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Increasing vertical mixing to reduce Southern Ocean deep convection in NEMO

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Abstract

Most CMIP5 models unrealistically form Antarctic Bottom Water by open ocean deep convection in the Weddell and Ross Seas. To identify the triggering mechanisms leading to Southern Ocean deep convection in models, we perform sensitivity experiments on the ocean model NEMO forced by prescribed atmospheric fluxes. We vary the vertical velocity scale of the Langmuir turbulence, the fraction of turbulent kinetic energy transferred below the mixed layer, and the background diffusivity and run short simulations from 1980. All experiments exhibit deep convection in the Riiser-Larsen Sea in 1987; the origin is a positive sea ice anomaly in 1985, causing a shallow anomaly in mixed layer depth, hence anomalously warm surface waters and subsequent polynya opening. Modifying the vertical mixing impacts both the climatological state and the associated surface anomalies. The experiments with enhanced mixing exhibit colder surface waters and reduced deep convection. The experiments with decreased mixing are warmer, open larger polynyas and have deep convection across the Weddell Sea until the simulations end. Extended experiments reveal an increase in the Drake Passage transport of 4 Sv each year deep convection occurs, leading to an unrealistically large transport at the end of the simulation. North Atlantic deep convection is not significantly affected by the changes in mixing parameters. As new climate model overflow parameterisations are developed to form Antarctic Bottom Water more realistically, we argue that models would benefit from stopping Southern Ocean deep convection, for example by increasing their vertical mixing.

1 Introduction

Full depth open ocean deep convection has been observed only once in the Southern Ocean, following the Weddell Polynya of 1974–1976 (Gordon, 1978). In state-of-the-art CMIP5 models, Antarctic Bottom Water (AABW) is formed via open ocean deep convection which occurs most winters in the seasonally ice covered southern subpo-

lar gyres (Heuzé et al., 2013) and continues throughout climate change simulations (Heuzé et al., 2015). Not only is this process relatively unrealistic, it also leads to spurious ocean properties, variabilities and drift (e.g. Latif et al., 2013; Megann et al., 2014) and can render reanalyses useless (Azaneu et al., 2014). Martin et al. (2012) and
5 Cheon et al. (2014) have shown that each event of open ocean deep convection accelerates the southern subpolar gyres and hence increases the Antarctic Circumpolar Current (ACC) transport to unrealistically high values. With the advent of new methods to form and export AABW more realistically (e.g. Briegleb et al., 2010; Bates et al., 2012; Adcroft, 2013), it would be desirable to eliminate unrealistic Southern Ocean
10 deep convection.

Here we investigate the role of ocean model parameterisations in triggering open ocean deep convection in the Southern Ocean. Observations and models alike show that a preconditioning mechanism, reducing the static stability of the water column, is necessary before deep convection can commence (Martinson et al., 1981). Accord-
15 ing to Losch et al. (2006), static instability arises because the current generation of ocean models mixes vertically too weakly when given heterogeneous surface conditions. Buoyancy loss and mixing imposed by sea ice concentration less than unity are averaged over the horizontal grid; with weak mixing, the stratification becomes unstable. Wind-induced mixing determines the amount of turbulent kinetic energy (TKE) pen-
20 etrating the mixed layer, which if underestimated can result in a disconnection between the surface and the rest of the water column, leading to a surface salt accumulation and subsequent loss of static stability (Timmermann and Beckmann, 2004). More recently, Calvert and Siddorn (2013) and Megann et al. (2014) have demonstrated the sensitivity of high latitude mixed layers to vertical mixing. All these studies suggest that, maybe
25 counter-intuitively, vertical mixing needs to be increased to reduce deep convection in models.

The ocean model NEMO (Madec, 2008) is used here in its configuration ORCA025 (nominal horizontal resolution of 0.25°, 75 vertical levels) and is coupled to the CICE sea ice model (Hunke and Lipscomb, 2008). Vertical mixing of momentum and tracers

is treated using a TKE scheme (Gaspar et al., 1990), with convection being parameterised by an enhancement of the calculated vertical eddy diffusivity and viscosity. The ocean is forced with a prescribed atmospheric forcing from CORE2 (Large and Yeager, 2009). The ocean model configuration used in this study exhibits deep convection in the Weddell Sea (Megann et al., 2014).

Here we investigate the role of vertical mixing parameterisations on Southern Ocean deep convection with sensitivity experiments on NEMO. We detail these experiments in Sect. 2, along with our methods. In Sect. 3, we briefly explain the mechanism triggering deep convection in the control experiment before showing how this mechanism is modified in the sensitivity experiments. We also study the impact of the changes of parameter on the ACC and North Atlantic deep convection in Sect. 3. Our results are summarised and discussed, along with their limitations, in Sect. 4.

2 Sensitivity experiments and methods

2.1 Parameters studied

The purpose of these experiments is to identify a method to reduce the occurrence of open ocean deep convection in the NEMO Southern Ocean. The experiments are based on similar work performed by Calvert and Siddorn (2013) who varied 13 parameters of the TKE scheme and studied their impact on the shallow bias in the Southern Ocean summer mixed layer depth. Note that their experiments were performed with a different ocean resolution (1° compared with 0.25° here). Their findings established which values to use in the most recent joint NERC-Met Office configuration of the ocean model, "GO5" (Megann et al., 2014), which provide the base settings used in our experiments. Following their findings as well as those of Megann et al. (2014), we increase and decrease just three parameters as detailed below. The background diffusivity experiments have run thoughout the period of the atmospheric forcing (27 years,

1980 to 2006) while the other experiments have run for 10 years (1980–1989). All the experiments are summarised in Table 1.

2.1.1 Langmuir turbulence velocity scale "c_{LC}"

The NEMO TKE scheme utilises the parameterisation of Langmuir turbulence of Axell (2002). If W is the maximum downwelling velocity of the Langmuir cell (assumed to be directly related to the wind forcing), we have $W = cV_{10}$, with c a constant coupling coefficient and V_{10} the 10 m wind magnitude (Leibovich, 1983). Then assuming that W decays sinusoidally with depth (idealised Langmuir cell), we obtain:

$$W = cV_{10} = c_{LC}V_s|_{z=0}\sin\left(-\frac{\Pi z}{L}\right) \quad \text{for} -z \le L,$$

$$W = 0 \quad \text{for} -z > L, \tag{1}$$

where L is the vertical extent of the Langmuir cell. The parameter that we change is the vertical velocity scale c_{LC}, which Axell (2002) suggests should be limited to between 0.15 and 0.2 following observations in the Baltic Sea.

Averaging over 1982–1985 and between 60–45° S in the Southern Ocean (i.e. north of the deep convection regions), Calvert and Siddorn (2013) show that increasing c_{LC} deepens the mixed layer throughout the year, decreasing the summer-autumn shallow bias but increasing the winter-spring deep bias. Here we test three values:

– $c_{LC} = 0$: case with no Langmuir turbulence ("LangmuirD")

– $c_{LC} = 0.15$: control experiment ("control")

– $c_{LC} = 0.20$: case with increased vertical velocity scale ("LangmuirI").

2.1.2 Near-inertial wave breaking TKE scaling "γ"

NEMO features an ad-hoc parameterisation of the mixing due to the breaking of near-inertial waves excited by high-frequency winds, $\bar{e}_{\text{inertial}}$, which is added to the time-

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integrated TKE (Rodgers et al., 2014):

$$\bar{e}(t + \Delta t, z) = \int_{t}^{t+\Delta t} \frac{\partial \bar{e}(z)}{\partial z} \mathrm{d}z + \bar{e}_{\text{inertial}}(t, z). \qquad (2)$$

$\bar{e}_{\text{inertial}}$ is defined as:

$$\bar{e}_{\text{inertial}}(t, z) = \gamma \bar{e}|_{z=0} \exp^{z/\lambda}, \qquad (3)$$

where λ is an exponential decay scale, set as 10 m globally in the GO5 configuration, and γ is the parameter varied here (fraction of TKE penetrating below the mixed layer). Calvert and Siddorn (2013) tested values of γ ranging from 0.005 to 0.095 (default 0.05). Considering the zonal average of the Southern Ocean between 60–45° S only, mixed layer depth biases decrease in NEMO in summer and winter as γ increases. As for c_{LC}, these averages do not include the Southern Ocean deep convection areas. The value of γ for the real ocean is unknown, hence the range of values suggested by Madec (2008) is purely numerical. In addition to the control, we consider only the two most extreme values:

- $\gamma = 0.005$: case with decreased extra TKE mixing ("GammaD")

- $\gamma = 0.050$: control experiment ("control")

- $\gamma = 0.095$: case with increased extra TKE mixing ("GammaI").

2.1.3 Background diffusivity profile and surface value

Unresolved and otherwise unparameterised vertical mixing processes are represented by a background vertical eddy diffusivity. In sensitivity experiments on ORCA025 (i.e. the same resolution of NEMO as used here), Megann et al. (2014) increased the background diffusivity from 1.0×10^{-5} to $1.2 \times 10^{-5}\,\mathrm{m^2\,s^{-1}}$ and found a significant surface

freshening and increased stratification in the Arctic. They did not detail their results for the Southern Ocean.

In NEMO, it is also possible to change the shape of the background diffusivity profile. Two shapes are implemented: the background diffusivity can either be constant through depth, or increase linearly with depth (diffusivity reaches 10 times the surface value at 4000 m depth, Madec, 2008). Unpublished experiments with HiGEM (Shaffrey et al., 2009) suggest that the open ocean deep convection area is reduced in both southern subpolar gyres when the background diffusivity increases linearly instead of being constant with depth. Here we test the effects of modifying the vertical profile and/or the surface value of the background diffusivity:

– constant profile, surface diffusivity = $1.0 \times 10^{-5}\,\mathrm{m^2\,s^{-1}}$ ("KnoprofD")

– constant profile, surface diffusivity = $1.2 \times 10^{-5}\,\mathrm{m^2\,s^{-1}}$, control experiment ("control")

– linear profile, surface diffusivity = $1.2 \times 10^{-5}\,\mathrm{m^2\,s^{-1}}$ ("Kprof")

– linear profile, surface diffusivity = $1.3 \times 10^{-5}\,\mathrm{m^2\,s^{-1}}$ ("Kprofl").

2.2 Methods

To assess the state of the ocean, we use the potential temperature (hereafter referred as temperature only), sea ice concentration and mixed layer depth (MLD) diagnostics from NEMO, in its configuration ORCA025 (horizontal resolution of 0.25°) and with the GO5.0 default settings (Megann et al., 2014). The MLD is determined in the model using a density σ_θ threshold of $0.01\,\mathrm{kg\,m^{-3}}$ from the 10 m depth value (Madec, 2008). Following Heuzé et al. (2015), we consider that there is open ocean deep convection in the Southern Ocean (latitude south of 50° S) if the local monthly mean MLD reaches at least 2000 m where the bathymetry is deeper than 3000 m. We compute the area of deep convection as the sum of the areas of the individual latitude-longitude grid cells where the MLD exceeds 2000 m.

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We define an open ocean polynya as a region where the sea ice area fraction is locally less than 0.15, surrounded by a zone where the sea ice area fraction is greater than 0.15 (i.e. not directly connected to the ice-free ocean). Likewise, we compute the area of the polynya as the sum of the areas of the individual latitude-longitude grid cells where the criteria for a polynya are met. Anomalies are calculated relative to the first half of the simulation (January 1980 to December 1984), before deep convection is initiated. Anomalies are considered significant if they are larger than the monthly temporal SD during 1980–1984.

To assess the impact of Southern Ocean deep convection on circulation, we calculate the Antarctic Circumpolar Current strength as the annual total mean volume transport through Drake Passage. We integrate the zonal velocity from the Antarctic Peninsula to South America, and then over depth from the sea floor to the surface. The ACC strength is the total resulting from these integrations. We also check if the MLD in the North Atlantic is modified by our sensitivity experiments. We consider that there is deep convection in the North Atlantic (north of 50° N, 70° W to 20° E) if the local monthly mean MLD exceeds 1000 m (following observations by Våge et al., 2009, for example).

3   Results

3.1   Mechanism for the control run

In order to assess how our parameter changes impact the model ocean, we first need to understand how model processes trigger Southern Ocean deep convection. In the control run, between 1980 and 1989, the maximum MLD exceeds 2000 m in the open ocean in the Riiser-Larsen Sea only (Fig. 1, blue box), in winter 1987.

The process leading to deep convection starts two years earlier with a decrease of the katabatic winds off the Antarctic coast, resulting in a positive and significant sea ice area anomaly in June 1985 (ice fraction increased by more than 0.2 locally, Fig. 2a).

Increased ice cover means less new ice formation and a relatively salinity-stratified underlying ocean, resulting in a shallow anomaly in MLD from July to October 1985 (Fig. 2b, anomaly in MLD of 40 m compared with the median MLD of 60 m). With a reduced mixed layer entrainment, a warm anomaly develops below the mixed layer (1 °C in September 1985, Fig. 2c). When ice-free in summer, the warm waters are incorporated in the mixed layer, which is no longer anomalously shallow. The surface waters become anomalously warm (Fig. 2c), remaining above the freezing point until June 1986 (Fig. 2d). The warm surface waters impede sea ice formation, resulting in the development of an open ocean polynya over August to October 1986 (maximum extent of 22 000 km^2 in September 1986, Fig. 2e). In agreement with observations (e.g. Killworth, 1979), the model polynya allows the formation of dense water at the surface, destabilising the water column and inducing an increased convection that reaches a maximum depth of 827 m in September 1986 at the centre of the polynya.

The polynya and increased convection both act towards a decrease in stratification, preconditioning the ocean for deep convection. Moreover, the increased convection reaches the layer of warm deep water that is sitting around 300 m depth. These relatively warm waters, now brought up, enhance the warm anomaly at the surface (from January 1987, Fig. 2c) so that again the ocean surface is above freezing temperature when the sea ice should grow (Fig. 2d). As the surface waters are anomalously warm, the polynya reopens in July 1987 and reaches a maximum extent of 68 000 km^2 in October 1987 (Fig. 2f). One can notice that the polynya reopens further south than it does in 1986. That is because, similar to observations (Smith and Barber, 2007), the 1986 polynya and increased convection had an impact over a large area, lifting isopycnals significantly even at 66° S where a lower than usual sea ice concentration can be seen (0.7 in the black contours of Fig. 2e).

The conjunction of low stratified waters and the second polynya results in open ocean deep convection from August to November 1987, reaching a maximum depth of 3200 m in October 1987. In the control run, deep convection does not restart in the subsequent years. As explained in the following section, surface waters are not warm enough to

sustain deep convection after 1987. We shall now examine how the process triggering deep convection is modified in the parameter sensitivity experiments.

3.2 Modifications induced by the experiments

All sensitivity experiments start from the same initial conditions as the control simulation. The parameter changes will consequently spin up a new ocean state over the first years of the simulation. Considering the mean 1980–1984 sea surface temperature, over the area that will convect (black contours on Fig. 2a), from November to March the decreased parameter experiments (reduced vertical mixing) are warmer than the control. The increased parameter experiments (increased mixing) in contrast are colder than the control. The largest difference from the control is found in February, over the region of deep convection, for the TKE experiments: on average GammaD is warmer than the control by 0.34 °C while GammaI is colder by 0.22 °C. There is no significant difference in the winter sea surface temperature, as all experiments are ice-covered and hence at freezing temperature. The winter mixed layer is modified by the Langmuir turbulence experiments (5 m shallower than the control in LangmuirI, 5 m deeper in LangmuirD), but not significantly by the other experiments. This is consistent with the results of Calvert and Siddorn (2013): increasing the Langmuir turbulence velocity scale deepened their winter MLD by 5 to 10 m, whereas Gamma had little effect on it.

In winter 1985, all experiments show the same sea ice anomaly, at the same location, constrained by the fixed anomaly in the atmospheric forcing. The resulting shallow anomaly in MLD is modified for the Langmuir experiments: it is 10 m shallower for LangmuirD than for the control and 10 m deeper for LangmuirI. The MLD of KprofI is also 10 m deeper than the control. The temperature anomaly of summer 1986 differs significantly from that of the control run only for the Langmuir experiments (anomaly in LangmuirD warmer by 0.3 °C than in the control and 0.1 °C colder in LangmuirI) and for KnoprofD (0.3 °C warmer than the control). The results of the Langmuir experiments are consistent with those of Calvert and Siddorn (2013); the results of the background diffusivity experiments are more surprising as it is thought that the background diffu-

sivity would not have any effect on such shallow MLD (e.g. Megann et al., 2014). The regional climatologies of the decreased parameter experiments are warmer than the control in the summer of 1986. Their vertical mixing is weaker than the control, and hence they accumulate warm anomalies at the surface.

The warmer the ocean is in summer-autumn 1986, the less sea ice can form, so the larger the polynya in winter 1986 (Fig. 3a). The larger the polynya, the deeper and more extensive the convection of 1986 (Fig. 3b). The deeper the convection, the warmer the surface of the ocean after the convection (Fig. 3c), and consequently the larger the polynya in winter 1987 (Fig. 3d). The larger the 1987 polynya, the deeper the mixed layer and the larger the area of deep convection in 1987 (Fig. 3e and f). All relationships are relatively linear and significant, albeit with some uncertainty in the surface properties.

After 1987, the LangmuirD, GammaD, KnoprofD and Kprof keep convecting. This is not surprising: the "temperature loop", sustained by warm water advection, found above has been responsible for decades of deep convection every winter in other models (e.g. Martin et al., 2012). As in other modelling studies (e.g. Martinson et al., 1981) the centre of the deep convection area is advected to the middle of the Weddell Gyre (Fig. 4). There is a very strong and significant across-run relationship between the deep convection area of 1987 in the Riiser-Larsen Sea and the area of the 1988 polynya and deep convection (correlation of 0.90). In fact, advected by a steady westward slope current, the warm anomalies induced by the 1987 deep convection in the Riiser-Larsen Sea propagate to the centre of the Weddell Gyre where deep convection occurs in 1988 and 1989 (Figs. 4 and 5a). The three increased parameter experiments (and the control) are not warm enough for their anomaly to trigger the opening of a polynya in 1988 (i.e. they do not have sufficient heat content to prevent sea ice formation), hence they stop convecting and their convection area is restricted to the Riiser-Larsen Sea (Fig. 4).

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3.3 ACC and North Atlantic deep convection

We have focused on the impact of changes to the vertical mixing parameters on a specific ocean characteristic, the MLD. If these changes are to be considered part of the solution to spurious deep convection, we must ensure that they have no adverse impact elsewhere in the global model ocean. We found that the ACC increases following deep convection in the Weddell Sea. Such a response had been hypothesised by Timmermann and Beckmann (2004). In an 18 year run of the model BRIOS2 with a specified atmospheric forcing, they found that Weddell Sea deep convection intensified the Weddell Gyre circulation, and suggested that it would then strengthen the ACC (theirs was prescribed). We found a significant correlation between the ACC transport and the area of deep convection (correlation of more than 0.6 for the three 27 year simulations, see Fig. S1 in the Supplement). For each year during which deep convection occurs, the ACC transport increases by 2 to 4 Sv.

We found an increase in the ACC tranport of nearly 25 Sv in the 27 year simulations. This result is consistent with the increase of more than 20 Sv found by Martin et al. (2012), and the increase of 20 % found by Cheon et al. (2014) while modelling the Weddell Polynya in GFDL-MOM4. Such an increase is an issue for the model, especially as in our experiments the atmospheric forcing is prescribed and does not react to deep convection. In a fully coupled climate model, where the large heat flux to the atmosphere by the polynya leads to reduced sea surface pressure and increased westerlies (Martin et al., 2012), the ACC would probably become even stronger. Even with a forced model, we obtain an unrealistically strong ACC: at the end of the run, all three long simulations have exceeded the observational range of 134–164 Sv (Griesel et al., 2012). These perturbed parameter experiments show that Southern Ocean deep convection disturbs the large scale oceanic circulation.

While deep convection in the Southern Ocean is a rather unrealistic process occurring in models, deep convection in the North Atlantic is a key driver of the global ocean circulation (e.g. Johnson, 2008). Ideally, we would like to minimalise Southern Ocean

deep convection whilst maintaining North Atlantic deep convection. Our results suggest that this is possible (Fig. 5). Unlike in the Southern Ocean where a preconditioning is necessary (Fig. 5a), deep convection in the North Atlantic is persistent throughout the simulation (Fig. 5b), with differences between the experiments being less than the inter-annual variability. The differences due to parameterisations are not altered throughout the 27 year simulations, consistent with the findings of Megann et al. (2014); changing the background diffusivity has no significant effect on the northern high latitude MLD. In general, the GammaD and GammaI simulations have respectively the largest and smallest areas of deep convection; the other experiments show no consistency in the North Atlantic from 1980 to 1989 (Fig. 5b). Our results suggest that at least up to a decade, deep convection in the North Atlantic is not significantly modified by the three vertical mixing parameters found to have a large impact on the Southern Ocean deep convection. It thus seems feasible to reduce Southern Ocean deep convection by tuning these three parameters without impacting North Atlantic deep convection.

4 Discussion, limitations and conclusions

We performed sensivity experiments on a global version of NEMO with prescribed atmospheric forcing to study the trigger of unrealistic Southern Ocean deep convection. A complex chain of events, which preconditioned the ocean, initiated open ocean deep convection in the Riiser-Larsen Sea in winter 1987. It began two years earlier with a weakening of katabatic winds from Antarctica resulting in a positive sea ice anomaly in winter 1985 (Fig. 2a). These results are consistent with those of Goosse and Fichefet (2001) that highlight a strong sensitivity of ocean models to sea ice anomalies leading to open ocean deep convection in the Southern Ocean.

Modifying the oceanic vertical mixing parameters greatly altered the extent of Southern Ocean open ocean deep convection and the process triggering it. The positive anomaly in sea ice induces an anomaly in mixed layer depth. In agreement with Calvert and Siddorn (2013), we found that the experiments with increased mixing parameters

(notably LangmuirI) resulted in a deeper mixed layer in winter 1985 than the experiments with decreased parameters, and hence smaller surface temperature anomalies the subsequent summer. The perturbed parameter ensemble revealed distinct relationships between the summer and autumn sea surface temperature and the area of the subsequent polynya in winter (Fig. 3). Once the polynya forms, the ocean enters a positive feedback loop: warmer surface leads to larger polynyas, which lead to deeper convection, which mixes up relatively warm water and leads to even warmer surface waters. A similar process occurs in the Kiel Climate Model (Martin et al., 2012) where winter deep convection persists for decades and reoccurs on centennial timescales. In our simulations, deep convection stops after winter 1987 in the control run and the increased parameter experiments. Longer simulations would be needed to determine if the frequency of occurrence of deep convection is generally diminished. Experiments which reduce the vertical mixing produce strong temperature anomalies, caused by the 1987 event, which are advected westward and result in renewed deep convection across the centre of the Weddell Gyre in 1988 (Fig. 4). These experiments exhibit convection from the surface to the sea bed over most of the Weddell Sea in 1989 when the simulations stop.

Our results are consistent with those of Martin et al. (2012) and Cheon et al. (2014): deep convection increases the ACC transport above the range indicated by observations. Whilst increasing the model vertical mixing dramatically alters the open ocean deep convection in the Southern Ocean, it does not significantly modify the North Atlantic deep convection (Fig. 5). This is the same conclusion reached by Timmermann and Beckmann (2004), although they went to the extent of devising a new vertical mixing scheme specifically for seasonal sea ice regions. Our conclusion is that a similar result can be achieved within the parameterisation of an existing, globally consistent, mixing scheme. Our results, however, need to be treated with some caution, as our simulations have been too short to determine if the parameter changes result in longer term (i.e. a spun-up model) damage to the ocean simulation.

It is clear that the current state-of-the-art models require Southern Ocean deep convection in order to form their Antarctic Bottom Water (Heuzé et al., 2013). New methods, which focus on shelf overflow parameterisations, are being developed and implemented to form this water in a more physically accurate manner. Some examples are a pipe from the shelf to the open ocean (Briegleb et al., 2010), porous barriers (Adcroft, 2013) and an embedded Lagrangian model (Bates et al., 2012). This study shows the general direction that models need to take to at least reduce spurious Southern Ocean deep convection: the vertical mixing needs to be increased, not decreased as one would intuitively think. This paper paves the way for further model improvement that could help NEMO (and the CMIP6 models that use it) not to form its bottom waters unrealistically, by reducing open ocean deep convection in the southern subpolar gyres.

The Supplement related to this article is available online at doi:10.5194/gmdd-8-2949-2015-supplement.

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Interactive Discussion

Table 1. Sensitivity experiments performed on NEMO, detailed in Sect. 2: "Langmuir" experiments look at Langmuir turbulence velocity scale, "Gamma" at the penetration of an additional turbulent kinetic energy term below the mixed layer, "Knoprof" and "KProf" at background diffusivity. "I" indicates that the parameter was increased compared to the reference value, "D" that it was decreased. The parameters column identifies the shorthand name used in the NEMO simulation name-list.

name	parameter	control value	run value	run length
LangmuirD	ln_lc	true	false	10 yr
LangmuirI	rn_lc	0.15	0.20	10 yr
GammaD	rn_efr	0.05	0.005	10 yr
GammaI	rn_efr	0.05	0.095	10 yr
KnoProfD	rn_avt0	1.2×10^{-5}	1.0×10^{-5}	27 yr
KProf	nn_avb	0	1	27 yr
KProfI	nn_avb	0	1	27 yr
	rn_avt0	1.2×10^{-5}	1.3×10^{-5}	
Control				10 yr
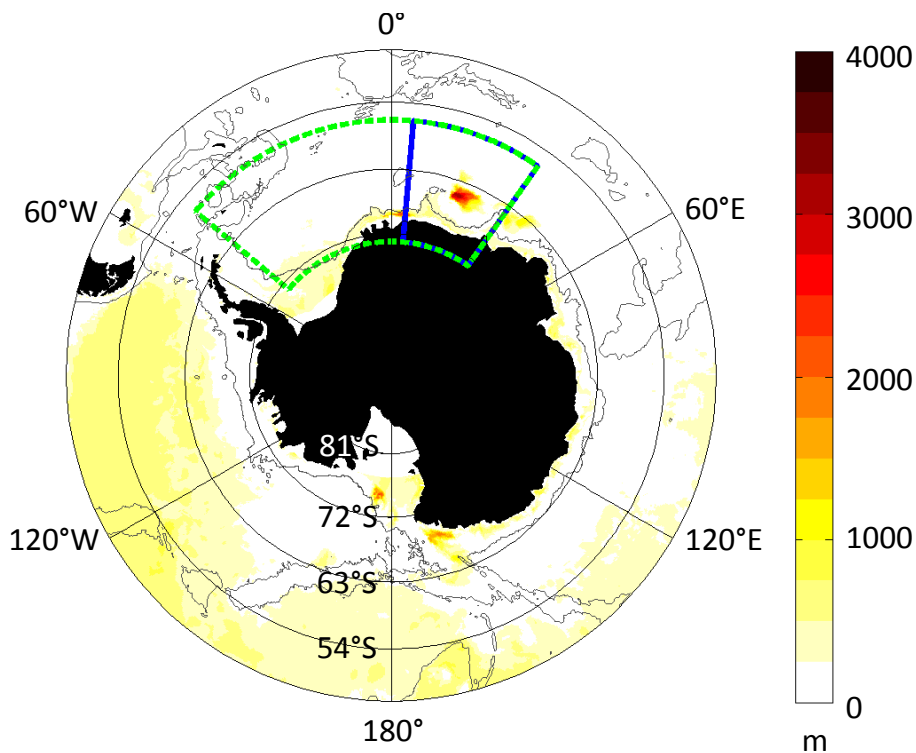
Figure 1. NEMO, control run, for each grid point, maximum monthly mixed layer depth between January 1980 and December 1989. The blue box indicates the Riiser-Larsen Sea region, which is overlaid by the dashed green box that indicates the Weddell Sea region studied in Sect. 3.2.
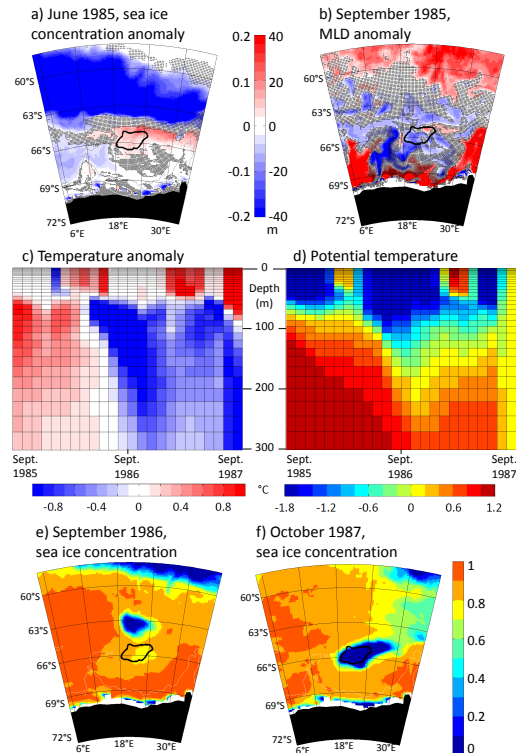
Figure 2. Chain of events leading to deep convection in 1987: **(a)** sea ice concentration anomaly in June 1985; **(b)** MLD anomaly in September 1985; **(c)** Hovmoller diagram of the potential temperature anomalies in the area where the maximum MLD from Fig. 1 exceeds 2000 m, as a function of depth and time; **(d)** same as **(c)** for the potential temperature; **(e, f)** sea ice concentration in September 1986 and October 1987. Stippling on **(a, b)** indicates areas where the anomalies are not significant. Black contours on (**a, b, e** and **f**) indicate the area where the maximum MLD from Fig. 1 exceeds 2000 m, while grey contours indicate the 3000 m isobath.
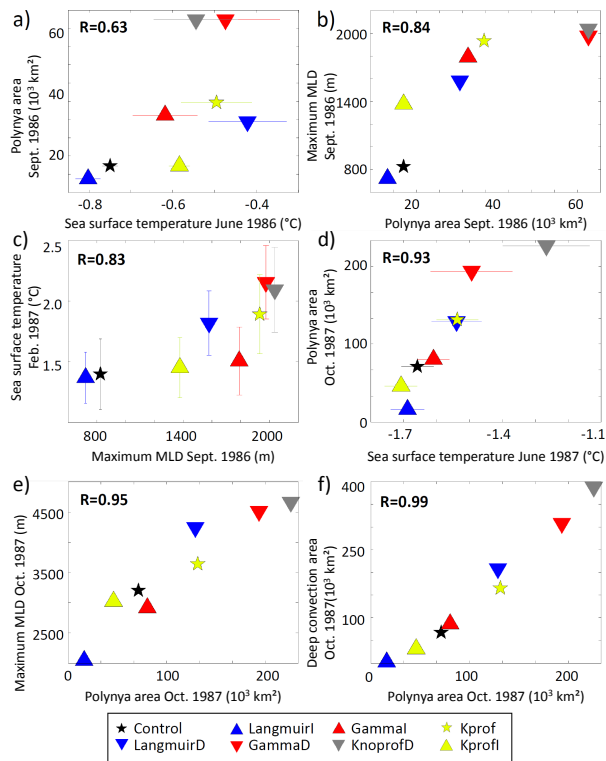
Figure 3. Across-run significant relationships between the steps leading to the deep convection event of winter 1987. **(a)** Sea surface temperature in June 1986 and polynya area in September 1986; **(b)** polynya area and maximum depth of the mixed layer, both in September 1986; **(c)** maximum depth of the mixed layer in September 1986 and sea surface temperature in February 1987; **(d)** sea surface temperature in June 1987 and polynya area in October 1987; **(e)** polynya area and maximum depth of the mixed layer, both in October 1987; **(f)** polynya area and deep convection area, both in October 1987. Horizontal (vertical) bars on a and **(d, c)** indicate the spatial SD.

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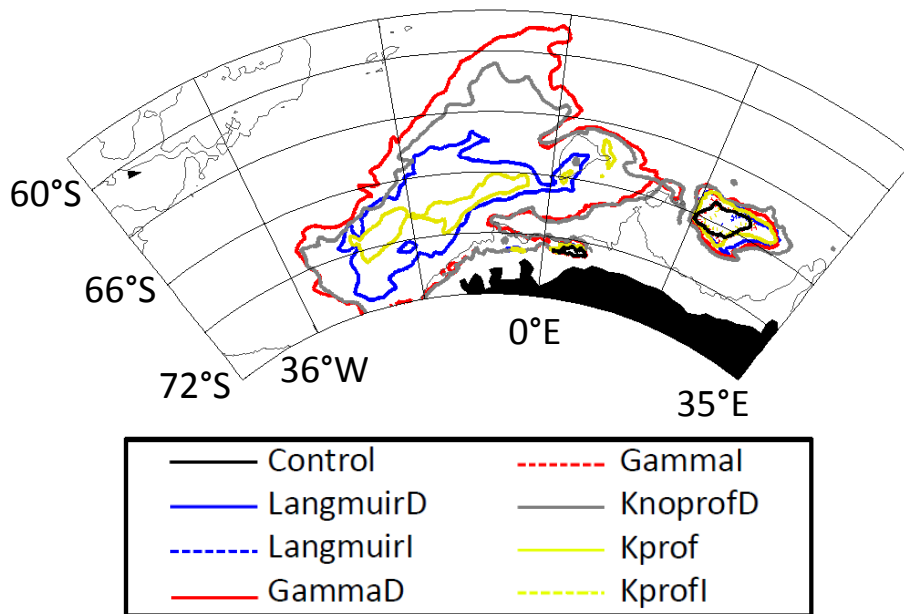Printer-friendly Version

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Figure 4. For each experiment, contours indicating where the maximum 1980–1989 MLD exceeds 2000 m in the Weddell Gyre region (green box in Fig. 1). Thin grey contours indicate the 3000 m isobath. Note that the contours for the increased-parameter experiments and the control mostly coincide.
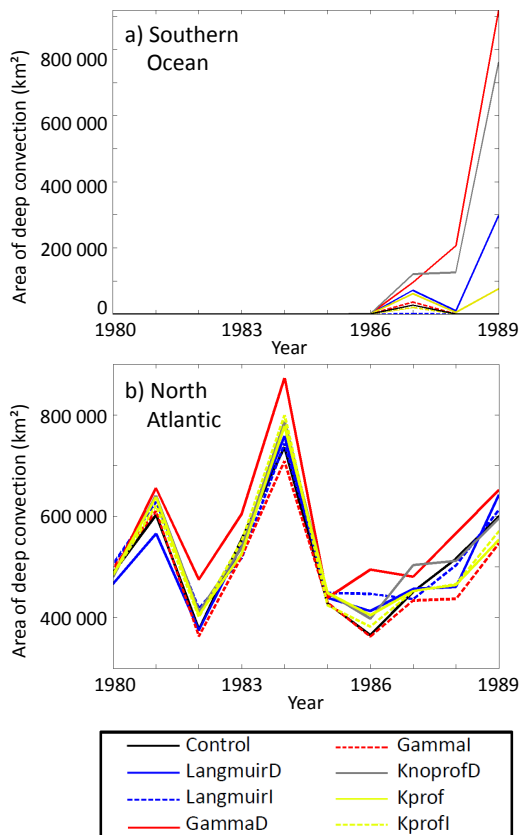
Figure 5. For each experiment, annual maximum area of deep convection in **(a)** the Southern Ocean and **(b)** the North Atlantic.