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## Can machine learning improve the model representation of turbulent kinetic energy dissipation rate in the boundary layer for complex terrain?

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**Abstract.** Current turbulence parameterizations in numerical weather prediction models at the mesoscale assume a local equilibrium between production and dissipation of turbulence. As this assumption does not hold at fine horizontal resolutions, improved ways to represent turbulent kinetic energy (TKE) dissipation rate  $(\epsilon)$  are needed. Here, we use a 6-week data set of turbulence measurements from 184 sonic anemometers in complex terrain at the Perdigão field campaign to suggest improved representations of dissipation rate. First, we demonstrate that the widely used Mellor, Yamada, Nakanishi, and Niino (MYNN) parameterization of TKE dissipation rate leads to a large inaccuracy and bias in the representation of  $\epsilon$ . Next, we assess the potential of machinelearning techniques to predict TKE dissipation rate from a set of atmospheric and terrain-related features. We train and test several machine-learning algorithms using the data at Perdigão, and we find that the models eliminate the bias MYNN currently shows in representing  $\epsilon$ , while also reducing the average error by up to almost 40 %. Of all the variables included in the algorithms, TKE is the variable responsible for most of the variability of  $\epsilon$ , and a strong positive correlation exists between the two. These results suggest further consideration of machine-learning techniques to enhance parameterizations of turbulence in numerical weather prediction models.

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## 1 Introduction

While turbulence is an essential quantity that regulates many phenomena in the atmospheric boundary layer (Garratt, 1994), numerical weather prediction models are not capable of fully resolving it. Instead, they rely on parameterizations to represent some of the turbulent processes. Investigations into model sensitivity have shown that out of the various parameterizations currently used in mesoscale models, that of turbulent kinetic energy (TKE) dissipation rate ( $\epsilon$ ) has the greatest impact on the accuracy of model predictions of wind speed at wind turbine hub height (Yang et al., 2017; Berg et al., 2018).

Current boundary layer parameterizations of  $\epsilon$  in mesoscale models assume a local equilibrium between production and dissipation of TKE. While this assumption is generally valid for homogeneous and stationary flow (Albertson et al., 1997), as the horizontal grid resolution of mesoscale models is constantly pushed toward finer scales thanks to the increase of the computing resource capabilities, the theoretical bases of this assumption are violated. In fact, turbulence produced within a model grid cell can be advected farther downstream in a different grid cell before being dis-

sipated (Nakanishi and Niino, 2006; Krishnamurthy et al., 2011; Hong and Dudhia, 2012).

The inaccuracy of the mesoscale model representation of  $\epsilon$  impacts a wide variety of processes that are controlled by the TKE dissipation rate. In fact, the dissipation of turbulence affects the development and propagation of forest fires (Coen et al., 2013), it has consequences on aviation meteorology and potential aviation accidents (Gerz et al., 2005; Thobois et al., 2015), it regulates the dispersion of pollutants in the boundary layer (Huang et al., 2013), and it affects wind energy applications (Kelley et al., 2006): for example, in terms of the development and erosion of wind turbine wakes (Bodini et al., 2017).

Several studies have documented the variability of  $\epsilon$  using observations from both in situ (Champagne et al., 1977; Oncley et al., 1996; Frehlich et al., 2006) and remotesensing instruments (Frehlich, 1994; Smalikho, 1995; Shaw and LeMone, 2003). Bodini et al. (2018, 2019b) showed how  $\epsilon$  has strong diurnal and annual cycles onshore, with topography playing a key role in triggering its variability. On the other hand, offshore turbulence regimes (Bodini et al., 2019a) are characterized by smaller values of  $\epsilon$ , with cycles mostly impacted by wind regimes rather than convective effects. Also,  $\epsilon$  greatly increases in the wakes of obstacles, for example, wind turbines (Lundquist and Bariteau, 2015; Wildmann et al., 2019) or whole wind farms (Bodini et al., 2019b).

This knowledge on the variability of TKE dissipation rate provided by observations lays the foundation to explore innovative ways to improve the model representation of  $\epsilon$  in the atmospheric boundary layer. In this study, we leverage the potential of machine-learning techniques to explore their potential application to improve the parameterizations of  $\epsilon$ . Machine-learning techniques can successfully capture the complex and nonlinear relationship between multiple variables without the need of representing the physical process that governs this relationship. They have been successfully used to advance the understanding of several atmospheric processes, such as convection (Gentine et al., 2018), turbulent fluxes (Leufen and Schädler, 2019), and precipitation nowcasting (Xingjian et al., 2015). The renewable energy sector has also experienced various applications of machinelearning techniques, in both solar (Sharma et al., 2011; Cervone et al., 2017) and wind (Giebel et al., 2011; Optis and Perr-Sauer, 2019) power forecasting. Applications have also been explored at the wind turbine level, for turbine power curve modeling (Clifton et al., 2013), turbine faults and controls (Leahy et al., 2016), and turbine blade management (Arcos Jiménez et al., 2018).

Here, we train and test different machine-learning algorithms to predict  $\epsilon$  from a set of atmospheric and topographic variables. Section 2 describes the Perdigão field campaign and how we retrieved  $\epsilon$  from the sonic anemometers on the meteorological towers. In Sect. 3, we then evaluate the accuracy of one of the most common planetary boundary layer

parameterization schemes used in numerical weather prediction: the Mellor, Yamada, Nakanishi, and Niino (MYNN) parameterization scheme (Nakanishi, 2001). Section 4 presents the machine-learning algorithms that we used in our analysis. The results of our study are shown in Sect. 5, and discussed in Sect. 6, where future work is also suggested.

#### 2 Data

# 2.1 Meteorological towers at the Perdigão field campaign

The Perdigão field campaign (Fernando et al., 2018), an international cooperation between several universities and research institutes, brought an impressive number of instruments to a valley in central Portugal to survey the atmospheric boundary layer in complex terrain. The Perdigão valley is limited by two mountain ridges running from northwest to southeast (Fig. 1), separated by  $\sim 1.5\,\mathrm{km}$ . The intensive operation period (IOP) of the campaign, used for this study, was from 1 May to 15 June 2017.

At Perdigão, 184 sonic anemometers were mounted on 48 meteorological towers, which provided an unprecedented density of instruments in such a limited domain (Fig. 1). Observations from the sonic anemometers (a mix of Campbell Scientific CSAT3, METEK uSonic, Gill WindMaster, and YOUNG Model 81000 instruments) were recorded at a 20 Hz frequency.

The height of the towers ranged from 2 to 100 m, with the sonic anemometers mounted at various levels on each tower, as detailed in Table 1 and summarized in the histogram in Fig. 2, allowing for an extensive survey of the variability of the wind flow in the boundary layer. Data from the sonic anemometers have been tilt-corrected following the planar fit method (Wilczak et al., 2001), and rotated into a geographic coordinate system.

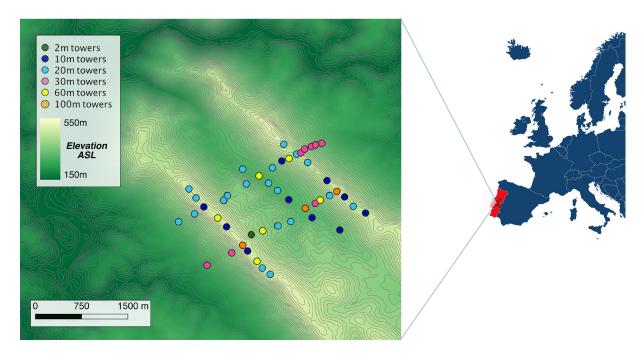
To classify atmospheric stability, we calculate the Obukhov length L from each sonic anemometer as

$$L = -\frac{\overline{\theta_{v}} \cdot u_{*}^{3}}{k \cdot g \cdot \overline{w'\theta_{v}'}}.$$
 (1)

 $\theta_{\rm v}$  is the virtual potential temperature (K, here approximated as the sonic temperature);  $u_*$  is the friction velocity (m s<sup>-1</sup>); k=0.4 is the von Kármán constant;  $g=9.81\,{\rm m\,s^{-2}}$  is the gravity acceleration; and  $\overline{w'\theta'_{\rm v}}$  is the kinematic buoyancy flux (m K s<sup>-1</sup>). For atmospheric stability, we classify unstable conditions as  $\zeta=z/L<-0.02$ ; and stable conditions as  $\zeta>0.02$ ; nearly neutral conditions as  $|\zeta|\leq0.02$ .

### 2.2 TKE dissipation rate from sonic anemometers

TKE dissipation rate from the sonic anemometers on the meteorological towers is calculated from the second-order structure function  $D_U(\tau)$  of the horizontal velocity U (Muñoz-



**Figure 1.** Map of the Perdigão valley showing the location and height of the 48 meteorological towers whose data are used in this study. Digital elevation model data are courtesy of the US Geological Survey.

Table 1. Heights where sonic anemometers were mounted on the meteorological towers at the Perdigão field campaign.

Tower height	Sonic anemometer heights (m a.g.l.)	Number of towers
2 m	2	1
10 m	10	5
	2, 10	5
20 m	10, 20	10
	2, 10, 20	6
	2, 4, 6, 8, 10, 12, 20	4
30 m	10, 30	3
	2, 4, 6, 8, 10, 12, 20, 30	5
60 m	10, 20, 30, 40, 60	5
	2, 4, 6, 8, 10, 12, 20, 30, 40, 60	1
100 m	10, 20, 30, 40, 60, 80, 100	3
Total number of	48	
Total number of	184	

Esparza et al., 2018):

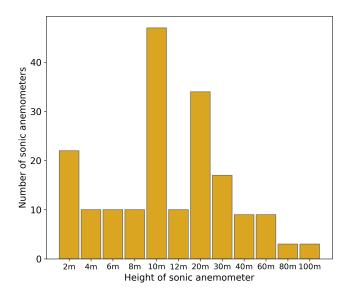
$$\epsilon = \frac{1}{U\tau} [aD_U(\tau)]^{3/2},\tag{2}$$

where  $\tau$  indicates the time lags over which the structure function is calculated, and a=0.52 is the one-dimensional Kolmogorov constant (Paquin and Pond, 1971; Sreenivasan, 1995). We calculate  $\epsilon$  every 30 s, and then average values at a 30 min resolution. At each calculation of  $\epsilon$ , we fit experimental data to the Kolmogorov model (Kolmogorov, 1941; Frisch, 1995) using time lags between  $\tau_1=0.1$  s and

 $\tau_2 = 2 \, \text{s}$ , which represent a conservative choice to approximate the inertial subrange (Bodini et al., 2018).

To account for the uncertainty in the calculation of  $\epsilon$ , we apply the law of combination of errors, which tracks how random errors propagate through a series of calculations (Barlow, 1989). We apply this method to Eq. (2) and quantify the fractional standard deviation in the  $\epsilon$  estimates (Piper, 2001; Wildmann et al., 2019) as

$$\sigma_{\epsilon} = \frac{3}{2} \frac{\sigma_{I}}{I} \epsilon,\tag{3}$$



**Figure 2.** Histogram of the heights a.g.l. of the 184 sonic anemometers considered in this analysis.

where I is the sample mean of  $\tau^{-2/3}D_U(\tau)$ , and  $\sigma_I^2$  is its sample variance. To perform our analysis only on  $\epsilon$  values with lower uncertainty, we discard dissipation rates characterized by  $\sigma_\epsilon > 0.05$ . About 3% of the data are discarded based on this criterion.

As additional quality controls, to exclude tower wake effects, data have been discarded when the recorded wind direction was within  $\pm 30^{\circ}$  of the direction of the tower boom. Data during precipitation periods (as recorded by a precipitation sensor on the tower "riSW06" on the southwest ridge) have also been discarded from further analysis. After all the quality controls have been applied, a total (from all sonic anemometers) of over 284 000 30 min average  $\epsilon$  data remains for the analysis.

## 3 Accuracy of current parameterization of TKE dissipation rate in mesoscale models

Before testing the performance of machine-learning algorithms in predicting TKE dissipation rates, we first assess the current accuracy of the parameterization of  $\epsilon$  in numerical models. In the Weather Research and Forecasting model (WRF; Skamarock et al., 2005), the most widely used numerical weather prediction model, turbulence in the boundary layer can be represented with several planetary-boundary-layer (PBL) schemes, most of which implicitly assume a local balance between turbulence production and dissipation. Among the different PBL schemes, the MYNN scheme is one of the most commonly chosen. Turbulence dissipation rate in MYNN is given (Nakanishi, 2001) as a function of TKE as

$$\epsilon = \frac{(2 \text{ TKE})^{3/2}}{B_1 L_{\text{M}}},\tag{4}$$

where  $B_1 = 24$ , and the master length scale,  $L_{\rm M}$ , is defined with a diagnostic equation, based on large-eddy simulations, as a function of three other length scales:

$$\frac{1}{L_{\rm M}} = \frac{1}{L_{\rm S}} + \frac{1}{L_{\rm T}} + \frac{1}{L_{\rm B}}.$$
 (5)

 $L_{\rm S}$  is the length scale in the surface layer, given by

$$L_{S} = \begin{cases} \kappa \ z/3.7 & \zeta \ge 1 \\ \kappa \ z \ (1 + 2.7 \ \zeta)^{-1} & 0 \le \zeta < 1 \\ \kappa \ z \ (1 - \alpha_4 \ \zeta)^{0.2} & \zeta < 0, \end{cases}$$
(6)

where  $\kappa = 0.4$  is the von Kármán constant,  $\zeta = z/L$  (with L the Obukhov length), and  $\alpha_4 = 100.0$ .

*L*<sub>T</sub> is the length scale depending upon the turbulent structure of the PBL (Mellor and Yamada, 1974), defined as

$$L_{\rm T} = \alpha_1 \frac{\int_0^\infty q \ z \ dz}{\int_0^\infty q \ dz},\tag{7}$$

where  $q = \sqrt{2 \text{ TKE}}$ , and  $\alpha_1 = 0.23$ .

 $L_{\rm B}$  is a length scale limited by the buoyancy effect, given by

$$L_{\rm B} = \begin{cases} \alpha_2 \ q/N & \partial\Theta/\partial z > 0 \text{ and } \zeta \ge 0\\ \frac{\alpha_2 \ q + \alpha_3 \ q \ (q_c/L_{\rm T} \ N)^{1/2}}{N} & \partial\Theta/\partial z > 0 \text{ and } \zeta < 0 \\ \infty & \partial\Theta/\partial z \le 0, \end{cases}$$
(8)

with N the Brunt-Väisälä frequency,  $\Theta$  the mean potential temperature,  $\alpha_2 = 1.0$ ,  $\alpha_3 = 5.0$ , and  $q_c = \left[ (g/\Theta_0 \overline{w'\theta'} L_{\rm T}) \right]^{1/3}$ .

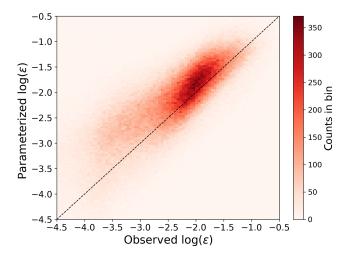
From the available observations from the meteorological towers at Perdigão, only  $L_{\rm S}$  can be determined, while the calculation of  $L_{\rm T}$  and  $L_{\rm B}$  would only be possible with critical assumptions about the vertical profile of TKE. Therefore, we decide to approximate  $L_{\rm M}$  as

$$\frac{1}{L_{\rm M}} \approx \frac{1}{L_{\rm S}}.\tag{9}$$

By doing so,  $L_{\rm M}$  is overestimated (proof shown in the Supplement), which in turn implies that  $\epsilon$  calculated using Eq. (4) will be underestimated.

To evaluate the accuracy of the MYNN parameterization of  $\epsilon$ , we calculated, using 30 min average data, the parameterized  $\epsilon$  using Eq. (4) (with the approximation in Eq. 9) from all of the 184 sonic anemometers considered in the study and compared with the observed values of TKE dissipation rate (Fig. 3) derived from the sonic anemometers with Eq. (2). Given the extremely large range of variability of  $\epsilon$ , we calculate all the error metrics using the logarithm of predicted and observed  $\epsilon$ .

The TKE dissipation rate predicted by the MYNN parameterization shows, on average, a large positive bias compared to the observed values, with a mean bias of +12% in terms



**Figure 3.** Scatter plot showing the comparison between observed and MYNN-parameterized  $\epsilon$  from the 184 sonic anemometers at Perdigão.

of the logarithm of  $\epsilon$ , +47% in terms of  $\epsilon$ . The root mean square error (RMSE) is 0.61, and the mean absolute error (MAE) is 0.46. The observed bias would be even larger if  $L_{\rm M}$  was calculated including all the contributions according to Eq. (5), and not  $L_{\rm s}$  only as in our approximation. Therefore, while the approximation in Eq. (9) is major and could be eased by making assumptions on the vertical profile of TKE at Perdigão, it does not affect the conclusion of a high inaccuracy in the MYNN parameterization of  $\epsilon$ .

Different atmospheric stability conditions give different biases. Figure 4 compares observed and parameterized  $\epsilon$  values for stable and unstable conditions, classified based on  $\zeta = z/L$ , measured at each sonic anemometer, according to the thresholds described in Sect. 2.1.

Stable cases show the largest bias (mean of +24% in terms of the logarithm of  $\epsilon$ , +101% in terms of  $\epsilon$ ), whereas for unstable conditions the bias is smaller (mean of +6% in terms of the logarithm of  $\epsilon$ , +19% in terms of  $\epsilon$ ). The MYNN parameterization of  $\epsilon$  is therefore especially inadequate to represent small values of  $\epsilon$ , which mainly occur in stable conditions.

Different heights also impact the accuracy of the parameterization of  $\epsilon$ .

As shown in Fig. 5, the mean bias in parameterized  $\log(\epsilon)$  decreases with height, while its spread (quantified in terms of the standard deviation of the bias at each height) does not show a large variability at different levels. Close to the surface (data from the sonic anemometers at 2 m a.g.l.), a mean bias (in logarithmic space) of about +25% is found, whereas for the sonic anemometers at 100 m a.g.l., we find a mean bias of just  $\sim$  +3%. This difference in bias with height becomes much larger if the bias is calculated on the actual  $\epsilon$  values (and not on their logarithm). We obtain comparable results when computing the bias in the MYNN pa-

rameterization only for the sonic anemometers mounted on the three 100 m meteorological towers (figure shown in the Supplement), thus confirming that the observed trend is not due to the larger variability of the conditions sampled by the more numerous sonics at lower heights. Therefore, our results show how the MYNN formulation fails in accurately representing atmospheric turbulence especially in the lowest part of the boundary layer.

## 4 Machine-learning algorithms

To test the power of machine learning to improve the numerical representation of the TKE dissipation rate, we consider three learning algorithms in this study: multivariate linear regression, multivariate third-order polynomial regression, and random forest. Given the proof-of-concept nature of this analysis in proving the capabilities of machine learning to improve numerical model parameterizations, we defer an exhaustive comparison of different machine-learning models to a future study and only consider relatively simple algorithms in the present work. The learning algorithms are trained and tested to predict the logarithm of  $\epsilon$  using 30 min average data. For all but the random forest algorithm, the data were scaled and normalized by removing the mean and scaling to unit variance. No data imputation was performed, and missing data were removed from the analysis.

For the purpose of machine-learning algorithms, the data set has to be divided into three subsets: training, validation, and testing sets (Friedman et al., 2001). The algorithms are first trained multiple times with different hyperparameters (model parameters whose values are set before the training phase and that control the learning process) on the training set, then the validation set is used to choose the best set of hyperparameters, and finally the predicting performance of the trained algorithm is assessed on the testing set. Usually, the data set is split randomly into training, validation, and testing sets. However, as the data used in this study consist of observations averaged every 30 min, data in contiguous time stamps are likely characterized by some autocorrelation. Therefore, the traditional random split between training and testing data would lead to an artificially enhanced performance of the machine-learning algorithms, which would be tested on data with a large autocorrelation with the ones used for the training. Therefore, here we use one concurrent week of the data for testing ( $\sim 17\%$  of the data), whereas the other 5 weeks are split between training (4 weeks, 66 % of the data) and validation (1 week, 17 % of the data). The 1-week testing period is shifted continuously throughout the considered 6 weeks of observations at Perdigão, so that each model is trained and its prediction performance tested six times. For each algorithm, we evaluate the overall performance based on the RMSE between the actual and predicted (logarithm of)  $\epsilon$ , averaged over the different week-long testing periods.

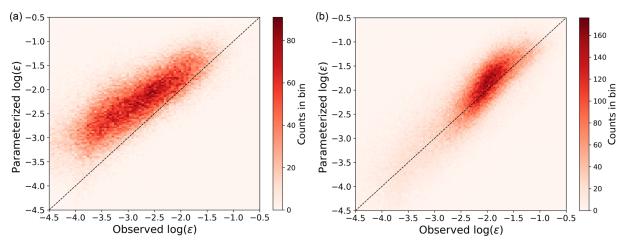
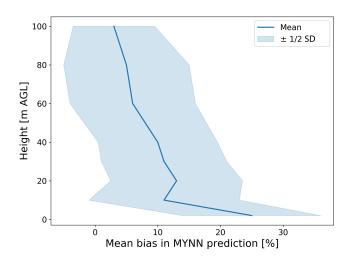


Figure 4. Scatter plot showing the comparison between observed and MYNN-parameterized  $\epsilon$  from the 184 sonic anemometers at Perdigão for stable conditions (a) and unstable conditions (b), as quantified by  $\zeta = z/L$  calculated at each sonic anemometer.



**Figure 5.** Bias in the MYNN-parameterized  $\log(\epsilon)$  at different heights, as calculated from the 184 sonic anemometers at Perdigão.

Before testing the models, however, it is important to avoid overfitting by setting the values of hyperparameters. Each learning algorithm has specific model-specific hyperparameters that need to be considered, as will be specified in the description of each algorithm. To test different combinations of hyperparameters and determine the best set, we use cross validation with randomized search, with 20 parameter sets sampled for each learning algorithm. For each set of hyperparameters, the RMSE between the actual and predicted  $\log(\epsilon)$ in the validation test is calculated. For each model, we select the hyperparameter combination (among the ones surveyed in the cross validation) that leads to the lowest mean (across the five validation sets) RMSE. We then use this set as the final combination for assessing the performance of the models on the testing set. Overall, the procedure is repeated six times by shifting the 1-week testing set (Fig. 6).

In the following paragraphs, we describe the main characteristics of the three machine-learning algorithms used in our study. A more detailed description can be found in machine-learning textbooks (Hastie et al., 2009; Géron, 2017).

## 4.1 Multivariate linear regression

To check whether simple learning algorithms can improve the current numerical parameterization of  $\epsilon$ , we test the accuracy of multivariate linear regression:

$$\log(\hat{\epsilon}) = \theta_0 + \theta_1 x_1 + \theta_2 x_2 + \ldots + \theta_n x_n, \tag{10}$$

where  $\hat{\epsilon}$  is the machine-learning predicted value of  $\epsilon$ , n is the number of features used to predict  $\epsilon$  (here, six; see Sect. 4.4),  $x_i$  is the ith feature value, and  $\theta_i$  is the jth model weight.

To avoid training a model that overfits the data, regularization techniques need to be implemented, so that the learning model is constrained: the fewer degrees of freedom the model has, the harder it will be for it to overfit the data. We use ridge regression (Hoerl and Kennard, 1970) (Ridge in Python's library Scikit-learn) to constrain the multivariate regression. Ridge regression constrains the weights of the model  $\theta_j$  to have them stay as small as possible. The ridge regression is achieved by adding a regularization term to the cost function (MSE):

$$J(\theta) = \text{MSE}(\theta) + \alpha \sum_{i=1}^{n} \theta_i^2,$$
(11)

where the hyperparameter  $\alpha$  controls how much the model will be regularized. The optimal value of the hyperparameter  $\alpha$  is determined by cross validation, as described earlier, with values sampled in the range from 0.1 to 10.

## 4.2 Multivariate third-order polynomial regression

Multivariate polynomial regression can easily be achieved by adding powers of each input feature as new features. The re-

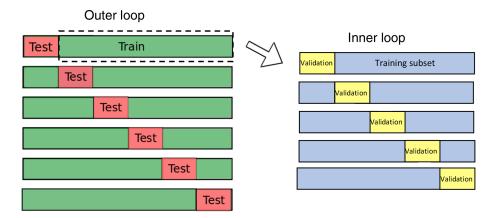


Figure 6. Cross-validation approach used to evaluate the performance of the machine-learning models considered in this study.

gression algorithm is then trained as a linear model on this extended set of features. For a third-order polynomial regression, the model becomes

$$\log(\hat{\epsilon}) = \theta_0 + \sum_{i=1}^n \theta_i \ x_i + \sum_{i=1}^n \theta_{ii} \ x_i^2 + \sum_{i=1}^{n-1} \sum_{j=i+1}^n \theta_{ij} \ x_i \ x_j$$
$$+ \sum_{i=1}^n \theta_{iii} \ x_i^3 + \sum_{i=1}^n \sum_{j\neq i} \theta_{iij} \ x_i^2 \ x_j$$
$$+ \sum_{i=1}^{n-2} \sum_{j=i+1}^{n-1} \sum_{k=j+1}^n \theta_{ijk} \ x_i \ x_j \ x_k.$$
 (12)

Ridge regression (Ridge in Python's Scikit-learn library) is used again to constrain the multivariate polynomial regression, with the hyperparameter  $\alpha$  in Eq. (11) determined via cross validation, with values sampled in the range from 1 to 2000.

### 4.3 Random forest

Random forests (RandomForestRegressor in Python's Scikit-learn library) combine multiple decision trees to provide an ensemble prediction. A decision tree can learn patterns and then predict values by recursively splitting the training data based on thresholds of the different input features. As a result, the data are divided into groups, each associated with a single predicted value of  $\epsilon$ , calculated as the average target value (of the observed  $\epsilon$ ) of the instances in that group.

As an ensemble of decision trees, a random forest trains them on different random subsets of the training set. Once all the predictors are trained, the ensemble (i.e., the random forest) can make a prediction for a new instance by taking the average of all the predictions from the single trees. In addition, random forests introduce some extra randomness when growing trees: instead of looking for the feature that, when split, reduces the overall error the most when splitting a node, a random forest searches for the best feature among a random subset of features.

Decision trees make very few assumptions about the training data. As such, if unconstrained, they will adapt their structure to the training data, fitting them closely, and most likely overfitting them, without then being able to provide accurate predictions on new data. To avoid overfitting, regularization can be achieved by setting various hyperparameters that insert limits to the structure of the trees used to create the random forests. Table 2 describes which hyperparameters we considered for the random forest algorithm. For each hyperparameters listed, we include the range of values that are randomly sampled in the cross-validation search to form the 20 sets of hyperparameters considered in the training phase.

## 4.4 Input features for machine-learning algorithms

Given the large variability of  $\epsilon$ , which can span several orders of magnitude (Bodini et al., 2019b), we apply the machine-learning algorithms to predict the logarithm of  $\epsilon$ . To select the set of input features used by the learning models, we take advantage of the main findings of the observational studies on the variability of  $\epsilon$  to select as inputs both atmospheric-and terrain-related variables to capture the impact of topography on atmospheric turbulence. For each variable, we calculate and use in the machine-learning algorithms 30 min average data, to reduce the high autocorrelation in the data and limit the impact of the high-frequency large variability of turbulent quantities. We use the following input features (calculated at the same location and height as  $\epsilon$ ) for the three learning algorithms considered in our study:

- Wind speed (WS). Bodini et al. (2018) found that WS has a moderate correlation with  $\epsilon$ .
- The logarithm of TKE. This quantity is calculated as

$$\log(\text{TKE}) = \log\left[\frac{1}{2}\left(\sigma_u^2 + \sigma_v^2 + \sigma_w^2\right)\right],\tag{13}$$

Hyperparameter	Meaning	Sampled values
Number of estimators	Number of trees in the forest	10–250
Maximum depth	Maximum depth of the tree	1–50
Maximum number of leaf nodes	Maximum number of leaf nodes in the decision tree	2000–500 000
Maximum number of features	Number of features to consider when looking for the best split	1–6
Minimum number of samples to split	Minimum number of samples required to split an internal node	1–200
Minimum number of samples for a leaf	Minimum number of samples required to be at a leaf node	1–50

**Table 2.** Hyperparameters considered for the random forest algorithm.

where the variances of the wind components are calculated over 30 min intervals. The choice of using the logarithm of TKE is justified by the fact Eq. (4) suggests this quantity is linearly related to the logarithm of  $\epsilon$ .

– The logarithm of friction velocity of  $u_*$ . This quantity is calculated as

$$u_* = (\overline{u'w'}^2 + \overline{v'w'}^2)^{1/4}.$$
 (14)

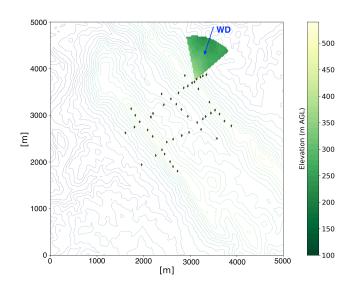
An averaging period of 30 min (De Franceschi and Zardi, 2003; Babić et al., 2012) has been used to apply the Reynolds decomposition and calculate average quantities and fluctuations.

- The log-modulus transformation (John and Draper, 1980) of the ratio  $\zeta = z_{\rm son}/L$ , where  $z_{\rm son}$  is the height above the ground of each sonic anemometer, and L is the 30 min average Obukhov length:

$$\operatorname{sign}(\zeta) \log(|\zeta| + 1). \tag{15}$$

The use of  $\zeta$  is justified within the context of the Monin–Obukhov similarity theory (Monin and Obukhov, 1954). The use of the logarithm of  $\zeta$  is consistent with the use of the logarithm of  $\epsilon$  as target variable. Finally, the log-modulus transformation allows for the logarithm to be calculated on negative values of  $\zeta$  and be continuous in zero.

The standard deviation SD(z<sub>terr</sub>) of the terrain elevation in a 1 km radius sector centered on the measurement point (i.e., the location of the sonic anemometer). The angular extension of the sector is set equal to ±30° from the recorded 30 min average wind direction (an example is shown in Fig. 7). While we acknowledge that some degree of arbitrariness lies in the choice of this variable to quantify the terrain influence, it represents a quantity that can easily be derived from numerical models, should our approach be implemented



**Figure 7.** Example of an upwind terrain elevation sector with a 1 km radius centered on the location of one of the meteorological towers at Perdigão.

for practical applications, to capture the influence of upwind topography to trigger turbulence. To compute this variable, we use Shuttle Radar Topography Mission (SRTM) 1 arcsec global data, at 30 m horizontal resolution.

- The mean vegetation height  $h_{\rm veg}$  in the upwind 1 km radius sector centered on the measurement point. Given the forested nature of the Perdigão region, we expect canopy to have an effect in triggering turbulence, especially at lower heights. To compute this variable, we use data from a lidar survey during the season of the field campaign, at a 20 m horizontal resolution.

The distribution of the input features and of  $\log(\epsilon)$  are shown in the Supplement.

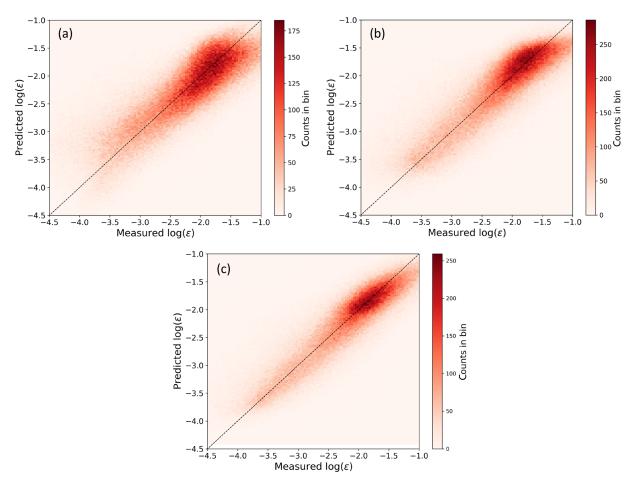


Figure 8. Scatter plot showing the comparison, performed on the testing set, between observed and machine-learning-predicted  $\epsilon$  from a multivariate linear regression (a), a multivariate third-order polynomial regression (b), and a random forest (c).

While we acknowledge that the input features are not fully uncorrelated, we found that including all these features provides a better predictive power for the learning algorithms, despite negatively affecting the computational requirements of the training phase. The application of principal component analysis can help reduce the number of dimensions in the input features while preserving the predictive power of each, but it is beyond the scope of the current work.

#### 5 Results

## **5.1** Performance of machine-learning algorithms

To evaluate the prediction performance of the three machine-learning algorithms we considered, we use, for each method, a scatter plot showing the comparison between observed and machine-learning-predicted  $\epsilon$  (Fig. 8).

The predictions from all the considered learning algorithms do not show a significant mean bias, as found in the MYNN representation of  $\epsilon$ . As specific error metrics, we compare RMSE and MAE of the machine-learning pre-

dictions with what we obtained from the MYNN parameterization, with the caveat that while the MYNN scheme is thought to provide a universal representation of  $\epsilon$ , the machine-learning models have been specifically trained on data from a single field campaign. Each machine-learning algorithm was tested on six 1-week long testing periods, as described in Sect. 4. For each method, we present the RMSE and MAE averaged across the different testing periods. Even the simple multivariate linear regression (Fig. 8a) improves, on average, on MYNN. Overall, the average RMSE (0.47) is 23 % smaller than the MYNN parameterization, and the average MAE (0.36) is 22 % lower than the MYNN prediction. The multivariate third-order polynomial regression provides an additional improvement (Fig. 8b) for the representation of  $\epsilon$ , with the average RMSE (0.44) over 28 % smaller than the MYNN parameterization, and the average MAE (0.33) 28 % lower than the MYNN representation. The additional input features created by the polynomial model allow for an accurate prediction of  $\epsilon$  even in the low turbulence regime. Finally, the random forest further reduces the spread in machine-learning predicted  $\epsilon$ , with the RMSE (0.40) re-

-28%

-37%

MYNN Linear Third-order Random parameterization regression polynomial regression forest **RMSE** 0.61 0.47 0.44 0.40 % change in RMSE -23%-28%-35%MAE 0.46 0.36 0.33 0.29

-22%

**Table 3.** Performance of the machine-learning algorithms trained and tested at Perdigão, measured in terms of RMSE and MAE between the logarithm of observed and MYNN-parameterized  $\epsilon$ .

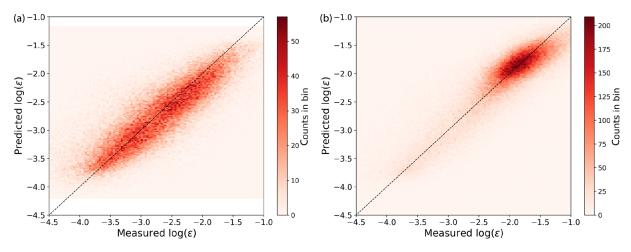


Figure 9. Scatter plot showing the comparison, performed on the testing set, between observed and machine-learning-predicted  $\epsilon$  from a random forest for stable conditions (a) and unstable conditions (b).

duced by about 35 % from the MYNN case, and the MAE (0.29) by 37 %, with no average bias between observed and predicted values of  $\epsilon$ .

% change in MAE

Table 3 summarizes the performance of all the considered algorithms. We note that, because the length scale approximation we made in calculating MYNN-predicted  $\epsilon$  led to a better agreement with the observed values compared to what would be obtained using the full MYNN parameterization, the RMSE and MAE for the MYNN case would in reality be higher than what we report here, and so the error reductions achieved with the machine-learning algorithms would even be greater than the numbers shown in the table.

Given the large gap in the performance of the MYNN parameterization of  $\epsilon$  between stable and unstable conditions, it is worth exploring how the machine-learning algorithms perform in different stability conditions. To do so, we train and test two separate random forests: one using data observed in stable conditions, the other one for unstable cases. We find that both algorithms eliminate the bias observed in the MYNN scheme (Fig. 9).

The random forest for unstable conditions provides, on average, more accurate predictions (RMSE of 0.37, MAE of 0.28) compared to the algorithm used for stable cases (RMSE of 0.44, MAE of 0.33), thus confirming the complexity in modeling atmospheric turbulence in quiescent condi-

tions. However, when the error metrics are compared to those of the MYNN parameterization, the random forest for stable conditions provides the largest relative improvement, with a  $50\,\%$  reduction in MAE, while for unstable conditions the reduction is of  $20\,\%$ .

## 5.2 Physical interpretation of machine-learning results

Not only do machine-learning techniques provide accuracy improvements to represent atmospheric turbulence, but additional insights on the physical interpretation of the results can – and should – be achieved. In particular, random forests allow for an assessment of the relative importance of the input features used to predict (the logarithm of)  $\epsilon$ . The importance of a feature is calculated by looking at how much the tree nodes that use that feature reduce the MSE on average (across all trees in the forest), weighted by the number of times the feature is selected. Table 4 shows the feature importance for the six input features we used in this study.

The feature importance results are affected by the correlation between some of the input features used in the models. We find how the logarithm of turbulence kinetic energy is the preferred feature for tree splitting, with the largest importance (47 %) in reducing the prediction error for  $\epsilon$  in the random forest. This result, which can be expected as both TKE

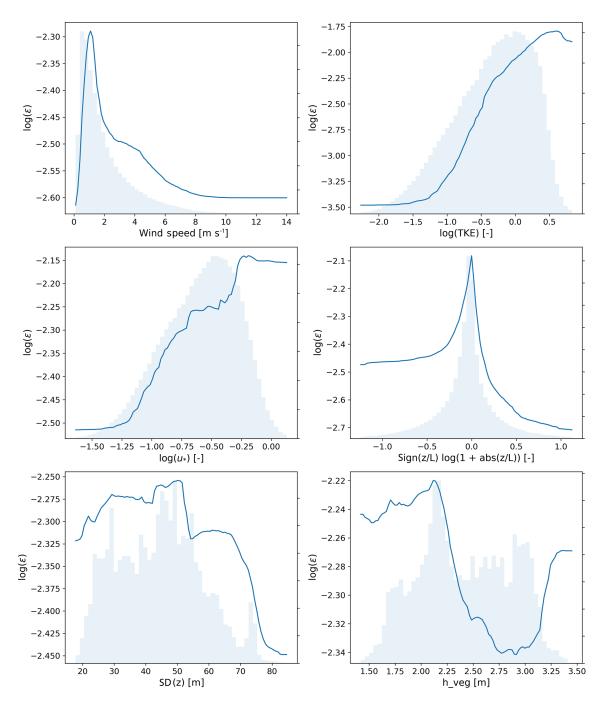


Figure 10. Partial dependence plots for the input features used in the analysis. Distributions of the considered features are shown in the background.

and  $\epsilon$  are variables connected to turbulence in the boundary layer, agrees well with the current formulation of the MYNN parameterization of  $\epsilon$ , which includes TKE as main term. As TKE is correlated to  $u_*$  and z/L, we find that the decision trees more often split the data based on TKE, so that the feature importance of its correlated variables is found to be lower. The limitations of the Monin–Obukhov similarity theory (Monin and Obukhov, 1954) in complex terrain

might also be an additional cause for the relatively low feature importance of the feature associated with L. The standard deviation of the upwind elevation and the mean vegetation height have the lowest importance of, respectively, 3% and 2%. Though not negligible, the importance of topography and canopy might increase by considering different parameters that could better encapsulate their effect. Also, the impact of topography and canopy might be hidden as it could

**Table 4.** Feature importance classification as derived from the random forest.

Input feature	Feature importance
log(TKE)	47 %
$\log(u_*)$	24 %
$sign(\zeta) log( \zeta  + 1)$	13 %
WS	11 %
$SD(z_{terr})$	3 %
$\overline{h_{ ext{veg}}}$	2 %

be already incorporated in the variability of parameters with larger relative importance, such as TKE. We have tested how the feature importance varies when considering several random forests, each trained and tested with data from all the sonic anemometers at a single height only, and did not find any significant variation of the importance of the considered variables in predicting  $\epsilon$  (plot shown in the Supplement).

Finally, to assess the dependence of TKE dissipation rate on the individual features considered in this study, Fig. 10 shows partial dependence plots for the input features considered in the analysis. These are obtained, for each input feature, by applying the machine-learning algorithm (here, random forests) multiple times with the other feature variables constant (at their means) while varying the target input feature and measuring the effect on the response variable (here,  $\log(\epsilon)$ ). In each plot, the values on the y axes have not been normalized, so that large ranges show a strong dependence of  $\log(\epsilon)$  on the feature, whereas small ranges indicate weaker dependence.

The strong relationship between  $\epsilon$  and TKE is confirmed, as the range shown on its y axis is the largest among all features. As TKE increases, so does  $\epsilon$ . A similar trend, tough with a weaker influence, emerges when considering the dependence of  $\epsilon$  on friction velocity. The relationship between  $\epsilon$  and wind speed shows a less clear trend, and with a weaker dependence:  $\epsilon$  increases for 30 min average wind speeds up to  $\sim 2 \,\mathrm{m \, s^{-1}}$ , and then decreases for stronger wind speed values. A more distinct trend could emerge when considering data averaged at shorter time periods. The dependence between TKE dissipation and atmospheric stability shows a moderate impact, with stable conditions (positive values of the considered metric) showing smaller  $\epsilon$  values compared to unstable cases (negative values of the considered metric). Interestingly, the largest  $\epsilon$  values seem to be connected to neutral cases. Finally, both terrain elevation and vegetation height show a weak impact on determining the values of  $\epsilon$ , as quantified by the narrow range of values sampled on the y axis for these two variables.

#### 6 Conclusions

Despite turbulence being a fundamental quantity for the development of multiple phenomena in the atmospheric boundary layer, the current representations of TKE dissipation rate  $(\epsilon)$  in numerical weather prediction models suffer from large inaccuracies. In this study, we quantified the error introduced in the MYNN parameterization of  $\epsilon$  by comparing predicted and observed values of  $\epsilon$  from 184 sonic anemometers from 6 weeks of observations at the Perdigão field campaign. A large positive bias (average +12 % in logarithmic space, +47 % in natural space) emerges, with larger errors found in atmospheric stable conditions. The need for a more accurate representation of  $\epsilon$  is therefore clearly demonstrated.

The results of this study show how machine learning can provide new ways to successfully represent TKE dissipation rate from a set of atmospheric and topographic parameters. Even simple models such as a multivariate linear regression can provide an improved representation of  $\epsilon$  compared to the current MYNN parameterization. More sophisticated algorithms, such as a random forest approach, lead to the largest benefits, with over a 35 % reduction in the average error introduced in the parameterization of  $\epsilon$ , and eliminate the large bias found in it, for the Perdigão field campaign. When considering stable conditions only, the reduction in average error reaches 50 %. Although the generalization gap between the universal nature of the MYNN parameterization of  $\epsilon$  and the campaign-specific training and testing of the machinelearning models needs to be acknowledged, the results of this study can be considered as a proof of concept of the potentialities of machine-learning-based representations of complex atmospheric processes.

Multiple opportunities exist to extend the work presented here. In the future, additional learning algorithms, such as support vector machines and extremely randomized trees, should be considered. Deep learning methods, such as recurrent neural networks, and specifically long- to short-term memory, which are well suited for time-series-based problems, could also be considered to obtain a more complete overview of the capabilities of machine-learning techniques for improving numerical representations of  $\epsilon$ . Moreover, additional input features could be added to the learning algorithms to possibly identify additional variables with a large impact on atmospheric turbulence. Finally, the learning algorithms developed here would need to be tested using data from different field experiments to understand whether the results obtained in this study can be generalized everywhere. Data collected in flat terrain and offshore would likely need to be considered to create a more universal model to predict dissipation in various terrains. Once the performance of a machine-learning representation of  $\epsilon$  has been accurately tested, its implementation in numerical weather prediction models, such as the Weather Research and Forecasting model, should be achieved.

Code and data availability. High-resolution data from anemometers on the meteorological towers (UCAR/NCAR, 2019) are available through the EOL project at https://doi.org/10.26023/8X1N-TCT4-P50X. Digital elevation model data are taken from the SRTM 1 arcsec global at https://doi.org/10.5066/F7PR7TFT (USGS, 2020). The vegetation height data are available upon request to Jose Laginha Palma at the University of Porto. The machine-learning code used for the analysis is stored at https://doi.org/10.5281/zenodo.3754710 (Bodini, 2020).

*Supplement.* The supplement related to this article is available online at: https://doi.org/10.5194/gmd-13-4271-2020-supplement.

Author contributions. NB and JKL designed the analysis. NB analyzed the data from the sonic anemometers and applied the machine-learning figures, in close consultation with JKL and MO. NB wrote the paper, with significant contributions from JKL and MO.

Competing interests. The authors declare that they have no conflict of interest.

*Disclaimer.* The views expressed in the article do not necessarily represent the views of the DOE or the US Government.

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### References

- Albertson, J. D., Parlange, M. B., Kiely, G., and Eichinger, W. E.: The average dissipation rate of turbulent kinetic energy in the neutral and unstable atmospheric surface layer, J. Geophys. Res.-Atmos., 102, 13423–13432, 1997.
- Arcos Jiménez, A., Gómez Muñoz, C., and García Márquez, F.: Machine learning for wind turbine blades maintenance management, Energies, 11, 13, 2018.
- Babić, K., Bencetić Klaić, Z., and Večenaj, Ž.: Determining a turbulence averaging time scale by Fourier analysis for the nocturnal boundary layer, Geofizika, 29, 35–51, 2012.
- Barlow, R. J.: Statistics: a guide to the use of statistical methods in the physical sciences, vol. 29, John Wiley & Sons, 1989.
- Berg, L. K., Liu, Y., Yang, B., Qian, Y., Olson, J., Pekour, M., Ma, P.-L., and Hou, Z.: Sensitivity of Turbine-Height Wind Speeds to Parameters in the Planetary Boundary-Layer Parametrization Used in the Weather Research and Forecasting Model: Extension to Wintertime Conditions, Bound.-Lay. Meteorol., 170, 507–518, 2018.
- Bodini, N.: Random forest for TKE dissipation rate gmd-2020-16 paper, Zenodo, https://doi.org/10.5281/zenodo.3754710, 2020.
- Bodini, N., Zardi, D., and Lundquist, J. K.: Three-dimensional structure of wind turbine wakes as measured by scanning lidar, Atmos. Meas. Tech., 10, 2881–2896, https://doi.org/10.5194/amt-10-2881-2017, 2017.
- Bodini, N., Lundquist, J. K., and Newsom, R. K.: Estimation of turbulence dissipation rate and its variability from sonic anemometer and wind Doppler lidar during the XPIA field campaign, Atmos. Meas. Tech., 11, 4291–4308, https://doi.org/10.5194/amt-11-4291-2018, 2018.
- Bodini, N., Lundquist, J. K., and Kirincich, A.: US East Coast Lidar Measurements Show Offshore Wind Turbines Will Encounter Very Low Atmospheric Turbulence, Geophys. Res. Lett., 46, 5582–5591, 2019a.
- Bodini, N., Lundquist, J. K., Krishnamurthy, R., Pekour, M., Berg, L. K., and Choukulkar, A.: Spatial and temporal variability of turbulence dissipation rate in complex terrain, Atmos. Chem. Phys., 19, 4367–4382, https://doi.org/10.5194/acp-19-4367-2019, 2019b.
- Cervone, G., Clemente-Harding, L., Alessandrini, S., and Delle Monache, L.: Short-term photovoltaic power forecasting using Artificial Neural Networks and an Analog Ensemble, Renew. Energ., 108, 274–286, 2017.
- Champagne, F. H., Friehe, C. A., LaRue, J. C., and Wynagaard, J. C.: Flux measurements, flux estimation techniques, and fine-scale turbulence measurements in the unstable surface layer over land, J. Atmos. Sci., 34, 515–530, 1977.
- Clifton, A., Kilcher, L., Lundquist, J., and Fleming, P.: Using machine learning to predict wind turbine power output, Environ. Res. Lett., 8, 024009, https://doi.org/10.1088/1748-9326/8/2/024009, 2013.
- Coen, J. L., Cameron, M., Michalakes, J., Patton, E. G., Riggan, P. J., and Yedinak, K. M.: WRF-Fire: coupled weather–wildland

- fire modeling with the weather research and forecasting model, J. Appl. Meteorol. Climatol., 52, 16–38, 2013.
- De Franceschi, M. and Zardi, D.: Evaluation of cut-off frequency and correction of filter-induced phase lag and attenuation in eddy covariance analysis of turbulence data, Bound.-Lay. Meteorol., 108, 289–303, 2003.
- Fernando, H. J., Mann, J., Palma, J. M., Lundquist, J. K., Barthelmie, R. J., Belo Pereira, M., Brown, W. O., Chow, F. K., Gerz, T., Hocut, C., Klein, P., Leo, L., Matos, J., Oncley, S., Pryor, S., Bariteau, L., Bell, T., Bodini, N., Carney, M., Courtney, M., Creegan, E., Dimitrova, R., Gomes, S., Hagen, M., Hyde, J., Kigle, S., Krishnamurthy, R., Lopes, J., Mazzaro, L., Neher, J., Menke, R., Murphy, P., Oswald, L., Otarola-Bustos, S., Pattantyus, A., Rodrigues, C. V., Schady, A., Sirin, N., Spuler, S., Svensson, E., Tomaszewski, J., Turner, D., van Veen, L., Vasiljević, N., Vassallo, D., Voss, S., Wildmann, N., and Wang, Y.: The Perdigão: Peering into Microscale Details of Mountain Winds, B. Am. Meteorol. Soc., 100, 799–819, https://doi.org/10.1175/BAMS-D-17-0227.1, 2018.
- Frehlich, R.: Coherent Doppler lidar signal covariance including wind shear and wind turbulence, Appl. Opt., 33, 6472–6481, 1994.
- Frehlich, R., Meillier, Y., Jensen, M. L., Balsley, B., and Sharman, R.: Measurements of boundary layer profiles in an urban environment, J. Appl. Meteorol. Climatol., 45, 821–837, 2006.
- Friedman, J., Hastie, T., and Tibshirani, R.: The elements of statistical learning, vol. 1, Springer series in statistics New York, 2001.
- Frisch, U.: Turbulence: the legacy of A.N. Kolmogorov, Cambridge University Press, 1995.
- Garratt, J. R.: The atmospheric boundary layer, Earth-Sci. Rev., 37, 89–134, 1994.
- Gentine, P., Pritchard, M., Rasp, S., Reinaudi, G., and Yacalis, G.: Could machine learning break the convection parameterization deadlock?, Geophys. Res. Lett., 45, 5742–5751, 2018.
- Géron, A.: Hands-on machine learning with Scikit-Learn and TensorFlow: concepts, tools, and techniques to build intelligent systems, O'Reilly Media, Inc., 2017.
- Gerz, T., Holzäpfel, F., Bryant, W., Köpp, F., Frech, M., Tafferner, A., and Winckelmans, G.: Research towards a wake-vortex advisory system for optimal aircraft spacing, C. R. Phys., 6, 501–523, 2005.
- Giebel, G., Brownsword, R., Kariniotakis, G., Denhard, M., and Draxl, C.: The state-of-the-art in short-term prediction of wind power: A literature overview, ANEMOS plus, 2011.
- Hastie, T., Tibshirani, R., and Friedman, J.: The elements of statistical learning: data mining, inference, and prediction, Springer Science & Business Media, 2009.
- Hoerl, A. E. and Kennard, R. W.: Ridge regression: Biased estimation for nonorthogonal problems, Technometrics, 12, 55–67, 1970.
- Hong, S.-Y. and Dudhia, J.: Next-generation numerical weather prediction: Bridging parameterization, explicit clouds, and large eddies, B. Am. Meteorol. Soc., 93, ES6–ES9, 2012.
- Huang, K., Fu, J. S., Hsu, N. C., Gao, Y., Dong, X., Tsay, S.-C., and Lam, Y. F.: Impact assessment of biomass burning on air quality in Southeast and East Asia during BASE-ASIA, Atmos. Environ., 78, 291–302, 2013.
- John, J. and Draper, N. R.: An alternative family of transformations, J. R. Stat. Soc. C-Appl., 29, 190–197, 1980.

- Kelley, N. D., Jonkman, B., and Scott, G.: Great Plains Turbulence Environment: Its Origins, Impact, and Simulation, Tech. rep., National Renewable Energy Laboratory (NREL), Golden, CO, available at: https://www.nrel.gov/docs/fy07osti/40176.pdf (last access: 3 September 2020), 2006.
- Kolmogorov, A. N.: Dissipation of energy in locally isotropic turbulence, Dokl. Akad. Nauk SSSR, 32, 16–18, 1941.
- Krishnamurthy, R., Calhoun, R., Billings, B., and Doyle, J.: Wind turbulence estimates in a valley by coherent Doppler lidar, Meteorol. Appl. 18, 361–371, 2011.
- Leahy, K., Hu, R. L., Konstantakopoulos, I. C., Spanos, C. J., and Agogino, A. M.: Diagnosing wind turbine faults using machine learning techniques applied to operational data, in: 2016 IEEE International Conference on Prognostics and Health Management (ICPHM), Ottawa, ON, Canada, 20–22 June 2016, https://doi.org/10.1109/ICPHM.2016.7542860, 2016.
- Leufen, L. H. and Schädler, G.: Calculating the turbulent fluxes in the atmospheric surface layer with neural networks, Geosci. Model Dev., 12, 2033-2047, https://doi.org/10.5194/gmd-12-2033-2019, 2019.
- Lundquist, J. K. and Bariteau, L.: Dissipation of Turbulence in the Wake of a Wind Turbine, Bound.-Lay. Meteorol., 154, 229–241, https://doi.org/10.1007/s10546-014-9978-3, 2015.
- Mellor, G. L. and Yamada, T.: A hierarchy of turbulence closure models for planetary boundary layers, J. Atmos. Sci., 31, 1791– 1806, 1974.
- Monin, A. S. and Obukhov, A. M.: Basic laws of turbulent mixing in the surface layer of the atmosphere, Contrib. Geophys. Inst. Acad. Sci. USSR, 151, 1954.
- Muñoz-Esparza, D., Sharman, R. D., and Lundquist, J. K.: Turbulence dissipation rate in the atmospheric boundary layer: Observations and WRF mesoscale modeling during the XPIA field campaign, Mon. Weather Rev., 146, 351–371, 2018.
- Nakanishi, M.: Improvement of the Mellor–Yamada turbulence closure model based on large-eddy simulation data, Bound.-Lay. Meteorol., 99, 349–378, 2001.
- Nakanishi, M. and Niino, H.: An improved Mellor–Yamada level-3 model: Its numerical stability and application to a regional prediction of advection fog, Bound.-Lay. Meteorol., 119, 397–407, 2006.
- Oncley, S. P., Friehe, C. A., Larue, J. C., Businger, J. A., Itsweire, E. C., and Chang, S. S.: Surface-layer fluxes, profiles, and turbulence measurements over uniform terrain under near-neutral conditions, J. Atmos. Sci., 53, 1029–1044, 1996.
- Optis, M. and Perr-Sauer, J.: The importance of atmospheric turbulence and stability in machine-learning models of wind farm power production, Renew. Sustain. Energ. Rev., 112, 27–41, 2019.
- Paquin, J. E. and Pond, S.: The determination of the Kolmogoroff constants for velocity, temperature and humidity fluctuations from second-and third-order structure functions, J. Fluid Mech., 50, 257–269, 1971.
- Piper, M. D.: The effects of a frontal passage on fine-scale nocturnal boundary layer turbulence, PhD thesis, University of Boulder, 2001.
- Sharma, N., Sharma, P., Irwin, D., and Shenoy, P.: Predicting solar generation from weather forecasts using machine learning, in: 2011 IEEE International Conference on Smart Grid Communications (SmartGrid-

- Comm), Brussels, Belgium, 17–20 October 2011, 528–533, https://doi.org/10.1109/SmartGridComm.2011.6102379, 2011.
- Shaw, W. J. and LeMone, M. A.: Turbulence dissipation rate measured by 915 MHz wind profiling radars compared with in-situ tower and aircraft data, in: 12th Symposium on Meteorological Observations and Instrumentation, available at: https://ams.confex.com/ams/pdfpapers/58647.pdf (last access: 3 September 2020), 2003.
- Skamarock, W. C., Klemp, J. B., Dudhia, J., Gill, D. O., Barker, D. M., Wang, W., and Powers, J. G.: A description of the advanced research WRF version 2, Tech. rep., National Center For Atmospheric Research, Boulder, CO, Mesoscale and Microscale Meteorology Div, 2005.
- Smalikho, I. N.: On measurement of the dissipation rate of the turbulent energy with a cw Doppler lidar, Atmos. Ocean. Opt., 8, 788–793, 1995.
- Sreenivasan, K. R.: On the universality of the Kolmogorov constant, Phys. Fluids, 7, 2778–2784, 1995.
- Thobois, L. P., Krishnamurthy, R., Loaec, S., Cariou, J. P., Dolfi-Bouteyre, A., and Valla, M.: Wind and EDR measurements with scanning Doppler LIDARs for preparing future weather dependent separation concepts, in: 7th AIAA Atmospheric and Space Environments Conference, AIAA 2015-3317, https://doi.org/10.2514/6.2015-3317, 2015.

- UCAR/NCAR: NCAR/EOL Quality Controlled High-rate ISFS surface flux data, geographic coordinate, tilt corrected, Version 1.1, Dataset, https://doi.org/10.26023/8x1n-tct4-p50x, 2019.
- USGS: USGS EROS Archive Digital Elevation Shuttle Radar Topography Mission (SRTM) 1 Arc-Second Global, https://doi.org/10.5066/F7PR7TFT, 2020.
- Wilczak, J. M., Oncley, S. P., and Stage, S. A.: Sonic anemometer tilt correction algorithms, Bound.-Lay. Meteorol., 99, 127–150, 2001.
- Wildmann, N., Bodini, N., Lundquist, J. K., Bariteau, L., and Wagner, J.: Estimation of turbulence dissipation rate from Doppler wind lidars and in situ instrumentation for the Perdigão 2017 campaign, Atmos. Meas. Tech., 12, 6401–6423, https://doi.org/10.5194/amt-12-6401-2019, 2019.
- Xingjian, S., Chen, Z., Wang, H., Yeung, D.-Y., Wong, W.-K., and Woo, W.-C.: Convolutional LSTM network: A machine learning approach for precipitation nowcasting, in: Advances in Neural Information Processing Systems, 802–810, 2015.
- Yang, B., Qian, Y., Berg, L. K., Ma, P.-L., Wharton, S., Bu-laevskaya, V., Yan, H., Hou, Z., and Shaw, W. J.: Sensitivity of turbine-height wind speeds to parameters in planetary boundary-layer and surface-layer schemes in the weather research and forecasting model, Bound.-Lay. Meteorol., 162, 117–142, 2017.