4 \\
 &= \varepsilon_s \varepsilon_a \sigma T_a^4 - \varepsilon_s \sigma T_s^4 \\
 &= \varepsilon'_a \sigma T_a^4 - \varepsilon_s \sigma T_s^4
 \end{aligned}$$

Then by assuming $\varepsilon \approx \varepsilon'_a \approx \varepsilon_s$, the derivation following equation 18 still holds. The sensitivity analysis suggests that the derived coefficients are insensitive to ε (cf. S for ε in Figure 2).

As such, we deem the omission of $(1 - \varepsilon_s)L_\downarrow$ will not qualitatively change the results of this work.

Appendix B: Comparison in OHM Coefficients between Different Modelling Approaches and Observation Regression

- 5 The comparison in OHM coefficients by different modelling and observational regression approaches (Figure A2) indicate AnOHM generally follows the results by observation regression, whereas the typical coefficient values adopted by OHM do not.

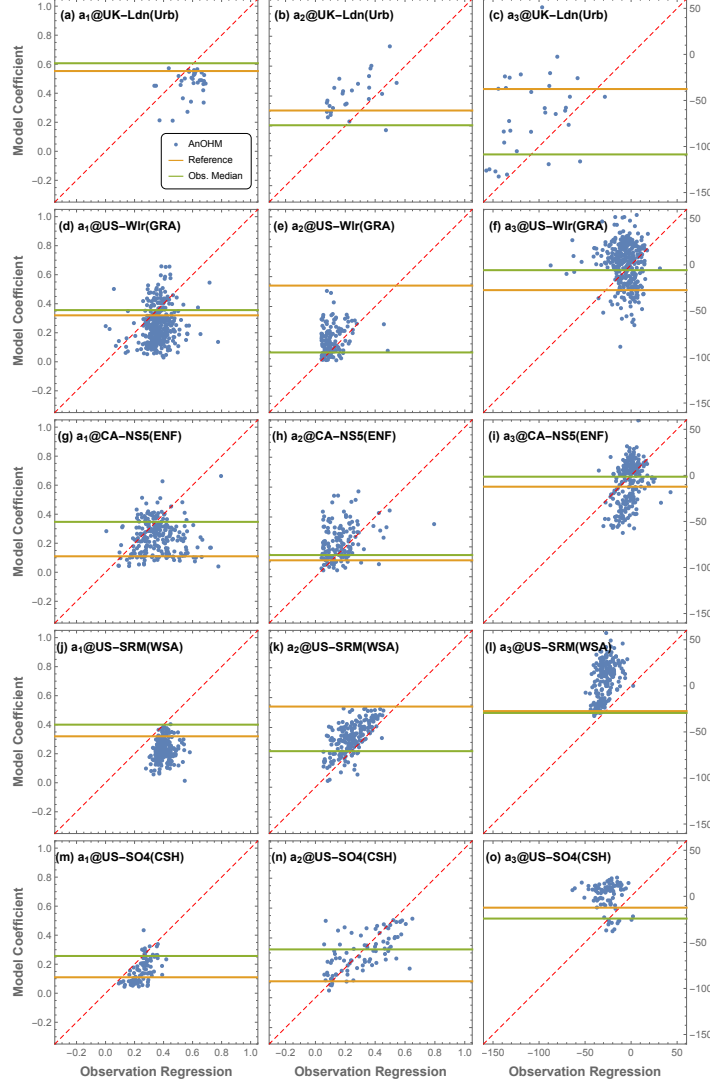


Figure A2 Comparison of OHM coefficients (left, central and right columns for a_1 , a_2 and a_3 , respectively) between different modelling approaches and observation regression at five sites: UK-Ldn (a, b, c), US-Wlr (d, e, f), CA-NS5 (g, h, i), US-SRM (j, k, l) and US-SO4 (m, n, o). The blue dots denote the paired values between AnOHM and observation regression. The orange lines represent the reference value used in OHM simulations for land covers of grass and tree (Grimmond and Oke, 1999), whereas the green lines shows median values derived from results by observation regression at corresponding sites.

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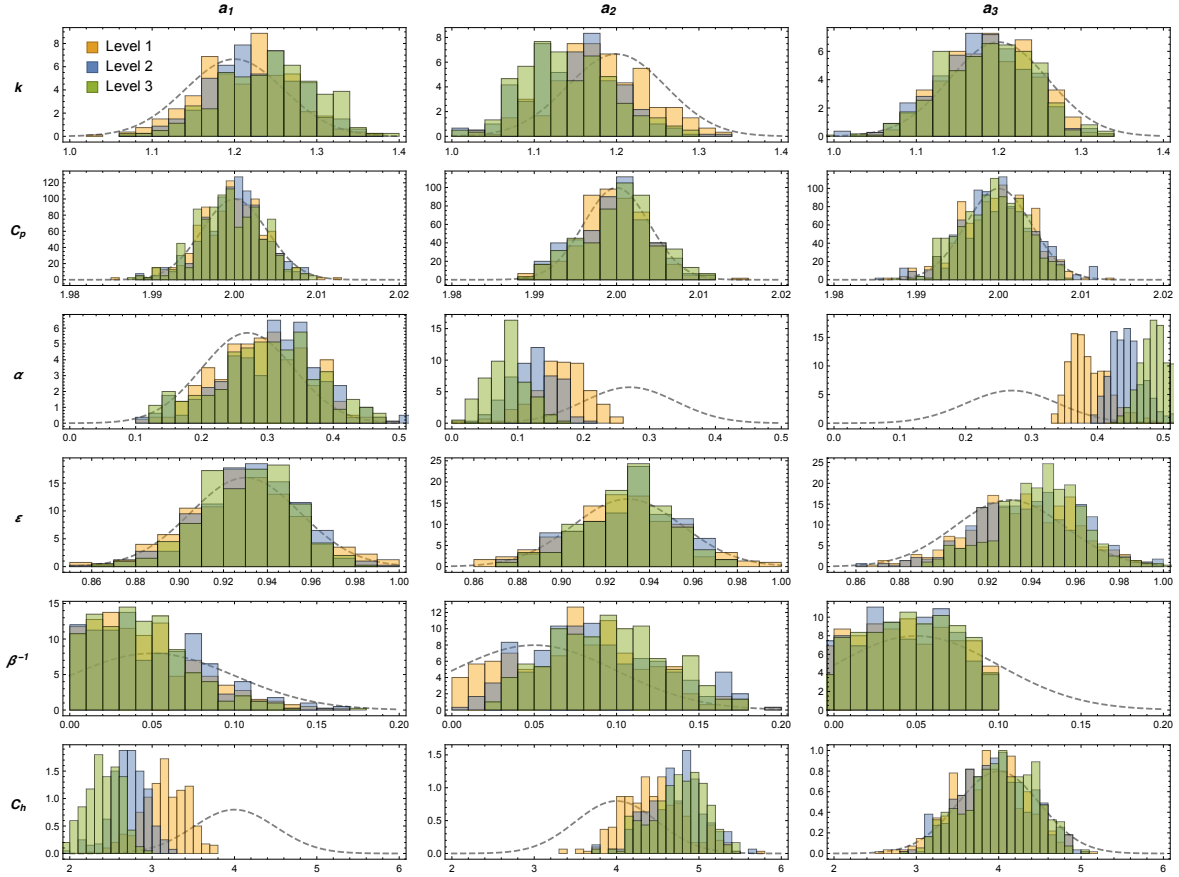


Figure 1 Histograms of conditional samples at different conditional levels for surface property parameters (rows from top: thermal conductivity k in $\text{W m}^{-1} \text{K}^{-1}$, heat capacity C_p in $\text{MJ m}^{-3} \text{K}^{-1}$, albedo α , emissivity ϵ , inverse Bowen ratio β^{-1} and bulk transfer coefficient C_h in $\text{J m}^{-3} \text{K}^{-1}$) with AnOHM coefficients as the model output (columns from left: a_1 , a_2 and a_3). Each subplot x-axis is the parameter value and y-axis is the PDF value. The original proposal distribution (dashed line) and simulation levels (different colours) are shown.

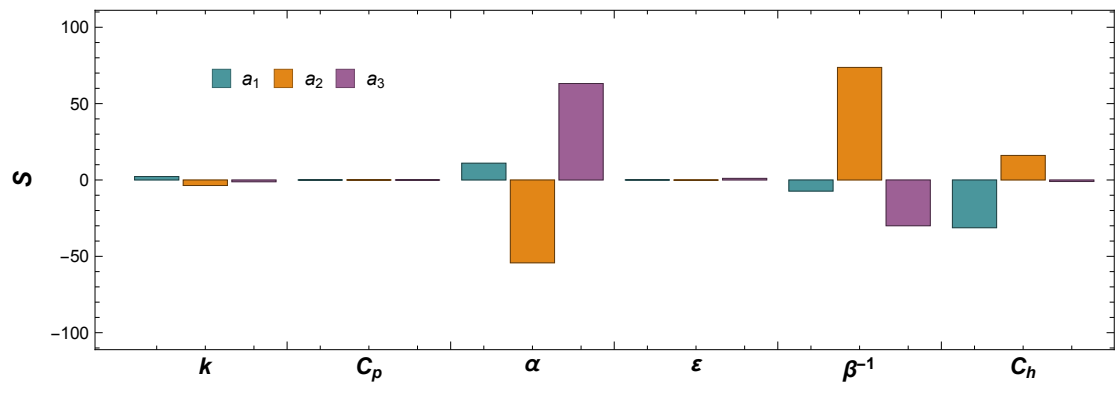


Figure 2 Relative variation in sensitivity (S , %, equation 28) to surface parameters. See Figure 1 for further details.

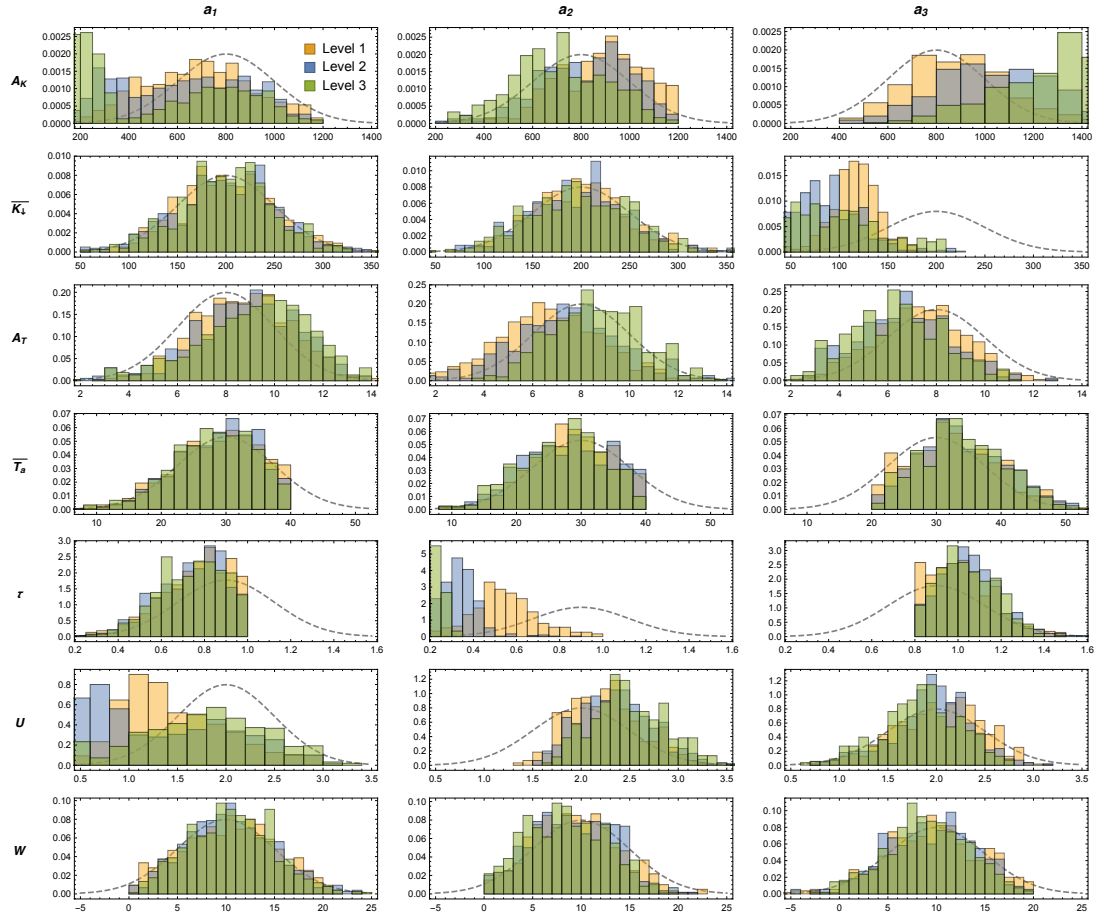


Figure 3 Histograms of conditional samples at different conditional levels for ambient forcing parameters (rows from top: incoming solar radiation amplitude A_K in W m^{-2} and its daytime mean $\overline{K_L}$ in W m^{-2} , air temperature amplitude A_T in $^{\circ}\text{C}$ and its daily mean $\overline{T_a}$ in $^{\circ}\text{C}$, the phase lag τ in rad between K_L and T_a , wind speed U in m s^{-1} and water flux density W in m s^{-1}) with AnOHM coefficients as the model output (columns from left: a_1 , a_2 and a_3). As Fig. 1.

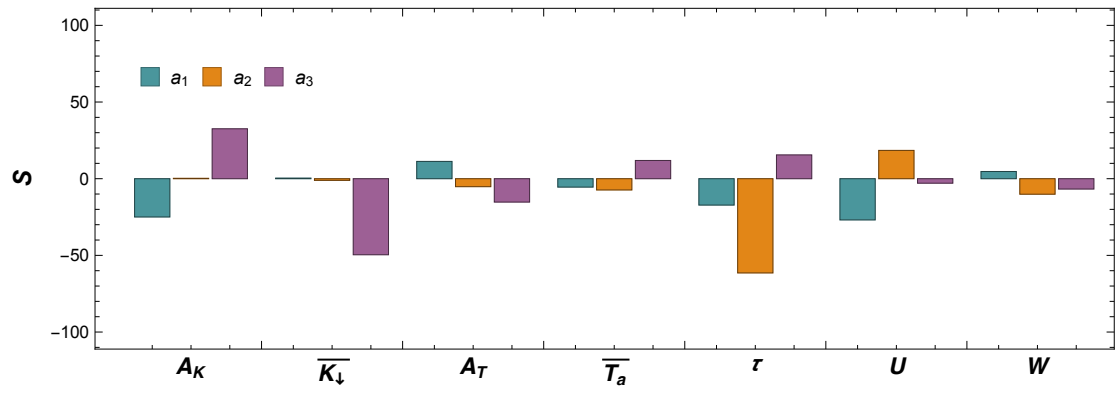


Figure 4 Relative variation in sensitivity (S , %, equation 28) to forcing parameters. See Figure 3 for further details.

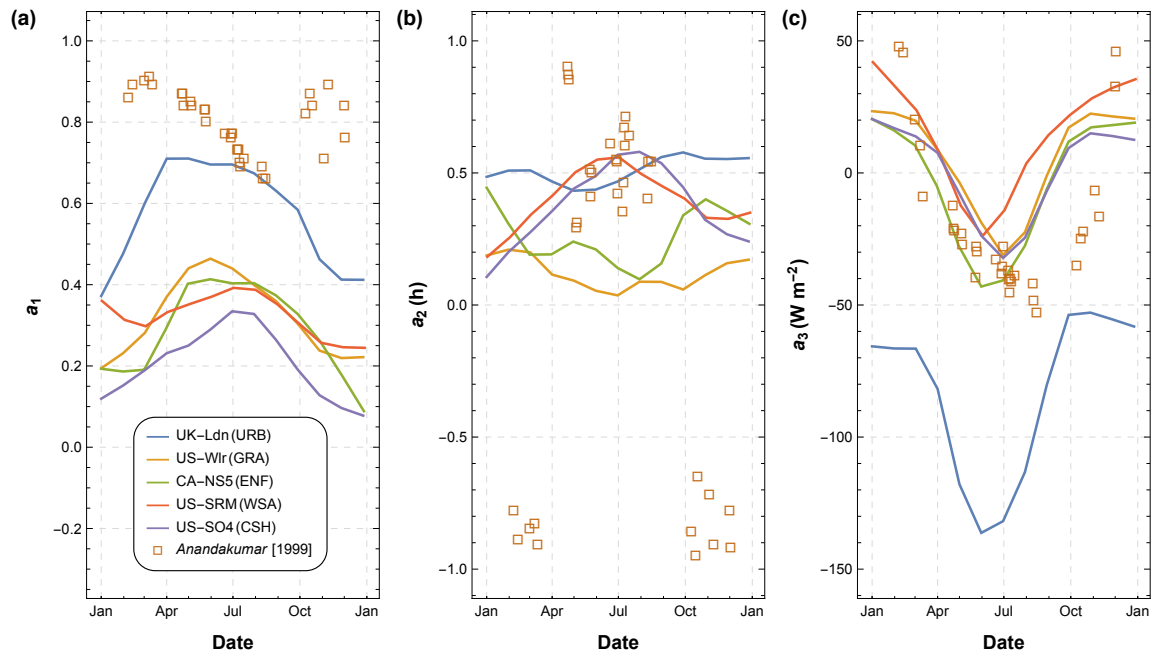


Figure 5 Intra-annual variations of OHM coefficients: (a) a_1 , (b) a_2 and (c) a_3 . LOESS fits (solid lines) through the daily values predicted by AnOHM and daily values (squares) measured at an asphalt road site (Anandakumar 1999) are shown. The LOESS (Cleveland and Devlin 2012) fitting is a locally weighted polynomial regression approach.

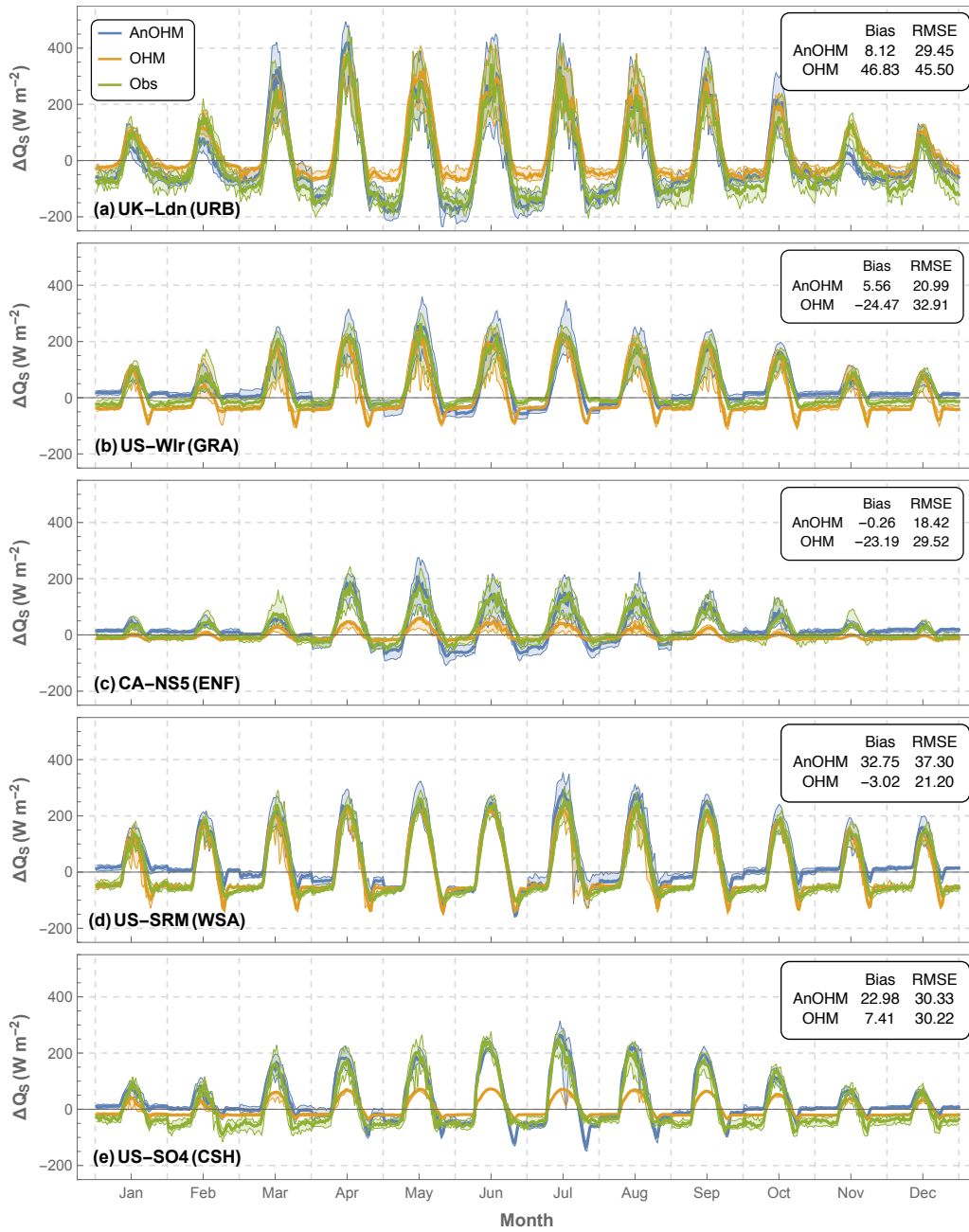


Figure 6 Monthly median (line) diurnal cycles and interquartile (shaded) values of ΔQ_s for AnOHM predictions (blue), OHM predictions (orange) and observations (green) at (a) UK-Ldn (URB), (b) US-Wlr (GRA), (c) CA-NS5 (ENF), (d) US-SRM (WSA), and (e) US-SO4 (CSH) (see Table 2 for site information). Statistics include average bias and RMSE (W m⁻²). The OHM coefficients a_1 , a_2 and a_3 used for different land covers are: 0.553, 0.303 and -37.6 at the urban site (UK-Ldn) (Ward et al., 2016), 0.32, 0.54 and -27.4 at the grass covered sites (US-Wlr and US-SRM) (Grimmond and Oke, 1999), and 0.11, 0.11 and -12.3 at the forest covered sites (CA-NS5 and US-SO4) (Grimmond and Oke, 1999).

Table 1 Range of values used as basis for the sensitivity analysis (a) surface parameters and (b) meteorological variables. All are assumed to have normal PDF. Values of surface parameters are based on values reported in Stull (Stull, 1988).

Parameter/Variable		Unit	Min	Max	Mean	Standard deviation
(a) Surface						
thermal conductivity	k	$\text{W m}^{-1} \text{K}^{-1}$	0	3	1.2	0.1
bulk material heat capacity	C_p	$\text{MJ m}^{-3} \text{K}^{-1}$	0	4	2.0	0.04
albedo	α	--	0	1	0.27	0.07
emissivity	ε	--	0.8	1.0	0.93	0.025
midday* mean Bowen ratio (inverse)	β^{-1}	--	0	20	0.05	0.05
bulk transfer coefficient	C_h	$\text{J m}^{-3} \text{K}^{-1}$	0	8	4	0.5
(b) Hydrometeorological						
Amplitude or range of the daily incoming shortwave radiation	A_K	W m^{-2}	0	1200	800	200
Mean daytime incoming shortwave radiation	$\overline{K_d}$	W m^{-2}	0	500	200	50
Amplitude or range of the daily air temperature	A_T	$^{\circ}\text{C}$	0	15	8	2
Mean daily air temperature	$\overline{T_a}$	$^{\circ}\text{C}$	0	40	30	7.5
Phase lag between radiation and air temperature	τ	rad	0	$\pi/2$	$\pi/4$	$\pi/10$
Mean daytime wind speed	U	m s^{-1}	0	4	2	0.5
Mean daily water flux density	W	$10^{-7} \text{m}^3 \text{s}^{-1} \text{m}^{-2}$	0	100	10	5

* midday period: 1000–1400 local standard time.

5 Table 2 Characteristics of the flux towers at the study sites.

Site	UK-Ldn	US-Wlr	CA-NS5	US-SRM	US-SO4
Location	51.50° N 0.12° W	37.52° N 96.86° W	55.86° N 98.49° W	31.82° N 110.87° W	33.38° N 116.64° W
Land cover classification	Urban/ Built-up	Grassland	Evergreen Needleleaf Forest	Woody Savannas	Closed Shrublands
Land cover code	URB	GRA	ENF	WSA	CSH
Study Year	2011	2003	2004	2004	2005
Reference	Kotthaus and Grimmond (2014a; 2014b)	Klazura et al. (2006), Coulter et al. (2006)	Goulden et al. (2006)	Scott et al. (2009)	Luo et al. (2007)

Table 3 Surface properties used in AnOHM simulation for the study sites based on calibration. The values of α and β are monthly climatology from January to December and are used when observations are not available (see Table 1 for notation definition).

Parameter	Unit	Site				
		UK-Ldn	US-Wlr	CA-NS5	US-SRM	US-SO4
k	$\text{W m}^{-1} \text{K}^{-1}$	2.8	0.43	0.51	0.41	0.56
C_p	$\text{MJ m}^{-3} \text{K}^{-1}$	2.4	0.31	0.36	0.56	0.27
α	--	0.24,0.24, 0.22,0.20, 0.14,0.13, 0.12,0.14, 0.18,0.24, 0.24,0.18	0.29,0.29, 0.17,0.18, 0.18,0.12, 0.11,0.10, 0.19,0.13, 0.24,0.35	0.30,0.29, 0.22,0.15, 0.10,0.10, 0.10,0.11, 0.22,0.24, 0.28,0.30	0.13,0.17, 0.16,0.14, 0.13,0.12, 0.13,0.15, 0.14,0.19, 0.13,0.18	0.22,0.11, 0.11,0.10, 0.11,0.10, 0.10,0.10, 0.11,0.10, 0.17,0.24
ε	--	0.92	0.93	0.95	0.95	0.92
β	--	6.1, 5.1, 8.3, 7.9, 5.4, 3.9, 5.3, 4.2, 5.2, 4.3, 4.8, 3.2	2.9, 0.8, 7.6, 2.7, 0.3, 0.3, 0.3, 0.8, 0.5, 0.7, 2.3, 2.3	6.1, 6.0, 8.7, 8.0, 1.9, 1.6, 0.7, 0.7, 1.3, 1.4, 3.1, 8.0	1.9, 5.5, 3.3, 2.0, 10.1, 9.7, 2.0, 0.9, 3.0, 4.3, 10.0, 3.3	1.5, 1.4, 1.9, 3.0, 1.4, 1.4, 2.1, 1.2, 2.8, 1.9, 2.1, 4.1
C_h	$\text{J m}^{-3} \text{K}^{-1}$	4.3	1.9	5.1	3.6	3.9