The format of this reply to Referee 1 is as follows:

In the first part, "Authors' Response to Anonymous Referee 1", we provide a point-by-point response to the referee's comments. We provide each of the referee's comments in **bold font**. After each comment that requires a response we provide a response (in regular font). If we modified the manuscript in response to a comment, we describe what the modification was, and indicate where it was made in the revised manuscript (the revised manuscript is provided at the end of this document).

In the second section, Additional modifications to the manuscript, we describe modifications to the manuscript not made in response to any specific comment of either referee. For the most part these modifications are minor, however we did fix two errors in the final analytical solution (errors in the text, not the code; so these did not affect any of the presented results).

In the third section, we provide the revised manuscript.

Authors' Response to Anonymous Referee 1

The authors derive an analytical solution of the two-dimensional steady Boussinesq equations in the limit of zero Reynolds number. To obtain this solution they transform the equations into a sixth-order equation for the streamfunction. Then they seek a solution in the form of a single-harmonic, A(z)cos(kx). Then the general solution is found as an infinite summation of single-harmonic solutions. The procedure is rigorous and well-described. In the second part of the paper the authors proceed to use this analytical solutions to test different boundary conditions for the pressure Poisson equation in a particular numerical implementation. They consider two cases in which they compare numerical results for homogeneous and inhomogeneous types boundary conditions for the pressure Poisson equation to the analytical solution. Only the inhomogeneous boundary condition passes both tests. In summary, this is an interesting paper that deserves publication. I have several minor comments:

1. Section 2.1: Please explain the physical meaning of equation (3), i.e. that the equation for b is based on the transport equation for the temperature.

Equation (2.3) is the thermal energy equation (differential form of the first law of thermodynamics). In the revised manuscript we now briefly discuss the governing equations, including (2.3) in the paragraph right after (2.4). We also give a reference to where the equations are described more fully (Chandrasekhar 1961). Further down in the paragraph we give a reference (Kundu 1990) for the Brunt-Väisälä frequency *N*, a parameter that appears in (2.3).

2. Please specify which condition has been substituted in which equation in order to obtain Eqs. (24-25).

Right before (2.24) we now state that: "In view of (2.7) and (2.17), the impermeability condition w(x,0) = 0 and no-slip condition u(x,0) = 0 yield"

3. Section 2.3: "The derivation of the u field requires considerable effort and is not pursued." I do not understand this. As far as I understand, the analytical solution of u is

simply the analytical partial derivative of Eq. (39) with respect to x, which can be explicitly written as an infinite sum. Do I miss anything?

The referee is not quite right. As can be seen from (7), one gets u by differentiating the streamfunction (specified in (39)) with respect to z, not x [One gets w by differentiating the streamfunction with respect to x, and that's a relatively easy calculation]. Moreover, as can be seen from (31), z is implicitly inside four of the factors in (39). This is why obtaining u is potentially complicated.

However, we decided to go ahead and take the z-derivative of (39) and see if we could obtain a relatively compact final form for u. It turns out that by making use of addition formulas for sines, the results do simplify considerably and the final form of u is not too bad. So, we now present the analytical solution for u: (2.32) for the single-harmonic wave and (2.41) for the square wave. The differences between the analytical u and the u obtained by finite differencing the streamfunction are visually imperceptible for tests A-1 and A-2. The quantitative changes are generally less than 1 % but are on the order of 1 % near the surface. Our reported values for the R ratios do not change for A-1 but there is a minor change in one of the R ratios for A-2: R_b was originally $R_b \cong 3.6 \times 10^{-3}$ but is now $R_b \cong 3.8 \times 10^{-3}$. The refined value is given in the revised manuscript. With the analytical solution for u now provided, there is now no reason to discuss obtaining the u field by finite-differencing the streamfunction. We have omitted such prose from the revised manuscript.

4. Section 3: "The surface condition on pressure is the inhomogeneous Neumann condition that arises from projecting the vertical equation of motion into the vertical, and imposing the impermeability condition (Vreman, 2014; also see our Appendix)." The sentence can be maintained, but a sentence should be added that, in addition, it is important that the discretized Poisson equation somehow incorporates the condition that delta = div u=0 on the wall or in the direct vicinity of the wall. This was also stressed by Vreman (and others) and is briefly mentioned in the appendix. It is good to include this requirement also in the main text. In the method of the authors the condition delta=0 near the wall is probably implicitly enforced via the alternative Poisson equation, specified in the Appendix, Eq. (A3b). After Eq. (A3b) the authors cite the pressure Poisson equation parodox using a sentence of Gresho and Sani. Please mention there that Vreman has revisited this paradox

and has shown that, at least for the standard staggered method, the discretized version of (A3b) (with appropriate Neumann condition) is equivalent to the discretized version of (A3a) supplemented with the condition that div(Laplacian(u))=0 in the direct vicinity of the wall. Equipped with the latter the condition, the diffusion equation d delta/dt = nu * Laplacian(delta) leads to delta=0 for all time.

Yes, the divergence-free condition near the wall is enforced via the alternative Poisson equation (A3b). We have modified the sentence right after (A3b) to emphasize that this alternative Poisson equation assures that (A2) (the divergence-free condition) is satisfied.

Since we want section 3 to focus on the verification tests, we did not want to divert too much attention in that section to the technical details of the pressure equation, pressure boundary condition, and related topics. However, we agree that these details are certainly important and should be clearly discussed. So, as a compromise, we modified a sentence in section 3 to read: "The pressure is diagnosed from a Poisson equation (equation (3b), discussed in the Appendix),..." Then, in the Appendix we modified the sentence about (A3b) enforcing the divergence-free condition (the modification described in the paragraph above) and added two sentences to the discussion of (A3b) along the lines suggested by the referee. Then, after those two new sentences, we end the paragraph with: "We note that (A3b) is the form adopted in our numerical code."

5. Section 3: The numerical solutions are obtained on an un-staggered grid. Please explain what was done to prevent odd-even decoupling of the pressure. Was the Rhie-Chow interpolation method used, for example?

The numerical solutions are obtained on a staggered (Arakawa C) grid so no decoupling of the pressure was occuring/noticed. We now mention "staggered (Arakawa C) grid" in the discussion of the DNS code in section 3.

6. Section 3: Please explain the meaning of the abbreviations HNC and INC (I guess homogeneous Neumann condition and inhomogeneous Neumann condition).

Yes, HNC is out acronym for homogeneous Neumann condition and INC is our acronym for inhomogeneous Neumann condition. We now introduce these acronyms near the end of the second paragraph of section 3.

Additional modifications to the manuscript (i.e., not made in response to the reviewers' comments)

Please note that the third author would like his middle initial "A" included in his name: Jeremy A. Gibbs.

We have added a new reference: Egger (1981). The Egger study was related to ours in that it was concerned with a linear analysis of the 2D Boussinesq governing equations for thermally driven flow. Egger's analysis was largely for slope flows, though with flat terrain (our focus) considered as a special case. However, Egger outlines how to get the analytical solution but does not actually provide the final analytical solution. We mention this Egger study in the second paragraph of Section 1. We also mention it in the paragraph right after (2.8): the restriction on acceptable surface buoyancies described in that paragraph was first noted by Egger, though without details.

A correction was made to the original equations (2.39) and (2.40) [these now appear as equations (2.40) and (2.42), respectively]. The factor $k^{1/3}$ in the denominator of the term in front of the summation in (2.39) and the factor $k^{2/3}$ in the numerator of the term in front of the summation in (2.40) should be kept inside the summations. These factors were treated correctly in the computer code, so none of the presented results were affected.

Section 3. We now make the number of points in the x and z direction unambiguous: instead of writing the number of points in test A-1 as (513, 1025) we write, "...consisted of 513 points in the x direction and 1025 points in the z direction,..." Similarly, for test A-2, we now write, "...was generated with 2049 points in the x direction and 513 points in the z direction,..."

In the Appendix we now write the time step as Δt instead of δt since the symbol δ has already been used to represent the divergence of the velocity field.

In several places in the manuscript we now use bold to indicate the vector $\mathbf{u}$ (formerly we used $\vec{u}$).

We have slightly modified the acknowledgements statement (we now thank the anonymous reviewer).

An analytical verification test for numerically simulated convective flow above a thermally heterogeneous surface

by Alan Shapiro, Evgeni Fedorovich, and Jeremy A. Gibbs

Abstract. An analytical solution of the Boussinesq equations for the motion of a viscous stably stratified fluid driven by a surface thermal forcing with large horizontal gradients (step changes) is obtained. This analytical solution is one of the few available for wall-bounded buoyancy-driven flows. The solution can be used to verify that computer codes for Boussinesq fluid system simulations are free of errors in formulation of wall boundary conditions and to evaluate the relative performances of competing numerical algorithms. Because the solution pertains to flows driven by a surface thermal forcing, one of its main applications may be for testing the no-slip, impermeable wall boundary conditions for the pressure Poisson equation. Examples of such tests are presented.

1 Introduction

Thermal disturbances associated with variations in underlying surface properties can drive local circulations in the atmospheric boundary layer (Atkinson, 1981; Briggs, 1988; Hadfield et al., 1991; Segal and Arritt, 1992; Simpson, 1994; Mahrt et al., 1994; Pielke, 2001; McPherson, 2007; Kang et al., 2012) and affect the development of the convective boundary layer (Patton et al., 2005; van Heerwaarden et al., 2014). Computational fluid dynamics (CFD) codes for modeling such flows commonly solve the Boussinesq equations of motion and thermal energy for a viscous/diffusive stably stratified fluid. In this paper we present an analytical solution of the Boussinesq equations for flows driven by a surface thermal forcing with large gradients (step changes) in the horizontal. The solution can be used to verify that CFD codes for Boussinesq fluid system simulations are free of errors, and to evaluate the relative performances of competing numerical algorithms. Such verification procedures are important in the development of CFD models designed for research, operational, and classroom applications.

We solve the linearized Navier-Stokes and thermal energy equations analytically for the case where the surface buoyancy varies laterally as a square wave (Fig. 1). Attention is restricted to the steady state. No boundary-layer approximations are made; the solution is non-hydrostatic, and both horizontal and vertical derivatives are included in the viscous stress and thermal diffusion terms. The solution is similar to that of Axelsen et al. (2010) for katabatic flow above a cold strip, but is easier to evaluate (no slope present) and applies to the more general scenario where the viscosity and diffusivity coefficients can differ. The flow is also similar to a special case (no slope) considered by Egger (1981), although a final analytical solution was not provided in that study. Strictly speaking, the linearized Navier-Stokes equations apply to a class of very low Reynolds number motions known as creeping flows. Such flows appear in studies of lubrication, locomotion of microorganisms, lava flow, and flow in porous media. Of course, for the task at hand, if our linear solution is to serve as a benchmark for a nonlinear numerical model solution, it is essential that the parameter space be restricted to values for which the model's nonlinear terms are negligible.

Because the solution pertains to flows driven by a surface thermal forcing, one of its main applications may be as a test for surface boundary conditions in the pressure Poisson equation. In models of atmospheric boundary layer flows, the buoyancy is a major contributor to the forcing term in the Poisson equation and also appears in the associated surface boundary condition. The pressure boundary condition on a solid boundary in incompressible (Boussinesq) fluid flows is an important and complex issue that has long been fraught with technical difficulties and controversies (Strikwerda, 1984; Orszag et al., 1986; Gresho and Sani, 1987; Gresho, 1990; Temam, 1991; Henshaw, 1994; Petersson, 2001; Sani et al., 2006; Rempfer, 2006; Guermond et al., 2006; Nordström et al., 2007; Shirokoff and Rosales, 2011; Hosseini and Feng, 2011; Vreman, 2014). Typical fractional-step solution methodologies and associated pressure (or pseudo-pressure) boundary-condition implementations are often verified using various prototypic flows such as Poiseuille flows, lid-driven cavity flows, flows over cylinders or bluff bodies, viscously decaying vortices, and dam-break flows. We are unaware of verification tests in which flows were driven by a heterogeneous surface buoyancy forcing. Our solution is designed to fill this gap.

The analytical solution is derived in Sect. 2. In Sect. 3, this solution is compared to numerically simulated fields in a steady state. Two versions of a numerical code are run: a version in which the correct surface pressure boundary condition is applied, and a version in which the pressure condition is mis-specified. A summary follows in Sect. 4.

2 Analytical solution

We derive the solution for steady flow over an underlying surface along which the buoyancy varies laterally as a single harmonic function. This single-harmonic solution is then used as a building block in a Fourier representation of the square-wave solution.

2.1 Governing equations

Consider the flow of a viscous stably stratified fluid that fills the semi-infinite domain above a solid horizontal surface (placed at $z = 0$). This surface undergoes a steady thermal forcing that varies periodically in the right-hand Cartesian x direction, but is independent of the y direction. The two-dimensional (x, z) flow is periodic in x, and satisfies the linearized (assuming the disturbance is of small amplitude) governing equations under the Boussinesq approximation,

$$0 = -\frac{\partial \Pi}{\partial x} + \nu \nabla^2 u \,, \tag{2.1}$$

$$0 = -\frac{\partial \Pi}{\partial z} + b + \nu \nabla^2 w \,, \tag{2.2}$$

$$0 = -N^2 w + \alpha \nabla^2 b \,, \tag{2.3}$$

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0 \,. \tag{2.4}$$

Apart from notational differences, (2.1)–(2.4) are the two-dimensional steady state versions of (55)–(57) of Sect. II of Chandrasekhar (1961). Equations (2.1) and (2.2) are the horizontal (x) and vertical (z) equations of motion, respectively, (2.3) is the thermal energy equation (differential form of the first law of thermodynamics) expressed in terms of the buoyancy variable (defined below), and (2.4) is the incompressibility condition. Here u and w are the horizontal and vertical velocity components, $\Pi \equiv [p - p_e(z)]/\rho_w$ is the kinematic pressure perturbation [p is pressure, $p_e(z)$ is pressure in a hydrostatic environmental state in which the density profile is $\rho_e(z)$, ρ_w is a constant reference density, say, $\rho_e(0)$], and $b \equiv -g[\rho - \rho_e(z)]/\rho_w$ is the buoyancy, where ρ is the actual density, and g is the acceleration due to gravity. The Brunt-Väisälä frequency $N \equiv \sqrt{-(g/\rho_w) d\rho_e/dz}$ of the ambient fluid (Kundu 1990), kinematic viscosity ν, and thermal diffusivity α are taken constant.

We obtain our solution using a standard vorticity/streamfunction formulation. Cross-differentiating (2.1) and (2.2) yields the vorticity equation,

$$0 = -\frac{\partial b}{\partial x} + \nu \nabla^2 \eta \,, \qquad (2.5)$$

where $\eta \equiv \partial u/\partial z - \partial w/\partial x$ is the vorticity. Eliminating b from (2.3) and (2.5) yields

$$\nabla^4 \eta = \frac{N^2}{\nu\alpha}\frac{\partial w}{\partial x} \,. \qquad (2.6)$$

Introducing a streamfunction ψ defined through

$$u = \partial\psi/\partial z, \qquad w = -\partial\psi/\partial x \,, \qquad (2.7)$$

guarantees that (2.4) is satisfied, and transforms (2.6) into a single equation for ψ,

$$\nabla^6 \psi + \frac{N^2}{\nu\alpha}\frac{\partial^2 \psi}{\partial x^2} = 0 \,. \qquad (2.8)$$

The dependent variables are assumed to vanish far above the surface ($z \to \infty$). On the surface we apply no-slip ($u = 0$) and impermeability ($w = 0$) conditions, and specify a periodic (in x) buoyancy distribution. As we will now see, restricting the dependent variables to steady periodic forms that vanish as $z \to \infty$ also restricts acceptable distributions of the surface buoyancy. The restriction was first noted by Egger (1981, Sect. 3c), though without details. Averaging (2.3) over one period (using $w = -\partial\psi/\partial x$) yields $d^2\overline{b}/dz^2 = 0$, which integrates to $\overline{b} = A + Bz$ ($\overline{b}$ is the average of b; A and B are constants). Taking $b \to 0$ as $z \to \infty$, implies that $\overline{b} \to 0$ as $z \to \infty$, in which case $A = B = 0$, and $\overline{b}(z) = 0$. In particular, at the surface, $\overline{b}(0) = 0$. If a surface distribution

$b(x,0)$ violates this condition, the ground acts as a net heat source/sink. In an unsteady model, such a source/sink would force a continually upward-developing disturbance, and a steady state could never be attained.

2.2 Single-harmonic forcing

For a surface buoyancy of the form $b(x,0) \propto \sin kx$, (2.3) indicates that ψ is of the form

$$\psi = A(z)\cos kx \,. \tag{2.9}$$

Application of (2.9) in (2.8) yields

$$\left(\frac{d^2}{dz^2} - k^2\right)^3 A - \frac{N^2k^2}{\nu\alpha} A = 0\,, \tag{2.10}$$

which has solutions of the form $A \propto e^{Mz}$ for M satisfying

$$(M^2 - k^2)^3 = \frac{N^2k^2}{\nu\alpha}\,. \tag{2.11}$$

Taking the one-third power of (2.11) yields a useful intermediate result:

$$M^2 - k^2 = \frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}} e^{2n\pi i/3}\,, \tag{2.12}$$

where n is an integer. Rearranging (2.12) and taking the square root yields

$$M = \pm\sqrt{k^2 + \frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}} e^{2n\pi i/3}}\,. \tag{2.13}$$

Equation (2.13) furnishes six roots, two for each of $n = 0$, 1, 2. To ensure that $A(z) \to 0$ as $z \to \infty$, we reject the roots with a positive real part. With the radicand of (2.13) expressed in polar form, the physically acceptable roots are

$$M_0 = -\sqrt{k^2 + \frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}}}, \quad (n = 0), \tag{2.14a}$$

$$M_1 = -r^{1/2}e^{i\phi/2}, \quad (n = 1), \tag{2.14b}$$

$$M_2 = -r^{1/2}e^{-i\phi/2}, \quad (n = 2), \tag{2.14c}$$

where the subscript on M denotes the associated value of n, and r and ϕ are defined by

$$r \equiv \sqrt{\left[k^2 + \frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}}\cos\left(\frac{2\pi}{3}\right)\right]^2 + \left[\frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}}\sin\left(\frac{2\pi}{3}\right)\right]^2}, \tag{2.15}$$

$$\cos\phi = \frac{1}{r}\left[k^2 + \frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}}\cos\left(\frac{2\pi}{3}\right)\right], \qquad \sin\phi = \frac{1}{r}\left(\frac{N^{2/3}k^{2/3}}{\nu^{1/3}\alpha^{1/3}}\right)\sin\left(\frac{2\pi}{3}\right) > 0. \tag{2.16}$$

While solving (2.16) for ϕ, care must be taken when evaluating arcsin or arccos functions that ϕ appears in the correct quadrant (ϕ should be in quadrant I or II so $\phi/2$ should always be in quadrant I). Also note from (2.14b) and (2.14c) that M_2 is the complex conjugate of M_1 ($M_2 = M_1{}^*$), a fact that will often be used below.

With the general solution for ψ written as

$$\psi = (B\,e^{M_0 z} + Ce^{M_1 z} + De^{M_2 z})\cos kx, \tag{2.17}$$

where B, C, and D are constants, the vorticity becomes,

$$\eta = \left[B(M_0^2 - k^2)\,e^{M_0 z} + C(M_1^2 - k^2)e^{M_1 z} + D(M_2^2 - k^2)e^{M_2 z}\right]\cos kx, \tag{2.18}$$

and the buoyancy follows from (2.3) as

$$b = \frac{kN^2}{\alpha}\left(\frac{B}{M_0^2 - k^2}e^{M_0 z} + \frac{C}{M_1^2 - k^2}e^{M_1 z} + \frac{D}{M_2^2 - k^2}e^{M_2 z}\right)\sin kx + b_h\,, \tag{2.19}$$

where $\nabla^2 b_h = 0$. In view of (2.12), equation (2.19) becomes

$$b = \frac{k^{1/3}\nu^{1/3}N^{4/3}}{\alpha^{2/3}}(B\,e^{M_0 z} + e^{-2\pi i/3}\,C\,e^{M_1 z} + e^{-4\pi i/3}\,D\,e^{M_2 z})\sin kx + b_h\,. \tag{2.20}$$

Applying (2.18) and (2.20) in (2.5) yields an equation for $\partial b_h/\partial x$, which upon use of (2.12) and $M_2 = M_1{}^*$ reduces to $\partial b_h/\partial x = 0$. So b_h is, at most, a function of *z*. Since $\nabla^2 b_h = 0$, b_h is, at most, a linear function of *z*, and since b should vanish as $z \to \infty$, that linear function must be 0. Thus, $b_h = 0$.

The pressure follows from (2.1) and (2.12) as

$$\Pi = \frac{\nu^{2/3}N^{2/3}}{k^{1/3}\alpha^{1/3}}\,(BM_0\,e^{M_0 z} + CM_1 e^{2\pi i/3}e^{M_1 z} + DM_2 e^{4\pi i/3}e^{M_2 z})\sin kx + G(z)\,, \tag{2.21}$$

where $G(z)$ is a function of integration. Applying (2.21) in (2.2), and using (2.11) yields $dG/dz = 0$, so G is constant. For Π to vanish as $z \to \infty$, this constant must be zero.

The surface conditions determine *B*, *C*, and *D*. The surface buoyancy is

$$b(x,0) = b_0 \sin kx\,, \tag{2.22}$$

where b_0 is a constant forcing amplitude. Application of (2.20) in (2.22) yields

$$B + e^{-2\pi i/3}\,C + e^{-4\pi i/3}\,D = \frac{b_0\alpha^{2/3}}{k^{1/3}\nu^{1/3}N^{4/3}}\,. \tag{2.23}$$

In view of (2.7) and (2.17), the impermeability condition $w(x,0) = 0$ and no-slip condition $u(x,0)=0$ yield

$$B+C+D=0\,, \tag{2.24}$$

$$BM_0+CM_1+DM_2=0\,. \tag{2.25}$$

Straightforward but lengthy manipulations yield the solution of (2.23)–(2.25):

$$B=-\left(\frac{b_0\,\alpha^{2/3}}{\sqrt{3}\,k^{1/3}\nu^{1/3}N^{4/3}}\right)\frac{2r^{1/2}\sin(\phi/2)}{M_0+2r^{1/2}\cos(\pi/3+\phi/2)}\,, \tag{2.26}$$

$$C=-i\left(\frac{b_0\,\alpha^{2/3}}{\sqrt{3}\,k^{1/3}\nu^{1/3}N^{4/3}}\right)\frac{M_2-M_0}{M_0+2r^{1/2}\cos(\pi/3+\phi/2)}\,, \tag{2.27}$$

$$D=i\left(\frac{b_0\,\alpha^{2/3}}{\sqrt{3}\,k^{1/3}\nu^{1/3}N^{4/3}}\right)\frac{M_1-M_0}{M_0+2r^{1/2}\cos(\pi/3+\phi/2)}\,. \tag{2.28}$$

Applying (2.26)–(2.28) in (2.17), (2.20), and (2.18), with (2.12) used in the latter equation, and noting that B is real, while $D=C^*$ (since $M_2=M_1^*$), we obtain

$$b=\frac{2b_0}{\sqrt{3}}\;\frac{e^{-Z_c}[\mu\cos(Z_s+\pi/6)+\cos(Z_s+\pi/6+\phi/2)]-e^{M_0 z}\sin(\phi/2)}{\mu+2\cos(\pi/3+\phi/2)}\sin kx\,, \tag{2.29}$$

$$\psi=\frac{2b_0\,\alpha^{2/3}}{\sqrt{3}\,k^{1/3}\nu^{1/3}N^{4/3}}\;\frac{e^{-Z_c}[\mu\sin Z_s+\sin(Z_s+\phi/2)]-e^{M_0 z}\sin(\phi/2)}{\mu+2\cos(\pi/3+\phi/2)}\cos kx\,, \tag{2.30}$$

where

$$Z_s\equiv z\,r^{1/2}\sin(\phi/2),\quad Z_c\equiv z\,r^{1/2}\cos(\phi/2),\quad \mu\equiv M_0/r^{1/2}\,. \tag{2.31}$$

Application of (2.30) in (2.7) yields the velocity components as

$$u = \frac{2b_0\, \alpha^{2/3} r^{1/2}}{\sqrt{3}\, k^{1/3} \nu^{1/3} N^{4/3}} \, \frac{e^{-Z_c}[\mu \sin(\phi/2 - Z_s) - \sin Z_s] - \mu e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \cos kx \qquad (2.32)$$

$$w = \frac{2b_0\, \alpha^{2/3} k^{2/3}}{\sqrt{3}\, \nu^{1/3} N^{4/3}} \, \frac{e^{-Z_c}[\mu \sin Z_s + \sin(Z_s + \phi/2)] - e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \sin kx \, . \qquad (2.33)$$

2.3 Piecewise constant (square wave) forcing

Next, consider the case where the surface buoyancy varies horizontally as a square wave, with a distribution over one period L given by

$$b(x,0) = \begin{cases} b_{\max}, & 0 < x < L/2, \\ -b_{\max}, & L/2 < x < L. \end{cases} \qquad (2.34)$$

Such a distribution can be expressed as the Fourier series:

$$b(x,0) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right), \qquad (2.35)$$

$$b_n = \frac{2}{L} \int_0^L b(x,0) \sin\left(\frac{n\pi x}{L}\right). \qquad (2.36)$$

Application of (2.34) in (2.36) yields

$$b_n = \frac{2 b_{\max}}{n\pi} \left[1 - 2\cos(n\pi/2) + \cos(n\pi)\right]. \qquad (2.37)$$

The solutions for b, ψ, u, and w can then be written as summations over the single-harmonic solutions (2.29), (2.30), (2.32), and (2.33), with k related to n by

$$k = \frac{n\pi}{L}, \qquad (2.38)$$

and with b_0 replaced by b_n:

$$b = \frac{2}{\sqrt{3}} \sum_{n=1}^{\infty} b_n \frac{e^{-Z_c}[\mu \cos(Z_s + \pi/6) + \cos(Z_s + \pi/6 + \phi/2)] - e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \sin\left(\frac{n\pi x}{L}\right), \qquad (2.39)$$

$$\psi = \frac{2\alpha^{2/3}}{\sqrt{3}\,\nu^{1/3} N^{4/3}} \sum_{n=1}^{\infty} \frac{b_n}{k^{1/3}} \frac{e^{-Z_c}[\mu \sin Z_s + \sin(Z_s + \phi/2)] - e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \cos\left(\frac{n\pi x}{L}\right), \qquad (2.40)$$

$$u = \frac{2\alpha^{2/3}}{\sqrt{3}\,\nu^{1/3} N^{4/3}} \sum_{n=1}^{\infty} b_n \frac{r^{1/2}}{k^{1/3}} \frac{e^{-Z_c}[\mu \sin(\phi/2 - Z_s) - \sin Z_s] - \mu e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \cos\left(\frac{n\pi x}{L}\right), \qquad (2.41)$$

$$w = \frac{2\alpha^{2/3}}{\sqrt{3}\,\nu^{1/3} N^{4/3}} \sum_{n=1}^{\infty} b_n k^{2/3} \frac{e^{-Z_c}[\mu \sin Z_s + \sin(Z_s + \phi/2)] - e^{M_0 z} \sin(\phi/2)}{\mu + 2\cos(\pi/3 + \phi/2)} \sin\left(\frac{n\pi x}{L}\right). \qquad (2.42)$$

3 Verification tests

A solution of the linearized equations may be used to verify a nonlinear code if the nonlinear terms are sufficiently small. Unfortunately, *a priori* estimates of such terms expressed, for example, through a Reynolds number, are not straightforward since the relevant velocity and length scales in our problem are only evident after a solution has been obtained. We thus seek an appropriate set of test parameters through trial and error, guided by *a posteriori* linear solution estimates of the terms $\mathbf{u}\cdot\nabla b$ and $\mathbf{u}\cdot\nabla\eta$ $[\mathbf{u} = (u, w)]$ present in nonlinear versions of (2.3) and (2.5), respectively. Specifically, for any computed candidate solution, we formed the ratios of the largest values of those nonlinear terms to the largest values of the corresponding linear terms, that is, the terms actually present in (2.3) and (2.5). We need only consider one such linear term per ratio since (2.3) and (2.5) are comprised of two terms of equal magnitude. A solution was deemed to be sufficiently linear if

$$R_\eta \equiv \frac{\max\left|\mathbf{u}\cdot\nabla\eta\right|}{\max\left|\partial b/\partial x\right|} < \varepsilon, \quad \text{and} \quad R_b \equiv \frac{\max\left|\mathbf{u}\cdot\nabla b\right|}{\max\left|\alpha\nabla^2 b\right|} < \varepsilon\,, \tag{3.1}$$

where ε ($<< 1$) is a prescribed threshold. The suitability of this approach was confirmed by the very close agreement between the analytical solutions and the numerical solutions obtained with the correct surface pressure condition.

The numerical model employed in our tests is a variant of a direct numerical simulation (DNS) code used in the boundary-layer and slope-flow studies of Fedorovich et al. (2001), Fedorovich and Shapiro (2009a,b), and Shapiro and Fedorovich (2013, 2014). The model solves the Boussinesq governing equations on a staggered (Arakawa C) grid. Although designed for three-dimensional simulations, the model was run in a two-dimensional (x, z) mode. The overall solution procedure is patterned on a fractional step method proposed by Chorin (1968). In our version, the prognostic equations are integrated using a filtered leapfrog scheme with explicit treatment of the viscous term. The pressure is diagnosed from a Poisson equation (equation (A3b), discussed in the Appendix), which is solved using a fast Fourier transform technique in horizontal planes, and a tridiagonal matrix inversion in the vertical. The surface condition on pressure is the inhomogeneous Neumann condition (INC) that arises from projecting the vertical equation of motion into the vertical, and imposing the impermeability condition (Vreman, 2014; also see the Appendix). We also run a version of the code in which the surface pressure condition is mis-specified as a homogeneous Neumann condition (HNC). We hasten to add, however, that our implementation of the HNC may be quite different from implementations described in the literature. We elaborate on these technical differences and review general aspects of the problem of surface pressure specification in the Appendix.

The analytical solution was evaluated on an un-staggered (x, z) grid extending over one period of the square wave ($x = 0$ to $x = L$). The series were truncated at 50000 terms. The governing parameters were adjusted so that the linearity criteria were satisfied in comparisons with $\varepsilon = 5\times10^{-3}$.

In the first test, we set $\nu = \alpha = 0.001\,\mathrm{m^2\,s^{-1}}$, $N = 0.02\,\mathrm{s^{-1}}$, $L = 5.12\,\mathrm{m}$, and $b_{\max} = 1\times10^{-5}\,\mathrm{m\,s^{-2}}$. For the analytical solution A-1, the (x, z) grid consisted of 513 points in the x direction and 1025 points in the z direction, with grid spacings $\Delta x = \Delta z = 0.01\,\mathrm{m}$. The linearity criteria (3.1) were satisfied with $R_\eta \cong 8.2\times10^{-5}$ and $R_b \cong 2.8\times10^{-3}$. The analytical b and w fields shown in Fig. 2 depict a broad zone of ascent above the warm surface and a compensating zone of descent over the cold surface, roughly for $z < 1.8\,\mathrm{m}$. In the upper part of these zones (at roughly $0.9\,\mathrm{m} < z < 1.8\,\mathrm{m}$), adiabatic expansion/compression has reversed the senses of the buoyancy fields. Surprisingly, the numerical fields in the inhomogeneous INC-1 and homogeneous HNC-1 cases are very similar to each other and to the A-1 fields. The u fields from A-1, INC-1, and HNC-1 shown in Fig. 3 are visually indistinguishable from one another.

To understand why the INC-1 and HNC-1 simulations are so similar, and to identify simulation parameters that might evince more substantial differences, we consider the idealized problem in which a specified buoyancy $b = b_0 e^{-\gamma z} \sin kx$ ($\gamma = h^{-1}$, where h is the e-folding depth scale) is the only forcing term in the Poisson equation $\nabla^2 \Pi = \partial b / \partial z$, with Neumann surface condition $\partial \Pi / \partial z \big|_0 = b(x,0)$. This idealized problem is solved as

$$\Pi^*_{\mathrm{INC}} = \frac{b_0}{\gamma^2 - k^2}\left(k\, e^{-kz} - \gamma\, e^{-\gamma z}\right) \sin kx \,. \tag{3.2}$$

The corresponding solution obtained with the homogeneous Neumann condition, $\partial \Pi / \partial z \big|_0 = 0$, is

$$\Pi^*_{\mathrm{HNC}} = \frac{b_0}{\gamma^2 - k^2}\left(\frac{\gamma^2}{k} e^{-kz} - \gamma\, e^{-\gamma z}\right) \sin kx \,. \tag{3.3}$$

The relative error (RE) in the vertical pressure gradient force associated with (3.2) and (3.3), defined as the local absolute error in that force divided by the local buoyancy, is calculated as

$$RE \equiv \left| \frac{\partial \Pi^*_{\mathrm{INC}} / \partial z - \partial \Pi^*_{\mathrm{HNC}} / \partial z}{b} \right| = e^{(a-1)kz}, \tag{3.4}$$

where $a \equiv \gamma / k$. Written in terms of the depth scale h and wavelength $\lambda = 2\pi / k$, a can be interpreted as an aspect ratio characterizing the width to depth scales of the disturbance, $a = \lambda / (2\pi h) \propto \lambda \gamma$. From (3.4) we see that RE decreases exponentially with z for disturbances characterized by small aspect ratios, $a < 1$ (which we refer to as deep disturbances) and increases exponentially with z for disturbances characterized by large aspect ratios, $a > 1$ (which we refer to as shallow disturbances). The buoyancy in Fig. 2 is suggestive of $a < 1$, which indicates that the first test could be classified as a deep (error-forgiving) simulation.

The preceding analysis suggests that simulations with shallow thermal disturbances ($a > 1$) might yield large differences between cases with inhomogeneous and homogeneous Neumann conditions. There did not appear to be a straightforward way to increase the effective a by systematically varying the parameters (e.g., increasing L tended to increase the effective h), but a set of suitable parameters were identified through trial and error and were used as the basis for the second test case.

In the second test, we set $\nu = \alpha =$ $0.0001\,\mathrm{m^2 s^{-1}}$, $N = 0.2\,\mathrm{s^{-1}}$, $L = 10.24\,\mathrm{m}$, and $b_{\max} = 5 \times 10^{-6}\,\mathrm{ms^{-2}}$. The analytical solution A-2 was generated with 2049 points in the x direction and 513 points in the z direction, with grid spacings of $\Delta x = \Delta z = 0.005\,\mathrm{m}$. The linearity criteria were satisfied with $R_\eta \cong 4.8 \times 10^{-5}$ and $R_b \cong 3.8 \times 10^{-3}$. In contrast to the counter-rotating convection rolls seen in the first test, the analytical b and w fields shown in Fig. 4 depict narrow updraft/downdraft pairs straddling the buoyancy discontinuities. Between the narrow updrafts is a broad region of relatively weak ascent. The w and b fields above the cold surface are mirror images of the fields above the warm surface. Note the change in the scales of the x and (especially) the z axes between Figs. 4 and 2: the low-level thermal disturbance in the second test is much shallower than the disturbance in the first test (and is suggestive of $a > 1$). In this second test case we find dramatic differences between the inhomogeneous INC-2 and homogeneous HNC-2 cases. Specifically, while the INC-2 and A-2 fields are in excellent agreement, the HNC-2 fields showed no signs of even approaching a steady state. Long after the INC-2 simulation had reached a steady state, the HNC-2 fields continued to amplify and develop asymmetric structures associated with flow nonlinearities. The very close agreement between the A-2 solution and the steady state in the INC-2 simulation is shown for the u field in Fig. 5. The u field in the disastrous HNC-2 simulation, at a time when a steady state had already been attained in the INC-2 simulation, is shown in Fig. 6.

4 Summary

The linearized Boussinesq equations for the motion of a viscous stably stratified fluid are solved analytically for a surface buoyancy that varies laterally as a square wave.

The solution describes two-dimensional laminar convective structures such as thermal convective rolls and updraft/downdraft pairs. The main applications of the solution may be in code verification and the evaluation of different implementations of the surface pressure condition for the pressure Poisson equation. Tests have been conducted for cases where the aspect ratios of the thermal disturbance have been large and small. With attention restricted to disturbances of sufficiently small amplitude, the linear solution and numerically simulated fields with the inhomogeneous Neumann condition for pressure (which is appropriate in the context of the particular fractional step procedure adopted in our DNS code) have been found to be in excellent agreement for both tests. However, in tests with a mis-specified Neumann condition, an excellent agreement with the analytical solution has been found only for the deep (small aspect ratio) disturbance case; errors in the shallow (large aspect ratio) disturbance case have been catastrophic.

Appendix A: Comment on the pressure condition at a lower solid surface

Consider a three-dimensional Boussinesq system with equation of motion,

$$\frac{\partial \mathbf{u}}{\partial t} = -\nabla\Pi + \nu\nabla^2\mathbf{u} + \mathbf{F}\,. \tag{A1}$$

Here $\mathbf{u} = (u, v, w)$ is the three-dimensional velocity vector, Π is a kinematic pressure perturbation, ν is the kinematic viscosity coefficient, and $\mathbf{F}$ is the sum of nonlinear acceleration and buoyancy terms. Applying the incompressibility condition,

$$\nabla\cdot\mathbf{u} = 0\,, \tag{A2}$$

in the equation that results from taking the divergence of (A1) (e.g., Orszag et al., 1986) yields the Poisson equation,

$$\nabla^2\Pi = \nabla\cdot\mathbf{F}\,. \tag{A3a}$$

Although (A1) and (A2) imply (A3a), the reverse statement is not generally true. Indeed, eliminating Π from between (A3a) and the equation arising from taking the divergence of (A1) yields the diffusion equation $\partial\delta/\partial t = \nu\nabla^2\delta$ for the velocity divergence $\delta \equiv \nabla\cdot\mathbf{u}$, whose solution is (A2) only if δ is zero initially and on all boundaries (Orszag et al., 1986; Gresho and Sani, 1987, Vreman 2014).

The same steps leading to (A3a) also lead to an alternative Poisson equation,

$$\nabla^2\Pi = \nabla\cdot\left(\nu\nabla^2\mathbf{u} + \mathbf{F}\right). \tag{A3b}$$

Although $\nabla\cdot\nu\nabla^2\mathbf{u}$ was omitted in (A3a) [this term is zero if (A2) is satisfied], without further constraints on δ (described above), (A2) may not be satisfied. Gresho and Sani

(1987) showed that the retention of $\nabla \cdot \nu \nabla^2 \mathbf{u}$ in (A3b) assures that (A2) is satisfied, and thus leads to the paradox: "If you include it, you don't need it; if you don't include it, you need it." Vreman (2014) revisited this paradox, and showed that for a standard staggered method, the discretized form of (A3b) is equivalent to that of (A3a) supplemented with the constraint that $\nabla \cdot \nabla^2 \mathbf{u} = 0$ $(\nabla^2 \delta = 0)$ on points adjacent to the solid boundary [with the same inhomogeneous Neumann boundary condition for Π implied for (A3a) and (A3b)]. When supplemented with this $\nabla^2 \delta = 0$ near-wall condition, the diffusion equation for δ led to $\delta = 0$ for all time. We note that (A3b) is the form adopted in our numerical code.

Evaluating the vertical component of (A1) on the surface, where the impermeability condition applies, yields the inhomogeneous Neumann condition,

$$\left.\frac{\partial \Pi}{\partial z}\right|_0 = \nu \left.\frac{\partial^2 w}{\partial z^2}\right|_0 + \left.F_z\right|_0, \tag{A4}$$

where $w \equiv \mathbf{k} \cdot \mathbf{u}$, $F_z \equiv \mathbf{k} \cdot \mathbf{F}$, $\mathbf{k}$ is the upward unit vector, and $(\ )|_0$ is a surface value. It has been argued that (A4), by itself, is not a proper boundary condition because it does not provide new information (it is not independent of the governing equations) and does not enforce the incompressibility condition (A2) at the boundary (Strikwerda, 1984; Henshaw, 1994; Sani et al., 2006). However, as pointed out by Henshaw (1994), many studies that impose (A4) (or a variant of it) also apply (A2) on the boundary.

In our numerical model, (A1) is integrated using a fractional step procedure with explicit treatment of the viscous term. First, a provisional velocity field $\tilde{\mathbf{u}}$ that does not satisfy (A2) is obtained by integrating a discretized form of (A1) in which the pressure gradient is omitted. The provisional velocity is equal to the velocity at the end of the previous time step plus the sum of the forcing terms (nonlinear acceleration, buoyancy, and viscous stress) multiplied by the time step Δt. With the forcing terms explicitly evaluated, $\tilde{\mathbf{u}}$ is readily computed throughout the flow domain, including on the surface, where, in surface-forced flows, the buoyancy will make a substantial contribution. In terms of $\tilde{\mathbf{u}}$ and its vertical component $\tilde{w}$, (A3b) and (A4) become,

$$\nabla^2 \Pi = \frac{\nabla \cdot \tilde{\mathbf{u}}}{\Delta t}, \tag{A5}$$

$$\left.\frac{\partial \Pi}{\partial z}\right|_0 - \frac{1}{\Delta t}\left.\tilde{w}\right|_0 = 0\,. \tag{A6}$$

In the second step, a velocity field that does satisfy (A2) is obtained by solving (A5) for Π and then adding the pressure gradient force associated with Π (multiplied by Δt) to $\tilde{\mathbf{u}}$.

In some explicit fractional step procedures (including the DNS code used in our study), the problem of solving (A5) subject to (A6) with $\left.\tilde{\mathbf{u}}\right|_0$ evaluated from model data is replaced by what appears to be an entirely different (but is actually equivalent) problem: solving (A5) subject to the homogeneous Neumann condition,

$$\left.\frac{\partial \Pi}{\partial z}\right|_0 = 0, \tag{A7}$$

in concert with $\tilde{\mathbf{u}}\big|_0$ being set to 0, obviating the need to calculate $\tilde{\mathbf{u}}\big|_0$ from model data. It can be shown that $\tilde{w}\big|_0$ and the discretized form of $\partial\Pi/\partial z\big|_0$ appear in the discretized form of (A5) valid half a grid point above the physical surface as $\partial\Pi/\partial z\big|_0 - \tilde{w}\big|_0/\Delta t$, that is, in the same combination as they appear in (A6). Thus, setting $\tilde{w}\big|_0$ and $\partial\Pi/\partial z\big|_0$ to 0, is equivalent to implementing (A6) with the model-computed values of $\tilde{w}\big|_0$: the discretized form of (A5) near the surface is the same in either case. Moreover, on the C grid, setting the tangential components $\tilde{u}\big|_0$ and $\tilde{v}\big|_0$ to 0 only affects the values of $\tilde{u}$ and $\tilde{v}$ half a grid point beneath the physical boundary. These values do not appear in the discretized form of (A5) at any z-level, and thus have no bearing on the solution. In essence, the errors associated with the conflation of the two physically unjustifiable specifications (homogeneous Neumann condition for pressure, and $\tilde{\mathbf{u}}\big|_0 = 0$) cancel out.

The homogeneous Neumann condition for pressure can be the source of confusion if the context in which the condition is applied is not made clear: it would be a correct condition if $\tilde{\mathbf{u}}\big|_0$ is set to zero (per the equivalence described above), but it would be an incorrect condition if the explicit model-computed values of $\tilde{\mathbf{u}}\big|_0$ are used. In the experiments with the mis-specified condition described in Sect. 3, the homogeneous condition is imposed in the latter context. Unfortunately, in many numerical model descriptions, the nature of the surface pressure condition is left vague, for example, by not indicating whether a Neumann condition is homogeneous or inhomogeneous, or, if a homogeneous Neumann condition is indicated, not mentioning how $\tilde{\mathbf{u}}\big|_0$ is treated.

Finally, we note that in fractional step procedures that treat the viscous term implicitly (e.g., Kim and Moin, 1985; Gresho, 1990; Armfield and Street, 2002; Guermond et al., 2006, and many others), the homogeneous Neumann condition is often applied as a surface condition for a Poisson equation, but it is again different from our implementation described in Sect. 3. In the implicit treatments, the provisional velocity is obtained as the solution of a boundary value problem ($\tilde{\mathbf{u}}\big|_0$ should be specified; often it is set to 0) in which the relevant Poisson equation resembles (A5) but applies to a scalar function (sometimes called a pseudo-pressure) that is not the real pressure. Temam (1991) refers to this scalar as, "... a technical quantity, a mathematical auxiliary..." and advocates that it should not even be considered as an approximation of the pressure. Interestingly, in the context of implicit treatments, the homogeneous Neumann condition on the pseudo-pressure has sometimes been implicated as corrupting solution accuracy through the development of spurious numerical boundary layers adjacent to solid boundaries (Gresho, 1990; Guermond et al., 2006; Hosseini and Feng, 2011).

Code availability

The Fortran program used to generate output data files from the analytical solution is available as a supplement to this article. That program (square.f) is configured for test A-1, but can be easily adjusted to run test A-2 or other tests. Running square.f automatically generates an output file for each dependent variable (e.g., u.dat) as well as an output file (square.out) that summarizes the test parameters and gives the computed values of the linearity ratios R_η and R_b defined in (3.1).

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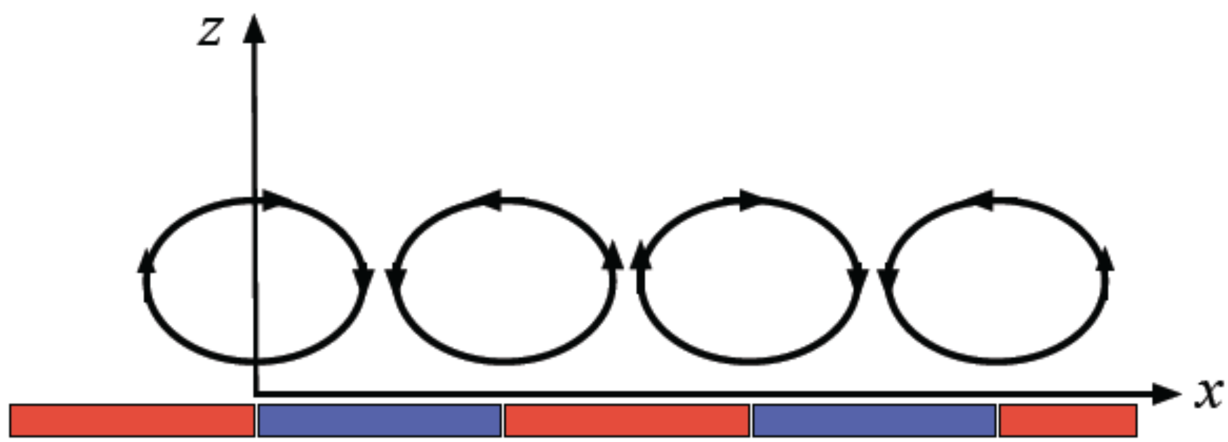

Figure 1. Schematic of two-dimensional (x, z) thermal convection induced by a surface buoyancy that varies horizontally (x) as a square wave. Red denotes positive surface buoyancy, blue denotes negative surface buoyancy.

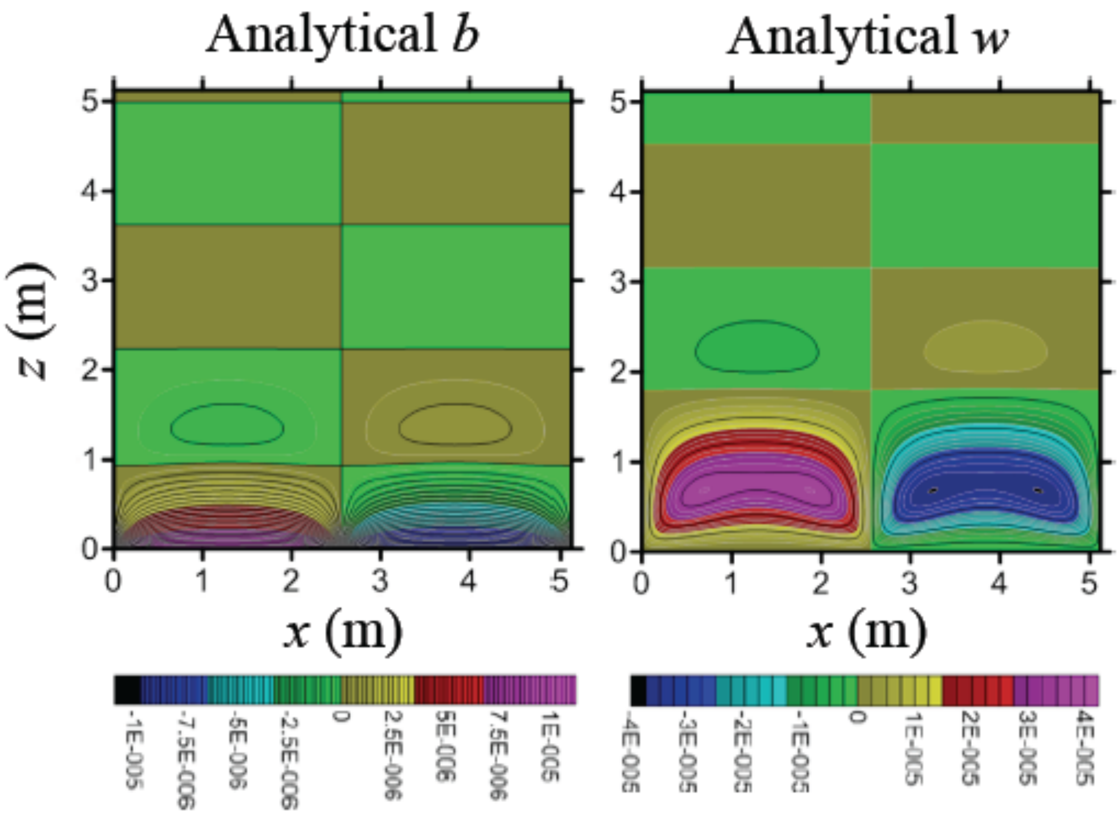

Figure 2. Vertical cross section of the analytical (A-1) buoyancy b and vertical velocity w fields from the first test case. Color bar units are $\mathrm{m\ s^{-2}}$ for b, and $\mathrm{m\ s^{-1}}$ for w.

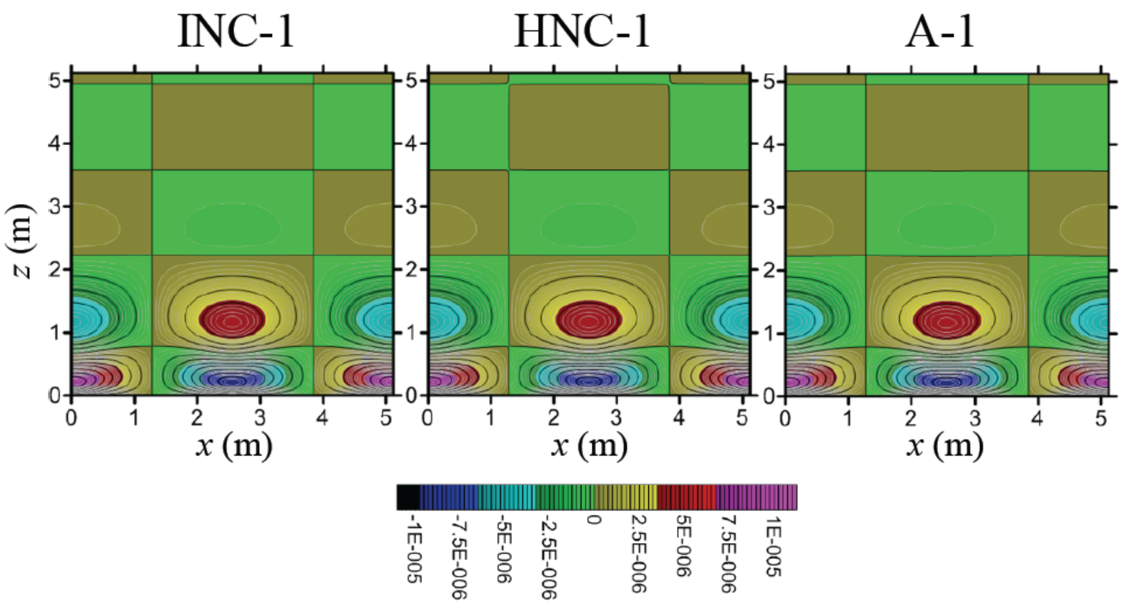

Figure 3. Vertical cross section of u from the first test case. A-1 is the analytical solution. INC-1 is the numerical simulation with inhomogeneous Neumann condition for pressure. HNC-1 is the numerical simulation with the homogeneous Neumann condition for pressure. Color bar units are m s^{-1}.

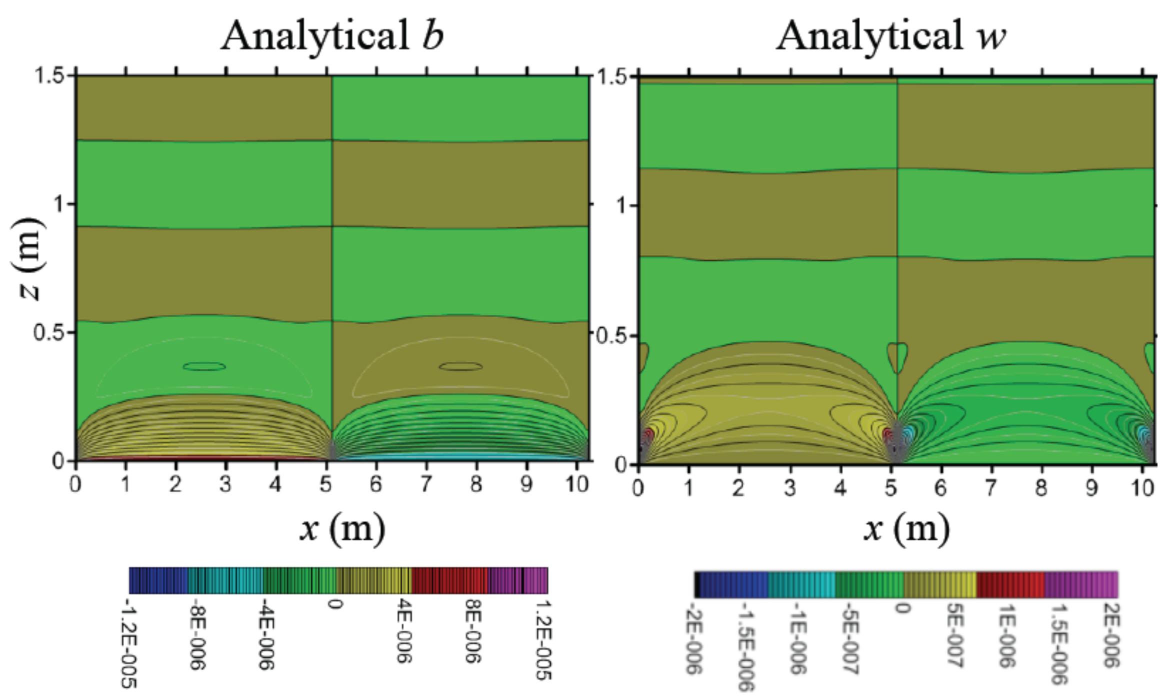

Figure 4. Vertical cross section of the analytical (A-2) buoyancy b and vertical velocity w fields from the second test case. Color bar units are $\mathrm{m\,s^{-2}}$ for b, and $\mathrm{m\,s^{-1}}$ for w.

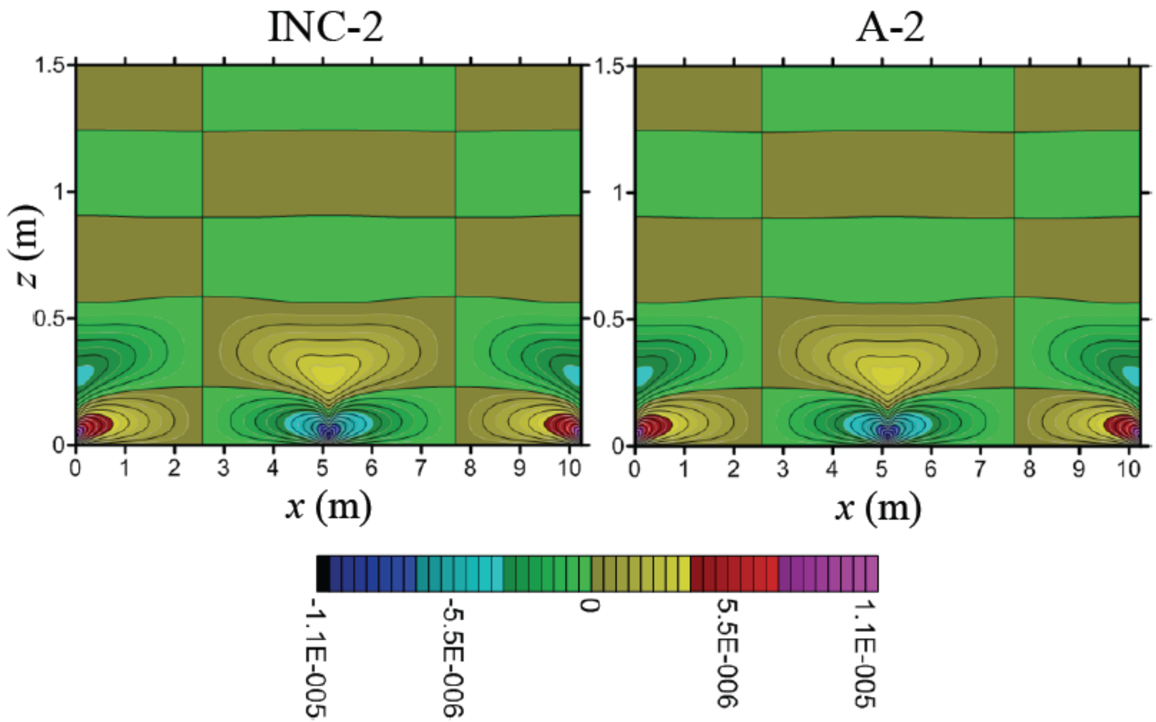

Figure 5. Vertical cross section of u from the second test case. A-2 is the analytical solution. INC-2 is the numerical simulation with inhomogeneous Neumann condition for pressure. Color bar units are $\mathrm{m\ s^{-1}}$.

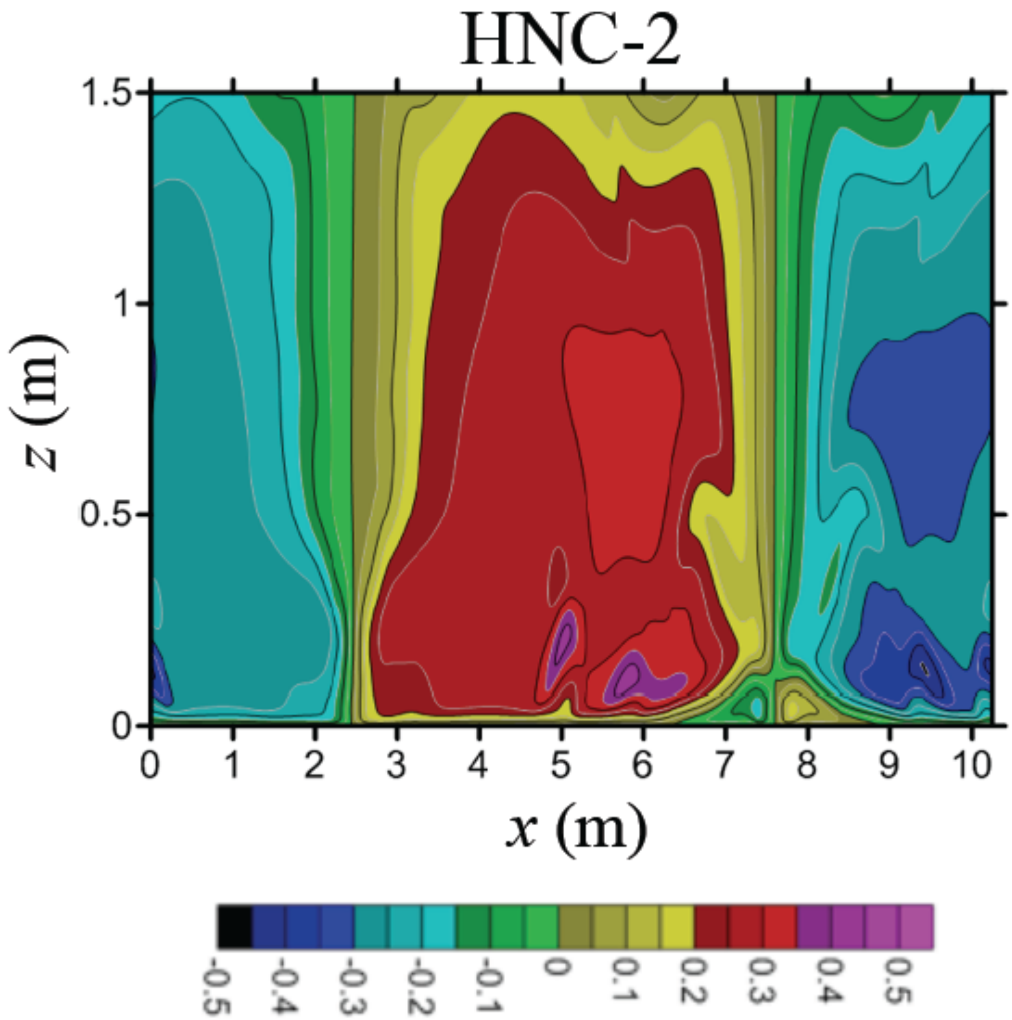

Figure 6. Vertical cross section of u from HNC-2, the numerical simulation with homogeneous Neumann condition for pressure in the second test case. Color bar units are m s^{-1}.