Development of a new semi-empirical parameterization for below-cloud scavenging of size-resolved aerosol particles by both rain and snow

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Abstract. A parameter called the scavenging coefficient Λ is widely used in aerosol chemical transport models (CTMs) to describe below-cloud scavenging of aerosol particles by rain and snow. However, uncertainties associated with available size-resolved theoretical formulations for Λ span one to two orders of magnitude for rain scavenging and nearly three orders of magnitude for snow scavenging. Two recent reviews of below-cloud scavenging of size-resolved particles recommended that the upper range of the available theoretical formulations for Λ should be used in CTMs based on uncertainty analyses and comparison with limited field experiments. Following this recommended approach, a new semi-empirical parameterization for size-resolved Λ has been developed for below-cloud scavenging of atmospheric aerosol particles by both rain (Λ_{rain}) and snow (Λ_{snow}). The new parameterization is based on the 90th percentile of Λ values from an ensemble data set calculated using all possible "realizations" of available theoretical Λ formulas and covering a large range of aerosol particle sizes and precipitation intensities (R). For any aerosol particle size of diameter d, a strong linear relationship between the 90th-percentile $\log_{10}(\Lambda)$ and $\log_{10}(R)$, which is equivalent to a power-law relationship between Λ and R, is identified. The log-linear relationship, which is characterized by two parameters (slope and y-intercept), is then further parameterized by fitting these two parameters as polynomial functions of aerosol size d. A comparison of the new parameterization with limited measurements in the literature in terms of the magnitude of Λ and the relative magnitudes of Λ_{rain} and Λ_{snow} suggests that it is a reasonable approximation. Advantages of this new semi-empirical parameterization compared to traditional theoretical formulations for Λ include its applicability to below-cloud scavenging by both rain and snow over a wide range of particle sizes and precipitation intensities, ease of implementation in any CTM with a representation of size-distributed particulate matter, and a known representativeness

based on the consideration in its development of all available theoretical formulations and field-derived estimates for $\Lambda(d)$ and their associated uncertainties.

1 Introduction

The removal of below-cloud aerosol particles by precipitation, either rain or snow, decreases the concentrations of particulate matter in the air and contributes to the wet deposition of toxic pollutants. This process has been identified as one of the most efficient removal mechanisms for atmospheric particles and is thus a key process in aerosol chemical transport models (CTMs) (Textor et al., 2006). Simulating this process with reasonable accuracy in CTMs has important impacts when model results from CTMs are used to assess air quality, climate, or ecosystem issues. This process, however, involves complex interactions between aerosol particles and falling hydrometeors and thus is commonly parameterized in CTMs (e.g., Zhang, 2008; Gong et al., 2011). A parameter called the scavenging coefficient Λ (s^{-1}) serves this purpose (Seinfeld and Pandis, 2006).

Various theoretical and empirical formulations for Λ exist in the literature to parameterize rain and snow scavenging of below-cloud aerosol particles. This choice matters because CTMs with different Λ formulations produce significantly different predictions of particulate matter concentrations and atmospheric deposition budgets (e.g., Rasch et al., 2000; Solazzo et al., 2012). To quantify the differences in the existing size-resolved formulations for Λ and to identify the dominant product terms causing these differences, we recently conducted detailed reviews of available parameterizations of below-cloud scavenging of size-resolved aerosol particles by rain (Λ_{rain}) and by snow (Λ_{snow}) (Wang et al., 2010, 2011; Zhang et al., 2013). The major conclusions from these review studies can be summarized as follows: (1) Different theoretical formulations for Λ can differ by one to two orders of magnitude for scavenging by rain (Λ_{rain}) and by up to three orders of magnitude for scavenging by snow (Λ_{snow}), depending on aerosol particle size. (2) Different formulas for hydrometeor-aerosol particle collection efficiency, which is one of the key product terms of the available theoretical formulations for Λ, can cause uncertainties of one order of magnitude or more for both Λ_{rain} and Λ_{snow} whereas different formulas for the three other product terms of Λ, i.e., the number size distribution, terminal velocity, and effective cross-sectional area of falling hydrometeors, can cause uncertainties of a factor of 2 to 5 in Λ. (3) The majority of field-derived estimates of Λ_{rain}, from which empirical Λ_{rain} formulas were developed, are one to two orders of magnitude larger than all theoretical Λ_{rain} formulas; the only exception is one controlled outdoor field experiment that obtained Λ_{rain} to a similar order of magnitude to the theoretical values (Sparmacher et al., 1993; Wang et al. 2010). A similar feature was also found for Λ_{snow}, although the differences between the few available field measurements and theoretical values are not as large as for Λ_{rain}. (4) The differences between empirical and theoretical Λ values can largely be explained by additional processes/mechanisms that influence field-derived estimates of Λ but that are not considered in the theoretical Λ formulas.

Based on the conclusions listed above, we provided some recommendations regarding the applications of Λ_{rain} and Λ_{snow} parameterizations in CTMs (Wang et al., 2010, 2011; Zhang et al., 2013) as follows: (1) Empirical Λ formulas should not be used in CTMs because some of the processes contributing to the field-derived estimates of Λ are treated in CTMs separately; (2) Upper-range values of available theoretical Λ formulations should be used in CTMs because they are closer to, while still smaller than, the field-derived estimates of Λ, and thus are thought to be more realistic than mid- to lower-range values from the available theoretical Λ formulations; (3) A simple semi-empirical formula for size-resolved Λ_{rain} and Λ_{snow} should be developed that takes into account the large range of Λ_{rain} and Λ_{snow} values that can be obtained from existing theoretical formulas, the many different possible choices for their product terms, and the upper-bound values provided by field-derived estimates. Note that certain physical processes that have potential to increase particle collection efficiency, e.g., storm dynamics (Chate, 2005) and rear capture of particles by falling drops (Quérel et al., 2013), are not explicitly or implicitly treated in any existing theoretical formulas. Thus, existing theoretical formulas are likely to be biased low for certain rain types.

The present study follows the above recommendations to develop a new semi-empirical formula for size-resolved Λ_{rain} and Λ_{snow}. The new parameterization is based on the existing theoretical framework for Λ_{rain} and Λ_{snow} (e.g., Slinn, 1984). Existing empirical Λ_{rain} and Λ_{snow} formulas purely based on field measurements are not used directly for the parameterization development; they are, however, used for comparison, selection, and evaluation purposes in this study. In the following sections, the methodology employed to develop the new parameterization is briefly described in Sect. 2. The development and resulting form of the parameterization is described in detail in Sect. 3. Next, a discussion on the new parameterization is presented in Sect. 4 followed by some conclusions in Sect. 5.

2 Methodology

In CTMs that simulate aerosol particle number concentrations, the time change of number concentration for aerosol particles undergoing below-cloud scavenging by falling hydrometeors is commonly described as (Seinfeld and Pandis, 2006):

$$\frac{\partial n(d,t)}{\partial t} = -\Lambda(d) \cdot n(d,t) \quad , \tag{1}$$

where $n(d, t)$ is the number concentration of aerosol particles with a diameter d at time t and $\Lambda(d)$ is the size-resolved scavenging coefficient (s^{-1}) for aerosol particles of size d. $\Lambda(d)$ can be described theoretically as (Slinn, 1984):

$$\Lambda(d) = \int_0^{\infty} A(d, D_p)(V_D - v_d)E(d, D_p)N(D_p)dD_p \ , \tag{2}$$

where D_p is the diameter of a hydrometeor (either raindrop or melted snow particle) and $N(D_p)$ is the number size distribution of hydrometeors, V_D and v_d are the terminal velocities of hydrometeors and aerosol particles, respectively, $E(d, D_p)$ is the collection efficiency (dimensionless) between an aerosol particle of size d and a hydrometeor of size D_p, and $A(d, D_p)$ is the effective cross-sectional area of a hydrometeor projected normal to the fall direction.

According to Eq. (2), if it is assumed that $V_D >> v_d$, then calculating Λ requires knowledge of four product terms: $E(d, D_p)$, $N(D_p)$, V_D, and A. Since raindrops are usually assumed to be spherical, the effective cross-sectional area A of a falling raindrop can be estimated as (e.g., Slinn, 1984)

$$A(d, D_p) = \frac{\pi}{4}(D_p + d)^2 \quad . \tag{3}$$

Extending the review of Wang et al. (2010), lists and references of available formulas for the other three product terms for the calculation of Λ_{rain} are provided in Tables 1, 2 and 3, respectively, while lists and references of available formulas for all four product terms for the calculation of Λ_{snow} are provided in Tables 4, 5, 6 and 7, respectively (Zhang et al., 2013). All symbols used in this study are defined in the appendices in Table 10 (Nomenclature).

As mentioned in the Introduction, different choices for these product terms give a large range of Λ values. To develop a new Λ parameterization, the following five-step approach was employed. The first step was to generate an ensemble of all potential Λ_{rain} values as a function of aerosol particle diameter d and a specified precipitation intensity R using all possible combinations of the product-term formulas listed in Tables 1 to 3, and to generate a second ensemble of all potential Λ_{snow} values using all possible combinations of the product-term formulas listed in Tables 4 to 7. In the second step, the ensembles of calculated Λ_{rain} and Λ_{snow} values were closely scrutinized and unrealistic values were modified or removed where it was possible to identify shortcomings in the formulation of any of the product-term parameterizations. In the third step, the 90th-percentile values of Λ_{rain} and Λ_{snow} were extracted from the reduced ensembles of Λ_{rain} and Λ_{snow} values for each aerosol particle diameter bin and precipitation intensity R. Note that the decision to choose 90th-percentile values was somewhat arbitrary, but it was based on the recommendations in Wang et al. (2010) and Zhang et al. (2013) that the upper range of theoretical Λ_{rain} and Λ_{snow} values should be used in CTMs and on the complementary evidence on upper bounds provided by field-derived estimates of Λ_{rain} and Λ_{snow}. Steps 1 to 3 were repeated many times in order to span a large range of precipitation intensity values, which resulted in a large data set of 90^{-th}-percentile $\Lambda_{rain}(d, R)$ and $\Lambda_{snow}(d, R)$ values. This 90^{-th}-percentile data set was then used as the basis for generating the new Λ_{rain} and Λ_{snow} parameterizations through a curve-fitting technique (Step 4) followed by an assessment of their relative errors (Step 5). The next section describes the application of the above approach to develop a new parameterization for the below-cloud scavenging of size-resolved aerosol particles by both rain and snow.

3 Development of the new parameterization

To solve Eq. (2) numerically for size-resolved Λ using selected product-term formulas, a number of size bins or sections need to be defined to describe both aerosol-particle and hydrometeor size distributions. A similar bin structure to that used previously in Wang et al. (2010) and Zhang et al. (2013) was also used here. Briefly, one set of 100 size bins was used to discretize the size distribution of raindrops (for Λ_{rain}) or snow particles (for Λ_{snow}) and a second set of 100 size bins was used to discretize the size distribution of aerosol particles. The size ranges considered were 1 μm to 10 mm in particle diameter for raindrops or snow particles (as liquid-water equivalent) and 0.001 μm to 100 μm in particle diameter for aerosol particles. A constant-volume ratio between successive size bins was used for both discretizations. The ambient temperature was assumed to be 15°C for rain cases and -10°C for snow cases and the ambient pressure was assumed to be 1013.5 hPa. Uncertainties associated with the choice of ambient temperature and pressure values are discussed in Sect. 4.3 below.

3.1 Λ_{rain}

Following step 1 of the approach described in Sect. 2, we calculated Λ_{rain} as a function of particle diameter for 100 size bins using Eq. (2) and 400 different combinations of formulas for $E(d, D_p)$, $N(D_p)$, and V_D (i.e., 5, 10, and 8 formulas, respectively, as listed in Tables 1, 2, and 3). Note that the product-term formulas were originally generated from a wide range of rain types such as "widespread", convective, thunderstorm and hurricane. Figure 1 shows the results for a precipitation intensity R of 1.0 mm h^{-1} as an example. The predicted Λ_{rain} values differ by one order of magnitude for ultrafine (e.g., <0.01 μm) and giant (e.g., >10 μm) aerosol particles and by nearly two orders of magnitude for particles in the diameter range from 0.01 μm to 10 μm.

Next, following step 2 from Sect. 2, we found that two groups of Λ_{rain} profiles had different shapes from the rest of the profiles for all of the precipitation intensities considered in this study. One group predicts much higher Λ_{rain} values for aerosol particles larger than 0.5 μm (see group of yellow lines in Fig. 1a) and the other group predicts much lower Λ_{rain} values for aerosol particles larger than 1.0 μm (see group of red lines in Fig. 1a). The first group was identified to be caused by the use of the $E(d, D_p)$ formula of Park et al. (2005) and the second group by the use of the $E(d, D_p)$ scheme of Ackerman et al. (1995).

Upon further investigation we found that the Park et al. (2005) formula neglects the critical Stokes number threshold in the inertial impaction mechanism, which leads to an additional contribution of inertial impaction to $E(d, D_p)$ for particles smaller than 3 μm in diameter. In fact, inertial impaction can only occur for particles with a Stokes number above the critical Stokes number, which is close to 1.2. The corresponding threshold diameter is close to 3 μm for a unit-density particle and a 1 mm raindrop (Phillips and Kaye, 1999; Loosmore and Cederwall, 2004). Thus, Λ_{rain} calculated using the $E(d, D_p)$ formula of Park et al. (2005) is believed to be an overestimation for particles with diameters from 0.5 μm to 3 μm. The $E(d, D_p)$ scheme of Ackerman et al. (1995), on the other hand, considers the collection mechanisms of Brownian diffusion, convective Brownian diffusion enhancement, and inertial impaction. In this scheme, the required collision efficiency values are interpolated from a look-up table from Hall (1980). The table, however, only covers collector (raindrop) sizes of 10 to 300 μm in radius colliding with aerosol particles (collected particles) with size ratios (the so-called p-ratio) from 0.05 to 1.0. There are no data available for collectors larger than 300 μm in radius, a size range that has appreciable concentrations in medium to heavy rain, or for particles with size ratios less than 0.05, which can include particles from 0.5 to 10 μm in radius. As well, collision efficiencies for collectors smaller than 30 μm were later found to be underestimated (Vohl et al., 2007). These deficiencies appear to be the main causes of the lower values of Λ_{rain} for particles in the diameter range from 1.0 μm to 10.0 μm compared to the rest of the Λ_{rain} formulas.

The above examination suggests that the two groups of Λ_{rain} profiles that used the $E(d, D_p)$ formulation of Park et al. (2005) and Ackerman et al. (1995) were not as realistic as the rest of the Λ_{rain} profiles. We thus removed the Λ_{rain} profiles based on the $E(d, D_p)$ formulation of Park et al. (2005) from further consideration since there was no easy way to fix the problem. We noticed, however, that Vohl et al. (2007) had updated the Hall (1980) table with new experimental results that provided more realistic collision efficiencies for wider size ranges for both collector and collected particles. Thus, we chose to keep the Λ_{rain} profiles based on the $E(d, D_p)$ scheme of Ackerman et al. (1995) for further analysis, but these were modified profiles based on the updated collision efficiency table of Vohl et al. (2007) in place of the Hall (1980) table.

With this finalized selection of the available $E(d, D_p)$ formulas (Table 1), there are 320 Λ_{rain} profiles based on different combinations of the product terms that are retained for further analysis (Fig. 1b). The use of the revised Ackerman et al. (1995) $E(d, D_p)$ scheme dramatically changed the corresponding 80 Λ_{rain} profiles (the red lines in Figure 1b), whose magnitudes increased by a factor of 2-3 for large particles (d > 10 μm) and over an order of magnitude for particles between 3.0 μm and 10.0 μm in diameter. The revised Λ_{rain} profiles were also comparable to the other 240 Λ_{rain}

profiles that used different $E(d, D_p)$ formulas (see the large group of black lines in Fig. 1b). Thus, it is recommended that the Hall (1980) table should be used with caution in the parameterization of Λ_{rain} in CTMs.

Using the 320 Λ_{rain} profiles shown in Fig. 1b, we identified a number of percentile values of Λ_{rain} for each aerosol particle diameter. These maximum, 95th-, 90th-, 80th-, 70th-, and 50th-percentile, and minimum Λ_{rain} profiles are shown in Fig. 1c. Note that the dots in this panel correspond to the original Λ_{rain} values shown in Fig. 1b and the lines are the calculated percentile Λ_{rain} profiles. Note also that the percentile profiles in Fig. 1c may not match exactly with any of the Λ_{rain} profiles shown in Fig. 1b, but they represent the range and distribution of the ensemble of all theoretical Λ_{rain} values across the range of different aerosol particle sizes.

In Fig. 1d the percentile Λ_{rain} profiles are compared with the available Λ_{rain} measurements and one empirical formula (Laakso et al., 2003: see Appendix A) that were summarized in Wang et al. (2010). Note that the blue solid triangles in this panel come from the controlled outdoor experiment of Sparmacher et al. (1993) while the other symbols come from *in situ* field measurements made by different researchers. Note that even the maximum theoretical Λ_{rain} values are smaller than the majority of field-experiment-derived values and those from the empirical formula of Laakso et al. (2003), and the differences can be larger than one order of magnitude for particles smaller than 3 μm. However, the 50th- to 90th-percentile theoretical Λ_{rain} profiles seem to agree reasonably well with the Λ_{rain} values estimated from the controlled outdoor experiment of Sparmacher et al. (1993). It is also worth noting that the Λ_{rain} profile from the parameterization of Henzing et al. (2006), which was developed using a three-parameter fit to a set of pre-calculated

Λ_{rain} values generated from a theoretical Λ_{rain} formulation (see Appendix B), falls into the lower range of the ensemble of available theoretical Λ_{rain} values.

The large differences in Λ_{rain} between the *in situ* field-derived values and those from the controlled outdoor experiment and between the field experiments and the theoretical formulations are caused by many different factors. Some of the differences might reflect the uncontrolled real-world situation while others are due to experimental errors and to errors in the theoretical formulations (Khain and Pinsky, 1997; Maria and Russell, 2005; Andronache et al., 2006; Wang et al., 2011; Quérel, 2012; Quérel et al., 2013). Choosing the upper range of theoretical Λ_{rain} values for applications in CTMs appears to be a reasonable choice because these values are only slightly higher than the corresponding values from the controlled outdoor experiment but are still lower than values from the majority of field experiments. Thus, the 90th percentile of the range of the ensemble of theoretical Λ_{rain} profiles was chosen for further analysis and parameterization development.

Moving to step 3 in Sect. 2, we repeated the calculation of Λ_{rain} with Eq. (2) for all of the 320 combinations of product-term formulas for each of 37 different precipitation intensities *R*, which covered the range of values from 0.01 to 100 mm h^{-1} and were uniformly distributed logarithmically (same as the tick values shown in x-axis of Fig. 2b). 90th-percentile Λ_{rain} values were then calculated from the ensemble of theoretical Λ_{rain} profiles for each aerosol particle diameter bin *d* and every precipitation intensity *R*. These 90th-percentile Λ_{rain} data are plotted against precipitation intensity in Fig. 2a as a set of 100 lines, with each line representing one aerosol particle diameter and in the form of Λ_{rain} vs. *R*.

Regression analysis suggests that for each aerosol particle diameter (i.e., each individual line in Fig. 2a), there exists a strong linear relationship between $\log_{10}(\Lambda_{rain})$ and $\log_{10}(R)$, or in other words a power-law relationship between Λ_{rain} and R, which can be expressed as:

$$\log_{10}(\Lambda(d,R)) = \log_{10}(A(d)) + B(d)(\log_{10} R) \quad , \tag{4}$$

$$\Lambda(d,R) = A(d)R^{B(d)} \quad . \tag{5}$$

Linear regression analysis based on Eq. (4) was performed for all 100 lines and the squares of the resulting correlation coefficients were very high, ranging from 0.9963 to 1.0. Figure 2b shows seven of these regression lines for seven selected aerosol particle diameters with the original data (the 90th-percentile Λ_{rain} values for 37 R values) shown as symbols. $B(d)$ values were obtained for all 100 aerosol sizes directly from the regression analysis. It is apparent from this panel that both the slope of the regression lines ($B(d)$) and its y-intercept ($\log_{10}A(d)$) may vary with aerosol particle diameter. Note, however, that the y-intercept does not cross the y-axis shown in Fig. 2b because the actual R value instead of $\log_{10}(R)$ is used for the x-axis. But according to Eq. (4), $A(d)$ equals Λ_{rain} (d, 1) (i.e., when R = 1.0 mm h^{-1}), so $A(d)$ values are also readily available. The resulting $A(d)$ and $B(d)$ values are plotted in Figs. 2c and 2d, respectively, for each of 100 aerosol particle diameters .

Since $A(d)$ and $B(d)$ correspond at this stage to sets of discrete data, a least-square polynomial curve-fitting technique was used to fit these power-law coefficient data and parameterize $A(d)$ and $B(d)$ as continuous functions of aerosol particle diameter. Due to the abrupt change of the values of both $A(d)$ and $B(d)$ at particle diameters between 1 and 2 µm, the particle diameter range of each of the two data sets was split into two contiguous segments for separate but more accurate fitting. After many tests, the separation point of the two segments was determined to be 1.97 µm for $A(d)$ (see Fig. 2c) and 1.94 µm for $B(d)$ (see Fig. 2d). We thus chose 2.0 µm to be the separation point for both the $A(d)$ and $B(d)$ curve fits. After some experimentation, the following polynomical functions (up to sixth order) were selected for fitting the four segments:

$$\log_{10}(A(d)) = \begin{cases} a_0 + a_1(\log_{10} d) + a_2(\log_{10} d)^2 + a_3(\log_{10} d)^3 & d \le 2.0\mu m \\ b_0 + b_1(\log_{10} d) + b_2(\log_{10} d)^2 + b_3(\log_{10} d)^3 + b_4(\log_{10} d)^4 + b_5(\log_{10} d)^5 + b_6(\log_{10} d)^6 & d > 2.0\mu m \end{cases} \quad (6)$$

$$B(d) = \begin{cases} c_0 + c_1(\log_{10} d) & d \le 2.0\mu m \\ e_0 + e_1(\log_{10} d) + e_2(\log_{10} d)^2 + e_3(\log_{10} d)^3 + e_4(\log_{10} d)^4 + e_5(\log_{10} d)^5 + e_6(\log_{10} d)^6 & d > 2.0\mu m \end{cases} \quad (7)$$

Note that the unit of d is µm, and the above equations should be applied to wet aerosol diameter. The empirical best-fit coefficients that were obtained for the above equations are listed in Table 8.

A comparison of Λ_{rain} values predicted by the new parameterization described by Eqs. (5), (6) and (7) with the data used for developing the parameterization (the 90th-percentile $\Lambda_{rain}(d, R)$ values) is shown in Fig. 3a for five different precipitation intensities. Very good agreement is evident for the full range of aerosol particle size and full range of precipitation intensity. To further examine the comparison shown in Fig. 3a, the relative error between Λ_{rain} values from the new parameterization and the original 90th-percentile values was also calculated (Fig. 3b). The relative error was within 10% for most of the aerosol particle sizes, except for the 2-6 µm diameter range for which the error could be larger than 30%. The largest relative errors corresponded to the aerosol particle diameters where Λ_{rain} increased abruptly with particle diameter. It should also be noted that various particle-size separation points were tested for the separate fits of Eqs. (6) and (7) (e.g., from 1.9 to 2.2 µm), and a separation point of 2.0 µm does lead to the minimum relative errors for most aerosol sizes.

To gain an idea of how $A(d)$ and $B(d)$ in Eqs. (6) and (7) would differ if $\Lambda_{rain}(d, R)$ values other than 90th-percentile ones had been used, a separate empirical fitting was performed using 50th-percentile values. It was found that $B(d)$ values did not change by very much whereas A(d) values differed by one order of magnitude. As noted above, B(d) represents the rate of change of $\Lambda_{rain}(d, R)$ for changes of R while A(d) represents the $\Lambda_{rain}(d, R)$ value when $R = 1.0$ mm h^{-1}. This means that $\Lambda_{rain}(d, R)$ for the 90th and 50th percentiles vary similarly with changes in R, but the magnitude of the 90th-percentile $\Lambda_{rain}(d, R)$ is much larger than the 50th-percentile $\Lambda_{rain}(d, R)$.

Overall, this new simple semi-empirical parameterization provides a good fit of the original Λ_{rain} data for all aerosol particle sizes and precipitation intensities. As well, uncertainties associated with the use of this new scheme in CTMs to parameterize Λ_{rain} should not be larger than those shown by Wang et al. (2010) to be associated with the existing theoretical formulas. The $\Lambda(d)$ profile generated from the new parameterization does not exactly match any of the existing theoretical profiles considered, but for all aerosol particle diameters its values will lie within the upper range of an ensemble of theoretical $\Lambda(d)$ values obtained from all possible combinations of existing product-term formulas. The new parameterization is designed for use in CTMs to describe below-cloud scavenging of size-resolved aerosol particles. We believe it to be a reasonable first-order approximation for any precipitation conditions, either stratiform or convective, considering that precipitation intensity and precipitation type (i.e., rain or snow) are likely to be the only precipitation information available in many CTMs (e.g., information on different rain types or droplet size distributions may not be available).

3.2 Λ_{snow}

The development of the new semi-empirical parameterization for Λ_{snow} follows the same approach described above for Λ_{rain}. The first step was to calculate an ensemble of theoretical Λ_{snow} profiles across the aerosol particle size spectrum using Eq. (2) for a precipitation intensity of 1.0 mm h^{-1} for all possible combinations of the product terms listed in Tables 4 to 7. There are three $E(d, D_p)$, four $N(D_p)$, eight V_D, and four A formulas available in the literature related to snow particles, but some of the V_D formulas were only applicable to specific snow types. Thus, a total of 168 combinations of these product-term formulas were used to calculate Λ_{snow} profiles (see Fig. 4a). Note that these formulas cover four habit types of snow crystals – spherical ice crystals, dendritic snow plates, columnar ice crystals, and graupel particles (see Table 7), all of which occur frequently in nature (e.g., Hobbs et al., 1972).

As discussed in Zhang et al. (2013), the range of the ensemble of available theoretical Λ_{snow} formulations is much larger than that for Λ_{rain} (compare Fig. 4a with Fig. 1b). It is likely that part of this larger range is due to real variability (e.g., different snow particle shapes and related properties affecting Λ_{snow}) while the other part is due to parameterization errors (e.g., improper formulation of related parameters). Examining the ensemble of Λ_{snow} profiles plotted in Fig. 4a (i.e., step 2), we did not find any obviously unrealistic profiles. The two clusters with distinct minima were caused by different formulas applying to different snow particle shapes and should not be considered as unrealistic (cf. Figs. 1, 2, and 8 of Zhang et al., 2013). Considering that information about snow particle shapes is not commonly available in CTMs, we chose to group all of the existing formulas together without explicit consideration of snow particle shape. Thus, all of the values in Fig. 4a were used for further analysis. Similar to Fig. 1c, the range and percentile values of Λ_{snow} were also generated as shown in Fig. 4b. Also plotted are two field-derived empirical formulas for Λ_{snow}, one from Paramonov et al. (2011) (Appendix C) and one from Kyrö et al. (2009) (Appendix D), but it should be noted that both formulas are more applicable to weaker snowfall intensities (e.g., 0.1-0.2 mm hr^{-1}) than the intensity assumed in Fig. 4b (1 mm hr^{-1}) and are only valid for aerosol particle sizes in 0.01-1.0 μm diameter range. Figure 4b shows that the upper range of the theoretical Λ_{snow} profiles calculated assuming a snowfall intensity of 1 mm hr^{-1} are of the same order of magnitude as the limited field data, which were observed under mostly weaker snowfall intensities. The theoretical Λ_{snow} profiles would be smaller than the experimental data if the same snowfall intensity as observed in the field were to be used for the calculation of Λ_{snow} using Eq. (2). To be consistent with the choice made for Λ_{rain}, the 90th percentile of the ensemble of all theoretical Λ_{snow} formulations at each aerosol particle diameter was also used to develop the new parameterization for Λ_{snow}. However, the evidence supporting this choice is somewhat weaker for Λ_{snow} than for Λ_{rain} due to the very limited field data for snow scavenging cases.

Theoretical size-resolved Λ_{snow} values were calculated in step 3 using the 168 combinations of product-term formulas for each of 37 precipitation intensities uniformly distributed logarithmically from 0.001 mm h^{-1} to 10 mm h^{-1} in liquid water equivalent. Given that 10 mm of snow is approximately equivalent to 1 mm of rain, a different range of precipitation intensities was used to generate the Λ_{snow} ensemble data set than that used in the Λ_{rain} case. 90th-percentile Λ_{snow} values for each aerosol particle diameter were then extracted for each precipitation intensity and are plotted in Fig. 5a, where again each line corresponds to a fixed aerosol particle diameter. The relationship between $\log_{10}(\Lambda_{snow})$ and $\log_{10}(R)$ can also be described by Eq. (4). Linear regressions were again calculated, and the squares of the correlation coefficients of the 100 regressions were again very high, ranging from 0.9736 to 0.9997. Seven of the 100 regression lines together with the data points being fit are plotted in Fig. 5b as examples.

The same approach described in Sect. 3.1 was also used here to generate $\log_{10}(A(d))$ and $B(d)$ values (Figs. 5c and d) and to conduct least-squares polynomial curve-fitting to parameterize $\log_{10}(A(d))$ and $B(d)$ for all d values. Again, the data sets were split into two contiguous segments for separate fitting. Multiple intersections between the two fitting functions were found for both the $\log_{10}(A(d))$ and $B(d)$ cases. This time a final separation point was chosen at a particle diameter of 1.44 μm because this value produced the minimum relative errors between the parameterized and the original theoretical Λ_{snow} values. The polynomial fitting formulas for the snow case are shown below and their corresponding empirical best-fit coefficients are listed in Table 8.

$$\log_{10}(A(d)) = \begin{cases} a_0 + a_1(\log_{10} d) + a_2(\log_{10} d)^2 + a_3(\log_{10} d)^3 + a_4(\log_{10} d)^4 + a_5(\log_{10} d)^5 + a_6(\log_{10} d)^6 & d \le 1.44\mu m \\ b_0 + b_1(\log_{10} d) + b_2(\log_{10} d)^2 + b_3(\log_{10} d)^3 + b_4(\log_{10} d)^4 + b_5(\log_{10} d)^5 + b_6(\log_{10} d)^6 & d > 1.44\mu m \end{cases} \quad (8)$$

$$B(d) = \begin{cases} c_0 + c_1(\log_{10} d) + c_2(\log_{10} d)^2 + c_3(\log_{10} d)^3 + c_4(\log_{10} d)^4 + c_5(\log_{10} d)^5 + c_6(\log_{10} d)^6 & d \le 1.44\mu m \\ e_0 + e_1(\log_{10} d) + e_2(\log_{10} d)^2 + e_3(\log_{10} d)^3 + e_4(\log_{10} d)^4 + e_5(\log_{10} d)^5 + e_6(\log_{10} d)^6 & d > 1.44\mu m \end{cases} \quad (9)$$

A comparison of the new parameterization described by Eqs. (5), (8) and (9) with the Λ_{snow} values from Fig. 5a is shown in Fig. 6a for five different precipitation intensities and the relative error from this comparison is shown in Fig. 6b. Reasonably good agreement was observed for the full range of aerosol particle size and full range of precipitation intensity. The relative error was within 30% for most aerosol particle sizes, except for the 1-4 μm diameter range, for which the error could be as large as 50%. Considering the very large range (i.e., two orders of magnitude or larger) of the existing theoretical Λ_{snow} values (cf. Fig. 4), an uncertainty of 50% or a factor of 2 in the parameterized Λ_{snow} values should be acceptable.

4 Discussion

4.1 Power-law relationship between Λ and *R*

A power-law relationship between the size-resolved Λ_{rain} or Λ_{snow} parameters and precipitation intensity *R* for each particle diameter *d* was identified in Sect. 3 and was used in the development of the new parameterization. The finding of such a power-law relationship is not surprising since many earlier theoretical and experimental studies also suggested the existence of such a relationship, although most of the earlier studies focused on bulk Λ instead of size-resolved Λ (Mircea et al., 1998; Andronache, 2003; Duhanyan and Roustan, 2011). A brief comparison of the results from the present study with earlier studies in terms of the power–law parameters is provided in Table 9 and presented below.

Early investigations reviewed by McMahon and Denison (1979) and more recent theoretical considerations (e.g., Scott, 1982; Mircea et al., 1998; Andronache, 2003) as well as field and experimental studies (Jylhä, 1991; Okita et al., 1996; Sparmacher et al.,1993) have suggested that the exponent *B* had values in the range of 0.59 to 0.94 for Λ_{rain} and 0.3 to 1.14 for Λ_{snow} (see Table 9 and the reviews of Sportisse (2007) and Duhanyan and Roustan (2011)). The field measurements by Jylhä (1991) and Okita et al. (1996) reported *B* values of 0.64-0.76. Sparmacher et al. (1993) fitted their experimental Λ data from their controlled outdoor study with a power-law relationship and obtained *B(d)* values of 0.59, 0.60, 0.94 and 0.61 for four selected aerosol particle diameters of 0.23, 0.46, 0.98 and 2.16 μm, respectively, for rain scavenging and values of 0.62, 0.89 and 1.09 for three selected aerosol particle diameters of 0.46, 0.98 and 1.66 μm, respectively, for snow scavenging. The *B* values obtained from theoretical derivations (Scott, 1982; Mircea et al., 1998;

Baklanov and Sorensen, 2001; Andronache, 2003; Feng, 2007) ranged from 0.59 to 0.86 for submicron particles and from 0.7 to 0.86 for coarse-mode particles for rain scavenging and from 0.31 to 1.14 for both submicron and for coarse-mode particles for snow scavenging with different habit types of snow crystals. However, the two most recent field studies on snow scavenging (Kyrö et al., 2009; Paramonov et al., 2011) did not identify a clear dependency of Λ_{snow} on R. As discussed in Zhang et al. (2013), we speculated that this might be due to the small range of snowfall intensities sampled in these experiments.

The values of $B(d)$ in the present study fall in the range of 0.64-0.91 for rain scavenging (Fig. 2d) and 0.53-0.86 for snow scavenging (Fig. 5d). More specifically, $B(d)$ has values in the ranges 0.64-0.67, 0.67-0.72, and 0.72-0.91 for ultrafine particles (d < 0.01 μm), mid-range particles (0.01 μm < d < 2 μm, and large particles (d > 2.0 μm), respectively, for rain scavenging and values in the ranges 0.66-0.77, 0.53-0.66, and 0.58-0.89, respectively, for the same particle diameter size ranges for snow scavenging. Thus, the results of the present study related to the exponent of the power-law relationship between Λ and R are comparable with most of the previous studies for both rain and snow scavenging.

As noted in Sect. 3.1 the parameter $A(d)$ equals $\Lambda(d)$ when R = 1.0 mm h^{-1}. Therefore, the values of $A(d)$ should be similar to the upper range of those in the theoretical formulas and lower than those in the field-data based empirical ones given the design decisions made in the development of the new parameterization. A comparison of $A(d)$ values from the new parameterization with those found in the literature (Table 9) supports this hypothesis.

4.2 Relative magnitudes of Λ_{rain} and Λ_{snow}

We briefly compared the relative magnitudes of Λ_{rain} and Λ_{snow} in one of our previous studies (Zhang et al., 2013) and concluded that snow scavenging seemed to be more effective than rain scavenging for equivalent precipitation amounts (i.e., liquid water equivalent) based on the median and upper-range theoretical Λ_{rain} and Λ_{snow} values. Since the 90th percentiles of the ensembles of both theoretical Λ_{rain} and Λ_{snow} formulations were used in this study to develop the new parameterizations for Λ_{rain} and Λ_{snow}, values of Λ_{snow} from the new scheme might be expected to be larger than values of Λ_{rain} from the new scheme for equivalent precipitation intensity. To obtain a quantitative measure of the relative magnitudes of Λ_{rain} and Λ_{snow} for the new parameterization, the ratios of Λ_{snow} to Λ_{rain} as a function of precipitation intensity were calculated for all 100 aerosol particle diameters.

Figure 7a shows that the magnitude of Λ_{snow} is higher than that of Λ_{rain} for the same precipitation intensity by a factor ranging from three to 300, depending on aerosol particle size and precipitation intensity. The ratio of Λ_{snow} to Λ_{rain} is the highest for medium particle sizes (i.e., $0.1 < d < 5.0$ µm; shown as yellow lines) and is the lowest for coarse and giant particles (e.g., $d > 5.0$ µm; shown as green lines). The largest ratios were found for a particle diameter of about 2.0 µm for all R values. However, the lowest ratios were found to occur for a particle diameter of 100 µm for small R values (lowest green line) and a particle diameter around 4.0 µm for large R values (lowest yellow line). The dependence of the Λ_{snow} to Λ_{rain} ratio on particle diameter can be better seen in Fig. 7b for selected R values. The ratio decreases with increasing R for medium-size particles (yellow lines in Fig. 7a), increases with increasing R for ultrafine particles (some of the blue lines in Fig. 7a), and only change slightly with increasing R for giant particles (e.g., $d > 10$ µm; some of the blue lines in

Fig. 7a).

It is possible to offer some explanation of the strong dependence of this ratio on aerosol particle diameter in terms of the physics of precipitation scavenging. Figures 1d and 4b show that the 90th-percentile Λ profiles are qualitatively similar for rain and snow scavenging. However, two significant differences exist between these two profiles. The first difference relates to the value of the aerosol particle diameter at which the minimum Λ value occurs. The Λ_{rain} minimum occurs at a particle diameter around 0.4 μm whereas the Λ_{snow} minimum occurs at a particle diameter around 0.1 μm (which corresponds to a local minimum in Fig. 7). For submicron particles, scavenging is mainly controlled by the interception mechanism and the contribution of this mechanism to scavenging increases with increasing particle diameter (e.g., see Fig. 1 of Wang et al., 2010). For snow scavenging, the increase of Λ_{snow} with particle diameter in this size range is faster than that for rain scavenging due to the larger cross-sectional areas of snow particles. Thus, the ratio between the snow and rain scavenging coefficients in Fig. 7 increases in the particle diameter range between 0.1 μm to 1.0 μm. The second significant difference relates to the abrupt transition of Λ_{rain} from an interception regime to an inertial-impaction regime at a particle diameter of about 2 μm (Fig. 1d). For particle diameters larger than 2 μm, Λ_{rain} increases more quickly with d than does Λ_{snow}. As a result, the Λ_{snow} to Λ_{rain} ratio decreases quickly with increasing d until leveling off for particle diameters close to 10 μm.

Some previous studies also support this result that snow scavenging is more effective than rain scavenging for equivalent precipitation amounts. Several field studies carried out before the 1980s found that snow scavenging of aerosols was 28 to 50 times more efficient than rain scavenging based on the equivalent water content of the precipitation (Reiter, 1964; Carnuth, 1967; Reiter and Carnuth, 1969; Graedel and Franey,1975). The average Λ_{snow} value obtained in the controlled outdoor experiment of Sparmacher et al. (1993) was five times higher than the average Λ_{rain} value obtained in similar controlled conditions for two aerosol particle diameters (0.46 and 0.98 μm). Tschiersch (2001) obtained values of Λ_{snow} up to two orders of magnitude higher than Λ_{rain} for particles in the size range of 0.5-3.5 μm for low precipitation intensities (water equivalent < 1 mm h^{-1}). Two recent field studies also claimed that snow is a better scavenger of aerosol particles than rain per equivalent water content (Kyrö et al., 2009; Paramonov et al., 2011). This limited experimental evidence suggests that the new parameterization is qualitatively correct in terms of the relative magnitudes of Λ_{rain} and Λ_{snow}, although it may not be quantitatively accurate.

4.3 Uncertainties in the new Λ parameterization related to the choice of ambient atmospheric conditions

The new parameterization for Λ_{rain} and Λ_{snow} was developed assuming the ambient temperature to be 15°C for rain scavenging and -10°C for snow scavenging and the ambient pressure to be 1013.5 hPa for both rain and snow scavenging. Such a choice may introduce uncertainties in Λ when the actual ambient atmospheric state differs from the assumed one. To investigate this issue a set of six sensitivity tests was performed covering the ambient temperature range of 5°C to 30°C for rain and -5°C to -30°C for snow and for a different ambient pressure (900 hPa) for both rain and snow. Figure 8 shows the percentage difference of the calculated 90^{th}-percentile Λ for the above mentioned temperature and pressure values relative to the Λ from the new parameterization scheme for different aerosol particle diameters and a precipitation intensity of 1.0 mm h^{-1}. The changes in Λ values due to different ambient temperature and pressure values are generally within 10% for all particle sizes for both rain and snow scavenging except for particle diameters from 0.1 μm to 2.0

μm for rain scavenging, where the differences can reach 30%. Of the four product terms needed to calculate Λ, only $E(d, D_p)$ and V_D might be impacted by changes in ambient temperature or pressure, and Λ is much more sensitive to $E(d, D_p)$ than to V_D (Wang et al., 2010; Zhang et al., 2013). Therefore, uncertainties in Λ due to ambient atmospheric condition are likely to arise mainly from the impact of different ambient temperatures and pressures on collection efficiency $E(d, D_p)$. The larger uncertainty at particle diameters of 0.1 to 2.0 μm for rain scavenging than for snow scavenging is due to the inclusion of thermophoresis and diffusiophoresis collection mechanisms in some of the theoretical formulas, since these two collection mechanisms are sensitive to the ambient atmospheric condition and have a large contribution to particle scavenging at this particular aerosol size range (Wang et al., 2010). Similar uncertainties were also found for other precipitation intensities considered in the present study.

5 Conclusions

The availability of a number of existing theoretical formulas for the size-resolved scavenging coefficient Λ(d) requires somewhat arbitrary choices to be made when selecting amongst these schemes and their product terms for implementation in a chemical transport model followed by the coding and run-time solution of often complex algorithms. The new semi-empirical Λ parameterization developed in the present study only requires input of precipitation intensity and precipitation type (rain or snow) – two routine output variables in any meteorological model used as a CTM driver. Thus, this new parameterization is readily implementable in any size-resolved aerosol CTM. The new parameterization produces Λ(d) values similar to the upper range (90^{th} percentile) of an ensemble of theoretical Λ(d) values generated using combinations of all available product-term formulas and is closer than the majority of theoretical Λ(d) formulas in terms of comparisons with field-derived Λ(d) values. The power-law relationship obtained in this study between $\Lambda(d)$ and precipitation intensity R appears to be comparable to empirical power-law relationships obtained from experimental measurements. The new parameterization produces faster removal of atmospheric aerosol particles by snow scavenging than by rain scavenging for equivalent precipitation intensity, a result in qualitative agreement with evidence from a limited number of field experiments. However, due to the large uncertainties in theoretical Λ formulations, the large gaps between theoretical and field-based Λ values, and the very limited existing data base of field measurements of below-cloud scavenging of size-resolved aerosol particles, especially for snow conditions, more experimental studies are needed at more locations under more climate regimes and for a wider range of aerosol particle sizes to improve our understanding of scavenging processes and to further improve Λ formulations.

Appendix A

Laakso et al. (2003) empirical parameterization for $\Lambda_{rain}(d)$

Laakso et al. (2003) suggested a parameterization for $\Lambda_{rain}(d)$ based on their analysis of six years of field measurements over forests in southern Finland:

$$\log_{10}\Lambda(d) = a_1 + a_2[\log_{10} d]^{-4} + a_3[\log_{10} d]^{-3} + a_4[\log_{10} d]^{-2} + a_5[\log_{10} d]^{-1} + a_6 R^{1/2} , \quad \text{(A1)}$$

where d is particle diameter (in m), a_1=274.35758, a_2=332839.59273, a_3=226656.57259, a_4=58005.91340, a_5=6588.38582, a_6=0.244984, R is rainfall intensity (in mm h^{-1}). The formula is valid only for limited ranges of particle diameters 0.01- 0.5 μm and for rain intensities 0-20 mm h^{-1}.

Appendix B

Henzing et al. (2006) $\Lambda_{rain}(d)$ formula fitted from comprehensive numerical simulation

Henzing et al. (2006) developed a simple Λ_{rain} parameterization that represents below-cloud scavenging coefficients as a function of aerosol particle size and rainfall intensity. The parameterization is a simple three-parameter fit through below-cloud scavenging coefficients calculated at high particle size resolution. The calculations were based on the concept of collection efficiency between polydisperse aerosol particles and raindrop distributions. Specifically, Slinn's semi-empirical formula was used for the raindrop-particle collection efficiency. The gamma-function fit of de Wolf (2001) and the empirical formula of Atlas et al. (1973) were applied to represent the raindrop size distribution and the terminal fall velocity, respectively. The parameterization has been applied in a global chemical transport model. The final fitting function has the form

$$\Lambda(d) = A_0\left(e^{A_1 R^{A_2}} - 1\right), \quad \text{(B1)}$$

where the parameters A_0, A_1 and A_2 are provided in a table that is available at http://www.knmi.nl/~velthove/wet_deposition/coefficients.txt.

Appendix C

The empirical $\Lambda_{snow}(d)$ formula from Paramonov et al. (2011)

Paramonov et al. (2011) proposed a Λ_{snow} parameterization from the empirical fit to field measurements from four winters (2006-2010) in an urban environment in Helsinki, Finland:

$$\Lambda(d) = 10^{a_1 + a_2[\log_{10} d]^{-2} + a_3[\log_{10} d]^{-1}} + g \cdot (RH) - h, \quad \text{(C1)}$$

where d is particle diameter (in m), a_1=28.0, a_2=1550.0, a_3=456.0, g=0.00015, h=0.00013, and RH is relative humidity. The formula is only valid for aerosol particles of 0.01–1.0 μm in diameter and snowfall intensities of 0.1 to 1.2 mm h^{-1} (as liquid water equivalent). Nevertheless, the formula is applicable to snowfall episodes of snowflakes, snow grains, snow crystals, ice pellets, as well as snow mixed with rain.

Appendix D

The empirical $\Lambda_{snow}(d)$ formula from Kyrö et al. (2009)

Kyrö et al. (2009) suggested a size-resolved Λ_{snow} parameterization from an empirical fit to four years (2005-2008) of field measurements in a rural background environment in Finland:

$$\Lambda(d) = 10^{a_1 + a_2[\log_{10} d]^{-2} + a_3[\log_{10} d]^{-1}} , \tag{D1}$$

where d is particle diameter (in m), a_1=22.7, a_2=1321.0, and a_3=381.0. The parameterization applies to snowfall types of light continuous snowfall and snow grains with intensities of the order of 0.1 mm h^{-1} (as liquid water equivalent) and to aerosol particles of 0.01–1.0 μm in diameter.

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Table 1. List of semi-empirical formulas for raindrop–aerosol particle collection efficiency $E(d, D_p)$. Symbols used in Tables 1-9 and their units are defined in Table 10.

Source	**Formula**
Slinn (1984)[(a)]	$E(d,D_p)=\frac{4}{\mathrm{Re}\,Sc}\left[1+0.4\,\mathrm{Re}^{1/2}\,Sc^{1/3}+0.16\,\mathrm{Re}^{1/2}\,Sc^{1/2}\right] + 4\frac{d}{D_p}\left[\frac{\mu_a}{\mu_w}+\left(1+2\,\mathrm{Re}^{1/2}\right)\frac{d}{D_p}\right]+\left(\frac{St-St^*}{St-St^*+2/3}\right)^{3/2}$
Andronache et al.(2006)[(b)]	$E(d,D_p)=\frac{4}{\mathrm{Re}\,Sc}\left[1+0.4\,\mathrm{Re}^{1/2}\,Sc^{1/3}+0.16\,\mathrm{Re}^{1/2}\,Sc^{1/2}\right] + 4\frac{d}{D_p}\left[\frac{\mu_a}{\mu_w}+\left(1+2\,\mathrm{Re}^{1/2}\right)\frac{d}{D_p}\right]+\left(\frac{St-St^*}{St-St^*+2/3}\right)^{3/2} + E_{th}(d,D_p)+E_{dph}(d,D_p)+E_{es}(d,D_p)$ $E_{th}(d,D_p)=\frac{4\alpha_{th}\left(2+0.6\,\mathrm{Re}^{1/2}\,\mathrm{Pr}^{1/3}\right)(T_a-T_s)}{V_D D_p}$ $E_{dph}(d,D_p)=\frac{4\beta_{dph}\left(2+0.6\,\mathrm{Re}^{1/2}\,Sc_w^{1/3}\right)\left(\frac{P_s^0}{T_s}-\frac{P_a^0 RH}{T_a}\right)}{V_D D_p}$ $E_{es}(d,D_p)=\frac{16KC_cQ_rq_p}{3\pi\mu_a V_D D_p^2 d}$
Park et al. (2005)	Brownian diffusion and interception from Jung and Lee (1998) Initial impaction from Calvert (1984)
Croft et al. (2009)	Brownian diffusion from Young (1993) Impaction from a modified Hall (1980) table
Ackerman et al. (1995)	Brownian diffusion from Fuchs (1964) Impaction from Hall (1980) table

(a) The formula takes into account the three most important collection mechanisms for below-cloud particle scavenging. The first term represents Brownian diffusion, the second term represents interception, and the third term represents inertial impaction. *Re* is the Reynolds number: $\mathrm{Re}=D_pV_D\rho_a/2\mu_a$. *Sc* is the Schmidt number: $Sc=\mu_a/\rho_aD_{diff}$, where $D_{diff}=k_bT_aC_c/(3\pi\mu_a d)$ with the Cunningham correction factor: $C_c=1+\frac{2\lambda_a}{d}\left(1.257+0.4\exp\left(-\frac{0.55d}{\lambda_a}\right)\right)$. *St* is the Stokes number: $St=2\tau(V_D-v_d)/D_p$ with the characteristic relaxation time of a particle: $\tau = (\rho_p - \rho_a)d^2 C_c / 18\mu_a$. St^* is the critical Stokes number expressed as: $St^* = \frac{1.2 + \ln(1+\mathrm{Re})/12}{1 + \ln(1+\mathrm{Re})}$.

(b) The formula takes into account three additional collection mechanisms due to thermophoresis $E_{th}(d, D_p)$, diffusiophoresis $E_{dph}(d, D_p)$, and electrostatic forces $E_{es}(d, D_p)$ based on Slinn (1984). The parameter α_{th} and the Prandtl number Pr in $E_{th}(d, D_p)$ are defined as, $\alpha_{th} = \frac{2C_c(k_a + 5\lambda_a / D_p k_p)k_a}{5P(1 + 6\lambda_a / D_p)(2k_a + k_p + 10\lambda_a / D_p k_p)}$, and $\mathrm{Pr} = c_p \mu_a / k_a$, respectively. The parameter β_{dph} and the Schmidt number for water Sc_w in $E_{dph}(d, D_p)$ are defined as $\beta_{dph} = \frac{T_a D_{diffwater}}{P}\sqrt{\frac{M_w}{M_a}}$, and $Sc_w = \mu_a / \rho_a D_{waterdiff}$, respectively. The parameter K in $E_{es}(d, D_p)$ is set as 9 x 10^9 (in $\mathrm{Nm}^2\ \mathrm{C}^{-2}$). Q_r and q_p are the mean charges on the raindrop and on the aerosol particle (in Coulomb, C), respectively, with opposite sign, and are parameterized as $Q_r = a\alpha D_p^2$ and $q_p = a\alpha d^2$ with a = 0.83 x 10^{-6} and α (C m^{-2}), an empirical parameter, in the range of 0-7 corresponding to cloud charges from neutral to highly electrified clouds.

Table 2. List of raindrop number size distribution ($N(D_p)$) formulas. The general forms of the (a) exponential, (b) gamma, and (c) lognormal distributions are commonly written as $N(D_p) = N_{0e}\exp(-\beta_e D_p)$, $N(D_p) = N_{0g} D_p^{\gamma}\exp(-\beta_g D_p)$, and $N(D_p) = \frac{N_{total}}{\sqrt{2\pi} D_p \ln(\sigma_D)} \exp\left[-\frac{(\ln(D_p) - \ln(D_{mean}))^2}{2(\ln(\sigma_D))^2}\right]$, respectively. See Table 10 for definitions of other symbols and units.

Raindrop number size spectrum	**Formula definition**	**Rain type**	**Source**
Exponential distributions(a)	$N_{0e} = 0.08$, $\beta_e = 41R^{-0.21}$	Widespread	Marshall and Palmer(1948)
	$N_{0e} = 0.30$, $\beta_e = 57R^{-0.21}$	Drizzle	Joss et al. (1968)
	$N_{0e} = 0.014$, $\beta_e = 30R^{-0.21}$	Thunderstorm	Joss et al. (1968)
	$N_{0e} = 0.07R^{0.37}$, $\beta_e = 38R^{-0.14}$	Thunderstorm	Sekhon and Srivastava (1971)
	$N_{0e} = 0.071M^{0.648}$, $\beta_e = \left(\frac{10^{-6}\rho_w \pi N_{0e}}{M}\right)^{0.25}$ $M = 0.0626R^{0.913}$	Convective	Zhang et al. (2008)
Gamma distributions(b)	$N_{0g} = 168.53R^{-0.384}$ $\gamma = 2.93$, $\beta_g = 53.8R^{-0.186}$	Widespread	de Wolf (2001)
	$N_{0g} = \frac{6.36\times10^{-4}M}{d_0^4}\left(\frac{1}{d_0}\right)^{2.5}$ $\gamma = 2.50$, $\beta_g = 5.57/d_0$ $d_0 = 0.157M^{0.168}$, $M = 0.062R^{0.913}$	Hurricane	Willis (1984)
	$N_{0g} = \frac{5.1285\times10^{-4}M}{d_0^4}\left(\frac{1}{d_0}\right)^{2.16}$ $\gamma = 2.16$, $\beta_g = 5.588/d_0$ $d_0 = 0.1571M^{0.1681}$, $M = 0.062R^{0.913}$	Hurricane	Willis and Tattelman (1989)

Lognormal distributions[c]	$N_{total} = 1.72 \times 10^{-4} R^{0.22}$, $D_{mean} = 0.072 R^{0.23}$ $\sigma_D = 1.43 - 3.0 \times 10^{-4} R$	Widespread	Feingold and Levin (1986)
	$N_{total} = 1.94 \times 10^{-4} R^{0.30}$, $D_{mean} = 0.063 R^{0.23}$ $\sigma_D = e^{\sqrt{0.191 - 1.1 \times 10^{-2} \cdot \ln(R)}}$	Widespread	Cerro et al. (1997)

Table 3. List of empirical and theoretical raindrop terminal velocity (V_D) formulas.

Type	**Formula**	**Source**
Empirical formulas	$V_D = 1300 D_p^{0.5}$	Kessler (1969)
	$V_D = 1767 D_p^{0.67}$	Atlas and Ulbrich (1977)
	$V_D = 4854 D_p \exp\left(-1.95 D_p\right)$	Willis (1984)
	$V_D = 958\left[1-\exp\left(-\left(\frac{D_p}{0.171}\right)^{1.147}\right)\right]$	Best (1950)
	$V_D = -10.21 + 4932 D_p - 9551 D_p^2 + 7934 D_p^3 - 2362 D_p^4$	Brandes et al. (2002)
	$V_D = \begin{cases} 0 & D_p \le 0.003 \\ 4323(D_p - 0.003) & 0.003 \le D_p \le 0.06 \\ 965 - 1030\exp\left(-6D_p\right) & D_p > 0.06 \end{cases}$	Henzing et al. (2006)
Theoretical formulas	Beard's scheme	Beard (1976)
	Feng's Scheme	Feng (2007)

Table 4. List of semi-empirical formulas for snow particle–aerosol particle collection efficiency E.

Source	Formula
Slinn (1984)[(a)]	$E(d,\lambda)=(\frac{1}{Sc})^{\alpha_\lambda}+\left[1-\exp(-(1+\mathrm{Re}_\lambda^{1/2}))\frac{(d/2)^2}{\lambda^2}\right]+\left(\frac{St-St^*}{St-St^*+2/3}\right)^{3/2}$
Murakami et al (1985)[(b)]	$E(d,D_m)=\frac{48D_{diff}}{\pi D_m V_D}(0.65+0.44Sc^{1/3}\,\mathrm{Re}^{1/2})+28.5I^{1.186}+\left(\frac{S_1-S_2}{S_2\exp(S_1t')-S_1\exp(S_2t')}\right)^2$
Dick (1990)[(c)]	$E(d,D_m)=\frac{2mV_D}{3\pi d\mu_a D_m}+\frac{4}{Pe}(1+0.4\,\mathrm{Re}^{1/6}\,Pe^{1/3})$

(a) λ is the characteristic capture length and α_λ is an empirical constant. Both λ and α_λ depend on the shape of snow particles (e.g., sleet/graupel, rimed crystals, powder snow, dendrite, tissue paper, and camera film). Re_λ is the Reynolds number corresponding to the specific λ, Sc is the Schmidt number: $Sc=\mu_a/\rho_a D_{diff}$, St is the Stokes number, and St^* is the critical Stokes number: $St^*=\frac{1.2+(1/12)\ln(1+\mathrm{Re}_\lambda)}{1+\ln(1+\mathrm{Re}_\lambda)}$.

(b) The formula is for snow aggregates. The Reynolds number of a snow particle is defined as: $\mathrm{Re}=D_m V_D \rho_a/\mu_a$, Sc is the Schmidt number, and I is the size ratio d/D_c with D_c the characteristic length of the snow particle. The third term is the theoretical solution of a simplified flow model by Ranz and Wong (1952), involving parameters S_1, S_2 and t', and it can be simplified to $\exp(\frac{-0.11}{St^{1/2}-0.25})$ if $St \geq 1/16$, or to 0 if $St < 1/16$ (Feng, 2009).

(c) Pe is the Peclet number: $Pe=D_m V_D/D_{diff}$ and Re is the Reynolds number: $\mathrm{Re}=D_m V_D \rho_a/2\mu_a$.

Table 5. List of exponential snow particle number size distribution ($N(D_p)$) formulas. Note that actual snow particle size D_m (cm) was used in Scott (1982) (see Appendix *A* in Zhang et al., 2013) whereas D_p was used in other formulas.

$$N(D_p) = N_{0e} \exp\left(-\beta_e D_p\right)$$

Source	N_{0e} **[cm^{-4}]**	β_e **[cm^{-1}]**
Marshall and Palmer (1948)	0.08	$\beta_e = 41R^{-0.21}$
Scott (1982)	0.5	$M = 0.37R^{0.94}$ $\beta_e = 20.7M^{-0.33} = 28.8R^{-0.31}$
Gunn and Marshall (1958)	$N_{0e} = 0.038R^{-0.87}$	$\beta_e = 25.5R^{-0.48}$
Sekhon and Srivastava (1970)	$N_{0e} = 0.025R^{-0.94}$	$\beta_e = 22.9R^{-0.45}$

Table 6. List of empirical and theoretical snow particle terminal velocity (V_D) formulas. X is the Best number: $X = \frac{2mg\rho_a D_m^2}{A\mu_a^2}$, α, β, δ and σ are empirical constants (see Table 7), and a_1 and b_1 are described as functions of X (see Mitchell and Heymsfield, 2005).

Source	**V_D formula**	**Particle shape**
Langleben (1954)	V_D=207.0$D_p^{0.310}$	plane dendrite
Jiusto and Bosworth (1971)	V_D=104.9$D_m^{0.206}$	plane dendrite
Locatelli and Hobbs (1974)	V_D=64.80$D_m^{0.257}$	plane dendrite
Molthan et al. (2010)	V_D=110.1$D_m^{0.145}$	plane dendrite
Jiusto and Bosworth (1971)	V_D=153.0$D_m^{0.206}$	column
Matson and Huggins (1980)	V_D=1145$D_p^{0.500}$	graupel
Mitchell (1996)	$V_D = \frac{\mathrm{Re}\,\mu_a}{D_m \rho_a}$ $\mathrm{Re} = \begin{cases} 0.04394X^{0.970}, 0.01 < X \le 10.0 \\ 0.06049X^{0.831}, 10.0 < X \le 585 \\ 0.2072X^{0.638}, 585 < X \le 1.56\times10^5 \\ 1.0865X^{0.499}, 1.56\times10^5 < X \le 10^8 \end{cases}$	any shape
Mitchell and Heymsfield (2005)	$V_D = a_v D_m^{b_v}$, $\mathrm{Re} = a_1 X^{b_1}$, $m = \alpha D_m^{\beta}$, $A = \delta D_m^{\sigma}$ $a_v = a_1 \left(\frac{\mu_a}{\rho_a}\right)^{(1-2b_1)} \left(\frac{2\alpha g}{\rho_\alpha \delta}\right)^{b_1}$, $b_v = b_1(\beta - \sigma + 2) - 1$	any shape

Table 7. Snow particle shapes considered in this study and their mass (m) and cross-sectional area (A) formulas.

Snow particle shape	**Mass** $m = \alpha D_m^{\beta}$ [g]	**Cross-sectional area** $A = \delta D_m^{\sigma}$ [cm^2]
Spheres	$m = 0.0524D_m^{3.00a}$	$A = 0.7854D_m^{2.00a}$
Dendrites	$m = 0.0022D_m^{2.19b}$	$A = 0.2285D_m^{1.88c}$
Columns	$m = 0.0450D_m^{3.00b}$	$A = 0.0512D_m^{1.41d}$
Graupel	$m = 0.0490D_m^{2.80e}$	$A = 0.5000D_m^{2.00e}$

[a] Obtained from $m=\rho_s(\pi/6)\,D_m^3$ and $A=(\pi/4)\,D_m^2$, with $\rho_s = 0.1$ g cm^{-3}.
[b] From Woods et al. (2008)
[c] From Mitchell (1996) for "Aggregates of side planes"
[d] From Mitchell (1996) for "Rimed long columns"
[e] From Mitchell (1996) for "Lump graupel"

Table 8. Empirical constants in the formulations of $\log_{10}$*(A(d))* and *B(d)* for Λ_{rain} and Λ_{snow} parameterizations.

Constants in Λ_{rain} parameterization							
$\log_{10}(A(d))$	a_0= -6.2609 x 10^{0}	a_1= 6.8200 x 10^{-1}	a_2= 8.6760 x 10^{-1}	a_3= 1.2820 x 10^{-1}			
	b_0= -1.4707 x 10^{1}	b_1= 5.1043 x 10^{1}	b_2= -9.7306 x 10^{1}	b_3= 9.7946 x 10^{1}	b_4= -5.3923 x 10^{1}	b_5= 1.5311 x 10^{1}	b_6 = -1.7510 x 10^{0}
B(d)	c_0= 7.2300 x 10^{-1}	c_1= 3.0300 x 10^{-2}					
	e_0= -6.4920 x 10^{-1}	e_1= 9.3483 x 10^{0}	e_2= -2.1929 x 10^{1}	e_3= 2.5317 x 10^{1}	e_4= -1.5395 x 10^{1}	e_5= 4.7242 x 10^{0}	e_6 = -5.7660 x 10^{-1}
Constants in Λ_{snow} parameterization							
$\log_{10}(A(d))$	a_0= -4.4260 x 10^{0}	a_1= 1.3940 x 10^{0}	a_2= -1.2020 x 10^{0}	a_3= -3.2942 x 10^{0}	a_4= -1.9521 x 10^{0}	a_5= -4.9040 x 10^{-1}	a_6= -4.5700 x 10^{-2}
	b_0= -4.3531 x 10^{0}	b_1= -7.8280 x 10^{-1}	b_2= 1.2768 x 10^{1}	b_3= -1.9864 x 10^{1}	b_4= 1.3618 x 10^{1}	b_5= -4.4350 x 10^{0}	b_6 = 5.5510 x 10^{-1}
B(d)	c_0= 5.6640 x 10^{-1}	c_1= 8.5000 x 10^{-3}	c_2= -1.9480 x 10^{-1}	c_3= -6.5320 x 10^{-1}	c_4= -5.462 x 10^{-1}	c_5= -1.7780 x 10^{-1}	c_6 = -2.0100 x 10^{-2}
	e_0= 5.6890 x 10^{-1}	e_1= -9.2300 x 10^{-2}	e_2= 4.0200 x 10^{-2}	e_3= 1.4523 x 10^{0}	e_4= -2.0780 x 10^{0}	e_5= 1.0500 x 10^{0}	e_6 = -1.8210 x 10^{-1}

Table 9. List of below-cloud Λ_{rain} and Λ_{snow} parameterizations from literature expressed as $\Lambda(d,R) = A(d)R^{B(d)}$ (where Λ is in units of s^{-1})

Λ (s^{-1})	Source	$A(d)$ (s^{-1})	$B(d)$	Aerosol diameter range (μm)	Calculation basis
Λ_{rain}	Jylhä (1991)	1.0×10^{-4}	0.64	0.3 - 0.9	Field measurements
	Okita et al. (1996)	1.0×10^{-4}	0.67 - 0.76	d > 2.0	
	Sparmacher et al. (1993)	2.34×10^{-7} 3.14×10^{-7} 2.56×10^{-7} 1.72×10^{-6}	0.59 0.60 0.94 0.61	0.23 0.46 0.98 2.16	Controlled experiment
	Scott (1982)	3.56×10^{-4}	0.78	10.0	Theoretical calculation
	Mircea et al. (1998)[a]	$2.43 \times 10^{-4}\,E - 7.41 \times 10^{-3}\,E$	0.78 - 0.86	Any sizes	
	Baklanov and Sorensen (2001)	8.40×10^{-5}	0.79	d < 2.8	
	Andronache (2003)	$2.78 \times 10^{-8} - 1.39 \times 10^{-6}$ $6.67 \times 10^{-5} - 2.44 \times 10^{-4}$	0.59 - 0.61 0.7	d < 2.0 d > 2.0	
	Feng (2007)	$1.19 \times 10^{-6} - 2.06 \times 10^{-6}$ $2.36 \times 10^{-7} - 3.69 \times 10^{-7}$ $2.11 \times 10^{-4} - 3.42 \times 10^{-4}$ $4.92 \times 10^{-4} - 5.06 \times 10^{-4}$	0.62 0.61 - 0.62 0.79 0.81 - 0.82	0.001 - 0.04 0.04 - 2.5 2.5 - 16.0 16.0 - 100.	

	This work	6.16×10^{-6} - 1.17×10^{-4} 3.83×10^{-7} - 6.16×10^{-6} 9.80×10^{-7} - 6.70×10^{-5} 6.75×10^{-5} - 6.89×10^{-4}	0.64 - 0.67 0.67 - 0.72 0.72 - 0.91 0.82 - 0.91	0.001 - 0.01 0.01 - 2.0 2.0 - 3.0 3.0 - 100.0	
Λ_{snow}	Sparmacher et al. (1993)	1.60×10^{-6} 8.10×10^{-7} 3.49×10^{-6}	0.62 0.89 1.09	0.46 0.98 1.66	Controlled experiment
	Mircea et al. (1998) [(a)]	$2.44 \times 10^{-3}\, E$ - $3.59 \times 10^{-2}\, E$	0.89 - 1.14	Any sizes	Theoretical calculation
	Baklanov and Sorensen (2001)	8.0×10^{-5}	0.31	Any sizes	
	Scott (1982)	2.44×10^{-4}	1.0	10.0	
	This work	6.84×10^{-5} - 1.80×10^{-3} 4.94×10^{-6} - 1.19×10^{-4} 1.19×10^{-4} - 5.13×10^{-4} 5.13×10^{-4} - 5.50×10^{-3}	0.66 - 0.77 0.53 - 0.66 0.58 - 0.61 0.61 - 0.86	0.001-0.01 0.01 - 2.0 2.0 -3.0 3.0 - 100.0	

[(a)] E is the collection efficiency and assumed to be a constant for a given precipitation distribution and aerosols types.

Table 10. Nomenclature. Note that CGS units are used in all of the equations and tables except Eqs. 6-9 because many empirical formulas in Tables 1-7 were developed based on CGS units.

A	hydrometeor-particle effective cross-sectional area projected normal to the fall direction (cm^2)
C_c	Cunningham correction factor
c_p	heat capacity of air (cm^2 s^{-2} K^{-1})
d	aerosol particle diameter (cm)
D_c	snow-particle characteristic length used in E formula of Murakami et al. (1985) (cm)
D_{diff}	aerosol-particle diffusivity coefficient (cm^2 s^{-1})
D_m	maximum dimension of a snow particle (cm)
D_{mean}	mean diameter of lognormal spectra (cm)
D_p	raindrop or melted snow-particle diameter (cm)
$D_{waterdiff}$	water vapour diffusivity in air (cm^2 s^{-1})
$E(d, D_p)$	overall hydrometeor-aerosol particle collection efficiency
$E_{dph}(d, D_p)$	collection efficiency due to diffusiophoresis
$E_{es}(d, D_p)$	collection efficiency due to charge effect
$E_{th}(d, D_p)$	collection efficiency due to thermophoresis
g	acceleration of gravity (cm s^{-2})
k_a	thermal conductivity of air (erg cm^{-1} s^{-1} K^{-1})
k_b	Boltzmann constant (erg K^{-1})
k_p	thermal conductivity of particle (erg cm^{-1} s^{-1} K^{-1})
m	particle mass (g)
M	precipitation water concentration (g m^{-3})
M_a	air molecular weight
M_w	water vapour molecular weight
$n(d, t)$	aerosol number concentration with diameters d at time t
$N(D_p)$	number size distribution of precipitation hydrometeors (cm^{-4})
N_{0e}	intercept parameter for exponential size distribution (cm^{-4})
N_{0g}	intercept parameter for gamma size distribution ($cm^{-\gamma-1}$ cm^{-3})
N_{total}	total number concentration of precipitation hydrometeors (cm^{-3})
P	atmospheric pressure (dyne)
Pe	Peclet number
Pr	Prandtl number for air
P_a^o	vapour pressure of water at temperature T_a (dyne)
P_s^o	vapour pressure of water at temperature T_s (dyne)
q_p	mean charge of a particle (C)
Q_r	mean charge of a raindrop (C)
R	precipitation intensity (mm h^{-1})
Re	Reynolds number
RH	relative humidity (%)
Sc	Schmidt number for aerosol particle
Sc_w	Schmidt number for water in air
St	Stokes number of aerosol particle
St^*	critical Stokes number of aerosol particle
T_a	air temperature (K)
T_s	raindrop surface temperature (K)
v_d	aerosol-particle terminal velocity (cm s^{-1})
V_D	raindrop or snow-particle terminal velocity (cm s^{-1})
X	Davies number
α, β	empirical constants in mass-diameter power-law relationships

δ, σ	empirical constants in area-diameter power-law relationships
β_e	slope parameter for exponential size distribution
β_a	slope parameter for gamma size distribution
γ	shape parameter for gamma size distribution
λ	snow-particle characteristic capture length used in *E* formula of Slinn (1984) (cm)
λ_a	mean free path of air molecules (cm)
$\Lambda(d)$	size-resolved aerosol-particle scavenging coefficient (s^{-1})
μ_a	dynamic air viscosity (g cm^{-1} s^{-1})
μ_w	water viscosity (g cm^{-1} s^{-1})
ρ_a	air density (g cm^{-3})
ρ_p	aerosol-particle density (g cm^{-3})
ρ_w	water density(g cm^{-3})
σ_D	standard deviation of lognormal size distribution
τ	characteristic relaxation time of a particle (s)

List of Figures

Figure 1. Size-resolved scavenging coefficients for rain conditions: (a) Λ_{rain} calculated using Eq. (2) from a total of 400 combinations of different $E(d, D_p)$, $N(D_p)$, and V_{D} formulas listed in Tables 1, 2 and 3, respectively. The yellow group uses the $E(d, D_p)$ formula of Park et al. (2005) and the red group uses the $E(d, D_p)$ formula of Ackerman et al. (1995); the black group includes all the other combinations; (b) same as in (a) but with the yellow group removed and the red group using the modified $E(d, D_p)$ formula of Ackerman et al. (1995) (reduced to a total of 320 combinations); (c) minimum, maximum, and five percentile Λ_{rain} profiles (coloured lines) based on ensemble of profiles from (b), where dots are the data from (b); (d) lines are the same as (c) and symbols are experimental data reviewed in Wang et al. (2010). Also shown in (d) are one empirical Λ_{rain} parameterization of Laakso et al. (2003) (denoted by LA; see Appendix A) and one semi-empirical Λ_{rain} parameterization of Henzing et al. (2006) (denoted by HS; see Appendix B), which is an empirical fit to theoretically calculated Λ_{rain} values.

Figure 2. (a) 90^{th}-percentile Λ_{rain} profiles as a function of precipitation intensity R derived from an ensemble of 320 Λ_{rain} realizations for 100 particle diameters (a total of 100 lines); (b) linear regression best-fit lines for the $90^{\text{-th}}$-percentile Λ_{rain} data (symbols) from (a) for seven aerosol particle diameters ; (c) Values (symbols) of y-intercept *A(d)* from the log-linear regressions for 100 particle diameters and their polynomial best-fit curves (lines); and (d) same as in (c) but for the slope *B(d)* of the log-linear regressions.

Figure 3. Parameterized size-resolved Λ_{rain} profiles using Eqs. (5), (6) and (7) (solid lines) and the original 90th percentile Λ_{rain} data (symbols) (a) and their percentage differences (b) for five different precipitation intensities.

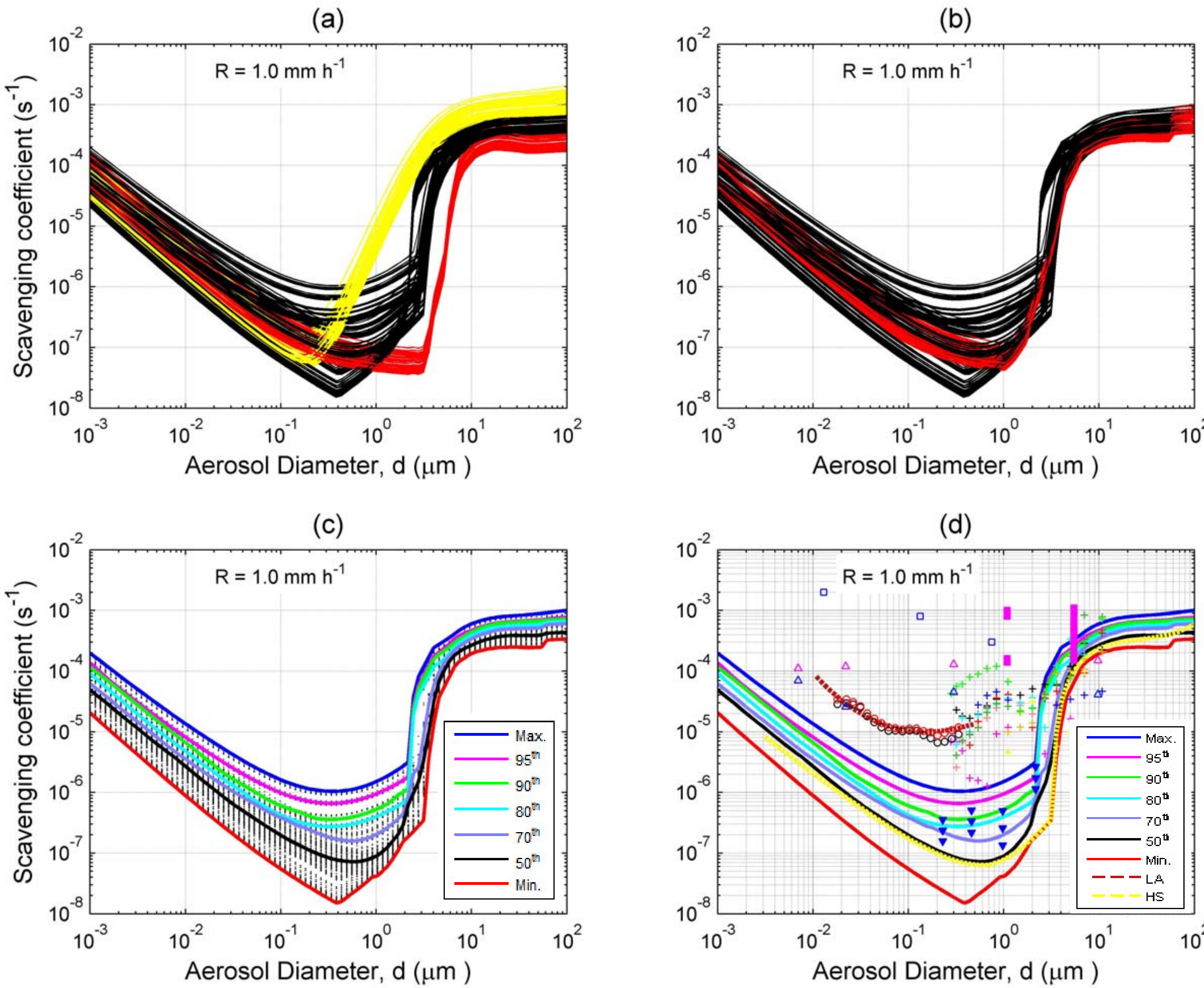
Figure 4. Size-resolved scavenging coefficient under snow conditions: (a) Λ_{snow} calculated using Eq. (2) from a total of 168 combinations of $E(d, D_p)$, $N(D_p)$, V_D and A listed in Tables 4, 5, 6 and 7, respectively; and (b) minimum, maximum, and five percentile Λ_{snow} profiles (coloured lines) based on ensemble of profiles from (a), where dots are the data from (a). Also shown in (b) are two empirical Λ_{snow} formulas of Paramonov et al. (2011) and Kyrö et al. (2009) (Appendices C and D, respectively).

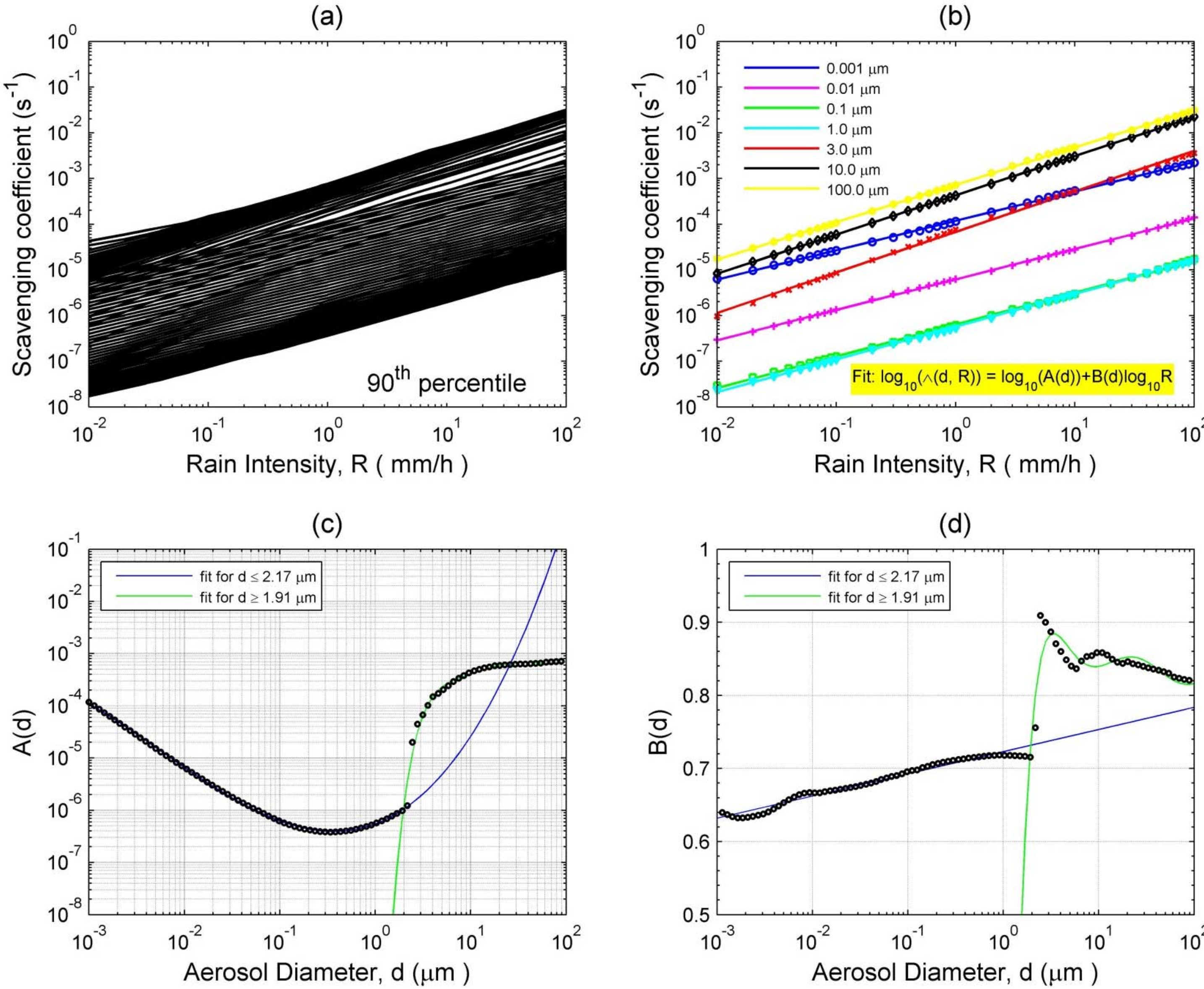
Figure 5. Same as in Fig. 2 except for Λ_{snow}

Figure 6. Same as in Fig. 3 except for Λ_{snow}

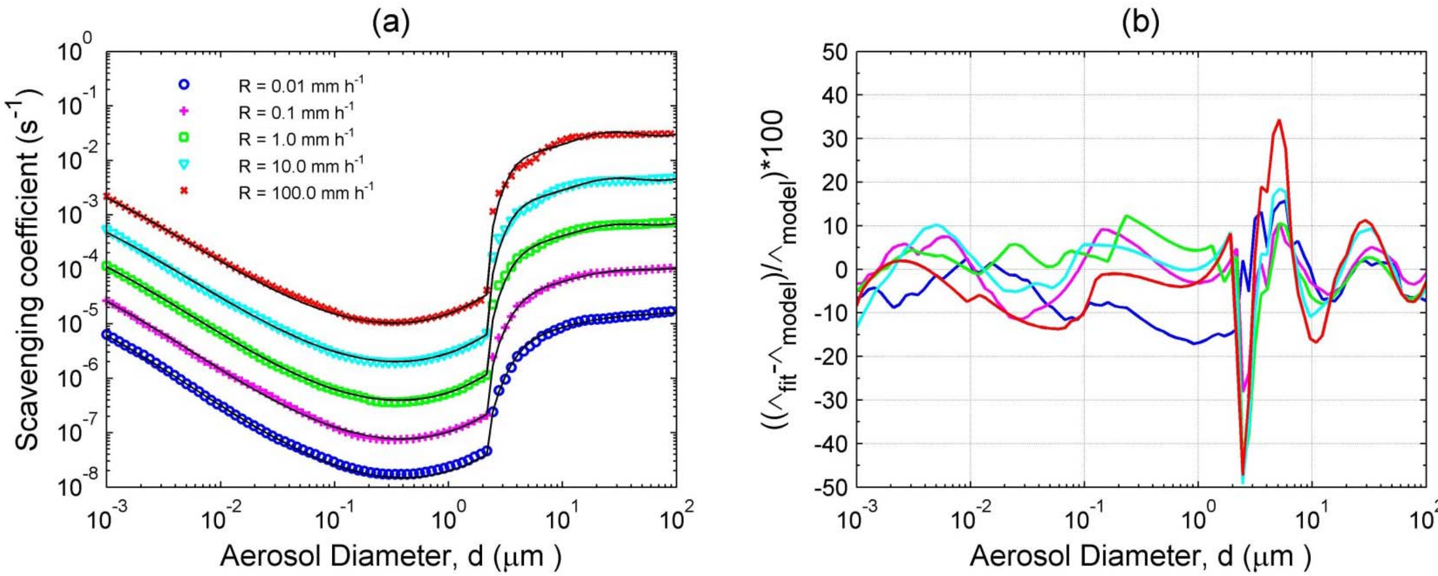
Figure 7. (a) The ratio of parameterized Λ_{snow} to Λ_{rain} as a function of precipitation intensity R (liquid water equivalent) for 100 aerosol particle diameters (100 lines in total). The groups of blue, yellow, and green lines correspond to aerosol particle diameters <0.1 µm, 0.1-5.0 µm, and >5.0 µm, respectively; (b) The ratio of parameterized Λ_{snow} to Λ_{rain} as a function of aerosol particle diameter d for four selected values R.

Figure 8. Percentage differences in Λ from the use of different temperature and pressure values for (a) rain and (b) snow scavenging versus base case. A precipitation intensity of 1.0 mm h^{-1} was assumed, and the base case refers to the ambient conditions used to develop the new Λ parameterization (i.e., p = 1013.25 hPa; T = 15°C for rain scavenging and T = -10°C for snow scavenging).

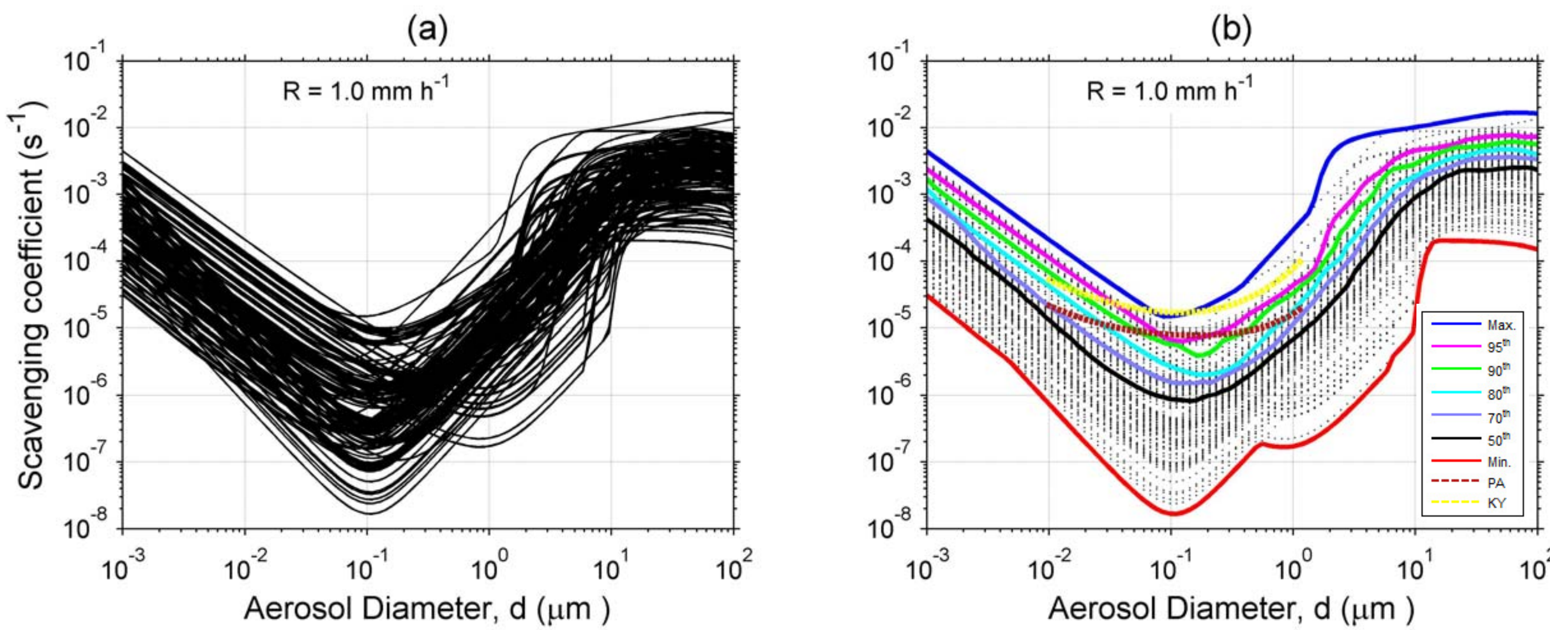

(a)
R = 1.0 mm h^{-1}
Scavenging coefficient (s^{-1})
Aerosol Diameter, d (μm)
(b)
R = 1.0 mm h^{-1}
Aerosol Diameter, d (μm)
(c)
R = 1.0 mm h^{-1}
Scavenging coefficient (s^{-1})
Max.
95th
90th
80th
70th
50th
Min.
Aerosol Diameter, d (μm)
(d)
R = 1.0 mm h^{-1}
Max.
95th
90th
80th
70th
50th
Min.
LA
HS
Aerosol Diameter, d (μm)

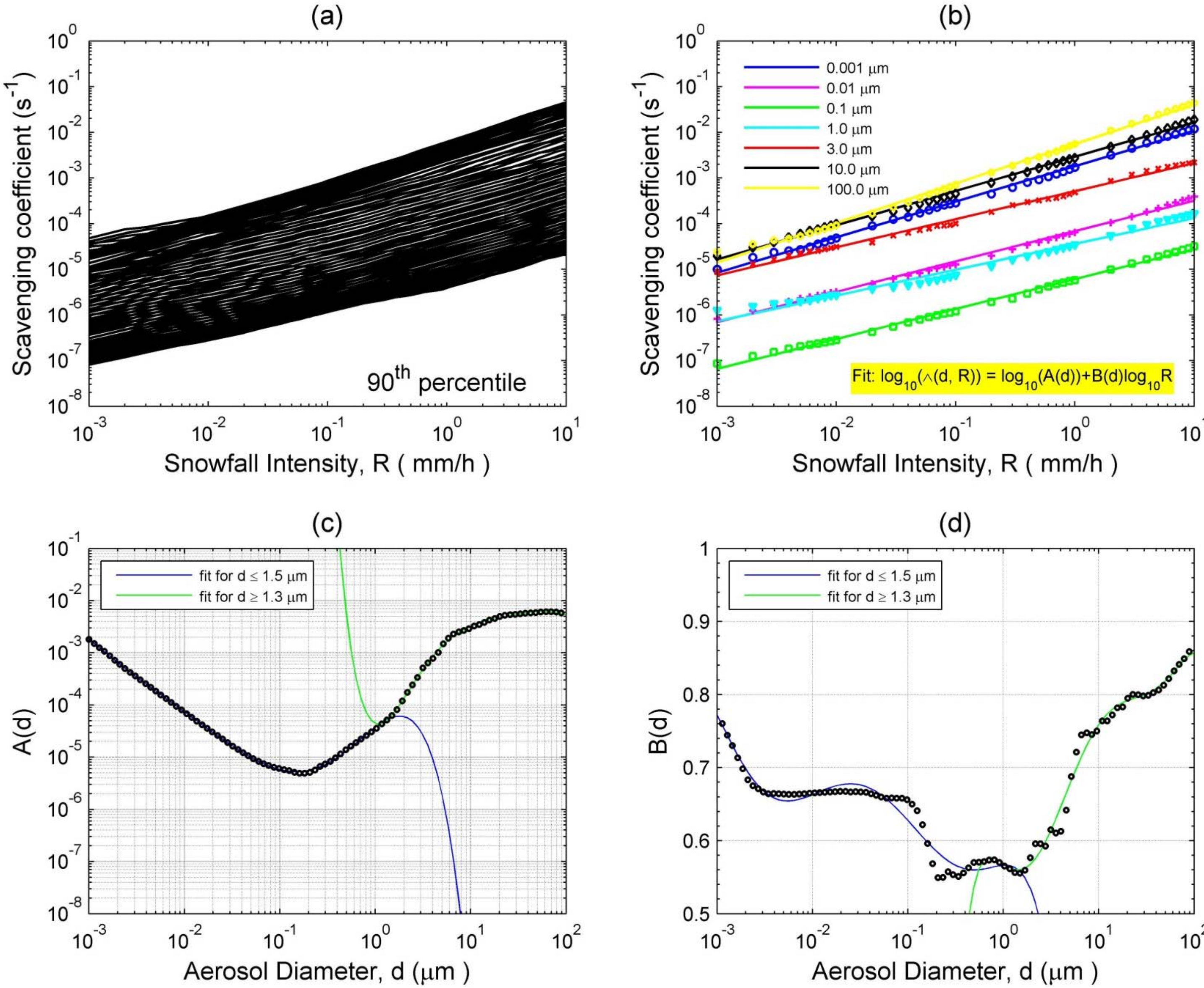


(a)
Scavenging coefficient (s^{-1})
90th percentile
Rain Intensity, R (mm/h)
(b)
0.001 μm
0.01 μm
0.1 μm
1.0 μm
3.0 μm
10.0 μm
100.0 μm
Fit: $\log_{10}(\wedge(d, R)) = \log_{10}(A(d))+B(d)\log_{10}R$
Scavenging coefficient (s^{-1})
Rain Intensity, R (mm/h)
(c)
fit for d ≤ 2.17 μm
fit for d ≥ 1.91 μm
A(d)
Aerosol Diameter, d (μm)
(d)
fit for d ≤ 2.17 μm
fit for d ≥ 1.91 μm
B(d)
Aerosol Diameter, d (μm)

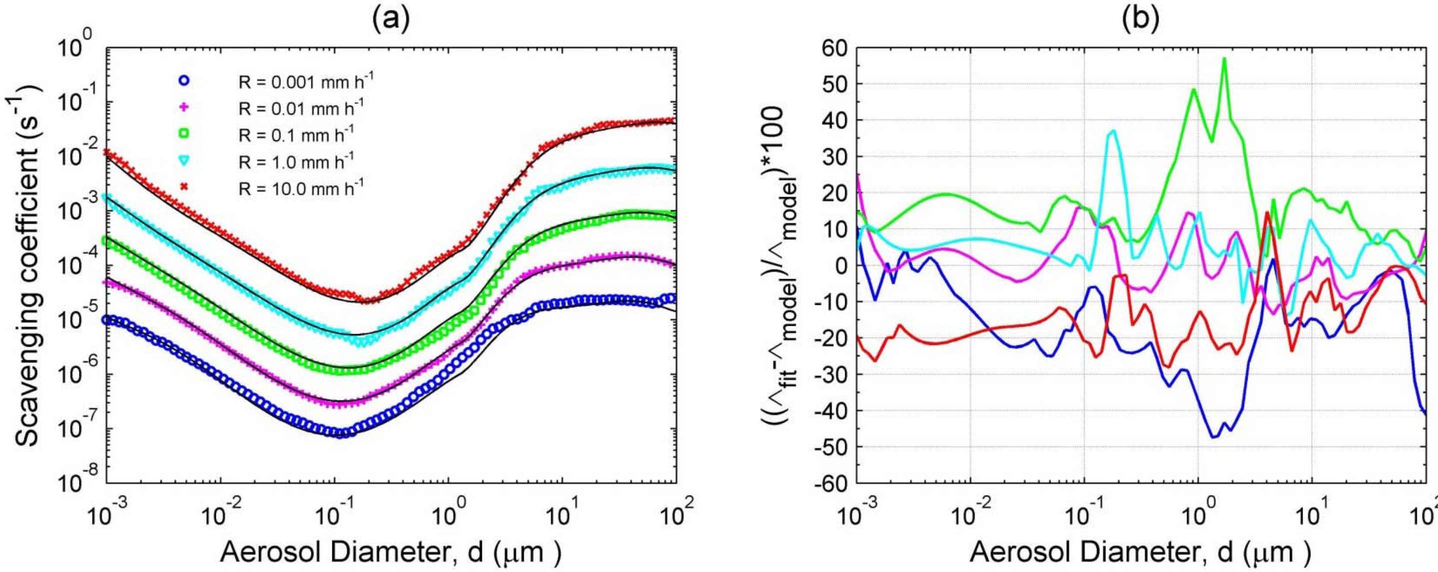


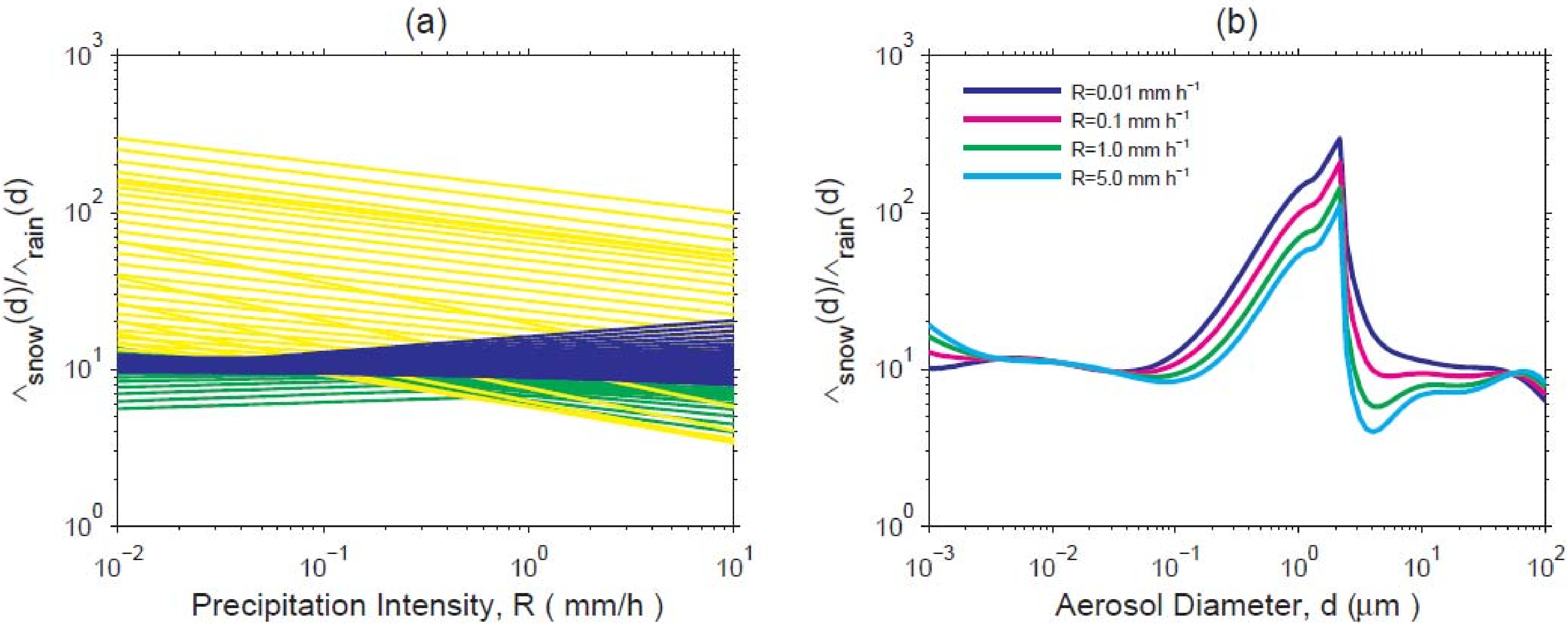
(a)
Scavenging coefficient (s^{-1})
R = 0.01 mm h^{-1}
R = 0.1 mm h^{-1}
R = 1.0 mm h^{-1}
R = 10.0 mm h^{-1}
R = 100.0 mm h^{-1}
Aerosol Diameter, d (μm)
(b)
$((\wedge_{fit}-\wedge_{model})/\wedge_{model})*100$
Aerosol Diameter, d (μm)


(a)
(b)
R = 1.0 mm h-1
Scavenging coefficient (s-1)
Aerosol Diameter, d (μm)
Max.
95th
90th
80th
70th
50th
Min.
PA
KY

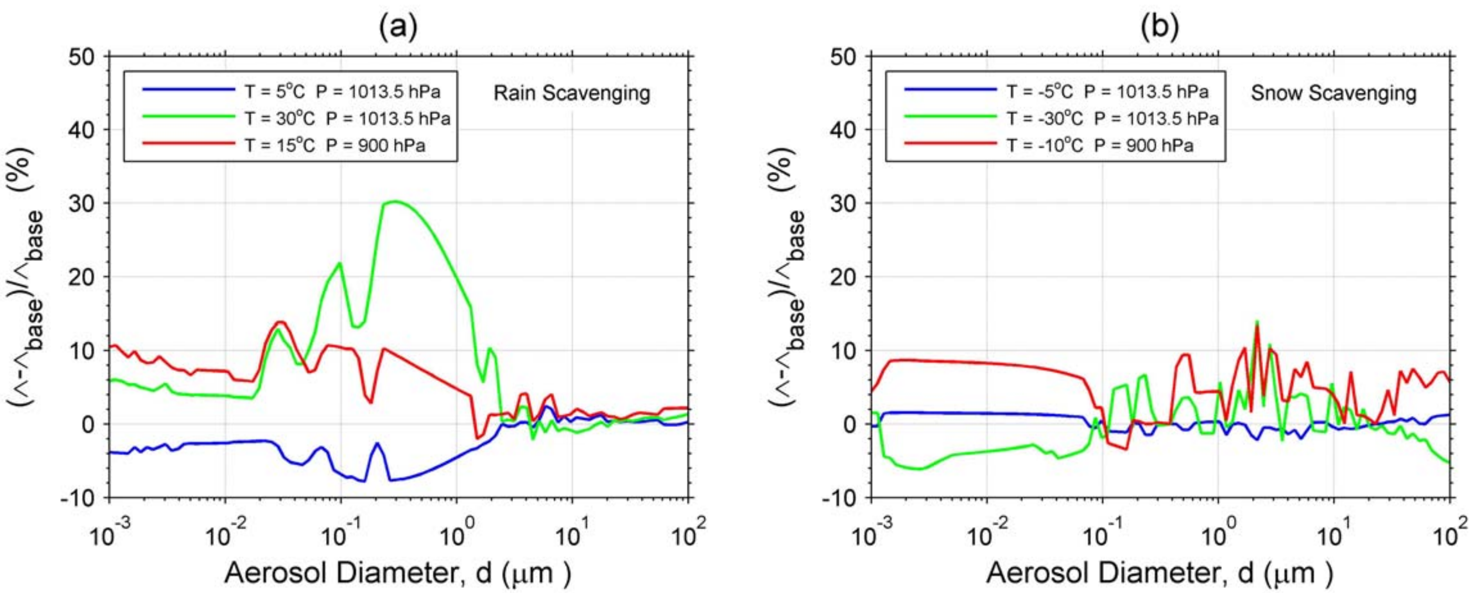


(a)
Scavenging coefficient (s^-1)
90th percentile
Snowfall Intensity, R (mm/h)
(b)
0.001 μm
0.01 μm
0.1 μm
1.0 μm
3.0 μm
10.0 μm
100.0 μm
Fit: log10(∧(d, R)) = log10(A(d))+B(d)log10R
(c)
fit for d ≤ 1.5 μm
fit for d ≥ 1.3 μm
A(d)
Aerosol Diameter, d (μm)
(d)
B(d)


(a)
(b)
Scavenging coefficient (s^{-1})
Aerosol Diameter, d (μm)
R = 0.001 mm h^{-1}
R = 0.01 mm h^{-1}
R = 0.1 mm h^{-1}
R = 1.0 mm h^{-1}
R = 10.0 mm h^{-1}
$((\Lambda_{fit}-\Lambda_{model})/\Lambda_{model})*100$
Aerosol Diameter, d (μm)



(a)
(b)
Λsnow(d)/Λrain(d)
Precipitation Intensity, R (mm/h)
Aerosol Diameter, d (μm)
R=0.01 mm h-1
R=0.1 mm h-1
R=1.0 mm h-1
R=5.0 mm h-1


(a)
Rain Scavenging
T = 5°C P = 1013.5 hPa
T = 30°C P = 1013.5 hPa
T = 15°C P = 900 hPa
(Λ-Λbase)/Λbase (%)
Aerosol Diameter, d (μm)
(b)
Snow Scavenging
T = -5°C P = 1013.5 hPa
T = -30°C P = 1013.5 hPa
T = -10°C P = 900 hPa
(Λ-Λbase)/Λbase (%)
Aerosol Diameter, d (μm)