Application of CCM SOCOL-AERv2-BE to cosmogenic beryllium isotopes: description and validation for polar regions

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Abstract. The short-living cosmogenic isotope \(^7\)Be, which is produced by cosmic rays in the atmosphere, is often used as a tracer for atmospheric dynamics, with precise and high-resolution measurements covering the recent decades. The long-living isotope \(^{10}\)Be, as measured in polar ice cores with an annual resolution, is a proxy for long-term cosmic-ray variability, whose signal can, however, be distorted by atmospheric transport and deposition that need to be properly modeled to be accounted for. While transport of \(^7\)Be can be modeled with high accuracy using the known meteorological fields, atmospheric transport of \(^{10}\)Be was typically modeled using case-study-specific simulations or simplified box models based on parameterizations. Thus, there is a need for a realistic model able to simulate atmospheric transport and deposition of beryllium with a focus on polar regions and (inter)annual timescales that is potentially able to operate in a self-consistent mode without the prescribed meteorology. Since measurements of \(^{10}\)Be are extremely laborious and hence scarce, it is difficult to compare model results directly with measurement data. On the other hand, the two beryllium isotopes are believed to have similar transport and deposition properties, being different only in production and lifetimes, and thus the results of \(^7\)Be transport can be generally applied to \(^{10}\)Be. Here we present a new model, called CCM SOCOL-AERv2-BE, to trace isotopes of \(^7\)Be and \(^{10}\)Be in the atmosphere based on the chemistry–climate model (CCM) SOCOL (SOlar Climate Ozone Links), which has been improved by including modules for the production, deposition, and transport of \(^7\)Be and \(^{10}\)Be. Production of the isotopes was modeled for both galactic and solar cosmic rays by applying the CRAC (Cosmic Ray Atmospheric Cascade) model. Transport of \(^7\)Be was modeled without additional gravitational settling due to the submicron size of the background aerosol particles. An interactive deposition scheme was applied including both wet and dry deposition. Modeling was performed using a full nudging to the meteorological fields for the period of 2002–2008 with a spin-up period of 1996–2001. The modeled concentrations of \(^7\)Be in near-ground air were compared with the measured ones at a weekly time resolution in four nearly antipodal high-latitude locations: two in the Northern (Finland and Canada) and two in the Southern (Chile and the Kerguelen Islands) Hemisphere. The model results agree with the measurements in the absolute level within error bars, implying that the production, decay, and lateral deposition are correctly reproduced. The model also correctly reproduces the temporal variability of \(^7\)Be concentrations on annual and sub-annual scales, including the presence and absence of the annual cycle in the Northern and Southern Hemispheres, respectively. We also modeled the production and transport of \(^7\)Be for a major solar energetic
particle event (SPE) on 20 January 2005, which appears insufficient to produce a measurable signal but may serve as a reference event for historically known extreme SPEs. Thus, a new full 3D time-dependent model, based on CCM SOCOL, of $^7$Be and $^{10}$Be atmospheric production, transport, and deposition has been developed. Comparison with real data on the $^7$Be concentration in the near-ground air validates the model and its accuracy.

1 Introduction

One of the most important outer-space factors affecting Earth is related to cosmic rays (highly energetic nuclei of extraterrestrial origin), which cause nucleonic–electromagnetic–muon cascade in the terrestrial atmosphere (Dorman, 2004). These cosmic-ray-induced atmospheric cascades form the main source of ionization in the troposphere and stratosphere (e.g., Mironova et al., 2015). Nuclear interactions between the nucleonic component of the cascade and nuclei of atmospheric gases such as nitrogen and oxygen can produce radioactive nuclides (e.g., Lal and Peters, 1967; Beer et al., 2012) including $^7$Be (half-life 53.22 d) and $^{10}$Be (half-life 1.4 $\times$ 10$^6$ years). They are called cosmogenic isotopes since galactic and solar cosmic rays form the dominant source of these nuclides in the terrestrial system. Production rates of $^7$Be and $^{10}$Be vary in time following the intensity of galactic cosmic rays (GCRs) modulated by solar magnetic activity (the 11-year solar cycle), geomagnetic field strength, and also sporadic solar energetic particle events (SPEs) (see, e.g., Usoskin and Kovaltsov, 2008; Kovaltsov and Usoskin, 2010).

The short-living cosmogenic isotope $^7$Be, which is relatively easy to measure via a proper collection system and $\gamma$ spectrometry, is routinely measured with daily, weekly, and monthly time resolution in near-ground air and precipitating water at different places around the globe in the framework of atmospheric radiation monitoring. Concentrations of long-living $^{10}$Be are measured (via acceleration mass spectrometry, AMS) in natural stratified archives, typically in polar ice cores, with pseudo-annual or longer time resolution (e.g., Beer et al., 2012). In particular, $^{10}$Be recorded in polar ice cores serves as a main proxy for long-term solar activity reconstructions in the past (e.g., Steinhilber et al., 2012; Usoskin, 2017; Wu et al., 2018).

In addition to the production pattern defined by cosmic-ray fluxes, concentrations of beryllium isotopes at any location are also affected by atmospheric transport and deposition processes, which may essentially distort the production signal. Thanks to the short life of $^7$Be, high-time-resolution data, and broad distribution of the measuring sites, this isotope has often been used as a tracer for atmospheric dynamics during the recent decades. Its atmospheric transport can be modeled with high accuracy using the known meteorological fields as summarized in Sect. 2. On the other hand, data on $^{10}$Be are mainly measured in polar regions with low temporal resolution and extend far back in the past to times when meteorological data are not available. Accordingly, the atmospheric transport of $^{10}$Be is typically modeled using simplified box models based on a parameterization of some specific simulations (e.g., Heikkilä et al., 2013). Thus, there is a need for a realistic model able to simulate atmospheric transport and deposition of beryllium focused in polar regions and (inter)annual timescales that is potentially able to operate in a self-consistent mode without the prescribed meteorology. Such a model needs to be able to reproduce the main pattern of beryllium transport and its deposition in polar regions on the annual timescale.

Since measurements of $^{10}$Be are laborious and expensive, their coverage for the last decades is low, making it difficult to compare model results directly with measured data. Moreover, they can be additionally distorted by deposition and post-depositional processes. However, since the two beryllium isotopes are believed to have similar transport and deposition, being different only in production and lifetime, a model applied and verified for $^7$Be transport can also be generally applied to $^{10}$Be (Heikkilä et al., 2008b). Here we develop such a model, CCM SOCOL-AERv2-BE, to trace isotopes of $^7$Be and $^{10}$Be in the atmosphere with modules for the production, deposition, and transport of $^7$Be and $^{10}$Be.

The main purpose of this work is to present a new combined model of beryllium production and transport and to confront its results with high-resolution (weekly) measurements of $^7$Be in near-ground air and precipitating water in polar regions. The isotope $^7$Be was chosen for the analysis because of the following reasons: (i) there are no systematic high-quality $^{10}$Be data in air samples since measurements of $^{10}$Be are much more difficult (AMS technology) than those of $^7$Be ($\gamma$ spectrometry); (ii) $^{10}$Be measured in ice cores contains an additional component (deposition), which also affects the measured data and needs to be modeled separately. On the other hand, the model eventually aims at the modeling of the $^{10}$Be isotope including its deposition. Since $^{10}$Be is typically measured in polar ice cores with pseudo-annual or longer resolution, we are primarily interested in validating the model for high-latitude regions and (inter)annual timescales. The new model version CCM SOCOL-AERv2-BE has been developed here for systematic modeling of $^7$Be and $^{10}$Be production, transport, and deposition in the atmosphere. This model set forms a methodological basis for a detailed study of the atmospheric dynamic and cosmogenic isotope distribution in the terrestrial system. The model computations were performed with full nudging (using ERA5) to the observed meteorological data to verify the capability of the model to trace beryllium isotopes in the atmosphere. The model computations were directly compared with measurements of $^7$Be isotopes in different high-latitude locations to validate the model.

The paper is organized as follows: a summary of the previous and existing models is presented in Sect. 2; data from
\[ ^{7}\text{Be} \] measurements are presented in Sect. 3; the model is described in great detail in Sect. 4; obtained results are discussed in Sect. 5; validation of the model by means of comparisons with measurements is presented in Sect. 6; Sect. 7 summarizes conclusions of the paper.

2 Overview of previous and existing models

Studies of the short-living \[ ^{7}\text{Be} \] isotope and its ratio to \[ ^{10}\text{Be} \] provide a precise way to probe large-scale atmospheric dynamical properties and validate the corresponding atmospheric models (e.g., Raisbeck et al., 1981; Jordan et al., 2003; Ioannidou and Paatero, 2014). For the long-living \[ ^{10}\text{Be} \] isotope, different transport and deposition assumptions as well as models were used (e.g., McCracken, 2004; Heikkilä et al., 2009; Sukhodolov et al., 2017; Usoskin et al., 2020a).

Earlier simulations with the transport model GLOMAC1 by Brost et al. (1991) yielded the concentration and deposition of \[ ^{7}\text{Be} \] isotopes in agreement with the data within 20%. However, the authors conclude that the application of this transport model and scavenging scheme for simulations of deposition and concentration was not well suited for polar regions.

Later, simulations of \[ ^{7}\text{Be} \] and \[ ^{10}\text{Be} \] transport for present and future climatic conditions were performed by Land and Feichter (2003) with the general circulation model MAECHAM4. The model successfully captured the maximum \[ ^{10}\text{Be}/^{7}\text{Be} \] ratio in the lower polar stratosphere.

The general ability of a full atmospheric–dynamic model to reproduce the pattern of beryllium transport and deposition even at regional spatial and synoptic temporal scales was demonstrated by Usoskin et al. (2009a), who used the GISS model (Field et al., 2006) to simulate \[ ^{7}\text{Be} \] production and transport for the 2 months of January–February 2005 and compared the model’s results with daily-resolution near-ground \[ ^{7}\text{Be} \] air concentrations measured at 11 sites worldwide (see Table 1 in Usoskin et al., 2009a). The results implied that the model captures \[ ^{7}\text{Be} \] transport at synoptic and longer timescales (longer than 4–10 d; see Fig. 5 in Usoskin et al., 2009a) in both absolute levels and time variability. However, as noted by Koch et al. (2006), the GISS model has a low bias (a factor of two) in the upper troposphere and lower stratosphere, while agreeing well with the surface concentrations. It was proposed that this can be related to model production that is too low and the fact that either the model delivers too much \[ ^{7}\text{Be} \] to the surface or does not scavenge efficiently enough. The GISS model was not further developed in this direction.

Modeling of beryllium isotopes in the Earth’s atmosphere was performed earlier using the ECHAM5-HAM atmospheric model (Heikkilä et al., 2008a, 2009) and zonal mean production rates. Cosmogenic radionuclides \[ ^{10}\text{Be} \] and \[ ^{7}\text{Be} \] were modeled during the Maunder Minimum and the present-day climate (Heikkilä et al., 2008b). ECHAM5-HAM simulates beryllium fluxes and concentrations in ice at Greenland sites well, in agreement with the data. In Antarctica, however, the model has some difficulties in reproducing the very low level of precipitation, which leads to an underestimation of the \[ ^{10}\text{Be} \] concentrations in ice. The precipitation rate modeled at the Concordia site was a factor 2–3 too high, leading to \[ ^{10}\text{Be} \] concentrations in ice that are too low, since, after full depletion of \[ ^{10}\text{Be} \] from the atmosphere, further rain or snow dilutes its concentration in the ice. The \[ ^{7}\text{Be} \] surface air concentrations generally agree within a factor of 2 with the observed concentrations. However, modeled concentrations at two Antarctic stations appeared too low. This model was very roughly parameterized (Heikkilä et al., 2013), being effectively reduced to a six-box model and not developed further in the application to beryllium transport modeling.

A simplified 2D model of beryllium stratospheric sedimentation was developed by Delaygue et al. (2015) but also not supported further.

Even though a full modeling can be done, transport of \[ ^{7}\text{Be} \] is typically considered, for application purposes, with simple box–reservoir exchange models (e.g., Pacini et al., 2015; Zheng et al., 2021) or with air-mass tracing codes (Pacini et al., 2011; Brattich et al., 2020).

In chemical transport models, wind and temperature values are taken from either (1) general circulation models or from (2) meteorological data. The advantage of the second approach is that chemical transport simulations are linked to real meteorology and the results are inherently compared to observations. Reanalysis extends meteorological data to the past. This allows us to perform long-term simulations which show temporal evolution of the atmosphere’s composition. The success of such a task relies on the reanalysis quality, i.e., on the accuracy of the indicator’s distributions. This also makes chemical transfer models a powerful tool for evaluating the reanalyses themselves (Monge-Sanz and Chipperfield, 2006).

The Global Modeling Initiative (GMI) chemistry and transport model (a model with prescribed meteorology) driven by NASA’s MERRA assimilated meteorological data was applied to understand the roles of transport and precipitation scavenging processes using the cosmogenic \[ ^{7}\text{Be} \] tracer (Brattich et al., 2017). The model results agree reasonably well with the \[ ^{7}\text{Be} \] data in general but diverge by a factor of 2 for warm periods. The discrepancy was ascribed to the model’s insufficiently detailed spatial sampling.

The approach based on prescribed meteorology is applicable only to recent years when the wind field is known; moreover, it includes neither stratospheric dynamics nor depositional processes and is therefore not suitable for \[ ^{10}\text{Be} \] in polar ice.

Thus, several models of different complexity and accuracy have been developed in the recent past to model transport and deposition of \[ ^{7}\text{Be} \] and \[ ^{10}\text{Be} \]. However, due to the lack of support for radionuclide simulations with the ECHAM-
HAM (Heikkilä et al., 2008b) and GISS (Field et al., 2006) climate models, a full working model capable of simulating production, transport, and deposition of \(^{7}\text{Be}\) and \(^{10}\text{Be}\) from all sources is presently missing. The model should be “user-friendly” (i.e., for a user without deep knowledge of climate modeling) and applicable for past or future projections when direct nudging to the meteorological reanalysis data is not possible.

Here we present a new version of CCM SOCOL – CCM SOCOL-AERv2-BE – for modeling of production, transport, and deposition of the cosmogenic isotopes of beryllium as well as its validation with the available measurements of \(^{7}\text{Be}\) at high-latitude locations. CCM SOCOL is potentially capable of simulating the beryllium transport and deposition and was already used (Sukhodolov et al., 2017; Usoskin et al., 2020a) for the estimation of \(^{10}\text{Be}\) deposition after major solar energetic particle (SPE) events, but it did not include the state-of-the-art interactive parameterization of the wet and dry deposition introduced recently (Feinberg et al., 2019). A recent model version of CCM SOCOL that simulates aerosols more realistically, however, does not include the treatment of all processes relevant to the beryllium life cycle, and its applicability has not been evaluated. We have further upgraded CCM SOCOL (Feinberg et al., 2019) here by adding the production, transport, and deposition of \(^{7}\text{Be}\) and \(^{10}\text{Be}\) isotopes from both GCRs and SPEs.

3 Data on \(^{7}\text{Be}\) measurements

For the model verification, we used data on the measured activity (concentration) of the \(^{7}\text{Be}\) isotope in near-ground air, with weekly sampling (corrected for the decay) performed in different hemispheres for the period 2002–2008. High-time-resolution data are not necessary for several reasons: (i) data availability, since weekly sampling is a standard resolution for most \(^{7}\text{Be}\)-measuring stations; (ii) inability of full atmospheric–dynamic models to catch the sub-synoptic timescales because of the rough grid size (see Heikkilä et al., 2009; Usoskin et al., 2009b); and (iii) focus on longer timescales (annual and interannual) for the application to \(^{10}\text{Be}\) in ice cores. On the other hand, rougher-resolution data (e.g., monthly or quarterly) are more uncertain because of the decay of short-living \(^{7}\text{Be}\). Thus, we found the weekly-resolution data to be optimal for our purposes.

Traditionally, for short-living isotopes, the measured quantity is not concentration but activity in units of becquerels (number of radioactive decays per second) per cubic meter of air (Bq m\(^{-3}\)), as directly measured by a \(\gamma\) spectrometer. Correction for decay between the times of sample collection and measurement is standardly performed (see Leppänen et al., 2012). The activity and concentration can be easily translated to each other using the known lifetime of the isotope.

Here we use data from four nearly antipodal high-latitude locations (see Table 1). Two locations are in the Northern Hemisphere in Finland (a set of stations) and Canada (Yellowknife), and two are in the Southern Hemisphere: Kerguelen island (France) and Punta Arenas (Chile). The longitudinal separation between the locations in each hemisphere is greater than 140\(^{\circ}\). The data series are shown in Fig. 1 and described below.

Finnish data on the measured \(^{7}\text{Be}\) activity were provided for the period 2002–2008 by the Finnish National Radiation and Nuclear Safety Authority (STUK) for four sites (see Table 1): Ivalo, Rovaniemi, Kajaani, and Kotka, spanning meridionally 8\(^{\circ}\) in latitude. The isotope of \(^{7}\text{Be}\) was collected on air filters using a high-volume air sampler on a weekly basis, and then its activity was measured with an HPGe (high-purity germanium) detector inside a 4 in. thick lead shield- ing in the STUK laboratory. The statistical measurement uncertainty is 3\%-5\% (see a detailed description of the sampling and measurement procedure in Leppänen et al., 2012). The data from the four stations were combined as the average of the four locations for each time interval to represent a regional-scale dataset, as shown in Fig. 1a.

We also used \(^{7}\text{Be}\) data from the Yellowknife site (Canada; see Fig. 1b) kindly provided by the Radiation Protection Bureau of Health Canada. The original data are with daily time resolution but have been combined into weekly samples (correction for decay applied) for consistency with other datasets.
samples for each year cover the period of January to March, together to form a quarterly sample. Hence, the four measured periods was 1 month, but the three samples were combined to-gether to form a quarterly sample. Hence, the four measured samples for each year cover the period of January to March, April to June, July to September, and October to December. The reference date was set to the middle of the sampling period (since the decay time of $^7$Be is shorter than the sampling interval) and the activities were decay-corrected to this date (see Leppänen, 2019). The $^7$Be deposition averaged over the two locations is shown in Fig. 2.

4 Model description

We used an extended version of CCM SOCOL (Stenke et al., 2013) with the aerosol module – AERv2 (Feinberg et al., 2019). This version consists of the general circulation model MA-ECHAM5 (Hommel et al., 2011) and the atmospheric chemistry module MEZON (Egorova et al., 2003) combined with an aerosol module AER (Weisenstein et al., 1997), interacting with each other every 2 h of modeling. In this study, MA-ECHAM5 is nudged with fifth-generation ECMWF atmospheric reanalysis data from the global climate ERA5 (era5T42L39) reanalyses. The chemical part of the MEZON core includes 56 chemical species of oxygen, hydrogen, nitrogen, carbon, chlorine, bromine, and sulfur groups, 295 gas-phase reactions, 64 photolysis reactions, and 16 heterogeneous reactions on stratospheric aqueous sulfuric acid aerosols as well as three types of polar stratospheric clouds. CCM SOCOL-AERv2-BE uses the horizontal resolution T42, where T refers to triangular truncation and the number denotes the series of spherical harmonics. The horizontal resolution is about $2.8^\circ \times 2.8^\circ$ ($\approx 300 \times 300$ km in the equatorial region). The vertical grid is defined using a sigma-hybrid coordinate system, which is calculated by a linear combination of surface pressure and constant coefficients that define the vertical coordinate. The model has 39 vertical levels between the Earth’s surface and the 0.01 hPa level ($\approx 80$ km). The real orography is smoothed over the model grid cells. The aerosol module exploits a sectional micro-physical scheme whereby the aerosol particle size distribution is represented by 40 discrete size bins. The model simulates the advective transport of stratospheric aerosol by the Brewer–Dobson large-scale circulation (Butchart, 2014) and gravitational sedimentation. We have made some adaptations

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Table 1. List of stations whose data were used for the present study. Data sources are the Finnish National Radiation and Nuclear Safety Authority (STUK), the Radiation Protection Bureau of Health Canada (RPBHC), and the Comprehensive Nuclear Test-Ban-Treaty Organization (CTBTO).

<table>
<thead>
<tr>
<th>Station name</th>
<th>Region</th>
<th>Location</th>
<th>Source</th>
<th>Time interval</th>
</tr>
</thead>
<tbody>
<tr>
<td>Ivalo</td>
<td>Finland</td>
<td>68° 39′ 09″ N 27° 32′ 47″ E</td>
<td>STUK</td>
<td>2002–2008</td>
</tr>
<tr>
<td>Rovaniemi</td>
<td>Finland</td>
<td>66° 30′ 05″ N 25° 44′ 05″ E</td>
<td>STUK</td>
<td>2002–2008</td>
</tr>
<tr>
<td>Kajaani</td>
<td>Finland</td>
<td>64° 13′ 30″ N 27° 44′ E</td>
<td>STUK</td>
<td>2002–2008</td>
</tr>
<tr>
<td>Kotka</td>
<td>Finland</td>
<td>60° 28′ N 26° 56′ 45″ E</td>
<td>STUK</td>
<td>2002–2008</td>
</tr>
<tr>
<td>Yellowknife</td>
<td>Canada</td>
<td>62° 26′ 32″ N 114° 23′ 51″ W</td>
<td>RPBHC</td>
<td>2003–2008</td>
</tr>
<tr>
<td>Kerguelen island</td>
<td>Port-aux-Français, France</td>
<td>49° 21′ S 70° 13′ E</td>
<td>CTBTO</td>
<td>2006–2008</td>
</tr>
<tr>
<td>Punta Arenas</td>
<td>Chile</td>
<td>53° 09′ S 70° 55′ W</td>
<td>CTBTO</td>
<td>2006–2008</td>
</tr>
</tbody>
</table>
of the main model to focus on beryllium transport, as described below.

4.1 Production rates of \(^{7}\text{Be}\)

Production of both \(^{7}\text{Be}\) and \(^{10}\text{Be}\) was computed using the CRAC : Be (Cosmic Ray Atmospheric Cascade: application to Beryllium) model (Poluianov et al., 2016), which is the most recent and accurate model of cosmogenic isotope production. This model simulates, using the GEANT4 Monte Carlo simulation tool (Geant4 collaboration, 2020), the full nucleonic–muon–electromagnetic cascade induced by primary cosmic-ray particles in the atmosphere. This model is an essential upgrade of the previous versions of the CRAC : Be model (Usoskin and Kovaltsov, 2008; Kovaltsov and Usoskin, 2010) based on the CORSIKA simulation tool (Heck et al., 1998). The CRAC model presents a set of accurately computed yield functions of production of cosmogenic isotopes for different primary particle types, energy, and atmospheric depths. By a numerical integration above the local geomagnetic rigidity cutoff (Cooke et al., 1991) of the product of the yield function and the known (or assumed) energy spectrum of primary cosmic rays (of either GCR or SPE origin), one can obtain the isotope’s production rate as a function of location (via the geomagnetic rigidity cutoff), atmospheric height (via the yield function), and time (via the time-variable energy spectrum). These production rates, computed for each location and with hourly temporal time resolution, were used as input for CCM SOCOL-AERv2-BE, providing a 3D + time source of the \(^{7}\text{Be}\) and \(^{10}\text{Be}\) in the atmosphere.

Production due to galactic cosmic rays (GCRs), which always bombard the Earth with slightly variable intensity due to solar modulation, was computed by applying the force-field model (Gleeson and Axford, 1968) with the modulation potential reconstructed from the World Neutron Monitor Network (Usoskin et al., 2017). The geomagnetic field was taken according to the International Geomagnetic Reference Field model (IGRF; Thébault et al., 2015) in the eccentric dipole approximation (Fraser-Smith, 1987; Usoskin et al., 2010). While GCRs are always present near Earth, sporadic solar particle events (SPEs) take place occasionally (Vainio et al., 2009; Desai and Giacalone, 2016) and can be sufficiently strong to be recorded by ground-based neutron monitors (Raukunen et al., 2018; Usoskin et al., 2020b). A large number of cosmogenic isotopes can be produced during such events (Usoskin et al., 2020a), which can be several orders of magnitude stronger on the multimillennial timescale than for the last decades covered by direct observations, with strong signals in the polar \(^{10}\text{Be}\) concentrations (e.g., Sukhodolov et al., 2017). To model the contribution of severe SPEs, we modeled the production of beryllium by the SPE on 20 January 2005 (GLE no. 69), which was the strongest event during the studied period (1996–2008) and the second strongest directly observed (1951–2020). The same yield function and geomagnetic cutoff approach as for GCR were applied along with the SPE energy spectrum obtained for this event by Koldobskiy et al. (2021). To study the possible effect of seasonality on the \(^{7}\text{Be}\) and \(^{10}\text{Be}\) transport we simulated the same SPE event as if it occurred not only on its actual date in midwinter (20 January 2005) but also mid-spring, summer, and autumn (20 April, 20 July, and 20 October 2005, respectively).

Figure 3 shows the zonal mean production of \(^{7}\text{Be}\) as a function of the atmospheric pressure and northern geographical latitude from two sources: daily GCRs (panel a) and SPEs (panel b). Since the duration of a strong SPE takes several hours, we compare it with the daily production of \(^{7}\text{Be}\) by GCR. One can see that the polar stratospheric production due to SPE grossly overwhelms that by the daily GCR flux, but it fades towards tropical latitudes. Since the production of \(^{7}\text{Be}\) is nearly symmetric between the global hemispheres, as defined by the geomagnetic shielding, only the Northern Hemisphere is shown. Even though the production rate is significantly higher in the polar region, its contribution to the global production is not dominant because of the small area of the polar regions. For the GCR-related production, the dominant production region is located in the middle- to high-latitude (40–70°) stratosphere and upper troposphere (about 400 hPa level or ≈ 7 km). Production of \(^{7}\text{Be}\) by SPEs is confined to the polar (latitude > 55°) upper stratosphere and mesosphere (above 50 hPa level or ≈ 35 km). Hardly any beryllium nuclei are produced by SPEs in the tropical region. In all cases, the

Figure 3. Zonal mean production of \(^{7}\text{Be}\) in the Northern Hemisphere as a function of geographical latitude and the model’s atmospheric pressure level. (a) Daily production by GCRs averaged over the period of 2002–2008. The color scale (on the right) is given in units of atoms per day per degree of latitude per gram per square centimeter (g cm\(^{-2}\)). (b) Total production by SPE on 20 January 2005. The color scale is in units of atoms per degree of latitude per gram per square centimeter (g cm\(^{-2}\)).
maximum of production lies at middle to high latitudes in the stratosphere, and the production decreases towards the surface because of the increasing atmospheric depth and equatorward due to the geomagnetic shielding. Because of the much softer energy spectrum, SPEs produce $^7$Be at shallower atmospheric depths and higher latitudes than GCRs do. The production pattern for $^{10}\text{Be}$ is similar to that for $^7\text{Be}$ but with smaller absolute values.

4.2 Decay of $^7\text{Be}$

$^7\text{Be}$ is a short-living isotope whose decay cannot be neglected in comparison with the typical transport and deposition time, especially in the stratosphere, in contrast to the long-living $^{10}\text{Be}$ isotope. Accordingly, a standard decay probability of 0.054 % per hour (corresponding to the isotope’s mean half-life of 53.22 d) was applied during tracing of $^7\text{Be}$ in a way similar to that used by Golubenko et al. (2020) for $^{222}\text{Rn}$. No decay needs to be applied to $^{10}\text{Be}$ for this kind of modeling. Figure 4 shows the modeled global content of the two beryllium isotopes in the atmosphere after an instant production by the SPE event on 20 January 2005 (day zero). One can see that the amount of $^{10}\text{Be}$ decreases nearly exponentially, reflecting a slow removal of the isotope through precipitation. The typical time of the isotope removal is $\tau_{^7\text{Be}} = 550 \pm 100$ d, implying a residence time of a few years (e.g., Beer et al., 2012). This time is comparable to the 1–2-year atmospheric residence time of other isotopes (e.g., $^{133}\text{Cs}$ and $^{244}\text{Am}$) defined after nuclear weapon tests (Haltia et al., 2021). The decrease in the $^7\text{Be}$ isotope concentration is much faster and nearly perfectly exponential with $\tau_{^7\text{Be}} = 72 \pm 3$ d, which includes both decay and removal. Since the removal process is assumed to be identical for both isotopes, the concentrations of $^{10}\text{Be}$ can be used to correct the $^7\text{Be}$ decay for the system effects. The decay time of $^7\text{Be}$ estimated in this way is $1/\tau = 1/\tau_{^7\text{Be}} - 1/\tau_{^{10}\text{Be}}$, as $\tau = 82 \pm 4$ d, which is generally consistent with the expected e-folding lifetime of $^7\text{Be}$ of 77 d, implying that the decay was accounted for correctly in the model. Thus, the $^7\text{Be}$ isotope is fully removed from the atmosphere, mostly due to decay, within several months.

4.3 Transport of cosmogenic isotopes $^7\text{Be}$ and $^{10}\text{Be}$

Beryllium atoms are captured by different aerosols in the air that can enhance or suppress its transport by gravitational sedimentation. Since the typical size of aerosol particles is less than 0.2 $\mu$m (Pierce et al., 2010), with a modal value of 0.05 $\mu$m, the corresponding gravitational sedimentation velocity ranges between 200 and 1000 m yr$^{-1}$ and does not play an important role in the lower stratosphere, where it is much smaller than the vertical wind speed (e.g., Weisenstein et al., 2015; Zhou et al., 2006). The role of gravitational sedimentation was evaluated by Delaygue et al. (2015) using numerical experiments with a 2D model, which includes detailed sub-

![Figure 4. Computed global mean column content of two beryllium isotopes, $^7\text{Be}$ and $^{10}\text{Be}$, produced by SPE on 20 January 2005 (day zero) as a function of time. The nearly exponential decay of the concentration ($\tau = 550 \pm 100$ d and $72 \pm 2$ d for $^{10}\text{Be}$ and $^7\text{Be}$, respectively) includes both decay and sink (deposition) of the isotopes.](https://doi.org/10.5194/gmd-14-7605-2021)
beryllium isotopes is parameterized as a function of surface properties, solubility, and reactivity of the considered species (Kerkweg et al., 2006a). These properties for the modeled beryllium are identical to those for small sulfate aerosols. It means that beryllium atoms are considered passive tracers only for the transport process, while with respect to the dry and wet deposition, they are treated as sulfates. The applied dry deposition scheme also considers the actual meteorological conditions and different surface types.

The interactive wet deposition scheme used in CCM SOCOL-AERv2-BE exploits the EASY2 version of the Scavenging (SCAV) submodule in the ECHAM/MESy Atmospheric Chemistry (EMAC) model. A detailed description of the interactive wet deposition scheme is presented and discussed elsewhere (Kerkweg et al., 2006b; Tost et al., 2007, 2010). The parameterization is based on the model-generated available liquid water in (cloud water content) and below (precipitating water) clouds as well as uptake and release form droplets, which depends on the concentration and solubility of the considered species. Scavenging coefficients for all tracers are calculated based on Henry’s law equilibrium constants. The Henry constants for beryllium are considered identical to those of small sulfate aerosols.

4.5 Setup of numerical experiments

Since $^{10}$Be is removed from the stratosphere via stratosphere–troposphere exchange and the stratospheric mixing time is long (a few years), we have performed a 6-year spin-up of the model for the period 1996–2001 to allow $^{10}$Be to reach equilibrium conditions. For $^7$Be, such a long spin-up is not needed since it is mostly removed from the atmosphere by decay, but we kept it similar to $^{10}$Be for consistency. After the spin-up period, we initiated a 7-year run (2002–2008) with full nudging (a linear relaxation of thermodynamic parameters: temperature, divergence, and vorticity of the wind field). In this way, we can validate the model by comparing the results with the measurements.

5 Modeling results

5.1 Global patterns of $^7$Be transport and deposition

Figure 5 shows the zonal mean concentration of $^7$Be averaged over the boreal summer and winter for 2002–2008, as modeled here. While the isotope production is symmetric between the hemispheres, the concentrations are not completely symmetric and depend on the large-scale circulation patterns, which are different in the summer and winter hemispheres (Terzi et al., 2020). The stratospheric meridional transport is driven by forces from the breaking planetary, synoptic-scale, and gravity waves. Stratospheric air ascends inside the so-called tropical pipe and then moves poleward, descending at middle and high latitudes, predominantly in the winter hemisphere where a more intensive stratosphere–troposphere exchange leads to slightly smaller concentrations of the $^7$Be in the stratosphere during the cold seasons. Lower-stratospheric air descends relatively fast over the middle latitudes at which the stratosphere–troposphere exchange is fast, while at high latitudes it resides longer in the lower stratosphere (Cohen et al., 2014).

Figure 6 presents distributions of the total (wet and dry) deposition of $^7$Be averaged over 2002–2008. The deposition is low in arid regions, e.g., the Sahara and the Middle East. It is also low to the west of the continents, following the low precipitation above the cold ocean currents. The precipitation rate is the highest in the tropical convection zone due to the highest temperature gradients and strong convection in the equatorial area. The deposition, however, is not very high at low latitudes because both the production rate and the downward transport of $^7$Be are low in the tropics. It is important that gradients of the deposition can be strong even on the regional scale, suggesting that meteorological processes can strongly influence beryllium deposition at any given location (e.g., Usoskin et al., 2009b). This pattern agrees well with a similar previous study (Heikkilä et al., 2009).

5.2 The effect of a strong SPE

While GCRs always bombard the Earth’s atmosphere with slightly variable intensity, major SPEs make a different impact, producing sporadic and short (typically several hours) but very intense fluxes of energetic particles entering the atmosphere mostly in polar regions. Potentially, extreme SPEs that are orders of magnitude stronger than those observed during the recent decades can be recorded in cosmogenic isotope data (Usoskin et al., 2006, 2020a; Miyake
et al., 2012), and a proper model is needed to study them (Sukhodolov et al., 2017). Here we modeled transport of such SPE-produced beryllium based on the strong SPE on 20 January 2005. In order to distinguish between different sources, we traced five beryllium tracers for the $^7$Be isotope: one for GCR-produced beryllium and four tracers for the SPE scenarios as described above.

Production of the $^7$Be isotope during this SPE (Fig. 3b) appears mostly in the polar stratosphere and lower mesosphere. After production, the isotope starts decaying and being transported by the air dynamic. Figure 7 shows an example of the SPE-produced isotope concentration on the 30th day after the event (19 February 2005). One can see that the production pattern is smeared by the transport, leading in particular to an essential hemispheric difference. It is interesting that tropospheric concentrations are higher at middle and low latitudes of about $30^\circ$ than in polar regions because of the atmospheric circulation. As an example, the modeled activity of $^7$Be in near-ground air is shown in Fig. 8 as averaged over Finland. The locally (polar troposphere) produced beryllium dominates during the first 20 d after the event, but then the transport starts playing a role, leading to the very low concentration during the subsequent period. The level of activity for this event in the near-air in Finland is very low: a factor 30–100 lower than the typical level of activity due to GCRs during the winter season (see Fig. 1a).

Although the SPE on 20 January 2005 was a very strong one (the second-strongest directly observed), its imprint in cosmogenic isotopes is not detectable on the background of air transport dynamic and GCR variability. An SPE must be stronger by a factor of 10 or more than this one to become detectable in isotope records even on the daily scale (Sukhodolov et al., 2017; Usoskin et al., 2020a).

The transport of $^7$Be and $^{10}$Be after production may differ for different seasons. In order to study that, we simulated the same SPE on 20 January 2005 as if it occurred in mid-spring (20 April 2005), midsummer (20 July 2005), and mid-autumn (20 October 2005). Figure 9 shows the concentration of $^{10}$Be in the Northern Hemisphere ($85^\circ$ latitude) for these cases. Here we consider $^{10}$Be because of the fast decay of $^7$Be, which smears the time variability of the concentration. The most pronounced and long-lasting effect occurs from the SPE taking place during the autumn (panel d) when a strong concentration of $^{10}$Be persists in the stratospheric region for about 1 year and gradually decreases after that for a few more years. The patterns of the SPE occurring during summer (panel c) and spring (panel b) are similar to each other in that it has a shorter residence time (only a few months) and a smaller spread of $^{10}$Be. The winter scenario suggests somewhat faster removal of $^{10}$Be from the atmo-
sphere, but it also stays in the upper atmosphere for several years. The concentration in near-ground air is mostly defined by the seasonal cycle and is consistent between different scenarios of the SPE occurrence date.

6 Evaluation of the model by comparison with direct $^7$Be measurements

In this section, we validate the model by comparing the simulation results with actual measurements of the near-ground-air $^7$Be concentration in air filters performed in boreal and austral high-latitude regions (Sect. 3), as shown in Fig. 1.

First, we compared the modeled $^7$Be activity with the measured one for the all-Finland record, which was compiled from four stations (see Sect. 3) as shown in Fig. 1a. The linear Pearson correlation between the modeled and measured weekly series is highly significant at $r = 0.68 \pm 0.02$ ($p$ value $< 10^{-6}$). The standard correlation analysis does not distinguish between different frequencies and/or timescales. Therefore, we show in Fig. 10 the global wavelet coherence, which is an analog of the correlation coefficient but extending it into the frequency domain between the modeled and measured $^7$Be activities.

Figure 10 depicts the coherence for the Finnish stations. The coherence is highly significant (confidence level above 95%) at all timescales between 1 month and 2 years. Although there are discrepancies on the very short timescales related to synoptic noise, the model correctly reproduces some strong spikes in the beryllium activity, such as earlier in 2003, which are often related to sudden stratospheric warming (SSW) events (Pacini et al., 2015; Brattich et al., 2021). It is important that the long-term variability (at an annual timescale) is correctly reproduced by the model.

However, neither correlation coefficient nor coherence analyses provide information on the absolute level match between the series. In order to assess possible model biases we considered the difference between the modeled and measured weekly $^7$Be concentrations as shown in Fig. 11a.

The difference is insignificant at $0.16 \pm 0.7$ mBq m$^{-3}$, implying the absence of a systematic bias in the model. The mean (averaged over seven yearly values) level of the modeled activity (2.05 mBq m$^{-3}$) is close to the measured one (1.86 mBq m$^{-3}$), within 10% (the difference is $0.19 \pm 0.19$ mBq m$^{-3}$). The $z$ test suggests that this difference is insignificant. On the other hand, there is an essential difference between the model and the measurements regarding the magnitude of summer peaks in 2007 and 2008. It could potentially be related to the local atmospheric aerosol properties, but the AERONET (AErosol ROBotic NETwork) (Emery and Camps, 2017) station in Sodankylä (northern Finland) did not show any anomalies for that time. The discrepancies appear during warm periods, in agreement with other direct models (Brattich et al., 2017).

For the Canadian station, the agreement between the modeled and measured activities (Fig. 1b) is nearly perfect, including a correct reproduction of the very strong seasonal
cycle and sudden short spikes, e.g., in late 2003. The Pearson correlation coefficient is $r = 0.78 \pm 0.02$ ($p$ value $< 10^{-6}$). The coherence between the simulated and measured datasets (Fig. 10b) is highly significant at all timescales between 1 month and 1.5 years. The difference between the modeled and measured concentrations is $-0.25 \pm 1.05 \text{ mBq m}^{-3}$ (Fig. 11b). The overall mean level of modeled activity (2.15 mBq m$^{-3}$) is close to the measured one (2.4 mBq m$^{-3}$), within 10% (the difference is $0.24 \pm 0.22 \text{ mBq m}^{-3}$). The $z$ test suggests that the difference is insignificant, implying the absence of the model bias. The most pronounced difference between the model and the measurements was observed in 2008.

The Chilean data (Fig. 1c) depict reasonable agreement between the model and the measurements. Since the data contain no seasonal pattern, as correctly reproduced by the model, the formal correlation is insignificant at $r = 0.06$, but the coherence (Fig. 10d) is highly significant at monthly and annual timescales. The difference between the modeled and measured concentrations is $0.33 \pm 0.6 \text{ mBq m}^{-3}$ (Fig. 11d). The mean levels of the modeled and measured series are reasonably matched, with a difference of 18%: 2.02 vs. 1.65 mBq m$^{-3}$ for the modeled and measured activities, respectively, also indicating no significant bias. The observed small systematic discrepancy may be related to the peculiarity of this site, which is located on a small island (about 100 km across) in the middle of the ocean, while the model grid (about $300 \times 300$ km) may be too rough to catch the actual orography.

Therefore, we conclude that the model correctly (within 10%) calculates the mean levels of beryllium concentration for both northern and southern high-latitude locations and contains no biases, at least for the locations where the orographic scale is comparable to the model’s grid size, but it may have larger uncertainties for unresolvable spatial scales. It is important that the model correctly reproduces the seasonal variability or its absence as well as some sudden spikes, particularly during the late winter and spring seasons, likely related to SSW events.

In addition to the concentration of $^7$Be in the near-surface air (expressed in measured activity), we also compared the $^7$Be deposition modeled by CCM SOCOL-AERv2-BE with that measured in northern Finland, as shown in Fig. 2. The agreement is very good: the average levels (183 and 188 Bq m$^{-2}$ for the modeled and measured data, respec-
tively) agree within 2.5%, and the correlation is highly significant ($r = 0.86^{+0.04}_{-0.06}$, $p < 10^{-6}$). However, while the overall agreement is very good, the discrepancy can be significant for individual years. Thus, we conclude that the model also reproduces the depositional flux of $^7\text{Be}$ reasonably well on the yearly timescale.

7 Conclusions

A full 3D transient model of production, transport, and deposition of cosmogenic $^7\text{Be}$ and $^{10}\text{Be}$ isotopes in the atmosphere has been developed. The model, named CCM SOCOL-AERv2-BE, is based on CCM SOCOL, specifically tuned for the best performance in tracing beryllium, and the CRAC production model. Realistic modeling of the $^7\text{Be}$ isotope was performed for the years 2002–2008, with a 5-year model spin-up during 1996–2001. The measured weekly concentrations of $^7\text{Be}$ in near-ground air have been compared with the model prediction for four nearly antipodal high-latitude locations: two in the Northern (Finland and Canada) and two in the Southern (Chile and Kerguelen island) Hemisphere. The model results generally agree well with the measurements in the absolute level within error bars, implying that the production, decay, and lateral deposition are correctly reproduced by the model. However, a larger discrepancy (up to 20%) was observed at Kerguelen island where the orographic scale is much smaller than the model grid size. The model correctly reproduces the temporal variability of $^7\text{Be}$ concentrations on the annual and sub-annual scales, including a good reproduction of the annual cycle, which dominates data in the Northern Hemisphere, and the absence of such a cycle in the Southern Hemisphere. In addition, occasional anomalous events such as SSWs are well reproduced by the model. This validates the newly developed model to be able to correctly simulate the production, transport, and deposition of the $^7\text{Be}$ (and hence, $^{10}\text{Be}$) isotope on local or regional spatial (in high-latitude regions) and monthly or annual temporal scales. The modeled $^7\text{Be}$ distribution is also in general agreement with earlier computations based on a similar approach (e.g., Heikkilä et al., 2008a; Field et al., 2006).

We have also modeled the production and transport of $^7\text{Be}$ for a major solar energetic particle event on 20 January 2005, which was one of the largest directly observed events. It was shown that, at minimum, an order-of-magnitude stronger event is needed to become observable in the beryllium data.

Concluding, the new full 3D time-dependent CCM SOCOL-AERv2-BE model of $^7\text{Be}$ and $^{10}\text{Be}$ atmospheric production, transport, and deposition has been developed and validated using directly measured data. The model is recommended to be used in studies related to, e.g., atmospheric dynamical patterns, extreme solar particle storms, long-term solar activity reconstruction from cosmogenic proxy data, and solar–terrestrial relations.


Author contributions. All authors participated in the model development, verification, discussions about the results, and revisions to the article. IU, GK, and ER designed the main ideas of this study, and KG carried them out. KG was supervised directly by ER and TS during the model code improvement work. APL found and prepared the real beryllium data for the model verification. IU and KG were responsible for data processing and figure preparation. IU prepared the paper.

Competing interests. The contact author has declared that neither they nor their co-authors have any competing interests.

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