

# A Semi-Lagrangian advection scheme for radioactive tracers in the NCEP regional spectral model (RSM)

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## Abstract

In this study, the non-iteration dimensional-split semi-Lagrangian (NDSL) advection scheme is applied to the National Centers for Environmental Prediction (NCEP) regional spectral model (RSM) to alleviate the Gibbs phenomenon. The Gibbs phenomenon is a problem wherein negative values of positive-definite quantities (e.g., moisture and tracers) are generated by the spectral space transformation in a spectral model system. To solve this problem, the spectral prognostic specific humidity and radioactive tracer advection scheme is replaced by the NDSL advection scheme, which considers advection of tracers in a grid system without spectral space transformations. A regional version of the NDSL is developed in this study and is applied to the RSM. Idealized experiments show that the regional version of the NDSL is successful. The model runs for an actual case study suggest that the NDSL can successfully advect radioactive tracers (iodine-131 and cesium-137) without noise from the Gibbs phenomenon. The NDSL can also remove negative specific humidity values produced in spectral calculations without losing detailed features.

## 1 Introduction

The spectral method is a numerical method applied to primitive meteorological equations to achieve high-order accuracy (Robert, 1966). The spectral method has advantages over the finite difference method: the absence of truncation error in linear terms, the absence of aliasing and phase errors, the absence of pole problems in global models, and the efficient use and simple programming of semi-implicit time integration (Bourke, 1974). Applying the

1 spectral method in a regional model can be difficult due to time-dependent lateral boundary  
2 conditions. However, several approaches for solving this problem have been successful (e.g.,  
3 Tatsumi, 1986; Fulton and Schubert, 1987; Hoyer, 1987; Segami et al., 1989; Chen and Kuo,  
4 1992). Juang and Kanamitsu (1994) presented the National Centers for Environmental  
5 Prediction (NCEP) regional spectral model (RSM). The RSM has advantages in accuracy for  
6 a regional high-resolution domain. In addition, the spectral representation of the RSM is two-  
7 dimensional perturbation method, which can eliminate the error due to reevaluation of the  
8 linear forcing from the base fields by the regional model (Juang et al., 1997). This is one of  
9 the reasons that the RSM can be easily used for long-range climate simulations. The RSM has  
10 been widely used for regional downscaling research (e.g., Kang and Hong, 2008; Kanamitsu  
11 et al., 2010; Chang and Hong, 2011; Li et al., 2012). Kang and Hong (2008) assessed impact  
12 of the land surface parameters on the regional climate circulations. Kanamitsu et al. (2010)  
13 presented a refined spectral nudging technique for regional dynamical downscaling. Chang  
14 and Hong (2011) used the RSM to produce regional future scenarios by dynamical  
15 downscaling. Li et al. (2012) showed that the fully coupled RSM and regional ocean  
16 modeling system (ROMS) can produce detailed oceanic circulations over the California coast.

17 Although the RSM has advantages, it also has a problem common in spectral model systems:  
18 the Gibbs phenomenon. This phenomenon is “overshooting” in the convergence of the partial  
19 sums of particular Fourier series in the neighborhood of a discontinuity of the function being  
20 expanded. In the case of spectral techniques, the Gibbs phenomenon can introduce negative  
21 values in positive-definite variables (i.e., hydrometeors and tracers). Yoshimura (2011)  
22 presented a tracer simulation for the Fukushima Dai-ichi nuclear power plant accident during  
23 the massive earthquakes and tsunami on 11 March 2011 in eastern Japan by utilizing the  
24 stable isotope mode of the RSM (IsoRSM; Yoshimura et al., 2010). General features, e.g., the  
25 tracers from Fukushima reached the metropolitan area and significant amounts of tracers were  
26 precipitated, are well captured in this simulation. However, these simulations also clearly  
27 show that computational noise is produced by the Gibbs phenomenon near the emission point.  
28 This problem is quite severe in this simulation because radioactive materials are emitted from  
29 a single grid-point source at the surface level, which creates a significant discontinuity in the  
30 wave space transformation. Thus, an advection method that does not require a spectral  
31 transformation for tracers is needed to resolve the Gibbs phenomenon.

1 The semi-Lagrangian method is an alternative advection scheme that can replace the spectral  
2 calculation of hydrometeors and tracers. Semi-Lagrangian advection schemes have long been  
3 preferred in numerical weather prediction because they are more accurate and efficient than  
4 traditional Eulerian schemes when large time steps are considered (Williamson, 2007).  
5 Staniforth and Côté (1991) reviewed semi-Lagrangian literatures for atmospheric models.  
6 They concluded that the semi-Lagrangian framework facilitates the incorporation of shape-  
7 preserving and monotonic schemes for moisture advection, because of the relatively small  
8 dispersion errors in the presence of discontinuities or near discontinuities. Juang (2007, 2008)  
9 proposed a non-iteration dimensional-split semi-Lagrangian (NDSL) scheme, which is simple  
10 and economical compared with the conventional semi-Lagrangian method. Although the  
11 NDSL scheme has not yet been applied to the regional spectral model, some versions of the  
12 NDSL have been added to the global spectral model system, e.g., NCEP Global Forecast  
13 System (GFS; Moorthi et al., 2001), the Global/Regional Integrated Model System (GRIMs;  
14 Hong et al., 2013) Global Model Program (GMP). For regional model system, the flux  
15 through the boundaries is needed to apply the mass restoration, whereas there are no  
16 boundaries for global domains. Aranami et al. (2015) applied a mass restoration scheme for  
17 limited-area models (LAMs) with semi-Lagrangian advection. As such, the boundary  
18 treatment is required to apply the NDSL advection scheme for the RSM.

19 The objective of this study is to remove the Gibbs phenomenon for hydrometeors and tracers  
20 in the regional spectral model by replacing the spectral tracer advection scheme with the  
21 semi-Lagrangian advection scheme in a real simulation. Detailed features of the semi-  
22 Lagrangian version of the regional spectral model and the experimental design are described  
23 in Section 2. Section 3 provides results from the original IsoRSM and the semi-Lagrangian  
24 version of IsoRSM for radioactive tracers and humidity fields in the Fukushima nuclear  
25 power plant accident case study. A summary and conclusion are provided in Section 4.

26

## 27 **2 Method**

### 28 **2.1 Regional spectral model for radioactive tracers**

29 Yoshimura et al. (2010) presented the IsoRSM, which includes the isotopic species for water  
30 vapor ( $HDO$  and  $H_2^{18}O$ ) as tracers in the latest version of the Scripps Experimental Climate  
31 Prediction Center's regional spectral model (Kanamitsu et al., 2005). In the IsoRSM, the

1 tracers can be incorporated into raindrops or cloud particles at every integration time step.  
2 Therefore, the interactions between the tracers and precipitation processes can be considered.  
3 In contrast, a general chemical transport model uses precipitation as an external forcing.  
4 Yoshimura (2011) and Saya et al. (2013) modified IsoRSM to enable the simulation of the  
5 transport of radioactive tracers. Isotopic variables were replaced with radioactive tracers (i.e.,  
6  $^{131}\text{I}$  and  $^{137}\text{Cs}$ ), and dry and wet deposition processes caused by gravity and precipitation  
7 processes, respectively, were introduced. Saya et al. (2013) proposed a wet deposition process  
8 wherein the deposition is proportional to the ratio of the amount of condensed water to the  
9 total amount of water, whereas a traditional method (e.g., Maryon, 1991) considers only the  
10 amount of condensed water. In this study, the radioactive tracer mode of IsoRSM is used as  
11 an original framework.

12

## 13 **2.2 NDSL scheme for RSM**

14 In this study, the NDSL advection scheme replaces the spectral prognostic calculation of the  
15 tracers in the RSM. Once each tracer field is provided from the initial field on the regular grid  
16 space, the advection of these fields is calculated by the NDSL on the grid space without any  
17 spectral transformation during the model integration. This process prevents the Gibbs  
18 phenomenon. The NDSL has two characteristics: 1) non-iteration to compute the trajectories  
19 of each tracer; and 2) a dimensional-splitting method. Figure 1 shows the basic concept of  
20 two-dimensional advection for both the traditional semi-Lagrangian scheme and the NDSL  
21 scheme (from Juang 2008). The traditional backward (forward) semi-Lagrangian scheme  
22 assumes that the arrival (departure) points are located on the regular model grid (Fig. 1a). In  
23 this case, an initial guess and iterations to compute the trajectories are required, which means  
24 finding mid-point wind and transferring the fluid particles from the departure points to the  
25 arrival points. These iterations make the semi-Lagrangian scheme expensive and inefficient.  
26 The NDSL scheme is a central scheme, which assumes a mid-point wind at the regular model  
27 grid point at time  $t$  to find the departure point at time  $t - \Delta t$  and the arrival point at time  
28  $t + \Delta t$  (Fig. 1b, also see Fig. 1 of Zhang and Juang, 2012). No initial guess or iteration is  
29 needed to find trajectories. Only one interpolation at the departure point and one remapping at  
30 the arrival point are needed. The wind in the central scheme can reduce numerical error better  
31 than the upstream scheme because no estimation or iteration occurs. During the advection

1 process, the quantity of the tracer is assumed to be steady. The NDSL uses the dimensional-  
2 splitting method for 2D advection, which simply splits the 2D advection into a sequence of  
3 1D advections. Because it only needs 1D interpolation and remapping, the method easily  
4 attains mass conservation. Another important advantage is that the 1D method can be easily  
5 coded; thus, it is compatible with most numerical models.

6 For the implementation of the regional version of the NDSL in a real-data case, the boundary  
7 treatment must be considered. The global NDSL uses a cyclic boundary condition, which  
8 places a duplicated global domain at the western and eastern boundaries. This allows for the  
9 arrival or departure point to be located, even when it lies outside of the model domain.  
10 However, cyclic boundary conditions are not suitable for regional domains because the  
11 domain edge points differ in both longitude and latitude. Thus, the domain is treated as three  
12 sections: the boundary zone, the buffer zone, and inner domains. Figure 2 shows the structure  
13 of these sections in the lower-left corner of the model domain as an example; it assumes that  
14 the boundary and buffer zones are defined as three grid points each. The boundary zone (dark  
15 gray in Fig. 2a) is the area where the semi-Lagrangian advection calculation is not applied and  
16 where the values of the global base field are specified. This area is necessary to prevent the  
17 calculated departure point of tracers located outside of the model domain. The values in the  
18 inner domain are calculated entirely from the regional model. In the buffer zone (light gray in  
19 Fig. 2a), the global base field and the regional model field are combined, with the weighting  
20 determined by an inverse exponential function (Eq. 1 and Fig. 2b), to smooth the gap between  
21 the boundary zone and the inner domain.

$$22 \quad W_G = \frac{1}{e^k} \quad (1)$$
$$W_R = 1 - W_G$$

23 Here,  $W_G$  is the weighting for the global base field,  $W_R$  is the weighting for the regional  
24 model field, and  $k$  is the grid point of the buffer zone toward the inner domain.  $W_G$  is defined  
25 according to Eq. 1 in the buffer zone but is specified as 1 and 0 in the boundary zone and the  
26 inner domain, respectively. Finally, the result over the model domain is defined by Eq. 2:

$$27 \quad F = W_G F_G + W_R F_R \quad (2)$$

28 where  $F$  is the final result of a particular tracer,  $F_G$  is the global base field, and  $F_R$  is the  
29 field calculated by the regional model.

1

### 2 **3 Idealized experiment**

3 An idealized experiment is performed to examine the feasibility of the regional version of the  
4 NDSL advection scheme. A horizontal advection scheme is examined in the 2D domain,  
5 which has 400 east-west grid points and 400 north-south grid points, with a 10-km resolution.  
6 A uniform wind field of  $10 \text{ m s}^{-1}$  in both the x and y directions is used for horizontal  
7 advection. The integration time interval ( $\Delta t$ ) is set to 10 seconds. Figure 3a shows the result  
8 of this horizontal advection every 1000 time steps. An ideal perturbation is imposed at the  
9 initial step, and the shape of the perturbation is maintained during the integration up to the  
10 3000th time step. The transported disturbance by the NDSL scheme at the 3000th time step is  
11 compared to the analytic solution (Fig. 3b); these two results are nearly identical. The shaded  
12 values in Fig. 3b indicate differences between the analytic solution and the result calculated  
13 from the NDSL. The differences are less than 1% of the maximum disturbance. The ratio of  
14 the mass change due to the NDSL to the initial mass (Eq. 3) is verified to confirm that the  
15 NDSL satisfies mass conservation.

$$16 \quad R_{mass} = \frac{(total\ mass) - (initial\ total\ mass)}{(initial\ total\ mass)} \quad (3)$$

17  $R_{mass} \approx 10^{-15}$  up to the 3000th time step, which shows that the NDSL scheme satisfies mass  
18 conservation during the integration.

19 The idealized vertical advection experiment is performed in one dimension. The vertical  
20 dimension is the sigma coordinate, which has 100 layers from 1 (bottom) to 0 (top) with equal  
21 spacing (0.01). A uniform vertical velocity of  $10^{-4} \text{ sigma s}^{-1}$  is prescribed. The integration  
22 time interval ( $\Delta t$ ) is 100 seconds. Figure 4 shows that the virtual concentration at the initial  
23 time is conserved during the transport process. The mass conservation ratio for vertical  
24 advection is  $R_{mass} \approx 10^{-15}$ .

25 The idealized experiments in the horizontal and vertical directions verify that the regional  
26 version of the NDSL scheme can accurately calculate the advection of these specific  
27 perturbations in the horizontal and vertical directions while conserving mass.

28

## 1 4 Real-case experiments

### 2 4.1 Experiment design

3 For the Fukushima case study, two experiments are performed. The first is the ORG run,  
4 which is identical to the control experiment of Saya et al. (2013), wherein tracer fields are  
5 calculated in the spectral space, as is performed in the original IsoRSM. The second is the SL  
6 run, wherein the specific humidity and radioactive material ( $^{131}\text{I}$  and  $^{137}\text{Cs}$ ) tracer fields are  
7 calculated by the NDSL scheme. All configurations except the advection method for tracers  
8 are identical in the two experiments. Figure 5 shows the experimental domain for the case  
9 study, with horizontal grid spacing of 10 km. The number of grid points is 161 (east-west) by  
10 200 (north-south). The number of vertical layers is 28 in terrain-following sigma coordinates;  
11 the lowest and highest sigma levels are 0.995 and 0.002, respectively. The red circle in Fig. 5  
12 indicates the emission point, which is the location of the Fukushima Dai-ichi nuclear power  
13 plant. The physical processes used are the relaxed Arakawa-Schubert deep convection scheme  
14 (Moorthi and Suarez, 1992), the Noah land surface model (Ek et al., 2003), the Chou radiation  
15 scheme (Chou and Suarez, 1994), and a non-local planetary boundary scheme (Hong and Pan,  
16 1996). The model simulation is integrated from 00 UTC 12 March 2011 to 00 UTC 28 March  
17 2011 (16 days). Atmospheric initial and lateral boundary conditions are provided by the  
18 NCEP-Department of Energy (DOE) reanalysis (Kanamitsu et al., 2002). In the case that  
19 negative values are introduced in the initial field, correction is performed in regional  
20 interpolation process by replacing negative values with zero. If negative tracer quantities are  
21 produced by the physical parameterizations, those negative values are transported to above  
22 layer and original values are replaced by zero. The emission rate of the radioactive tracers  
23 from Chino et al. (2011) is used in this study. To determine the size of the boundary and  
24 buffer zones (Fig. 2) for this case study, the fastest wave is assumed to be a sound wave,  
25 which moves at a speed of  $300 \text{ m s}^{-1}$ . The estimated longest traveling distance of this wave in  
26 one time step ( $\Delta t = 40 \text{ sec}$ ) is 12 km; this value is less than  $2\Delta x$ . Therefore, the fastest-  
27 moving tracer is confined within 2 grid points over one time step. Thus, it is impossible for  
28 tracers to enter from outside the domain when the boundary zone is larger than  $2\Delta x$ . In this  
29 study, the boundary zone is set to  $5\Delta x$  for safety. The buffer zone is the same size as the  
30 boundary zone.

1 When the NDSL advection used (the SL experiment), there occurs extra computation cost of  
2 35% with respect to the ORG run. However, the current version of the NDSL in the RSM still  
3 calculates spectral tracer advections even the result is not used any more. It means that the  
4 computational burden with the current NDSL in the IsoRSM is purely the increased  
5 computational cost which is required for the NDSL tracer advections. In updated release, this  
6 inefficiency will be solved.

7

## 8 **4.2 Radioactive tracer field**

9 Figure 6 shows column-integrated atmospheric radioactive tracer ( $^{137}\text{Cs}$ ) from the ORG and  
10 the SL experiments at 12 UTC 15 March 2011, when the maximum emission rates occurred.  
11 The ORG run produces very distinct noise in the zonal and meridional directions from the  
12 emission point (Fig. 6a). Additionally, a ring-shaped signal surrounds the emission point. This  
13 signal is the pattern of high-concentration, empty values; the ring shape extends from the  
14 center of the emission to outside the domain. These signals, also known as the ringing artifact,  
15 are typical of the Gibbs phenomenon. In the Fukushima case study, the discontinuity of the  
16 tracer field is pronounced because the tracers are emitted from a single grid point, which leads  
17 to significant noise from the spectral transformation processes in the ORG experiment.  
18 Widespread distributions of tracers with small concentrations occur over the domain.  
19 However, the SL experiment does not have any noise from the Gibbs phenomenon (Fig. 6b).  
20 Compared with the ORG run, the SL experiment produces a generally similar pattern of  
21 tracers, which advect northeast from the emission point. This result is the clearest advantage  
22 of the semi-Lagrangian advection scheme for tracer transport in the spectral model system.  
23 The simulated  $^{131}\text{I}$  tracers have almost same patterns with the  $^{137}\text{Cs}$  in both the ORG and the  
24 SL runs (figure not shown). Figure 7 shows longitudinal-vertical cross sections of tracers  
25 averaged in the meridional direction over  $36.5^{\circ}$ - $37.5^{\circ}\text{N}$ , which includes the emission point.  
26 This figure shows that the ringing noises extend above the surface in the ORG experiment  
27 (Fig. 7a). The SL experiment shows transport in the vertical and horizontal directions without  
28 computational noise. Clearly, the semi-Lagrangian advection method can calculate the  
29 transport of tracers without computational noise, whereas the spectral representation of  
30 advection produces severe errors.

31



### 1 **4.3 Humidity field**

2 The SL experiment applies the semi-Lagrangian method to radioactive tracers and the  
3 humidity field. The humidity field is also a positive-definite field and spectrally exhibits the  
4 Gibbs phenomenon, although the discontinuity is not as strong as it is in the radioactive tracer  
5 fields. Figure 8 shows that the specific humidity in the ORG run exhibits some negative  
6 values in the lower troposphere (Figs. 8a and 8b). The SL experiment simulates detailed  
7 humidity distributions that are similar to those of the ORG experiment but without any  
8 negative values (Figs. 8c and 8d). Negatives in the ORG run indicate that the original RSM  
9 has a systematic problem representing the positive-definite field, even though the RSM is  
10 widely used and evaluated for regional downscaling. The 16-days accumulated precipitation  
11 fields from each experiment are quite similar (Figs. 9a and 9b). General rainfall patterns  
12 observed in the tropical rainfall measuring mission (TRMM) multi-satellite precipitation  
13 analysis (TMPA) are well captured in both experiments (Fig. 9c). The spatial correlation  
14 coefficient of precipitation between the ORG run and the TMPA is 0.616 whereas the  
15 correlation coefficient between the SL run and the TMPA is 0.622. It means that the corrected  
16 humidity field by the NDSL scheme can slightly improve precipitation or keep the simulation  
17 skill of the original IsoRSM in the rainfall simulation. When we consider that the ORG  
18 experimental set have been widely used for various downscaling researches, it is possible to  
19 understand that the regional NDSL can successfully calculate the transport and distribution of  
20 humidity in the RSM. One possible reason why the improvement of the rainfall simulation by  
21 the NDSL scheme is not much significant is that the selected case in this study is not a heavy  
22 rainfall case. For a heavy rainfall case, the large discontinuity of humidity field is expected,  
23 which means higher possibility of negative value occurrences in the original IsoRSM. Further  
24 study will be continued to examine how the NDSL can improve skills for the precipitation  
25 simulation in a heavy rainfall cases.

26

### 27 **5 Summary and discussion**

28 In previous studies, the RSM has been utilized to simulate the Fukushima Dai-ichi nuclear  
29 power plant accident. The results exhibit severe noise in the simulated radioactive tracer fields  
30 (i.e., iodine-131 and cesium-137). This noise is due to the Gibbs phenomenon, wherein  
31 discontinuities in positive-definite fields create negative values after spectral transformations.  
32 This problem is common in spectral model systems. The spectral tracer advection is replaced

1 with a semi-Lagrangian advection scheme to prevent the Gibbs phenomenon in the tracer  
2 output. Prognostic tracer fields are calculated using the semi-Lagrangian method and are only  
3 considered in grid space. The Gibbs phenomenon does not occur when the spectral space  
4 transformation is not performed.

5 The semi-Lagrangian method used in this study is the NDSL scheme, which has the  
6 advantages of efficiency and simplicity. Because the NDSL has been previously applied in a  
7 global model system only, a regional version of the NDSL is developed in this study. For this  
8 application, the boundary conditions are applied by defining simple weighting functions. The  
9 regional version of the NDSL scheme is verified by performing idealized experiments using  
10 horizontal and vertical advection. These idealized experiments transport a particular  
11 disturbance to the uniform wind field. The results show that the shape is well maintained and  
12 the mass conservation is satisfied during advection. Therefore, the regional version of the  
13 NDSL is successfully applied in the RSM.

14 Two experiments are performed for the Fukushima case study to evaluate the NDSL  
15 advection scheme in the RSM. The ORG experiment is performed by the original RSM, and  
16 the SL experiment is produced by the NDSL version of the RSM. The ORG run shows severe  
17 errors in tracer fields induced by the Gibbs phenomenon. Errors appear as a ringing signal that  
18 extends zonally and meridionally from the emission point. Additionally, relatively strong  
19 ring-shaped noise is captured around the emission point. This noise is clearly removed when  
20 the tracer advection component is replaced by the NDSL scheme. The SL run shows that the  
21 NDSL advection scheme can capture the major transport of tracers without any noise from the  
22 Gibbs phenomenon. This finding is the clearest advantage of the NDSL scheme in the tracer  
23 field simulation. In the humidity field, the ORG experiment produces some negative values in  
24 the lower troposphere. However, the SL experiment does not exhibit such negatives; both SL  
25 and ORG capture the detailed distribution of the humidity field. The precipitation fields from  
26 the ORG and SL experiments are similar, which means that the NDSL properly calculates the  
27 humidity field.

28 This study reveals that replacing the tracer advection scheme with a semi-Lagrangian scheme  
29 can eliminate the Gibbs phenomenon in a regional spectral model. However, the simulated  
30 surface depositions of radioactive tracers are still deviate from the observation and  
31 precipitation from the SL experiment does not show significant improvement, even though  
32 the NDSL removes severe errors. Note that the objective of this study is to determine the

1 feasibility of the NDSL advection scheme in a regional spectral model. Thus, some  
2 quantitative validations from experiments are not included in this study. These results may be  
3 improved upon by applying enhanced physical parameterizations and advanced formulas for  
4 tracer surface deposition processes.

5

## 6 **Code availability**

7 One can access to the IsoRSM code through the concurrent versions system (CVS) server at  
8 the Center for Ocean-Atmospheric Prediction Studies (COAPS). Detailed descriptions how to  
9 get the code and install the model are located at the G-RSM homepage ([http://g-](http://g-rsm.wikispaces.com/Installation)  
10 [rsm.wikispaces.com/Installation](http://g-rsm.wikispaces.com/Installation)). For further information or requests on the model, please  
11 contact to E.-C. Chang ([echang@kongju.ac.kr](mailto:echang@kongju.ac.kr)) or K. Yoshimura ([kei@aori.u-tokyo.ac.jp](mailto:kei@aori.u-tokyo.ac.jp)).

12

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## 1 **References**

- 2 Aranami, K., Davies, T., and Wood, N.: A mass restoration scheme for limited-area models  
3 with semi-Lagrangian advection, *Q. J. R. Meteor. Soc.*, 141, 1795-1803, 2015.
- 4 Bourke, W.: An Efficient, One-Level, Primitive-Equation Spectral Model, *Mon. Wea. Rev.*,  
5 100(9), 683-689, 1972.
- 6 Chang, E.-C. and Hong, S.-Y.: Projected climate change scenario over East Asia by a regional  
7 spectral model, *J. Korean Earth Sci. Soc*, 32, 770-783, 2011.
- 8 Chen, Q.-S. and Kuo, Y.-H.: A harmonic-sine series expansion and its application for  
9 partitioning and reconstruction problems in a limited area. *Mon. Wea. Rev.*, 120, 91-112,  
10 1992.
- 11 Chino, M., Nakayama. H., Nagai. H., Terada. H., Katata. G., and Yamazawa. H.: Preliminary  
12 estimation of release amounts of 131I and 137Cs accidentally discharged from the Fukushima  
13 Daiichi Nuclear Power Plant into the atmosphere, *J. Nuclear Sci. Tech.*, 48(7), 1129-1134,  
14 2011.
- 15 Chou, M.-D., and Suarez, M. J.: An efficient thermal infrared radiation parameterization for  
16 use in general circulation models, NASA Tech. Rep. TM-1994-104606, Series on Global  
17 Modeling and Data Assimilation, NASA, Houston, Tex., 1994.
- 18 Ek, M. B., Mitchell, K. E., Lin, Y., Rogers, E., Grunmann, P., Koren, V., Gayno, G., and  
19 Tarpley, J. D.: Implementation of Noah land surface model advances in the National Centers  
20 for Environmental Prediction operational mesoscale Eta model, *J. Geophys. Res.*, 108(D22),  
21 8851, doi:10.1029/2002JD003296, 2003.
- 22 Fulton, S. R. and Schubert, W. H.: Chebyshev spectral methods for limited-area models. Part  
23 I: Model problem analysis, *Mon. Wea. Rev.*, 115, 1940-1965. 1987.
- 24 Hong, S. Y., and Pan, H. L.: Nonlocal boundary layer vertical diffusion in a medium-range  
25 forecast model, *Mon. Weather Rev.*, 124, 2322–2339, 1996.
- 26 Hong, S.-Y., Park, H., Cheong, H.-B., Kim, J.-E. E., Koo, M.-S., Jang, J., Ham, S., Hwang,  
27 S.-O., Park, B.-K., Chang, E.-C., and Li, H.: The Global/Regional Integrated Model System  
28 (GRIMs), *Asia-Pacific, J. Atmos. Sci.*, 49, 219–243, doi:10.1007/ s13143-013-0023-0, 2013.

1 Hoyer, J. M.: The ECMWF spectral limited area model, ECMWF Workshop Proc. on  
2 Techniques for Horizontal Discretization in Numerical Weather Prediction Models, 343-359,  
3 1987.

4 Juang, H.-M. H.: Semi-Lagrangian advection without iteration, In Proceedings of the  
5 Conference on Weather Analysis and Forecasting, Central Weather Bureau, Longtan, Taoyan,  
6 Taiwan, 277, 2007.

7 Juang H.-M. H.: Mass conserving and positive-definite semi-Lagrangian advection in NCEP  
8 GFS: decomposition of massively parallel computing without halo. In Proceedings of the  
9 Thirteenth Workshop on Use of High Performance Computing in Meteorology, European  
10 Centre for Medium-Range Weather Forecasts, Reading, United Kingdom, 3-7 November  
11 2008.

12 Juang, H.-M. H. and Kanamitsu, M.: The NMC nested regional spectral model, *Mon. Wea.*  
13 *Rev.*, 122, 3-26, 1994.

14 Juang, H.-M. H., Hong, S.-Y., and Kanamitsu, M.: The NCEP Regional Spectral Model: An  
15 Update, *Bull. Amer. Meteor. Soc.*, 78, 2125-2143, 1997.

16 Kanamitsu, M., Ebisuzaki, W., Woollen, J., Yang, S.-K., Hnilo, J. J., Fiorino, M., and Potter,  
17 G. L.: NCEP-DOE AMIP-II reanalysis (R-2), *Bull. Amer. Meteor. Soc.*, 83, 1631-1643, 2002.

18 Kanamitsu, M., Kanamaru, H., Cui, Y., and Juang, H.: Parallel implementation of the regional  
19 spectral atmospheric model, CEC Report CEC-500-2005-014, California Energy  
20 Commission, Sacramento, 2005.

21 Kanamitsu, M., Yoshimura, K., Yhang, Y.-B., and Hong, S.-Y.: Errors of Interannual  
22 variability and Trend in Dynamical Downscaling of Reanalysis, *J. Geophys. Res.*, 115,  
23 D17115, doi:10.1029/2009JD013511, 2010.

24 Kang, H.-S. and Hong, S.-Y.: An assessment of the land surface parameters on the simulated  
25 regional climate circulations: The 1997 and 1998 east Asian summer monsoon cases, *J.*  
26 *Geophys. Res.*, 113, D15121, doi:10.1029/2007JD009499, 2008.

27 Li, H., Kanamitsu, M., and Hong, S.-Y.: California reanalysis downscaling at 10 km using an  
28 ocean-atmosphere coupled regional model system, *J. Geophys. Res.*, 117, D12118,  
29 doi:10.1029/2011JD017372, 2012.

1 Mayron, R. H., Smith, F. B., Conway, B. J., and Goddard, D. M.: The U. K. Nuclear Accident  
2 Model, *Progress in Nuclear Energy*, 26(2), 85-104, 1991.

3 Moorthi, S. and Suarez. M. J.: Relaxed Arakawa-Schubert: A parameterization of moist  
4 convection for general circulation models, *Mon. Wea. Rev.*, 120, 978-1002, 1992

5 Mootrhi, S., Pan, H.-L., and Caplan, P.: Changes to the 2001 NCEP operational MRF/AVN  
6 global analysis/forecast system, *NWS Tech. Procedures Bull.* 484, 14 pp  
7 [<http://www.nws.noaa.gov/om/tpb/484.htm>], 2001.

8 Robert, A. J.: The Integration of a Low Order Spectral Form of the Primitive Meteorological  
9 Equations, *J. Meteor. Soc. Japan*, 44, 237-245, 1966.

10 Saya, A., Yoshimura, K., and Oki, T.: Simulation of radioactive tracer transport using  
11 IsoRSM and uncertainty analyses, *Japan Soc. Civil Engineers*, 69, I\_1765-I\_1770, 2013.

12 Segami, A., Kurihara, K., Nakamura, H., Ueno, M., Takano, I., and Tatsumi, Y.: Operational  
13 mesoscale weather prediction with Japan spectral model, *J. Meteor. Soc. Japan*, 67, 907-923,  
14 1989.

15 Staniforth, A. and Côté, J.: Semi-Lagrangian Integration Schemes for Atmospheric Models –  
16 A Review, *Mon. Wea. Rev.*, 119, 2206-2223.

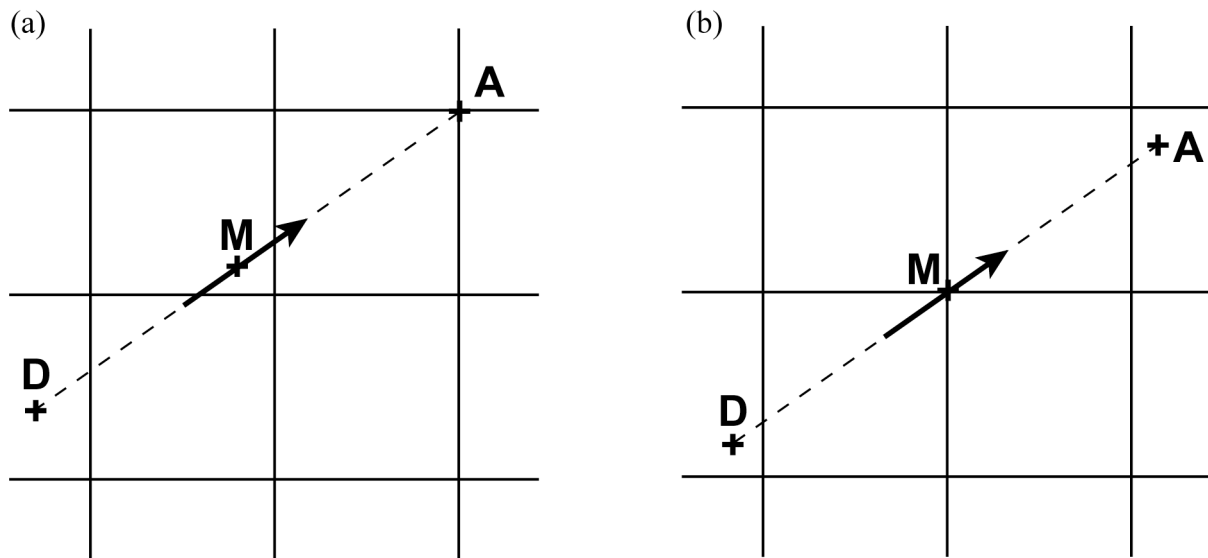
17 Tatsumi, Y.: A spectral limited-area model with time dependent lateral boundary conditions  
18 and its application to a multi-level primitive equation model, *J. Meteor. Soc. Japan*, 64, 637-  
19 663, 1986.

20 Williamson, D. L.: The Evolution of Dynamical Cores for Global Atmospheric Models, *J.*  
21 *Meteor. Soc. Japan*, 85B, 241-269, 2007.

22 Yoshimura, K.: A tracer simulation with IsoRSM on the issue of Fukushima Nuclear  
23 Accident, In proceedings of the 11th International RSM Workshop, National Central  
24 University, Jhongli, Taiwan, 15-19 August 2011.

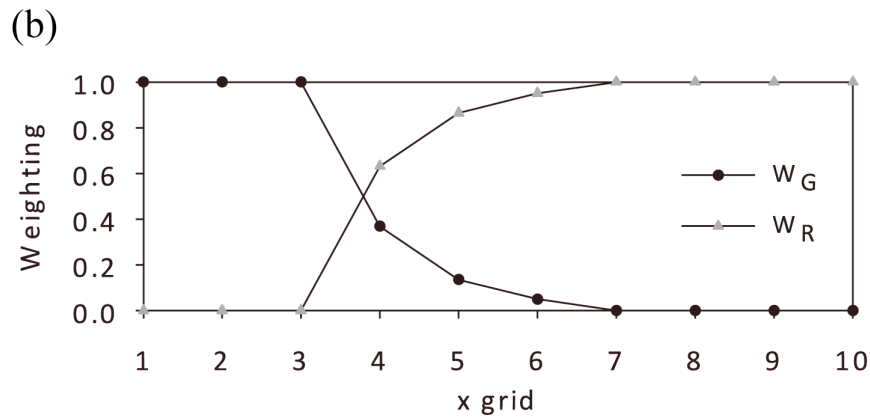
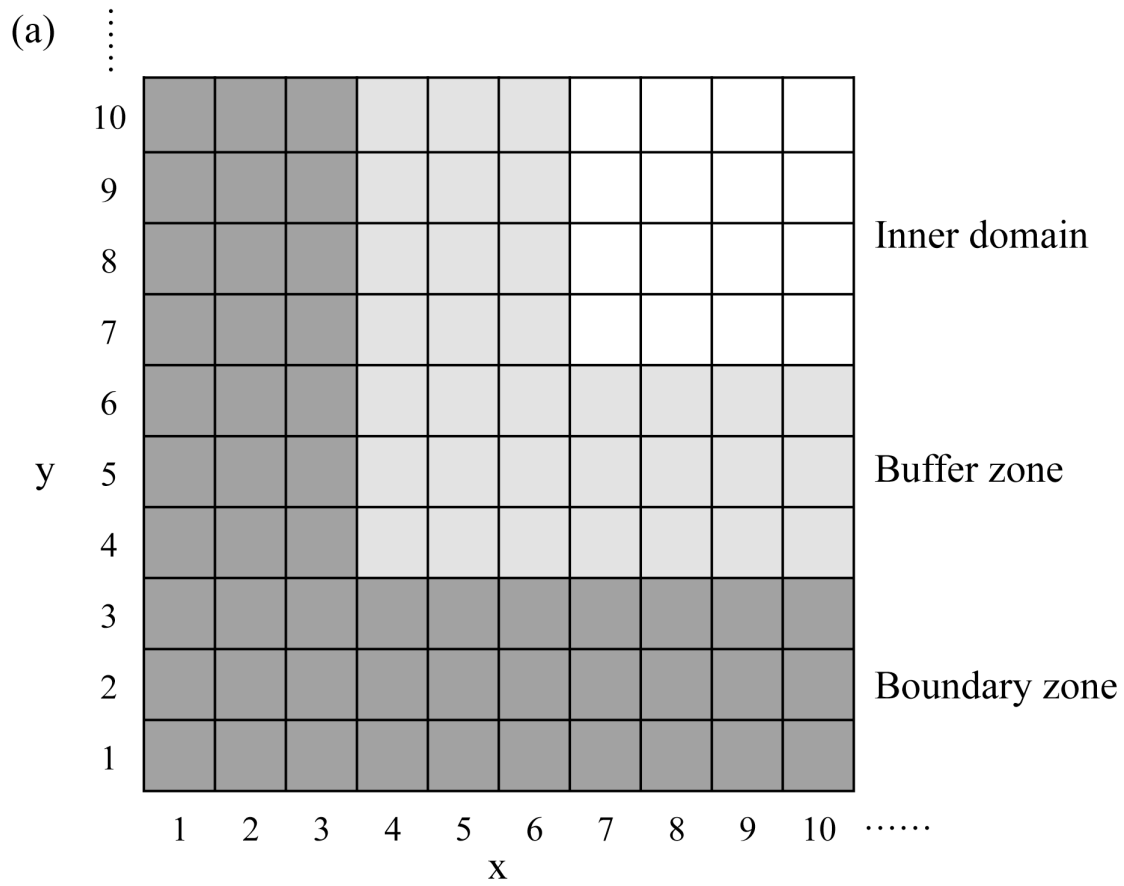
25 Yoshimura, K., Kanamitsu, M., and Dettinger, M.: Regional downscaling for stable water  
26 isotopes: A case study of an atmospheric river event, *J. Geophys. Res.*, 115, D18114,  
27 doi:10.1029/2010JD014032, 2010.

28 Zhang, Y. and Juang, H.-M. H.: A mass-conserving non-iteration-dimensional-split semi-  
29 Lagrangian advection scheme for limited-area modelling, *Q. J. R. Meteor. Soc.*, 138, 2118-  
30 2125, DOI:10.1002/qj.1938, 2012.



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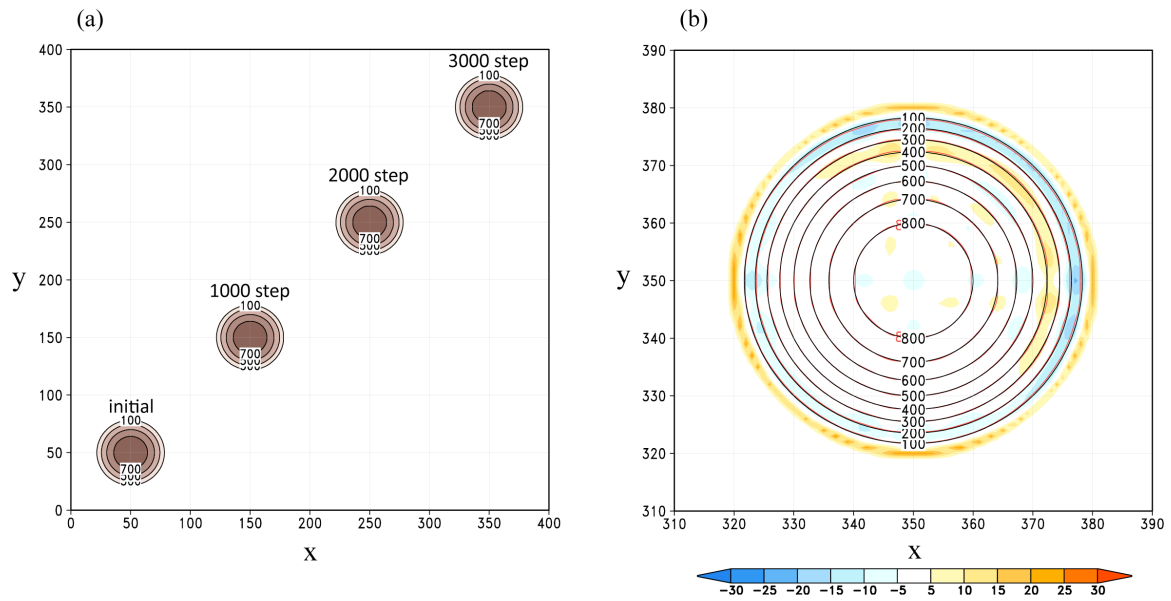
Figure 1. Schematic of two-dimensional advection for (a) the traditional backward semi-Lagrangian scheme and (b) the NDSL scheme. “D,” “M,” and “A” indicate the departure point, mid-point, and arrival point, respectively. The bold arrow indicates the wind vector at the mid-point.



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Figure 2. (a) The structure of the lower-left corner in the regional domain for boundary treatment. (b) Weighting function of the global base field (black) and regional model field (gray) along the x-axis, which includes the inner domain. This example assumes that the grid size of the boundary and buffer zones is 3.

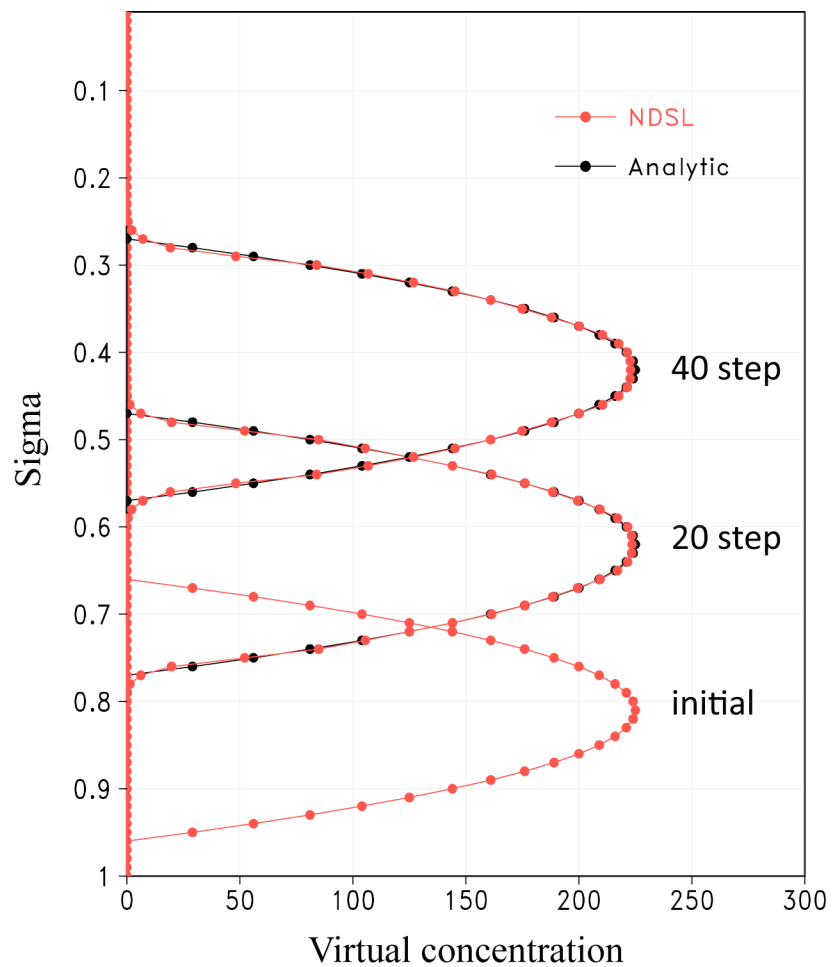




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2 Figure 3. Result from the idealized horizontal advection experiment. (a) Virtual concentration  
 3 every 1000 time steps by the NDSL advection scheme and (b) results at the 3000th time step  
 4 from the NDSL (red contour) and the analytic solution (black contour). The shaded values  
 5 indicate the differences between the NDSL results and the analytic solution.

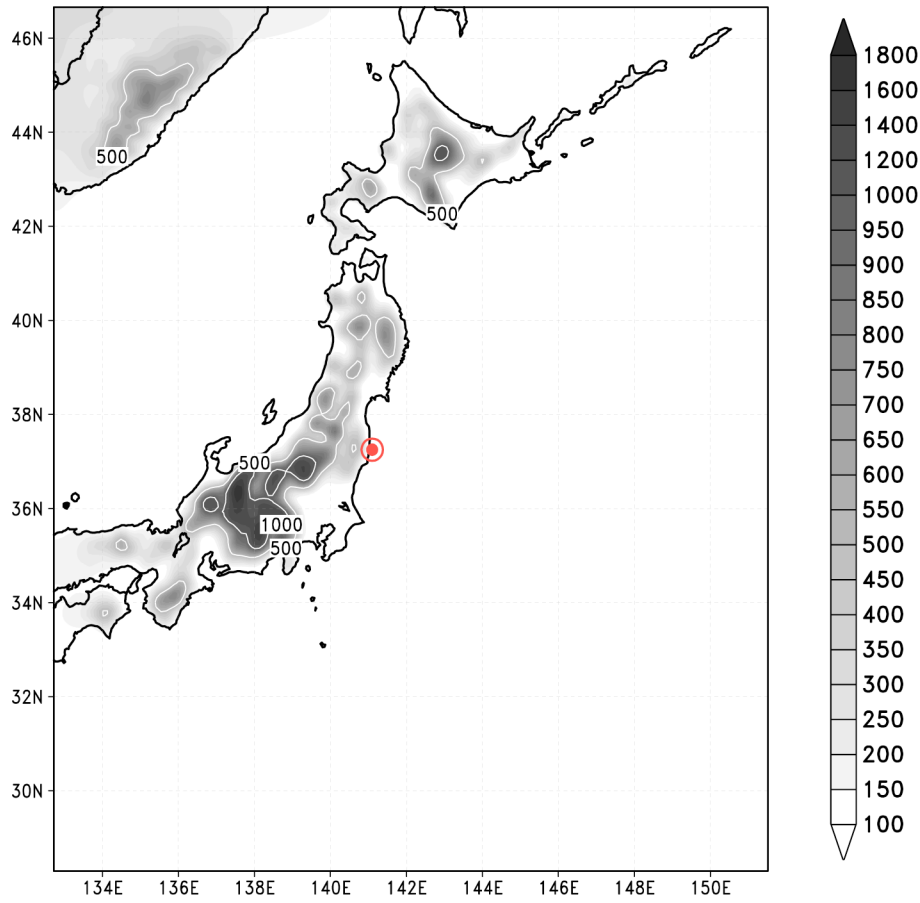
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2 Figure 4. Virtual concentration from the idealized vertical advection experiment. The red and  
 3 black lines indicate the transported concentration by the NDSL and the analytic solution,  
 4 respectively.

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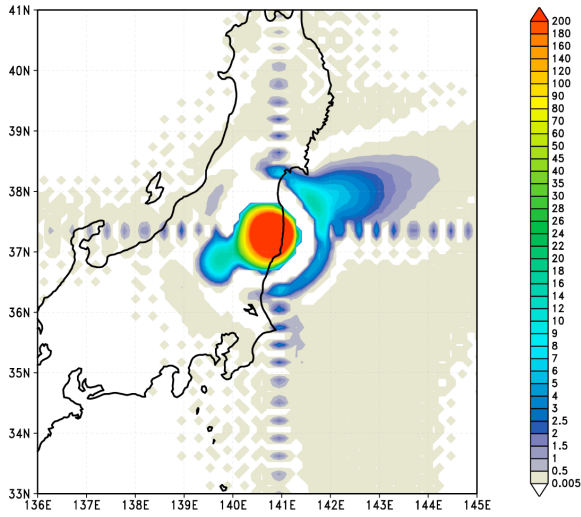


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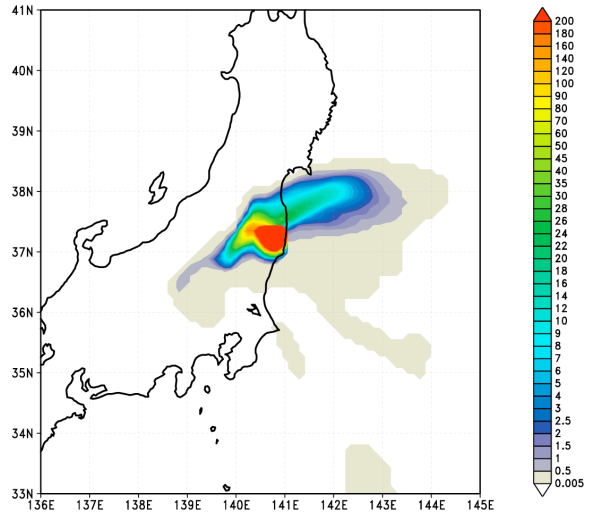
2 Figure 5. Experimental domain for the case study. The shaded values indicate orography (m).  
 3 The red circle is the emission point, which is the location of the nuclear power plant in  
 4 Fukushima.

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(a) ORG ( $^{137}\text{Cs}$ )



(b) SL ( $^{137}\text{Cs}$ )

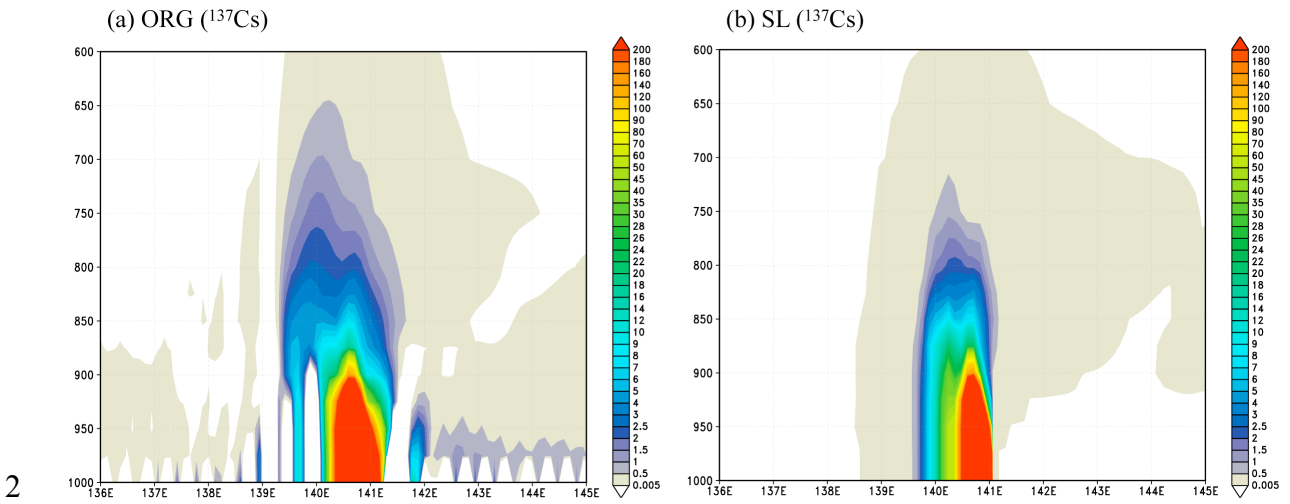


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2 Figure 6. Simulated column-integrated atmospheric radioactive tracer (cesium-137,  $\text{kBq m}^{-2}$ )  
3 at 12 UTC 15 March 2011.

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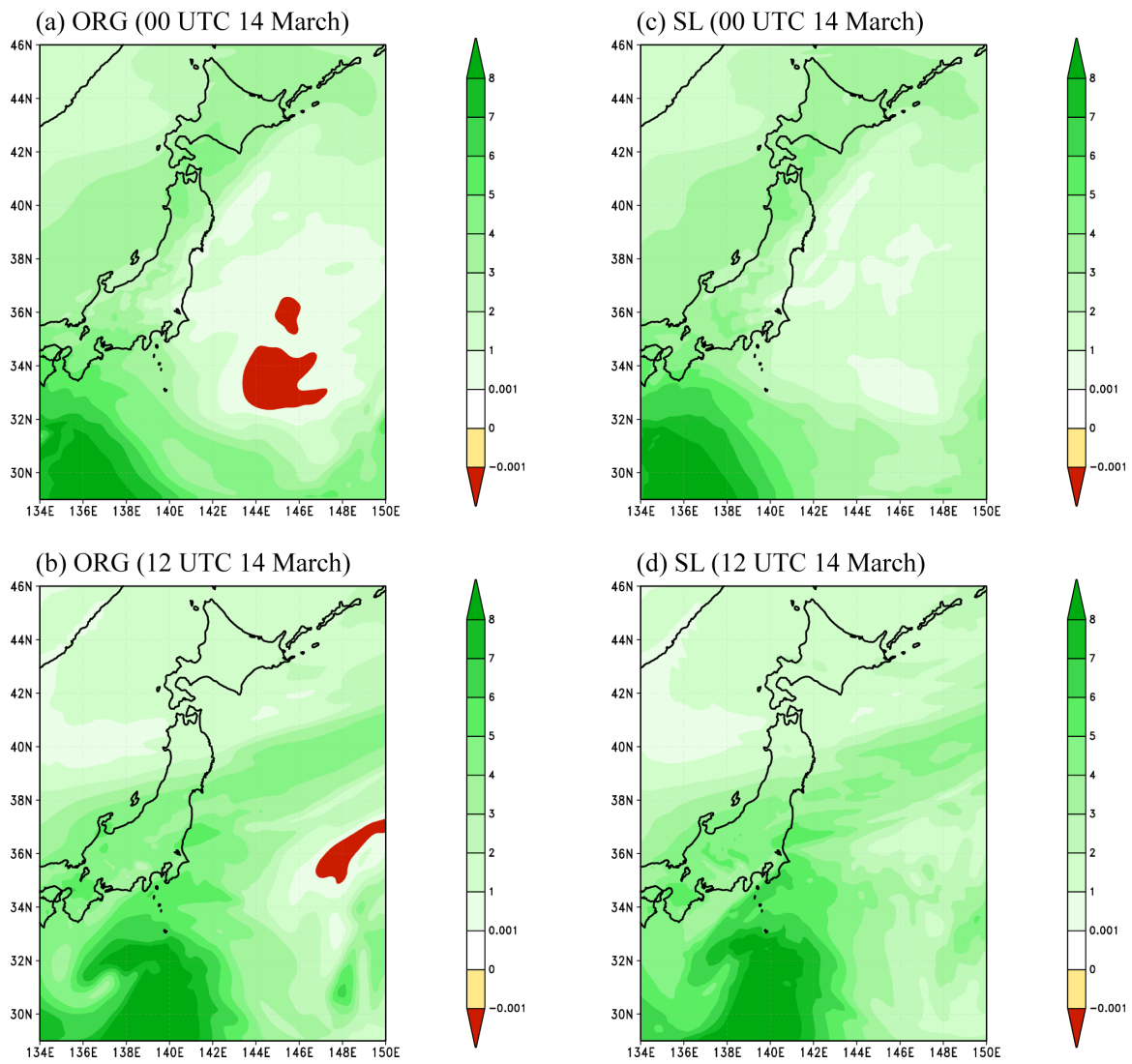


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3 Figure 7. Simulated mixing ratio of radioactive tracers (cesium-137,  $\text{Bq kg}^{-1}$ ) averaged over  
4  $36.5^\circ\text{N}$ - $37.5^\circ\text{N}$  at 12 UTC 15 March 2011.

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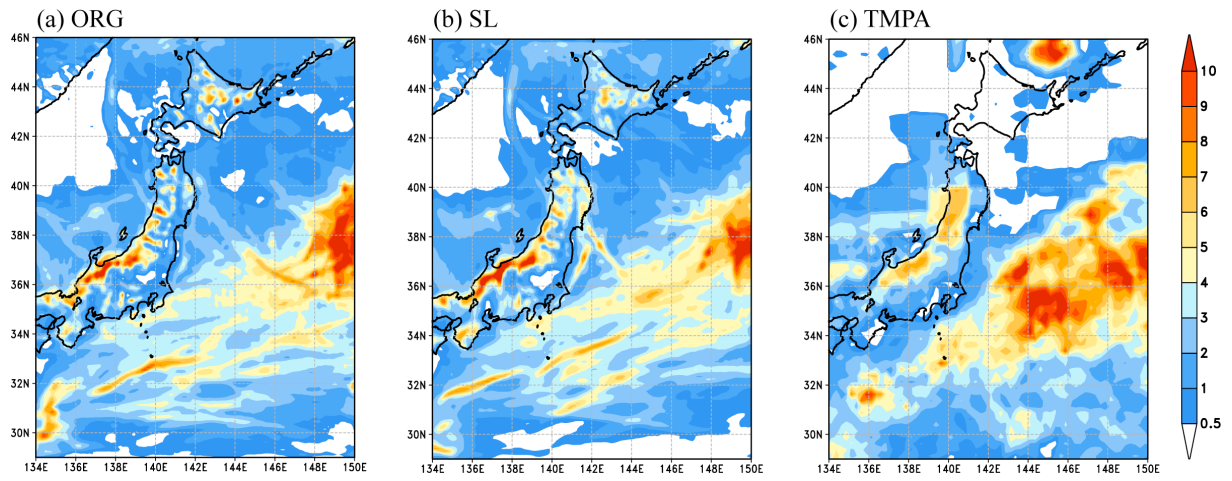
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3 Figure 8. Specific humidity ( $\text{g kg}^{-1}$ ) at 850 hPa from the (a, b) ORG and (c, d) SL runs at 00  
4 and 12 UTC 14 March 2011, respectively.

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2 Figure 9. Average precipitation over the total integration period (16 days) from the (a) ORG  
 3 run, (b) SL run, and (c) TMPA.